Robotics 1

January 13, 2025 [students with midterm]

Exercise 1

For the PRR planar robot in Fig. 1, consider first the task vector r made by the position $p \in \mathbb{R}^2$ of the end-effector and by its angle $\phi \in \mathbb{R}$ with respect to the x_w axis. Compute the corresponding Jacobian $J_r(q)$ and find all its singularities. With the robot in a generic singular configuration q_s :

- a) provide the expression of all joint velocities \dot{q} that produce no task velocity \dot{r} ;
- b) determine all task velocities \dot{r} that cannot be instantaneously realized.

Next, consider only the two-dimensional task vector r = p for the same robot and find all singularities of the corresponding Jacobian $J_p(q)$. When the robot is in a configuration q_s with all strictly positive joint values and such that the matrix $J_p(q_s)$ loses rank:

- c) provide the expression of all forces $f \in \mathbb{R}^2$ applied to the end-effector that need no joint force/torque $\tau \in \mathbb{R}^3$ to be balanced;
- d) determine the τ that statically balances a force $f = \begin{pmatrix} 3 & 1 \end{pmatrix}^T$ [N] applied to the end-effector.

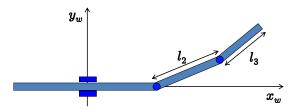


Figure 1: A PRR planar robot with the last two links of length l_2 and $l_3 \neq l_2$.

Exercise 2

A cylindrical robot has the direct kinematics of its end-effector position expressed by

$$m{p}(m{q}) = \left(egin{array}{c} q_3 \cos q_2 \ q_3 \sin q_2 \ q_1 \end{array}
ight).$$

When the desired position is $p_d = \begin{pmatrix} 1 & -1 & 3 \end{pmatrix}^T$ [m], provide the first few iterations of a Newton algorithm for the numerical solution of the inverse kinematics problem in the following two cases:

- a) starting from the initial guess $\mathbf{q}_a^{[0]} = \begin{pmatrix} -2 & 0.7\pi & \sqrt{2} \end{pmatrix}^T$ [m,rad,m]; b) starting from the initial guess $\mathbf{q}_b^{[0]} = \begin{pmatrix} 2 & \pi/4 & \sqrt{2} \end{pmatrix}^T$ [m,rad,m].

In case of convergence, the algorithm should stop as soon as $\|e^{[k]}\| = \|p_d - p(q^{[k]})\| \le \epsilon = 0.1 \text{ mm}$.

Exercise 3

A 2R planar robot with link lengths $l_1 = 1.2$, $l_2 = 0.8$ [m] is at rest at t = 0 in the configuration $q_0 = 0$ (stretched along the x_0 axis). A pointwise target moves at constant speed v = 1.5 m/s on a straight line with an angle $\delta = 15^{\circ}$ from the \boldsymbol{x}_0 axis, being in $\boldsymbol{p}_0 = \begin{pmatrix} -2 & 1 \end{pmatrix}^T$ [m] at t = 0 and entering after in the robot workspace. Solve the following rendez-vous problem:

- a) define a trajectory that will bring the robot end-effector on the target when the latter crosses the y_0 axis; the end-effector should have then the same velocity $v_t \in \mathbb{R}^2$ of the target;
- b) provide the rendez-vous time $t_{rv} > 0$ and the expression of the command $\dot{q}(t) \in \mathbb{R}^2$, $t \in [0, t_{rv}]$. How would you modify the velocity command $\dot{q}(t)$ as a function of q(t) so as to reach the target at the rendez-vous position if the robot starts from a configuration close but different from q_0 ?

[180 minutes (3 hours); open books]

Solution

January 13, 2025 [students with midterm]

Exercise 1

The kinematics of the three-dimensional task vector $\mathbf{r} = (\mathbf{p}, \phi)$ for the PRR robot of Fig. 1 is

$$m{r} = \left(egin{array}{c} q_1 + l_2 \cos q_2 + l_3 \cos \left(q_2 + q_3
ight) \ l_2 \sin q_2 + l_3 \sin \left(q_2 + q_3
ight) \ q_2 + q_3 \end{array}
ight) = m{f}(m{q}),$$

with the corresponding (analytic) Jacobian

$$\boldsymbol{J_r(q)} = \frac{\partial \boldsymbol{f}}{\partial \boldsymbol{q}} = \begin{pmatrix} 1 & -l_2 \sin q_2 - l_3 \sin (q_2 + q_3) & -l_3 \sin (q_2 + q_3) \\ 0 & l_2 \cos q_2 + l_3 \cos (q_2 + q_3) & l_3 \cos (q_2 + q_3) \\ 0 & 1 & 1 \end{pmatrix}.$$

Since det $J_r(q) = l_2 \cos q_2$, singular configurations are obtained when $q_2 = \pm \pi/2$. For $q_2 = \pi/2$, we have

$$\bar{\boldsymbol{J}} = \boldsymbol{J_r}|_{q_2 = \pi/2} = \begin{pmatrix} 1 & -l_2 - l_3 \cos q_3 & -l_3 \cos q_3 \\ 0 & -l_3 \sin q_3 & -l_3 \sin q_3 \\ 0 & 1 & 1 \end{pmatrix}$$

with the third column being the sum of the first column multiplied by l_2 and the second column. Therefore, all joint velocities that produce $\dot{r} = 0$ are of the form

$$\dot{\boldsymbol{q}} = \nu \begin{pmatrix} l_2 \\ 1 \\ -1 \end{pmatrix} \in \mathcal{N}\{\bar{\boldsymbol{J}}\}, \quad \text{for any } \nu \in \mathbb{R}.$$

On the other hand, task velocities \dot{r} that cannot be instantaneously realized lie outside the range of \bar{J} . Taking then as basis for $\mathcal{R}\{\bar{J}\}$, e.g., its first two columns, one should find vectors $\dot{r} \in \mathbb{R}^3$ that are linear independent from this basis, i.e., such that

$$\det \begin{pmatrix} 1 & -l_2 - l_3 \cos q_3 & \dot{r}_1 \\ 0 & -l_3 \sin q_3 & \dot{r}_2 \\ 0 & 1 & \dot{r}_3 \end{pmatrix} \neq 0 \qquad \Rightarrow \qquad \dot{r}_3 \, l_3 \sin q_3 + \dot{r}_2 \neq 0 \quad (\text{and no condition on } \dot{r}_1).$$

There are many instances of vectors that satisfy the above inequality, e.g., $\dot{\boldsymbol{r}}=(0,\lambda,0)$ for any $\lambda\neq 0$, and thus such that $\dot{\boldsymbol{r}}\not\in\mathcal{R}\{\bar{\boldsymbol{J}}\}$. However, the most elegant way to supplement the range space of $\bar{\boldsymbol{J}}$ is to build the null space of $\bar{\boldsymbol{J}}^T$, since $\mathcal{R}\{\bar{\boldsymbol{J}}\}\oplus\mathcal{N}\{\bar{\boldsymbol{J}}^T\}=\mathbb{R}^3$. Being

$$\dot{\boldsymbol{r}}_0 = \mu \begin{pmatrix} 0 \\ 1 \\ l_3 \sin q_3 \end{pmatrix} \in \mathcal{N}\{\bar{\boldsymbol{J}}^T\}, \quad \text{for any } \mu \neq 0,$$

we can conclude that any task velocity \dot{r} that has a component along \dot{r}_0 cannot be realized. Conversely, any \dot{r} that has no component along \dot{r}_0 belongs to $\mathcal{R}\{\bar{J}\}$ and thus can be realized.

For the other type of singular configurations having $q_2 = -\pi/2$, it is

$$\bar{\bar{J}} = J_r|_{q_2 = -\pi/2} = \begin{pmatrix} 1 & l_2 + l_3 \cos q_3 & l_3 \cos q_3 \\ 0 & l_3 \sin q_3 & l_3 \sin q_3 \\ 0 & 1 & 1 \end{pmatrix}.$$

Following the same procedure as above, we obtain

$$\dot{\boldsymbol{q}} = \nu \begin{pmatrix} l_2 \\ -1 \\ 1 \end{pmatrix} \in \mathcal{N}\{\bar{\bar{\boldsymbol{J}}}\}, \quad \text{for any } \nu \in \mathbb{R}$$

and

$$\dot{\boldsymbol{r}}_0 = \mu \begin{pmatrix} 0 \\ -1 \\ l_3 \sin q_3 \end{pmatrix} \in \mathcal{N}\{\bar{\bar{\boldsymbol{J}}}^T\}, \qquad \text{for any } \mu \neq 0.$$

As for the two-dimensional task vector r = p, the corresponding Jacobian reduces to the 2×3 matrix

$$\boldsymbol{J_p(q)} = \frac{\partial \boldsymbol{f}}{\partial \boldsymbol{q}} = \begin{pmatrix} 1 & -l_2 \sin q_2 - l_3 \sin \left(q_2 + q_3\right) & -l_3 \sin \left(q_2 + q_3\right) \\ 0 & l_2 \cos q_2 + l_3 \cos \left(q_2 + q_3\right) & l_3 \cos \left(q_2 + q_3\right) \end{pmatrix},$$

obtained by simply eliminating the third row from $J_r(q)$. The singular configurations are found by zeroing the three minors obtained by deleting one column at a time from $J_p(q)$ (resulting in three 2×2 matrices $J_{-i}(q)$, i = 1, 2, 3). One has

$$\det \mathbf{J}_{-1} = l_2 l_3 s_3 = 0 \qquad \det \mathbf{J}_{-2} = l_3 c_{23} = 0 \qquad \det \mathbf{J}_{-3} = l_2 c_2 + l_3 c_{23} = 0,$$

with the shorthand notation used for trigonometric functions (e.g., $c_{23} = \cos(q_2 + q_3)$). From these, the singular configurations are characterized by

$$c_2 = 0 \cap s_3 = 0$$
 \Rightarrow $(q_2, q_3) = \{(\pi/2, 0), (\pi/2, \pi), (-\pi/2, 0), (-\pi/2, \pi)\}.$

As it is always the case, the first variable q_1 never matters in the definition of kinematic singularities.

Since it is asked to use the Jacobian only in a singular configuration q_s characterized by all strictly positive joint values, we choose $q_s = (q_1, \pi/2, \pi)$, with a generic value $q_1 > 0$. The matrix

$$oldsymbol{J}_s^T = oldsymbol{J}_{oldsymbol{p}}^T(oldsymbol{q}_s) = \left(egin{array}{ccc} 1 & 0 \ l_3 - l_2 & 0 \ l_3 & 0 \end{array}
ight)$$

has clearly rank 1. All Cartesian forces $f \in \mathbb{R}^2$ applied to the end-effector that need no joint force/torque $\tau \in \mathbb{R}^3$ to be balanced have the simple form

$$m{f} = arphi \left(egin{array}{c} 0 \\ 1 \end{array}
ight) \in \mathcal{N}\{m{J}_s^T\}, \qquad ext{for any } arphi \in \mathbb{R}.$$

Finally, to balance the force $\mathbf{f} = (3,1)$ [N] applied to the robot end-effector so that the robot remains in static conditions, the motors need to provide the vector of force (at the first joint) and torques (at the second and third joint)

$$oldsymbol{ au} = -oldsymbol{J}_s^T oldsymbol{f} = \left(egin{array}{c} -3 \ 3(l_2 - l_3) \ -3l_3 \end{array}
ight) ext{ [N,Nm,Nm]},$$

where the minus sign in the general formula stands for the need of having zero balance of generalized forces at all joints. The resulting vector $\boldsymbol{\tau}$ would be the same for any applied Cartesian force of the form $\boldsymbol{f} = (3, \varphi)$ [N], with arbitrary φ .

Exercise 2

The Jacobian associated to the positional direct kinematics of the cylindrical robot is

$$\boldsymbol{J}(\boldsymbol{q}) = \frac{\partial \boldsymbol{p}}{\partial \boldsymbol{q}} = \begin{pmatrix} 0 & -q_3 \sin q_2 & \cos q_2 \\ 0 & q_3 \cos q_2 & \sin q_2 \\ 1 & 0 & 0 \end{pmatrix}$$

This matrix is singular when det $J = -q_3 = 0$. It is simple to write a code that performs a number of times the Newton iteration

$$m{q}^{[k]} = m{q}^{[k-1]} + m{J}^{-1}(m{q}^{[k-1]}) \left(m{p}_d - m{p}(m{q}^{[k-1]})
ight)$$

for solving the inverse kinematics problem, starting from a given initial guess $q^{[0]}$.

For the problem data $\boldsymbol{p}_d=(1,-1,3)$ [m], Tab. 1 summarizes the obtained numerical results when starting from $\boldsymbol{q}_a^{[0]}=(-2,0.7\pi,\sqrt{2})$ [m,rad,m]). The algorithm stops at the third iteration, having reached the requested accuracy of $\epsilon=10^{-4}$ m on the norm of the Cartesian position error. The inverse kinematics solution found is $\boldsymbol{q}_{\mathrm{sol},1}=(3,2.3562,-1.4142)$ [m,rad,m].

iteration k	$oldsymbol{q}^{[k]}$	$oldsymbol{p}^{[k]}$	$\ oldsymbol{e}^{[k]}\ $	$\det oldsymbol{J}(oldsymbol{q}^{[k]})$
0	(-2, 2.1991, 1.4142)	(-0.8312, 1.1441, 2)	5.74027	-1.4142
1	(3, 2.0426, -1.3968)	(0.6349, 1.2441, 3)	0.43918	1.3968
2	(3, 2.3549, -1.3452)	(0.9500, -0.9524, 3)	0.06895	1.3452
3	(3, 2.3562, -1.4142)	(1.0000, -0.9999, 3)	0.00009	1.4142

Table 1: Steps of the Newton method for the inverse kinematics problem when starting from $q_a^{[0]}$.

When starting from $\boldsymbol{q}_b^{[0]}=\left(2,\pi/4,\sqrt{2}\right)$ [m,rad,m], where $\boldsymbol{p}^{[0]}=(1,1,2)$ [m] and $\|\boldsymbol{e}^{[0]}\|=2.2361$, the next configuration computed by Newton's method is

$$m{q}^{[1]} = \left(egin{array}{c} 3 \ -0.2146 \ 0 \end{array}
ight), \qquad ext{with } m{p}^{[1]} = \left(egin{array}{c} 0 \ 0 \ 3 \end{array}
ight), \; \left\|m{e}^{[1]}
ight\| = 1.4142, \; ext{ and } \det m{J}(m{q}^{[1]}) = 0.$$

Having reached a singular configuration, the algorithm should be stopped at this first iteration to prevent divergence (NaN in MATLAB).

Note that if one replaces the (numerical) pseudoinverse of the Jacobian to its inverse, i.e.,

$$\boldsymbol{q}^{[k]} = \boldsymbol{q}^{[k-1]} + \boldsymbol{J}^{\#}(\boldsymbol{q}^{[k-1]}) \, \Big(\boldsymbol{p}_d - \boldsymbol{p}(\boldsymbol{q}^{[k-1]}) \Big),$$

the modified algorithm will tolerate the singularity after the first iteration, and eventually converge in six iterations to the other inverse kinematics solution $q_{\text{sol},2} = (3, -0.7854, 1.4142)$ [m,rad,m], as shown in Tab. 2.

iteration k	$oldsymbol{q}^{[k]}$	$oldsymbol{p}^{[k]}$	$\ oldsymbol{e}^{[k]}\ $	$\det oldsymbol{J}(oldsymbol{q}^{[k]})$
0	(2, 0.7854, 1.4142)	(1, 1, 2)	2.23607	-1.4142
1	(3, -0.2146, 0)	(0,0,3)	1.41421	0
2	(3, -0.2146, 1.1900)	(1.1627, -0.2534, 3)	0.76410	-1.1900
3	(3, -0.8567, 1.1900)	(0.7794, -0.8993, 3)	0.24252	-1.1900
4	(3, -0.7720, 1.4106)	(1.0107, -0.9840, 3)	0.01921	-1.4106
5	(3, -0.7854, 1.4141)	(0.9999, -0.9999, 3)	0.00013	-1.4141
6	(3, -0.7854, 1.4142)	(1, -1, 3)	0.00000	-1.4142

Table 2: Steps of the modified Newton method (with pseudoinverse) for the inverse kinematics problem when starting from $q_b^{[0]}$.

Exercise 3

The trajectory of the target is

$$p_T(t) = p_0 + v t \begin{pmatrix} \cos \delta \\ \sin \delta \end{pmatrix}, \qquad t \in [0.t_{rv}],$$

and the time instant of rendez-vous (appointment), i.e., when the target crosses the \boldsymbol{y}_0 axis, is obtained from

$$p_{Tx}(t_{rv}) = p_{0x} + v t_{rv} \cos \delta = 0$$
 \Rightarrow $t_{rv} = \frac{-p_{0x}}{v \cos \delta} = 1.380 \text{ s.}$

Accordingly, the y-component of the Cartesian rendez-vous position is given by

$$p_{Ty} = p_{0y} + v t_{rv} \sin \delta = 1.5359$$
 \Rightarrow $p_{rv} = \begin{pmatrix} 0 \\ 1.5359 \end{pmatrix}$ [m].

For rendez-vous position p_{rv} , the two (regular) inverse kinematics solutions for the 2R robot are computed from the usual formulas as

$$\boldsymbol{q}_a = \left(\begin{array}{c} 1.0294 \\ 1.4250 \end{array}\right) \text{ [rad]} \quad \text{(right arm)} \qquad \boldsymbol{q}_b = \left(\begin{array}{c} 2.1122 \\ -1.4250 \end{array}\right) \text{ [rad]} \quad \text{(left arm)}.$$

The velocity of the target is constant and equal to

$$\dot{\boldsymbol{p}}_T = v \begin{pmatrix} \cos \delta \\ \sin \delta \end{pmatrix} = \begin{pmatrix} 1.4489 \\ 0.3882 \end{pmatrix} \text{ [m/s]}.$$

Therefore, the robot joint velocity for matching the target velocity at the rendez-vous is given by

$$\dot{\boldsymbol{q}}_a = \boldsymbol{J}^{-1}(\boldsymbol{q}_a)\dot{\boldsymbol{p}}_T = \begin{pmatrix} -0.7359\\ -0.6278 \end{pmatrix} [\text{rad/s}] \quad (\text{right arm})$$

$$\dot{\boldsymbol{q}}_b = \boldsymbol{J}^{-1}(\boldsymbol{q}_b)\dot{\boldsymbol{p}}_T = \begin{pmatrix} -1.1508\\ 0.6278 \end{pmatrix} [\mathrm{rad/s}[\quad (\mathrm{left\ arm}).$$

A suitable trajectory for solving the problem is computed in the joint space, ¹ as cubic polynomials interpolating the following alternative boundary conditions

$$\mathbf{q}(0) = \mathbf{0}$$
 $\dot{\mathbf{q}}(0) = \mathbf{0}$ \Rightarrow $\mathbf{q}(t_{rv}) = \mathbf{q}_a$ $\dot{\mathbf{q}}(t_{rv}) = \dot{\mathbf{q}}_a$ (right arm) $\mathbf{q}(t_{rv}) = \mathbf{q}_b$ $\dot{\mathbf{q}}(t_{rv}) = \dot{\mathbf{q}}_b$ (left arm).

The interpolating cubics are conveniently written in normalized time as

$$\mathbf{q}_d(\tau) = (3\mathbf{q}_{rv} - t_{rv}\dot{\mathbf{q}}_{rv})\tau^2 + (-2\mathbf{q}_{rv} + t_{rv}\dot{\mathbf{q}}_{rv})\tau^3, \qquad \tau = \frac{t}{t_{rv}} \in [0, 1],$$

which applies to both cases, just replacing a or b for rv in q_{rv} and \dot{q}_{rv} . The corresponding joint velocity command is

$$\dot{\boldsymbol{q}}_{d}(\tau) = \frac{2}{t_{rv}} \left(3\boldsymbol{q}_{rv} - t_{rv}\dot{\boldsymbol{q}}_{rv} \right) \tau + \frac{3}{t_{rv}} \left(-2\boldsymbol{q}_{rv} + t_{rv}\dot{\boldsymbol{q}}_{rv} \right) \tau^{2}, \qquad \tau = \frac{t}{t_{rv}} \in [0, 1].$$

Substituting the numerical values of the boundary conditions, we obtain the joint trajectories in normalized time

$$\mathbf{q}_{d.a}(\tau) = \begin{pmatrix} 4.1040 \\ 5.1415 \end{pmatrix} \tau^2 + \begin{pmatrix} -3.0746 \\ -3.7165 \end{pmatrix} \tau^3, \qquad \tau = \frac{t}{t_{rv}} \in [0, 1] \quad \text{(right arm)}$$

or

$$\boldsymbol{q}_{d,b}(\tau) = \left(\begin{array}{c} 7.9251 \\ -5.1415 \end{array} \right) \tau^2 + \\ \left(\begin{array}{c} -5.8129 \\ 3.7165 \end{array} \right) \tau^3, \qquad \tau = \frac{t}{t_{rv}} \in [0,1] \quad (\text{left arm}).$$

When using the original time t, the equivalent expressions are

$$\mathbf{q}_{d.a}(t) = \begin{pmatrix} 2.1539 \\ 2.6984 \end{pmatrix} t^2 + \begin{pmatrix} -1.1690 \\ -1.4130 \end{pmatrix} t^3, \quad t \in [0, t_{rv}] \quad \text{(right arm)}$$

or

$$\mathbf{q}_{d,b}(t) = \begin{pmatrix} 4.1593 \\ -2.6984 \end{pmatrix} t^2 + \begin{pmatrix} -2.2101 \\ 1.4130 \end{pmatrix} t^3, \qquad t = \in [0, t_{rv}] \quad \text{(left arm)}.$$

Accordingly, the joint velocity commands in the two cases are

$$\dot{q}_{d.a}(\tau) = \begin{pmatrix} 5.9462 \\ 7.4495 \end{pmatrix} \tau + \begin{pmatrix} -6.6821 \\ -8.0773 \end{pmatrix} \tau^2, \qquad \tau = \frac{t}{t_{rv}} \in [0, 1] \quad \text{(right arm)}$$

or

$$\dot{\boldsymbol{q}}_{d,b}(\tau) = \left(\begin{array}{c} 11.4826 \\ -7.4495 \end{array}\right) \tau + \\ \left(\begin{array}{c} -12.6334 \\ 8.0773 \end{array}\right) \tau^2, \qquad \tau = \frac{t}{t_{rv}} \in [0,1] \quad (\text{left arm}),$$

and in terms of the original time

$$\dot{q}_{d.a}(t) = \begin{pmatrix} 4.3077 \\ 5.3967 \end{pmatrix} t + \begin{pmatrix} -3.5069 \\ -4.2391 \end{pmatrix} t^2, \quad t \in [0, t_{rv}] \quad \text{(right arm)}$$

or

$$\dot{\boldsymbol{q}}_{d,b}(t) = \left(\begin{array}{c} 8.3185 \\ -5.3967 \end{array}\right) t + \\ \left(\begin{array}{c} -6.6303 \\ 4.2391 \end{array}\right) t^2, \qquad t \in [0,t_{rv}] \quad (\text{left arm}).$$

¹In this way, we avoid any problem related to singularities in the inversion of a Cartesian trajectory. The initial position of the robot end-effector is singular, but the joint configuration is already specified.

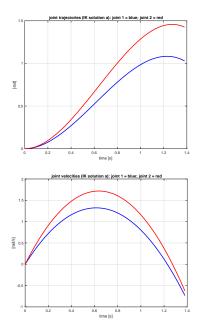


Figure 2: The planned trajectory for the 2R robot when choosing the **right** arm configuration at the rendez-vous: joint positions (top) and velocities (bottom).

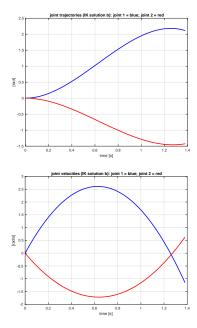


Figure 4: The planned trajectory for the 2R robot when choosing the **left** arm configuration at the rendez-vous: joint positions (top) and velocities (bottom).

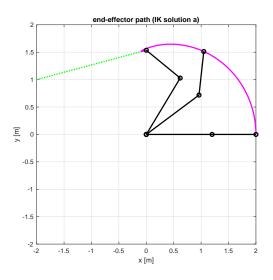


Figure 3: Cartesian path of the end-effector for the joint trajectory in Fig. 2.

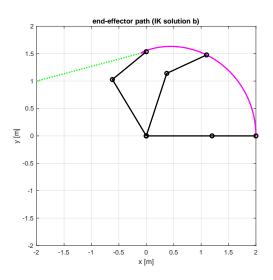


Figure 5: Cartesian path of the end-effector for the joint trajectory in Fig. 4.

Figure 2 shows the joint positions and velocity for case a (right arm configuration at rendez-vous), while the resulting path of the end-effector is given in Fig. 3. Therein, the path of the target is traced as a dotted green line. Similarly, Figs. 4–5 refer to case b (left arm configuration at rendez-vous). Note that in both cases the end-effector reaches approximately the rendez-vous position and then slightly moves back to the right to reach the exact desired position while matching also the velocity of the pointwise target.

If the robot starts from a perturbed initial configuration that is different from q_0 , the planned trajectory (and the corresponding joint velocity command $\dot{q}_d(t)$, for $t \in [0, t_{rv}]$) will fail to solve the rendez-vous task. Rather than replanning a trajectory from the new initial configuration (a different one for every new initialization!), it is more convenient to include a control action with a feedback term proportional to the joint position error, i.e.,

$$\dot{\boldsymbol{q}}_{c}(t) = \dot{\boldsymbol{q}}_{d}(t) + \boldsymbol{K}_{p} \left(\boldsymbol{q}_{d}(t) - \boldsymbol{q}(t) \right),$$

with a positive definite, 2×2 (diagonal) gain matrix \mathbf{K}_p . This additional action will recover the error with respect to the originally planned joint trajectory in an exponentially fast way, depending on the choice of the control gains. For a diagonal \mathbf{K}_p , by choosing sufficiently large gains $k_{p,1}$ and $k_{p,2}$, the error will vanish in practice well before the planned rendez-vous time instant t_{rv} , thus allowing successful completion also for the perturbed task. Indeed, the larger is the initial perturbation from \mathbf{q}_0 , the larger will have to be the control gains. If the required gains become too large (i.e., when the initial configuration is too far away from \mathbf{q}_0), risking a saturation of the applied command, then it would be appropriate to replan the rendez-vous trajectory $\mathbf{q}_d(t)$.

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