

Formulation

The goal is to express quantum mechanics in a new way with geometric algebra in order to get an understanding for space and spin.

The wavefunction can be expressed as

$$\Psi = \sum_i \frac{1}{\sqrt{2}} (e_i + J f_i) (\Re \psi_i + J \Im \psi_i)$$

where e_i, f_i are orthonormal vectors which square to $+1$, $J = e_0 f_0$ is an independent bivector, and ψ_i are the complex valued components of the wave vector. This makes the wave function an odd-grade multivector.

For a spin-1/2 particle space coordinates will come out as

$$\begin{aligned} X &= \frac{1}{2} (e_1 f_2 + e_2 f_1) \\ Y &= \frac{1}{2} (e_1 e_2 + f_1 f_2) \\ Z &= \frac{1}{2} (e_2 f_2 - e_1 f_1) \end{aligned}$$

which means that space rotations can be derived in this basis.

Observation in quantum mechanics

The probability $P = |\langle \psi_1 | \psi_2 \rangle|^2$ to measure a state ψ_1 in a state ψ_2 can be calculated from

$$\begin{aligned} P(1 \rightarrow 2) &= \langle (1 - \Psi_1 \Psi_1^\dagger) (1 - \Psi_2 \Psi_2^\dagger) \rangle \\ &= \langle \Psi_1 \Psi_1^\dagger \Psi_2 \Psi_2^\dagger \rangle - 1 \end{aligned}$$

which is an inner product between two state multivectors of the form $\Psi \Psi^\dagger$.

Unitary transformation

A unitary transformation of the wavefunction can be represented as a rotor in geometric algebra. Note that the same rotor can be applied to Ψ or $\Psi \Psi^\dagger$.

Single spin

The wavefunction for a single spin-up in a direction given by Euler angles θ, ϕ is usually written as

$$\psi = \begin{pmatrix} \cos \frac{\theta}{2} \\ \sin \frac{\theta}{2} e^{i\phi} \end{pmatrix}$$

up to an arbitrary phase.

Written in geometric algebra this is

$$\begin{aligned} J\Psi\Psi^\dagger &= J + e_1 f_1 \cos^2 \frac{\theta}{2} + e_2 f_2 \sin^2 \frac{\theta}{2} \\ &\quad + (e_1 f_2 + e_2 f_1) \cos \frac{\theta}{2} \sin \frac{\theta}{2} \cos \phi \\ &\quad + (e_1 e_2 + f_1 f_2) \cos \frac{\theta}{2} \sin \frac{\theta}{2} \sin \phi \end{aligned}$$

(see derivation in appendix).

$$\begin{aligned} J\Psi\Psi^\dagger &= J + e_1 f_1 \frac{1 - \cos \theta}{2} + e_2 f_2 \frac{1 + \cos \theta}{2} \\ &\quad + (e_1 f_2 + e_2 f_1) \frac{1}{2} \sin \theta \cos \phi \\ &\quad + (e_1 e_2 + f_1 f_2) \frac{1}{2} \sin \theta \sin \phi \\ &= J + \frac{1}{2} (e_1 f_1 + e_2 f_2) \\ &\quad + \frac{1}{2} (e_2 f_2 - e_1 f_1) \cos \theta \\ &\quad + \frac{1}{2} (e_1 f_2 + e_2 f_1) \sin \theta \cos \phi \\ &\quad + \frac{1}{2} (e_1 e_2 + f_1 f_2) \sin \theta \sin \phi \end{aligned}$$

Remembering that we have Euler angles, we can identify the multivectors for space coordinates from this expression for a single spin

$$\begin{aligned} X &= \frac{1}{2} (e_1 f_2 + e_2 f_1) \\ Y &= \frac{1}{2} (e_1 e_2 + f_1 f_2) \\ Z &= \frac{1}{2} (e_2 f_2 - e_1 f_1) \end{aligned}$$

They are orthogonal (e.g. $\langle XY \rangle = 0$) and obey

$\begin{aligned} XY &= Z & YZ &= X & ZX &= Y \end{aligned}$ Note that these multivectors

$$\Psi = \sum_i \frac{1}{\sqrt{2}} (e_i + J f_i) (\Re \psi_i + J \Im \psi_i)$$

$$\text{With } z_i = \Re \psi_i + J \Im \psi_i$$

$\begin{aligned}$

$$\begin{aligned} \Psi \Psi^\dagger &= \frac{1}{2} \sum_{ij} (e_i + J f_i) z_{iz_j}^\dagger (e_j - J f_j) \\ &= \sum_i (1 - e_{if_i J}) z_{iz_i}^\dagger \\ &\quad + \frac{1}{2} \sum_{i < j} \left((e_i + J f_i)(e_j - J f_j) z_{iz_j}^\dagger + (e_j + J f_j)(e_i - J f_i) z_{iz_i}^\dagger \right) \\ &= \sum_i (1 - e_{if_i J}) z_{iz_i}^\dagger \\ &\quad + \frac{1}{2} \sum_{i < j} \left((e_i e_j + f_{if_j} - (e_{if_j} + e_{jf_i}) J) z_{iz_j}^\dagger \right. \\ &\quad \left. + (e_j e_i + f_{jf_i} - (e_{jf_i} + e_{if_j}) J) z_{iz_i}^\dagger \right) \end{aligned}$$

With the real and imaginary parts

$\begin{aligned}$

$$\begin{aligned} R_{ij} &= \frac{1}{2} (z_{iz_j}^\dagger + z_{jz_i}^\dagger) \\ J I_{ij} &= \frac{1}{2} (z_{iz_j}^\dagger - z_{jz_i}^\dagger) \\ z_{iz_j}^\dagger &= R_{ij} + J I_{ij} \\ z_{jz_i}^\dagger &= R_{ij} - J I_{ij} \end{aligned}$$

where R_{ij}, I_{ij} are scalar this becomes

$\begin{aligned}$

$$\begin{aligned} \Psi \Psi^\dagger &= \sum_i (1 - e_{if_i J}) z_{iz_i}^\dagger \\ &\quad + \frac{1}{2} \sum_{i < j} \left((e_i e_j + f_{if_j} - (e_{if_j} + e_{jf_i}) J) (R_{ij} + J I_{ij}) \right. \\ &\quad \left. + (-e_i e_j - f_{if_j} - (e_{if_j} + e_{jf_i}) J) (R_{ij} - J I_{ij}) \right) \\ &= \sum_i z_{iz_i}^\dagger - \sum_i e_{if_i J} z_{iz_i}^\dagger \\ &\quad + \sum_{i < j} (e_{if_j} + e_{jf_i}) J R_{ij} \\ &\quad + \sum_{i < j} (e_{ie_j} + f_{if_j}) J I_{ij} \end{aligned}$$

For normalized wave vectors

$$\langle \Psi | \Psi \rangle = 1$$

Therefore

$\begin{aligned}$

$$J \Psi \Psi^\dagger = J + \sum_i e_{if_i} z_{iz_i}^\dagger + \sum_{i < j} (e_{if_j} + e_{jf_i}) R_{ij} -$$

$$\sum_{i < j} (e_{ie_j} + f_{if_j}), l_{ij}$$

$$\end{align}$$

is a grade -2 multivector. The probability in quantum mechanics can be calculated from the

$$(\psi_i \psi_i^*, \dots, \sqrt{2} \operatorname{Re}(\psi_i \psi_j^*), \dots, \sqrt{2} \operatorname{Im}(\psi_i \psi_j^*), \dots)$$

(shown elsewhere). Therefore in our case and for normalized wave vectors the probability

$$P(1 \rightarrow 2) = \langle (1 - \Psi_1 \Psi_1^\dagger) (1 - \Psi_2 \Psi_2^\dagger) \rangle = \langle \Psi_1 \Psi_1^\dagger \Psi_2 \Psi_2^\dagger \rangle - 1$$

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