

## Formulation

The goal is to express quantum mechanics in a new way with geometric algebra in order to get an understanding for space and spin.

The wavefunction can be expressed as

$$\Psi = \sum_{i=1}^n \frac{1}{\sqrt{2}} (e_i + J f_i) (\Re \psi_i + J \Im \psi_i)$$

where  $e_i, f_i$  are orthonormal vectors which square to  $+1$ ,

$$J = e_0 f_0$$

is an independent bivector, and  $\psi_i$  are the complex valued components of the wave vector. In total there are  $2n + 2$  orthonormal vectors for an  $n$ -dimensional original wave vector.

Alternatively, the bivector

$$\Psi f_0 = \frac{1}{\sqrt{2}} \sum_{i=1}^n ((e_i f_0 + f_i e_0) \Re \psi_i + (e_i e_0 - f_i f_0) \Im \psi_i)$$

is a more natural choice, but for the derivation I will use the first version. The difference to other treatments of spin with geometric algebra is that here the basis is of the form  $ab + cd$  with 2 additional basis vectors, instead of  $\sigma_{12} = e_1 e_2$ .

There is a possibility that the bivectors  $e_i f_0, f_i e_i, e_i e_0, -f_i f_0$  are related to the Dirac matrices and they come into action when they do not share the same coefficient like here.

One advantage is that the full probability calculation becomes

$$P(1 \rightarrow 2) = \langle \Psi_1 \Psi_1^\dagger \Psi_2 \Psi_2^\dagger \rangle - 1$$

which can take the same rotor for  $\Psi \Psi^\dagger$  and  $\Psi$ .

For a spin-1/2 particle the observable  $\Psi \Psi^\dagger$  and the space coordinates will come out as

$$J\Psi\Psi^\dagger = J + N + Xx + Yy + Zz$$

$$X = \frac{1}{2}(e_1 f_2 + e_2 f_1)$$

$$Y = \frac{1}{2}(e_1 e_2 + f_1 f_2)$$

$$Z = \frac{1}{2}(e_2 f_2 - e_1 f_1)$$

$$N = \frac{1}{2}(e_1 f_1 + e_2 f_2)$$

where coordinates do not depend on  $e_0, f_0$ .  $N$  belongs to  $Z$  somehow. Space rotations can be derived in this basis and the same rotor can be applied to  $\Psi$  too. They anti-commute and obey

$$XY = Z \quad YZ = X \quad ZX = Y$$

which has been derived from the spin wave vector alone.

## Observation in quantum mechanics

The probability  $P = |\langle\psi_1|\psi_2\rangle|^2$  to measure a state  $\psi_1$  in a state  $\psi_2$  can be calculated from (see appendix)

$$\begin{aligned} P(1 \rightarrow 2) &= \langle(1 - \Psi_1\Psi_1^\dagger)(1 - \Psi_2\Psi_2^\dagger)\rangle \\ &= \langle\Psi_1\Psi_1^\dagger\Psi_2\Psi_2^\dagger\rangle - 1 \\ &= \Omega_1 \cdot \Omega_2 - 1 \end{aligned}$$

with the state vector

$$\Omega = J\Psi\Psi^\dagger$$

This is an inner product between two state multivectors. The observable state vector can be expanded into a state bivector

$$\begin{aligned} \Omega &= J\Psi\Psi^\dagger \\ &= J + \sum_i e_i f_i |\psi_i|^2 + \sum_{i<j} (e_i f_j + e_j f_i) \Re(\psi_i \psi_j^*) - \sum_{i<j} (e_i e_j + f_i f_j) \Im(\psi_i \psi_j^*) \end{aligned}$$

(see appendix) for an easier translation from the usual representation with a complex wave vector  $\psi_i$ . Note that  $J$  appears only once at the front.  $\Psi\Psi^\dagger$  takes the role of the density matrix.

## Unitary transformation

A unitary transformation of the wavefunction can be represented as a rotor in geometric algebra. Note that the same rotor can be applied to  $\Psi$  or  $\Psi\Psi^\dagger$ .

## State vector of a single spin

The wavefunction for a single spin-up in a direction given by Euler angles  $\theta, \phi$  is usually written as

$$\psi = \begin{pmatrix} \cos \frac{\theta}{2} \\ \sin \frac{\theta}{2} e^{i\phi} \end{pmatrix}$$

up to an arbitrary phase.

The state vector written in geometric algebra then is

$$\begin{aligned} \Omega &= J\Psi\Psi^\dagger \\ &= J + \sum_i e_i f_i |\psi_i|^2 + \sum_{i<j} (e_i f_j + e_j f_i) \Re(\psi_i \psi_j^*) - \sum_{i<j} (e_i e_j + f_i f_j) \Im(\psi_i \psi_j^*) \\ &= J + e_1 f_1 \cos^2 \frac{\theta}{2} + e_2 f_2 \sin^2 \frac{\theta}{2} \\ &\quad + (e_1 f_2 + e_2 f_1) \cos \frac{\theta}{2} \sin \frac{\theta}{2} \cos \phi \\ &\quad + (e_1 e_2 + f_1 f_2) \cos \frac{\theta}{2} \sin \frac{\theta}{2} \sin \phi \\ &= J + e_1 f_1 \frac{1 - \cos \theta}{2} + e_2 f_2 \frac{1 + \cos \theta}{2} \\ &\quad + (e_1 f_2 + e_2 f_1) \frac{1}{2} \sin \theta \cos \phi \\ &\quad + (e_1 e_2 + f_1 f_2) \frac{1}{2} \sin \theta \sin \phi \\ &= J + \frac{1}{2} (e_1 f_1 + e_2 f_2) \\ &\quad + \frac{1}{2} (e_2 f_2 - e_1 f_1) \cos \theta \\ &\quad + \frac{1}{2} (e_1 f_2 + e_2 f_1) \sin \theta \cos \phi \\ &\quad + \frac{1}{2} (e_1 e_2 + f_1 f_2) \sin \theta \sin \phi \end{aligned}$$

Remembering that we have Euler angles, we can identify the multivectors for space coordinates from this expression for a single spin

$$\begin{aligned}
J\Psi\Psi^\dagger &= J + N + Xx + Yy + Zz \\
X &= \frac{1}{2}(e_1f_2 + e_2f_1) \\
Y &= \frac{1}{2}(e_1e_2 + f_1f_2) \\
Z &= \frac{1}{2}(e_2f_2 - e_1f_1) \\
N &= \frac{1}{2}(e_1f_1 + e_2f_2)
\end{aligned}$$

Space rotations will be derived soon, but from  $XY = Z$  we can already see the rotation.

## Interpretation

A vague idea why quantum mechanics is this way, is because due to the rules of probability, a state should be a (multi)vector and probabilities be calculated from an inner product. The probabilities rules are that they should sum to 1 and redoing a measurement yields the same results.

This state is constantly being rotated looking like  $\Omega = \cdots R_3 R_2 R_1 \Omega_0 R_1^\dagger R_2^\dagger R_3^\dagger \cdots$  and for some reason we get, that actually  $\Omega$  should split into  $\Omega = \Psi\Psi^\dagger$  - as if all particles start with the same state  $\Omega_0$ .

$\Psi$  has only roughly the square root number of element from the observable  $\Omega$ .

## Other ideas

Looking at

$$\Psi f_0 = \frac{1}{\sqrt{2}} \sum_{i=1}^n ((e_i f_0 + f_i e_0) \Re \psi_i + (e_i e_0 - f_i f_0) \Im \psi_i)$$

one may also consider what happens if the terms  $e_i f_0, f_i e_0, e_i e_0, f_i f_0$  have independent coefficients. Maybe this is related to Dirac matrices like

$$\Psi f_0 = \frac{1}{\sqrt{2}} \sum_{i=1}^n \left( e_i f_0 \alpha_i^{(1)} + f_i e_0 \alpha_i^{(2)} + e_i e_0 \alpha_i^{(3)} - f_i f_0 \alpha_i^{(4)} \right)$$

with real  $\alpha$  (and  $n = 2$  for spin-1/2), but I have not checked that yet. Basically, the system is that all basis vectors are combined with  $e_0$  and  $f_0$ .

## Multiple spins

For multiple spins one could take the wavefunction of dimension  $n = 2^k$  and apply the above procedure. However, it may also be insightful to generate a new wavefunction from multiplying two single spin wave functions

$$\Psi = \sum_{\substack{i=1,2 \\ j=3,4}} \frac{1}{\sqrt{2}} (e_i + J f_i) \left( \Re \psi_i^{(1)} + J \Im \psi_i^{(1)} \right) \frac{1}{\sqrt{2}} (e_j + J f_j) \left( \Re \psi_j^{(2)} + J \Im \psi_j^{(2)} \right)$$

This is for independent spins. For entangled spins, the cross-terms will have separate coefficients.

## Appendix

### Derivation

$$\Psi = \sum_i \frac{1}{\sqrt{2}} (e_i + J f_i) (\Re \psi_i + J \Im \psi_i)$$

With  $z_i = \Re \psi_i + J \Im \psi_i$

$$\begin{aligned} \Psi \Psi^\dagger &= \frac{1}{2} \sum_{ij} (e_i + J f_i) z_i z_j^\dagger (e_j - J f_j) \\ &= \sum_i (1 - e_i f_i J) z_i z_i^\dagger \\ &\quad + \frac{1}{2} \sum_{i < j} \left( (e_i + J f_i) (e_j - J f_j) z_i z_j^\dagger + (e_j + J f_j) (e_i - J f_i) z_j z_i^\dagger \right) \\ &= \sum_i (1 - e_i f_i J) z_i z_i^\dagger \\ &\quad + \frac{1}{2} \sum_{i < j} (e_i e_j + f_i f_j - (e_i f_j + e_j f_i) J) z_i z_j^\dagger \\ &\quad + \frac{1}{2} \sum_{i < j} (e_j e_i + f_j f_i - (e_j f_i + e_i f_j) J) z_j z_i^\dagger \end{aligned}$$

With the real and imaginary parts

$$\begin{aligned} R_{ij} &= \frac{1}{2} (z_i z_j^\dagger + z_j z_i^\dagger) \\ J I_{ij} &= \frac{1}{2} (z_i z_j^\dagger - z_j z_i^\dagger) \\ z_i z_j^\dagger &= R_{ij} + J I_{ij} \\ z_j z_i^\dagger &= R_{ij} - J I_{ij} \end{aligned}$$

where  $R_{ij}, I_{ij}$  are scalars this becomes

$$\begin{aligned}
\Psi\Psi^\dagger &= \sum_i (1 - e_i f_i J) z_i z_i^\dagger \\
&\quad + \frac{1}{2} \sum_{i < j} (e_i e_j + f_i f_j - (e_i f_j + e_j f_i) J) (R_{ij} + J I_{ij}) \\
&\quad + \frac{1}{2} \sum_{i < j} (-e_i e_j - f_i f_j - (e_i f_j + e_j f_i) J) (R_{ij} - J I_{ij}) \\
&= \sum_i z_i z_i^\dagger - \sum_i e_i f_i J z_i z_i^\dagger \\
&\quad - \sum_{i < j} (e_i f_j + e_j f_i) J R_{ij} \\
&\quad + \sum_{i < j} (e_i e_j + f_i f_j) J I_{ij}
\end{aligned}$$

For normalized wave vectors

$$\sum_i z_i z_i^\dagger = 1$$

Therefore

$$J\Psi\Psi^\dagger = J + \sum_i e_i f_i z_i z_i^\dagger + \sum_{i < j} (e_i f_j + e_j f_i) R_{ij} - \sum_{i < j} (e_i e_j + f_i f_j) I_{ij}$$

is a grade-2 multivector.

The probability in quantum mechanics can be calculated from the dot product of two real vectors where states have the components

$$(\psi_i \psi_i^*, \dots, \sqrt{2} \Re(\psi_i \psi_j^*), \dots, \sqrt{2} \Im(\psi_i \psi_j^*), \dots)$$

due to

$$\begin{aligned}
P(\psi \rightarrow \phi) &= |\langle \psi | \phi \rangle|^2 \\
&= \sum_i \psi_i \phi_i^* \sum_j \phi_j \psi_j^* \\
&= \sum_i \psi_i \phi_i^* \phi_i \psi_i^* + \sum_{i < j} (\psi_i \phi_i^* \phi_j \psi_j^* + \psi_j \phi_j^* \phi_i \psi_i^*) \\
&= \sum_i \psi_i \psi_i^* \phi_i \phi_i^* + \sum_{i < j} 2 \Re(\psi_i \psi_j^* \phi_i^* \phi_j) \\
&= \sum_i \psi_i \psi_i^* \phi_i \phi_i^* + \sum_{i < j} 2 (\Re(\psi_i \psi_j^*) \Re(\phi_i \phi_j^*) + \Im(\psi_i \psi_j^*) \Im(\phi_i \phi_j^*))
\end{aligned}$$

being a dot product of vectors with components

$$(\psi_i \psi_i^*, \sqrt{2}\Re(\psi_i \psi_j^*), \sqrt{2}\Im(\psi_i \psi_j^*)), (i < j).$$

Therefore in our case and for normalized wave vectors the probability can also be calculated from

$$P(1 \rightarrow 2) = \langle (1 - \Psi_1 \Psi_1^\dagger)(1 - \Psi_2 \Psi_2^\dagger) \rangle = \langle \Psi_1 \Psi_1^\dagger \Psi_2 \Psi_2^\dagger \rangle - 1$$

## Relations between single spin vectors

$$\begin{array}{lll} XY = Z & YZ = X & ZX = Y \\ XN = 0 & YN = 0 & ZN = 0 \end{array}$$

$$\begin{aligned} \Pi &= e_1 f_1 e_2 f_2 \\ XX &= -\frac{1}{2}(1 + \Pi) \\ YY &= -\frac{1}{2}(1 + \Pi) \\ ZZ &= -\frac{1}{2}(1 + \Pi) \\ NN &= -\frac{1}{2}(1 - \Pi) \end{aligned}$$

## Wave vector basis

$$\begin{aligned} \Psi &= \sum_{i=1}^n \frac{1}{\sqrt{2}} (e_i + J f_i) (\Re \psi_i + J \Im \psi_i) \\ &= \sum_{i=1}^n \Lambda_i (\Re \psi_i + J \Im \psi_i) \end{aligned}$$

with

$$\Lambda_i = \frac{1}{\sqrt{2}} (e_i + J f_i)$$

$$\begin{aligned} X_{ij} &= \frac{1}{2} (e_i f_j + e_j f_i) & \tilde{X}_{ij} &= \frac{1}{2} (e_i f_j - e_j f_i) \\ Y_{ij} &= \frac{1}{2} (e_i e_j + f_i f_j) & \tilde{Y}_{ij} &= \frac{1}{2} (e_i e_j - f_i f_j) \\ Z_{ij} &= \frac{1}{2} (e_j f_j - e_i f_i) & \tilde{Z}_{ij} &= \frac{1}{2} (e_j f_j + e_i f_i) \\ T_{ij} &= \frac{1}{2} (e_i f_i + e_j f_j) & \tilde{T}_{ij} &= \frac{1}{2} (e_i f_i - e_j f_j) \\ e_i f_i &= T_{ij} - Z_{ij} = \tilde{T}_{ij} + \tilde{Z}_{ij} \\ e_j f_j &= T_{ij} + Z_{ij} = \tilde{T}_{ij} - \tilde{Z}_{ij} \end{aligned}$$

Neither of these  $ab+cd$  terms has an inverse.

$$\Lambda_i^2=0\qquad \Lambda_i\Lambda_i^\dagger=1-J(T_{ij}-Z_{ij})$$

$$\Lambda_i\Lambda_j=$$

$$e_1f_1=T-Z$$

$$e_2f_2=T+Z$$

$$TT+XX=-1$$