

# Celestial Mechanics - Homework n. 1

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March 3 2024

## 1 Problem 1

The canonical equation of an ellipse is

$$\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1$$

with the two foci at  $(c, 0)$ ,  $(-c, 0)$ . The eccentricity is defined as  $c = ea$ ,  $0 \leq e < 1$ , and  $a^2 = b^2 + c^2$  which gives  $b = a\sqrt{1 - e^2}$ .

Now I apply the transformation in polar coordinates centered on the focus at  $(c, 0)$ .

$$\begin{cases} x = c + r \cos \theta \\ y = r \sin \theta \end{cases}$$

Substituting  $x$  and  $y$  in the equation gives

$$\frac{(c + r \cos \theta)^2}{a^2} + \frac{(r \sin \theta)^2}{a^2(1 - e^2)} = 1$$

that after some manipulating becomes

$$\frac{1 - e^2 \cos^2 \theta}{1 - e^2} + 2rc \cos \theta - a^2(1 - e^2) = 0$$

This quadratic equation gives two solutions

$$r_{1,2} = \frac{-ae \cos \theta \pm a}{(1 - e \cos \theta)(1 + e \cos \theta)}(1 - e^2) = \pm \frac{a(1 - e^2)}{1 \mp e \cos \theta}$$

Now because  $r > 0$ , the correct solution is

$$r = \frac{a(1 - e^2)}{1 - e \cos \theta}$$

## 2 Problem 2

The Stokes theorem states that the integral of a vectorial field  $\mathbf{F}$  calculated over a closed path  $C$  is equal to the integral of the curl of the field over the surface  $\Sigma$  enclosed by  $C$

$$\int_{\Sigma} \nabla \times \mathbf{F} \cdot d\mathbf{\Sigma} = \oint_C \mathbf{F} \cdot d\mathbf{\Gamma}$$

In this case, I will use  $d\mathbf{\Sigma} = \hat{\mathbf{n}} dx dy$  with  $\hat{\mathbf{n}} = \mathbf{e}_z$ , so  $\nabla \times \mathbf{F} \cdot d\mathbf{\Sigma} = (\partial_x \mathbf{F}_y - \partial_y \mathbf{F}_x) \hat{\mathbf{n}}$ . To calculate the area inside  $C$  I choose  $\mathbf{F} = (-y, x, 0)$ , so

$$\int_{\Sigma} (\partial_x \mathbf{F}_y - \partial_y \mathbf{F}_x) dx dy = \int_{\Sigma} 2 dx dy = 2A$$

By doing the change of coordinates

$$\begin{cases} x = r \cos \theta \\ y = r \sin \theta \end{cases}$$

I get the differentiation

$$\begin{cases} dx = dr \cos \theta - r \sin \theta d\theta \\ dy = dr \sin \theta + r \cos \theta d\theta \end{cases}$$

and by substituting in the path integral  $\oint_C \mathbf{F}_x dx + \mathbf{F}_y dy = \oint_C -y dx + x dy$  I obtain

$$\oint_C -r dr \sin \theta \cos \theta + r^2 \sin^2 \theta d\theta + r dr \cos \theta \sin \theta + r^2 \cos^2 \theta d\theta = \int_0^{2\pi} r^2 d\theta$$

Now I choose as a coordinate basis  $(\mathbf{e}_r, \mathbf{e}_\theta, \mathbf{e}_z)$ . In this basis the radius  $\mathbf{r}$  has representation  $(r, 0, 0)$ ,  $d\mathbf{r} = (0, r d\theta, 0)$  and the cross product between the two gives indeed  $r^2 d\theta \mathbf{e}_z$ , which proves that

$$A = \frac{1}{2} \int_0^{2\pi} \mathbf{r} \times d\mathbf{r}$$

For the ellipse  $A = \frac{1}{2} \int_0^{2\pi} \frac{a^2(1-e^2)^2}{(1-e \cos \theta)^2} d\theta = \frac{a^2(1-e^2)^2}{2} \int_0^{2\pi} \frac{d\theta}{(1-e \cos \theta)^2}$ . The indefinite integral  $\int \frac{d\theta}{(1-e \cos \theta)^2}$  has been solved by Wolfram Alpha by substituting  $u = \tan \frac{\theta}{2}$  and the solution is

$$F(\theta) = \frac{2e \tan\left(\frac{\theta}{2}\right)}{(1-e^2) \left(e \tan^2\left(\frac{\theta}{2}\right) + \tan^2\left(\frac{\theta}{2}\right) - e + 1\right)} + \frac{2 \arctan\left(\frac{\sqrt{e+1} \tan\left(\frac{\theta}{2}\right)}{\sqrt{1-e^2}}\right)}{\sqrt{1-e^2} (1-e^2)} + C$$

This integral function is not continue in  $\theta = \pi$ , so the value of the definite integral is  $F(2\pi) - \lim_{\theta \rightarrow \pi^+} F(\theta) + \lim_{\theta \rightarrow \pi^-} F(\theta) - F(0)$ . In 0 and  $2\pi$   $\tan \frac{\theta}{2} = 0$  and so  $\arctan(\dots) = 0$ , while  $\lim_{\theta \rightarrow \pi^+} \tan \frac{\theta}{2} = -\infty$ : the first part goes to zero while the second one gives  $\frac{-\pi}{(1-e^2)^{1.5}}$ . Similarly,  $\lim_{\theta \rightarrow \pi^-} \tan \frac{\theta}{2} = +\infty$  and the result is  $\frac{\pi}{(1-e^2)^{1.5}}$ . Combining all of this together gives

$$A = \frac{a^2(1-e^2)^2}{2} \left( \frac{2\pi}{(1-e^2)^{1.5}} \right) = \pi a \cdot a \sqrt{1-e^2} = \pi ab$$

### 3 Problem 3

Given the position at a given time  $s(t) = \frac{g}{2\omega} (\sinh \omega t - \sin \omega t)$  and the value  $s(1) = 1$ , to find the value of  $\omega$  a root finding algorithm is used with function  $f(\omega) = \frac{\sinh \omega - \sin \omega}{\omega} - \frac{2}{g}$ . The calculation are done in different ways in the python file attached. The result, with an accuracy of  $10^{-8}$ , is  $\omega = 0.782286418$ .

### 4 Problem 4

Given the ODE

$$m \frac{d^2 x}{dt^2} + c \frac{dx}{dt} + kx = 0$$

with  $m, c, k > 0$ , this is solved by a linear combination of solutions of the type  $x = e^{zt}$ , where  $z$  is a complex number.

The conditions on  $z$  are found by substituting the solution inside the ODE, giving

$$mz^2 + cz + k = 0 \implies z_{1,2} = \frac{-c \pm \sqrt{c^2 - 4mk}}{2m}$$

Using  $\gamma = \frac{c}{2m}$  and  $\omega_0^2 = \frac{k}{m}$ , the solution becomes

$$z_{1,2} = -\gamma \pm \sqrt{\gamma^2 - \omega_0^2}$$

Now there are three cases:

1.  $\gamma > \omega_0$  ( $c^2 > 4mk$ ): the roots are two real numbers, and the general solution is

$$x = Ae^{(-\gamma+\omega)t} + Be^{(-\gamma-\omega)t}$$

with  $\omega = \sqrt{\gamma^2 - \omega_0^2}$  and  $A, B$  to be determined from the initial conditions  $x(t_0)$  and  $\dot{x}(t_0)$ . The system is in the overdamped regime and quickly tends to zero.

2.  $\gamma = \omega_0$  ( $c^2 = 4mk$ ): the solution has one root of multiplicity two, and the general solution is

$$x = (A + Bt)e^{-\gamma t}$$

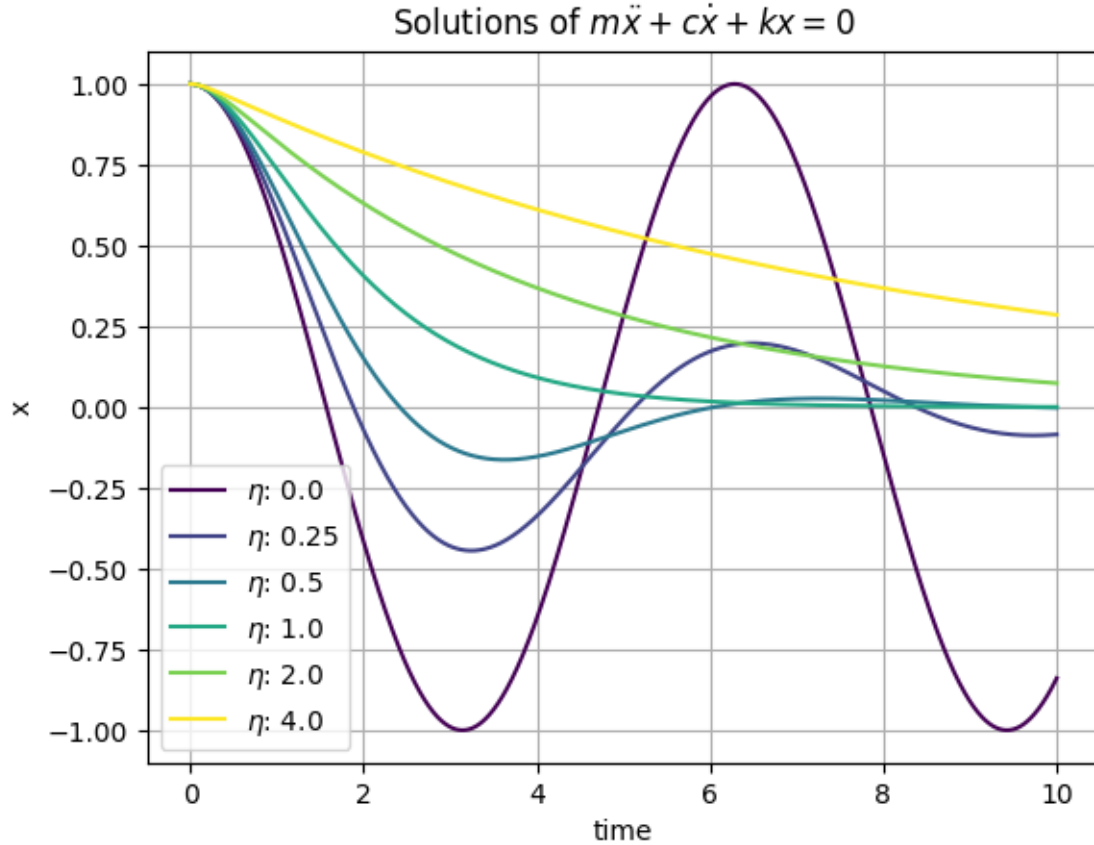
The system is said to be critically damped.

3.  $\gamma < \omega_0$  ( $c^2 < 4mk$ ): the roots are complex numbers  $z_{1,2} = -\gamma \pm i\omega$ , with  $\omega = \sqrt{\omega_0^2 - \gamma^2}$ . Grouping together the imaginary parts in a trigonometric function, the general solution is

$$x = Ae^{-\gamma t} \cos(\omega t + \phi)$$

The system is in the underdamped regime and oscillates around zero with increasingly smaller amplitude. If  $c = 0$  the system reduces to a lossless harmonic oscillator.

In the graph are plotted several solution with initial conditions  $x(0) = 1$ ,  $\dot{x}(0) = 0$  for different values of  $\eta = \frac{\gamma}{\omega_0} = \frac{c}{2\sqrt{mk}}$ : for  $\eta > 1$  the system is overdamped, for  $\eta = 1$  is critically damped, for  $\eta < 1$  is underdamped. In this plot I used  $m = k = 1$ , so  $\eta = \frac{c}{2}$ .



## 5 Problem 5

The formula describing the relationship between the angles and the sides of a spherical triangle is known as the **law of cosines**.

Because the sides of the triangle are circle segments of circles with radius 1, the length of the segment is equal to the angle between the vectors  $e_i$  (expressed in radians). So  $e_1 \cdot e_2 = \cos b$ ,  $e_2 \cdot e_3 = \cos a$ ,  $e_1 \cdot e_3 = \cos c$ . Then the vectors  $e_1 \times e_2$  and  $e_1 \times e_3$  are each normal to the plane described by the two vectors and have norms  $\sin b$  and  $\sin c$ . The angle  $\alpha$  is the angle between the two planes found before, so

$$\cos \alpha = \frac{(e_1 \times e_2) \cdot (e_1 \times e_3)}{\sin b \sin c}$$

and by applying the rule  $(a \times b) \cdot (c \times d) = (a \cdot c)(b \cdot d) - (a \cdot d)(b \cdot c)$ , this gives  $\cos \alpha \sin b \sin c = 1 \cdot \cos a - \cos b \cos c$  which can be arranged in the final form

$$\cos a = \cos b \cos c + \sin b \sin c \cos A$$