# Numerical Resolution of the Convection-Diffusion Equation The Smith-Hutton Case

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#### Abstract

This report presents a numerical study of the Smith-Hutton convection-diffusion problem, where the steady-state transport of a scalar property  $\phi$  is investigated under a predefined velocity field. The numerical solution is obtained using the finite volume method with Upwind Difference Scheme (UDS) and Central Difference Scheme (CDS). The effects of varying the  $\rho/\Gamma$  ratio on the solution stability and accuracy are analyzed.

### 1 Introduction

The Smith-Hutton case is a reference problem for testing numerical schemes in convectiondiffusion problems. It consists of a solenoidal flow in a 2D domain, where a transported property  $\phi$  is studied under different transport regimes.

The governing convection-diffusion equation is given by:

$$\frac{\partial(\rho\phi)}{\partial t} + \nabla \cdot (\rho v\phi) = \nabla \cdot (\Gamma \nabla \phi) + \dot{S}_{\phi} \tag{1}$$

where v = (u, v) is the known velocity field and  $\Gamma$  is the diffusion coefficient.

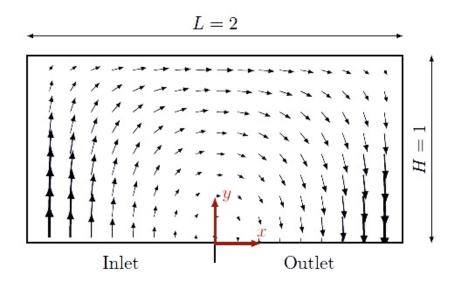


Figure 1: Smith-Hutton case

The boundary conditions for the present problem are:

- Inlet:  $\phi = 1 + \tanh[(2x+1)\alpha]$ , for  $x \in [-1,0], y = 0$
- Outlet:  $\frac{\partial \phi}{\partial y} = 0$ , for  $x \in [0, 1], y = 0$
- Walls:  $\phi = 1 \tanh[\alpha]$

where  $\alpha = 10$ .

The initial condition for the map is assumed to be  $\phi = 0 \ \forall (x, y)$ . The problem is solved using explicit temporal scheme and central difference scheme (CDS).

# 2 Theoretical Background

The convection-diffusion equation 1 provides a unified framework that can represent the Navier-Stokes, energy, and species transport equations by appropriately selecting the physical variable  $\phi$ . Regardless of the specific application, these equations share a common structure consisting of transient (unsteady) terms, convective terms, diffusive terms, and additional source or sink terms. In its general form, the equation models the unsteady transport of a scalar (or vector) quantity, where:

- $\phi$  is the scalar or vector variable being transported.
- $\Gamma_{\phi}$  is the rate of diffusion for the transported quantity  $\phi$ .
- $\dot{S}_{\phi}$  accounts for the generation or destruction of the quantity  $\phi$  in the system.

Using the mass conservation equation,  $\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0$ , the previous convection-diffusion equation can be rewritten as follows:

$$\rho \frac{\partial \phi}{\partial t} + \rho \vec{v} \cdot \nabla \phi = \nabla \cdot (\Gamma \nabla \phi) + \dot{S}_{\phi}$$
 (2)

Taking a small control volume, integrating it using the finite volume approach, and applying Gauss's theorem, the following equation is obtained:

$$\rho \frac{\partial \phi}{\partial t} \Delta V + \rho \sum_{i} (\vec{v}\phi)_f \cdot \Delta S = \sum_{i} (\Gamma \nabla \phi)_f \cdot \Delta S + \dot{S}_\phi \cdot \Delta V$$
 (3)

where the subscript f denotes that the properties are evaluated at the control volume faces rather than at the cell centroids.

Analyzing the terms individually, we obtain:

• Convective term:

$$conv = \rho \sum_{i} (\vec{v}\phi)_f \cdot \Delta S = \rho \left( u_e \phi_e \Delta S_e - u_w \phi_w \Delta S_w + v_n \phi_n \Delta S_n - v_s \phi_s \Delta S_s \right)$$
 (4)

• Diffusive term:

$$diff = \sum (\Gamma \nabla \phi)_f \cdot \Delta S = \Gamma_e \frac{\phi_E - \phi_P}{d_{PE}} \Delta S_e + \Gamma_w \frac{\phi_W - \phi_P}{d_{PW}} \Delta S_w + \Gamma_n \frac{\phi_N - \phi_P}{d_{PN}} \Delta S_n + \Gamma_s \frac{\phi_S - \phi_P}{d_{PS}} \Delta S_s$$
(5)

• Source/Sink term:

$$S = \dot{S}_{\phi} \cdot \Delta V = 0 \tag{6}$$

this term is neglected in the present exercise.

• Transient term:

$$\rho \frac{\partial \phi}{\partial t} \Delta V = \rho \frac{\phi_P^{n+1} - \phi_P^n}{\Delta t} \Delta V \tag{7}$$

Applying explicit Forward Euler time discretization, for simplicity, the final form of equation 1 is:

$$\frac{\rho \phi_P^{n+1} - \phi_P^n}{\Delta t} = \frac{1}{\Delta V} (-\text{conv}^n + \text{diff}^n)$$
 (8)

where both the convective and diffusive terms are evaluated at time instant  $t^n$ .

## 3 Numerical Implementation

The numerical implementation is structured using modular functions, with matrices handling data storage and processing. An iterative solver is used to advance the solution in time until steady-state convergence is reached. To streamline visualization, results are automatically plotted using Gnuplot directly within the C++ code. The following sections describe the key steps of the implementation

#### 3.1 Computational Mesh and Velocity Field

The computational domain is a rectangular region which has been discretized using a structured mesh of  $N_x = 100$  and  $N_y = 50$  control volumes. The mesh is semi-uniform, consisting of  $N_x + 2$  and  $N_y + 2$  nodes, which are positioned at the center of the control volumes and along the boundaries. The mesh is generated in the function MeshDefinition(), where the positions of nodes, faces, and control volume centers are computed:

- Control volume positions are stored in matrices  $x_{cv}[i][j]$  and  $y_{cv}[i][j]$ .
- Nodes' locations are derived from the control volume face coordinates and stored in  $x_P[i][j]$  and  $y_P[i][j]$ .
- The control volume's volume is computed and stored in  $v_P[i][j]$ .
- Boundary nodes are explicitly initialized to ensure correct positioning.

For the purpose of this exercise, the velocity field is assumed to be known and is assigned in the function VelocityField(), where the velocity components are defined at the control volume faces as:

$$u = 2y(1 - x^2), \quad v = -2x(1 - y^2)$$
 (9)

where the coordinates x and y refer, as specified earlier, to the coordinates of the corresponding face.

## 3.2 Discretization of the Governing Equation

The governing equation has been spatially discretized using the finite volume method and temporally using the explicit Forward Euler time integration method. The resulting equation from these implementations is as follows:

$$\frac{\rho \phi_P^{n+1} - \phi_P^n}{\Delta t} = \frac{1}{\Delta V} \left[ -\rho \left( u_e \phi_e^n \Delta S_e - u_w \phi_w^n \Delta S_w + v_n \phi_n^n \Delta S_n - v_s \phi_s^n \Delta S_s \right) + \Gamma \left( \frac{\phi_E^n - \phi_P^n}{d_{PE}} \Delta S_e + \frac{\phi_W^n - \phi_P^n}{d_{PW}} \Delta S_w + \frac{\phi_N^n - \phi_P^n}{d_{PN}} \Delta S_n + \frac{\phi_S^n - \phi_P^n}{d_{PS}} \Delta S_s \right) \right]$$
(10)

The transported property at each node/face in the right-hand side of the equation is evaluated at the time instant  $t^n$ . Since an explicit time discretization is employed, the updated value  $\phi_P^{n+1}$  can be directly computed at each time step without requiring the solution of a system of equations. This approach ensures a straightforward time advancement, as  $\phi_P^{n+1}$  depends only on known quantities from the previous time step, making the method computationally less expensive and easier to implement.

The convective term requires the transported property to be evaluated at the control volume face rather than at the node. Since the exact value at the face is unknown, an interpolation method must be applied. Initially, the Central Difference Scheme (CDS) was chosen as the first approach to solve the exercise. In the following, the solution using CDS will be explained in detail. Subsequently, the Upwind Difference Scheme (UDS) will be introduced to analyze the stability and convergence of the code.

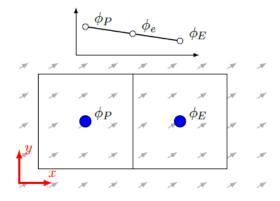


Figure 2: Central Difference Scheme (CDS)

The fundamental idea behind CDS is to approximate the value at the face using the known values of the surrounding nodes. This is achieved through linear interpolation. Following the example shown in Figure 2, and applying this method to each face, the interpolation is given by:

$$\phi_e^n = \frac{\phi_P^n + \phi_E^n}{2} \tag{11}$$

This scheme is not particularly dissipative but tends to be more unstable and may produce spurious oscillations if the chosen time step is not sufficiently small.

By introducing the CDS scheme into Equation 10, the resulting discretized equation is:

$$\frac{\rho\phi_P^{n+1} - \phi_P^n}{\Delta t} = \frac{1}{\Delta V} \left[ -\rho \left( u_e \frac{\phi_P^n + \phi_E^n}{2} \Delta S_e - u_w \frac{\phi_P^n + \phi_W^n}{2} \Delta S_w \right) + v_n \frac{\phi_P^n + \phi_N^n}{2} \Delta S_n - v_s \frac{\phi_P^n + \phi_S^n}{2} \Delta S_s \right) + \Gamma \left( \frac{\phi_E^n - \phi_P^n}{d_{PE}} \Delta S_e + \frac{\phi_W^n - \phi_P^n}{d_{PW}} \Delta S_w \right) + \frac{\phi_N^n - \phi_P^n}{d_{PN}} \Delta S_n + \frac{\phi_S^n - \phi_P^n}{d_{PS}} \Delta S_s \right) \right] (12)$$

Although an explicit time discretization is used, the coefficients  $a_F$  (F = N, S, E, W) are explicitly written to preserve consistency with the structured approach followed in previous exercises. This notation ensures a clearer representation of the finite volume formulation while explicitly distinguishing the contribution of neighboring nodes. Furthermore, the coefficients  $a_P$  and  $a_F$  naturally arise from the discretization process, making the transition between different schemes and numerical methods more intuitive.

The coefficients for each internal node of the mesh are then computed using the function InternalNodesCoefficients(), and they are defined as follows.

$$a_P = \frac{\rho}{\Delta t} \tag{13}$$

$$a_E = \frac{\Delta S_e}{\Delta V} \left( -\rho \frac{u_e}{2} + \frac{\Gamma}{d_{PE}} \right) \tag{14}$$

$$a_W = \frac{\Delta S_w}{\Delta V} \left( \rho \frac{u_w}{2} + \frac{\Gamma}{d_{PW}} \right) \tag{15}$$

$$a_N = \frac{\Delta S_n}{\Delta V} \left( -\rho \frac{v_n}{2} + \frac{\Gamma}{d_{PN}} \right) \tag{16}$$

$$a_S = \frac{\Delta S_s}{\Delta V} \left( \rho \frac{v_s}{2} + \frac{\Gamma}{d_{PS}} \right) \tag{17}$$

$$b_{P} = \frac{\rho}{\Delta t} + \frac{\Delta S_{e}}{\Delta V} \left( -\rho \frac{u_{e}}{2} - \frac{\Gamma}{d_{PE}} \right) + \frac{\Delta S_{w}}{\Delta V} \left( \rho \frac{u_{w}}{2} - \frac{\Gamma}{d_{PW}} \right) + \frac{\Delta S_{s}}{\Delta V} \left( -\rho \frac{v_{n}}{2} - \frac{\Gamma}{d_{PN}} \right) + \frac{\Delta S_{s}}{\Delta V} \left( \rho \frac{v_{s}}{2} - \frac{\Gamma}{d_{PS}} \right)$$
(18)

#### 3.3 Boundary Conditions

The function BoundaryConditions() imposes Dirichlet and Neumann boundary conditions at the domain boundaries:

• Inlet 
$$(y = 0, \text{ for } x < 0)$$
:  

$$\phi = 1 + \tanh((2x + 1)\alpha)$$
(19)

resulting in:

$$a_E = a_W = a_S = a_N = 0 (20)$$

$$a_P = 1, \quad b_P = 1 + \tanh((2x+1)\alpha)$$
 (21)

• Outlet (y = 0, for x > 0):

$$\frac{\partial \phi}{\partial y} = 0 \tag{22}$$

resulting in:

$$a_E = a_W = a_S = 0, \quad a_N = 1$$
 (23)

$$a_P = 1, \quad b_P = 0 \tag{24}$$

• Remaining boundaries (y = H, x = -L/2, x = L/2):

$$\phi = 1 - \tanh(\alpha) \tag{25}$$

resulting in:

$$a_E = a_W = a_S = a_N = 0 (26)$$

$$a_P = 1, \quad b_P = 1 - \tanh(\alpha)$$
 (27)

### 3.4 Solution Algorithm

The numerical solution is obtained following this procedure:

- Definition of global variables and matrices outside of **main()** to be used in multiple functions.
- Definition of the necessary functions, as described earlier.

- Inside main():
  - Definition of physical and numerical parameters.
  - Computation of the mesh using **MeshDefinition()**.
  - Computation of the velocity field using **VelocityField()**.
  - Initialization of the property map, setting  $\phi = 0$ .
  - Computation of the internal node coefficients for the initial map using **InternalNodesCoefficients()**.
  - Application of boundary conditions using **BoundaryConditions()**.

The steady-state map of the property  $\phi$  is then obtained iteratively in time using a solver, implemented in the function Solver():

- Initialization of the time counter to zero.
- Entry into the steady-state convergence loop:
  - Solve the discretized equation:
    - \* for internal nodes:

$$a_P \phi_P^{n+1} = a_W \phi_W^n + a_E \phi_W^n + a_N \phi_N^n + a_S \phi_S^n + b_P \phi_P^n$$
 (28)

\* for boundary nodes:

$$a_P \phi_P^{n+1} = a_W \phi_W^n + a_E \phi_W^n + a_N \phi_N^n + a_S \phi_S^n + b_P \tag{29}$$

– Assessment of convergence based on the maximum residual, ensuring that the changes in  $\phi$  remain below a predefined tolerance:

$$\max(\mathrm{Res}) < 10^{-6}$$

- Update of the time counter by adding  $\Delta t$ .
- If convergence is not reached, the previous map is updated:

$$\phi^n = \phi^{n+1}$$

• The iterative process is repeated until steady-state conditions are met, meaning that the difference between consecutive iterations remains below the prescribed tolerance.

### 4 Results

The solution was obtained for three different values of  $\rho/\Gamma$ : 10, 1000, and 1,000,000 and compared to the results given with the exercise. To solve the exercise,  $\Gamma$  was set to 1 and  $\rho$  was changed to fit the ratio  $\rho/\Gamma$ .

Position $x$	$\rho/\Gamma = 10$		$\rho/\Gamma = 1000$		$\rho/\Gamma = 1\ 000\ 000$	
	Given	Obtained	Given	Obtained	Given	Obtained
0.0	1.989	1.818	2.0000	1.9990	2.000	1.863
0.1	1.402	1.382	1.9990	1.9993	2.000	1.907
0.2	1.146	1.134	1.9997	1.9998	2.000	1.995
0.3	0.946	0.938	1.9850	1.9916	1.999	2.055
0.4	0.775	0.768	1.8410	1.8213	1.964	2.006
0.5	0.621	0.616	0.9510	0.9671	1.000	0.917
0.6	0.480	0.476	0.1540	0.1515	0.036	0.095
0.7	0.349	0.347	0.0010	0.0062	0.001	-0.020
0.8	0.227	0.225	0.0000	0.0001	0.000	0.001
0.9	0.111	0.111	0.0000	0.0000	0.000	0.000
1.0	0.000	0.000	0.0000	0.0000	0.000	0.000

Table 1: Table of provided and obtained results for different values of  $\rho/\Gamma$ .

Table 1 presents a comparison between the obtained and expected results. During the computation, some issues related to convergence and stability were encountered. For lower values of  $\rho/\Gamma$ , the algorithm exhibited convergence difficulties when using a time step of  $\Delta t = 10^{-3}$  s, whereas it successfully converged with  $\Delta t = 10^{-4}$  s. For the case of  $\rho/\Gamma = 10^3$ , no additional convergence issues were observed and the same time step was maintained. However, when dealing with high  $\rho/\Gamma$  values such as  $10^6$ , the program required significantly more computational time to reach convergence. Ultimately, it was necessary to slightly relax the convergence tolerance to  $10^{-5}$  and further reduce the time step to  $\Delta t = 10^{-5}$  s to achieve better results.

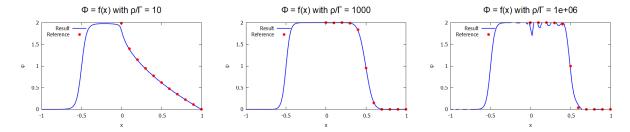


Figure 3: Plots of  $\phi$  at the outlet for different  $\rho/\Gamma$  ratios compared to the reference data.

The images in Figure 3 show the plot of the property  $\phi$  as a function of x, with the given reference values of  $\phi$  at the outlet represented by red dots to assess the accuracy of the results

It is of particular interest that in the case with  $\rho/\Gamma=10^6$ , spurious oscillations appeared. These errors are caused by the dominance of convection over diffusion, leading to numerical instabilities. In high  $\rho/\Gamma$  regimes, central differencing schemes tend to produce non-physical oscillations due to the lack of sufficient numerical dissipation. Additionally, rough grids and large time steps amplify aliasing errors. To mitigate this issue, a more stable convective scheme, such as UDS or QUICK can be used. Additionally, reducing the time step  $\Delta t$  improves stability, refining the mesh in high-gradient regions enhances spatial resolution, and ensuring consistent boundary conditions prevents errors at the domain outlet.

For a complete visualization of the results, color maps depicting the steady-state distribution of the property have been plotted using Gnuplot.

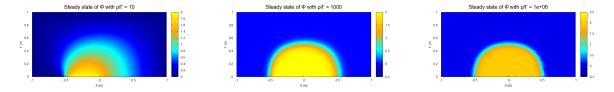


Figure 4: Color maps of  $\phi$  at steady state of the Smith-Hutton case for different  $\rho/\Gamma$  ratios.

The color map for  $\rho/\Gamma = 10^6$  shows some waviness on the right side, which corresponds to the spurious values observed in Figure 3. These are caused by numerical instabilities due to the central differencing scheme in convection-dominated cases. This highlights the need for more stable discretization methods or finer grid resolution to improve accuracy.

#### 5 Code Enhancements

To improve the numerical stability and flexibility of the implementation, some modifications were introduced. First, the code was restructured using classes, allowing easy interchangeability between the Central Differencing Scheme (CDS) and the Upwind Differencing Scheme (UDS). After this modification, the UDS was tested to compare its results with those obtained using the CDS.

#### 5.1 Class-Based Internal Nodes Coefficients Solver

The original InternalNodesCoefficients() function was implemented as a single global function, which made it harder to extend and modify when introducing different numerical methods such as CDS or UDS. To improve flexibility and modularity, the function was transformed into a class-based structure.

A base class, **InternalNodesCalculator**, was introduced to define a common interface for different numerical methods. This class provides a standardized framework for coefficient calculations and contains a pure virtual function, **computeCoefficients()**, which is implemented by specific subclasses.

Two main subclasses were created:

- CDSMethod: Implements the Central Differencing Scheme (CDS), which maintains the behavior of the original function.
- **UDSMethod**: Implements the Upwind Differencing Scheme (UDS), improving stability in convection-dominated problems.

This new approach facilitates the extension of the code, enabling the integration of additional discretization schemes without modifying the core solver structure.

To allow user selection of the method, the program asks the user to choose either 'CDS' or 'UDS' via command-line input. The selected method is created using the

createCalculator() function, which returns a pointer to the correct class. This method is then used to compute the internal coefficients before updating the solution over time until it reaches steady state.

In the **main()** function, the execution flow follows these steps:

- 1. The user inputs the desired discretization method ('CDS' or 'UDS').
- 2. The createCalculator() function returns an object of the corresponding class.
- 3. The selected method computes the internal node coefficients once, before time integration starts.
- 4. The explicit solver iteratively updates the transported property  $\phi$  over time until convergence is achieved.
- 5. The dynamically allocated memory for the class instance is cleared at the end of the execution.

# 5.2 Comparison Between UDS and CDS for $\rho/\Gamma = 10^6$

The Upwind Differencing Scheme (UDS) was then tested and compared to the previously used Central Differencing Scheme (CDS).

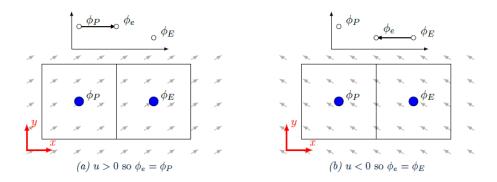


Figure 5: Upwind Difference Scheme (UDS) examples.

In the Upwind Difference Scheme (UDS), the value at a control volume face is assigned based on the direction of the velocity. Specifically, the face value is taken from the upstream node, meaning the node where the flow originates. This ensures that the transported quantity follows the natural direction of the fluid flow, improving numerical stability by introducing artificial diffusion. However, while UDS prevents oscillations and enhances robustness, it can also lead to excessive numerical diffusion, especially in high  $\rho/\Gamma$  regimes.

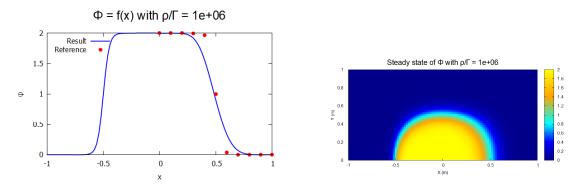


Figure 6: Plot at the outlet and color map at steady state of  $\phi$  with  $\rho/\Gamma = 10^6$ , using Upwind Difference Scheme (UDS).

Figure 6 presents results obtained using the Upwind Difference Scheme with  $\rho/\Gamma=10^6$ . The code turned out to be more stable and prone to convergence than when using CDS; a time step of  $10^{-4}$  and a time convergence tolerance of  $10^{-6}$  were sufficient to achieve convergence in less than 10 seconds.

Position $x$	$\rho/\Gamma = 1000000$				
	Given	CDS	UDS		
0.0	2.000	1.863	1.998		
0.1	2.000	1.907	2.000		
0.2	2.000	1.995	1.995		
0.3	1.999	2.055	1.919		
0.4	1.964	2.006	1.563		
0.5	1.000	0.917	0.906		
0.6	0.036	-0.095	0.330		
0.7	0.001	-0.020	0.068		
0.8	0.000	-0.001	0.007		
0.9	0.000	0.000	0.000		
1.0	0.000	0.000	0.000		

Table 2: Table of provided and obtained results for  $\rho/\Gamma = 1000000$  using CDS and UDS.

As shown in the comparison graph and in Table 2, no spurious oscillations were observed. However, despite its improved stability, UDS introduces significant numerical diffusion, particularly in regions with steep gradients. This effect is visible in the smoothing of the sharp profile, where the solution deviates from the expected reference values.

This result highlights the trade-off between numerical diffusion and stability: while UDS is more robust and ensures a smooth solution, it tends to over-diffuse sharp transitions, reducing accuracy in capturing steep gradients.

## 6 Conclusion

The numerical results show that increasing the  $\rho/\Gamma$  ratio leads to a sharper gradient in  $\phi$ , with possible oscillations when using the CDS scheme. The UDS scheme, while more

stable, introduces numerical diffusion. This highlights the trade-off between stability and accuracy when selecting convective schemes.

## A C++ Code

```
#include "Smith_Hutton_Upgraded.h"
#include <iostream>
#include <fstream>
#include <chrono>
#include <cmath>
#include <cstdlib>
\begin{array}{ll} \textbf{using namespace} & \text{std} \ ; \\ \textbf{using namespace} & \text{std} :: \text{chrono} \ ; \end{array}
  // Global Variables
double const L=2.0;
double const H=1.0;
double const W=1.0;
const int Nx=100, Ny=50;
                                                                                                                                 // Lenght
// Height
// Width
// Number of control volumes (100x50)
\label{eq:continuous} $ \begin{array}{ll} \textit{// Vectors Definition} \\ \textit{double phi} [Ny+2][Nx+2], \ phi-1[Ny+2][Nx+2], \ phi-g[Ny+2][Nx+2]; \\ \textit{double aP} [Ny+2][Nx+2], \ aE[Ny+2][Nx+2], \ aW[Ny+2][Nx+2], \\ aN[Ny+2][Nx+2], \ aS[Ny+2][Nx+2], \ bP[Ny+2][Nx+2]; \\ \textit{double } xcv[Ny+2][Nx+2], \ ycv[Ny+2][Nx+2], \ xP[Ny+2][Nx+2], \\ yP[Ny+2][Nx+2], \ vP[Ny+2][Nx+2]; \\ \textit{double } u[Ny+2][Nx+2], \ v[Ny+2][Nx+2]; \\ \end{aligned} 
                                                                                                                                                                                                                                                                  // Transported property
// Coefficients
                                                                                                                                                                                                                                                                   // Velocity components
xcv[i][j] = L / 2,
ycv[i][j] = H,
xP[i][j] = L / 2,
yP[i][j] = H,
vP[i][j] = 0;
else if (i == Ny+1 && j == 0)
    xcv[i][j] = -L / 2,
    ycv[i][j] = 0,
    xP[i][j] = 0,
    xP[i][j] = 0,
    vP[i][j] = 0;
else if (i == Ny+1 && j == Nx+1)
    xcv[i][j] = L / 2,
    ycv[i][j] = 0,
    xP[i][j] = L / 2,
    ycv[i][j] = L / 2,
    ycv[i][j] = 0,
    xP[i][j] = L / 2,
    yp[i][j] = H / 2,
    yP[i][j] = H,
    xv[i][j] = H,
    vP[i][j] = O;
else if (j == 0)
    xcv[i][j] = L / 2,
    ycv[i][j] = L / 2,
    yp[i][j] = U / 2,
    vP[i][j] = 0;
else if (j == 0)
    xcv[i][j] = L / 2,
    vP[i][j] = 0;
else if (j == Nx+1)
    xcv[i][j] = L / 2,
                                                 \begin{array}{c} \text{vP}[\,i\,][\,j\,] = 0; \\ \text{else if } (j = \text{Nx} + 1) \\ \text{xcv}[\,i\,][\,j\,] = \text{L} \ / \ 2, \\ \text{ycv}[\,i\,][\,j\,] = \text{H} - \text{Dy} \ * \ i\,, \\ \text{xP}[\,i\,][\,j\,] = \text{L} \ / \ 2, \\ \text{yP}[\,i\,][\,j\,] = \text{L} \ / \ 2, \\ \text{yP}[\,i\,][\,j\,] = 0; \\ \text{else if } (i = \text{Ny} + 1) \\ \text{xcv}[\,i\,][\,j\,] = -\text{L} \ / \ 2 + \text{Dx} \ * \ j\,, \\ \text{ycv}[\,i\,][\,j\,] = -\text{L} \ / \ 2 + \text{Dx} \ * \ j\,, \\ \text{ycv}[\,i\,][\,j\,] = 0, \\ \text{xP}[\,i\,][\,j\,] = 0, \\ \text{xP}[\,i\,][\,j\,] = 0, \\ \text{yP}[\,i\,][\,j\,] = 0; \\ \end{array} 
                                 }
                }
```

```
// Function to assign velocity value at each face
void VelocityField() {
    for (int i = 0; i < Ny+2; i++) {
        for (int j = 0; j < Nx+2; j++) {
            u[i][j] = 2 * yP[i][j] * (1 - pow(xcv[i][j],2)); // ue of node P, if I want uw I'll
            write u[i][j-1]
        v[i][j] = -2 * xP[i][j] * (1 - pow(ycv[i][j],2)); // vs of node P, if I want vn I'll
            write u[i-1][j]
                            }
             }
}
 // Base Class to implement multiple methods {f class} InternalNodesCalculator {
 public:
              virtual void computeCoefficients(double rho, double dt, double gamma) = 0; // Pure virtual
              function
virtual ~InternalNodesCalculator() = default;
                                                                                                                                                                                                                                                                                      // Virtual destructor
// Central Differencing Scheme (CDS) method class CDSMethod: public InternalNodesCalculator {
// Now computeCoefficients() is a method of the class, which can be called with CDSMethod object public:
              void computeCoefficients (double rho, double dt, double gamma) override { for (int i=1;\ i< Ny+1;\ i++) { for (int j=1;\ j< Nx+1;\ j++) {
                                                        \begin{array}{l} \textbf{double} \ \ DSe = (ycv \, [\, i \, - \, 1\, ] \, [\, j\, ] \, - \, ycv \, [\, i\, ] \, [\, j\, ]) \ * \ L; \\ \textbf{double} \ \ DSw = (ycv \, [\, i \, - \, 1\, ] \, [\, j\, ] \, - \, ycv \, [\, i\, ] \, [\, j\, ]) \ * \ L; \\ \textbf{double} \ \ DSn = (xcv \, [\, i\, ] \, [\, j\, ] \, - \, xcv \, [\, i\, ] \, [\, j\, - \, 1\, ]) \ * \ L; \\ \textbf{double} \ \ DSs = (xcv \, [\, i\, ] \, [\, j\, ] \, - \, xcv \, [\, i\, ] \, [\, j\, - \, 1\, ]) \ * \ L; \\ \end{array} 
                                                        \begin{array}{lll} \mbox{double} \ dPE = xP[\,i\,][\,j\,+\,1] - xP[\,i\,][\,j\,]; \\ \mbox{double} \ dPW = xP[\,i\,][\,j] - xP[\,i\,][\,j\,-\,1]; \\ \mbox{double} \ dPN = yP[\,i\,-\,1][\,j\,] - yP[\,i\,][\,j\,]; \\ \mbox{double} \ dPS = yP[\,i\,][\,j\,] - yP[\,i\,+\,1][\,j\,]; \end{array} 
                                                      aE[i][j] = DSe / vP[i][j] * (-rho * u[i][j] / 2 + gamma / dPE);
aW[i][j] = DSw / vP[i][j] * (rho * u[i][j - 1] / 2 + gamma / dPW);
aN[i][j] = DSn / vP[i][j] * (-rho * v[i - 1][j] / 2 + gamma / dPN);
aS[i][j] = DSs / vP[i][j] * (rho * v[i][j] / 2 + gamma / dPS);
aP[i][j] = rho / dt;
bP[i][j] = rho / dt + DSe / vP[i][j] * (-rho * u[i][j] / 2 - gamma / dPE) +
DSw / vP[i][j] * (rho * u[i][j - 1] / 2 - gamma / dPW
                                                                                                                                                                          DSn / vP[i][j] * (-rho * v[i - 1][j] / 2 - gamma /
                                                                                                                                                                          dPN) +
DSs / vP[i][j] * (rho * v[i][j] / 2 - gamma / dPS);
             }
 // Upwind Differencing Scheme (UDS) method
class UDSMethod: public InternalNodesCalculator {
    // Now_computeCoefficients() is another method of the class, which can be called with UDSMethod
 public:
              void computeCoefficients(double rho, double dt, double gamma) override {
   for (int i = 1; i < Ny + 1; i++) {
      for (int j = 1; j < Nx + 1; j++) {
                                                       \begin{array}{l} \textbf{double} \ \ DSe = \big(ycv \begin{bmatrix} i - 1 \end{bmatrix} [j] - ycv \begin{bmatrix} i \end{bmatrix} [j] \big) \ * \ L; \\ \textbf{double} \ \ DSw = \big(ycv \begin{bmatrix} i - 1 \end{bmatrix} [j] - ycv \begin{bmatrix} i \end{bmatrix} [j] \big) \ * \ L; \\ \textbf{double} \ \ DSn = \big(xcv \begin{bmatrix} i \end{bmatrix} [j] - xcv \begin{bmatrix} i \end{bmatrix} [j - 1] \big) \ * \ L; \\ \textbf{double} \ \ DSs = \big(xcv \begin{bmatrix} i \end{bmatrix} [j] - xcv \begin{bmatrix} i \end{bmatrix} [j - 1] \big) \ * \ L; \\ \textbf{double} \ \ \ DSs = \big(xcv \begin{bmatrix} i \end{bmatrix} [j] - xcv \begin{bmatrix} i \end{bmatrix} [j - 1] \big) \ * \ L; \\ \end{array}
                                                        \begin{array}{lll} \mbox{double} \ dPE = xP[\,i\,][\,j\,+\,1] - xP[\,i\,][\,j\,]; \\ \mbox{double} \ dPW = xP[\,i\,][\,j\,] - xP[\,i\,][\,j\,-\,1]; \\ \mbox{double} \ dPN = yP[\,i\,-\,1][\,j\,] - yP[\,i\,][\,j\,]; \\ \mbox{double} \ dPS = yP[\,i\,][\,j\,] - yP[\,i\,+\,1][\,j\,]; \end{array} 
                                                      aE[i][j] = DSe / vP[i][j] * (max(-rho * u[i][j], 0.0) + gamma / dPE);
aW[i][j] = DSw / vP[i][j] * (max(rho * u[i][j-1], 0.0) + gamma / dPW);
aN[i][j] = DSn / vP[i][j] * (max(-rho * v[i-1][j], 0.0) + gamma / dPN);
aS[i][j] = DSs / vP[i][j] * (max(rho * v[i][j], 0.0) + gamma / dPS);
aS[i][j] = rho / dt;
bP[i][j] = rho / dt + DSe / vP[i][j] * (min(-rho * u[i][j], 0.0) - gamma / dPE) +

DSw / vP[i][j] * (min(rho * u[i][j-1], 0.0) - gamma / dPE) +
                                                                                                                                                                         DSS / vP[i][j] * (min(rho * u[i][j-1], 0.0) - gam / dPW) +
DSn / vP[i][j] * (min(-rho * v[i-1][j], 0.0) - gamma / dPN) +
DSS / vP[i][j] * (min(rho * v[i][j], 0.0) - gamma / dPS);
                      }
}:
    // Function to choose the method
 // runction to choose the method
InternalNodesCalculator* createCalculator(string method) {
   if (method == "CDS") {
      return new CDSMethod();
   } else if (method == "UDS") {
                           return new UDSMethod();
              } else {
   cerr << "Not-a-valid-method!" << std::endl;</pre>
```

```
return nullptr;
       }
}
// Function to apply boundary conditions
void BoundaryConditions(double alpha) {
    for (int i = 0; i < Ny+2; i++) {
        if (i = Ny+1 && x P[i][j] < 0)
            aE[i][j] = 0,
            aW[i][j] = 0,
            aN[i][j] = 0,
            aS[i][j] = 1,
            aP[i][j] = 1 + tanh((2 * xP[i][j] + 1) * alpha);
    else if (i = Ny+1 && xP[i][j] > 0)
        aE[i][j] = 0,
        aW[i][j] = 0,
        aW[i][j] = 1,
        aN[i][j] = 1,
        aS[i][j] = 0,
        aW[i][j] = 0,
        aN[i][j] = 1,
        aS[i][j] = 0,
        aP[i][j] = 1,
        bP[i][j] = 0;
    else if (i = 0 || j == 0 || j == Nx+1)
        aE[i][j] = 0,
        aW[i][j] = 0,
        aN[i][j] = 1,
        bP[i][j] = 1 - tanh(alpha);
}

 // Function to apply boundary conditions
                                                                                                                                                        // Inlet Condition
                                                                                                                                                         // Outlet Condition
                                                                                                                                                       // Rest of the boundaries
                                bP[i][j] = 1 - tanh(alpha);
             }
        }
}
 void Solver (double maxRes, double &t_count, double dt, double rho, double gamma) {
        // Time Loop
double res = maxRes + 1; // Condition to enter the loop
while (res > maxRes) {
    double maxDiff = 0.0;
               }
                                double diff = fabs(phi_1[i][j] - phi[i][j]);
if (diff > maxDiff) {
    maxDiff = diff;
                       }
                 t_count += dt;
                 res = maxDiff;
                // Update phi
for (int i = 0; i < Ny + 2; i++) {
    for (int j = 0; j < Nx + 2; j++) {
        phi[i][j] = phi-1[i][j];
                }
        std::cout << "Steady-state-reached-in-" << t_count << "-seconds" << std::endl;
}
int main() {
        // Physical Data
double phi_0 = 0.0;
double gamma = 1.0;
double ratio = 10000000;
                                                                                                // Initial condition in all the domain

// Assumed equal to 1

// To be changed with 10, 1.000, 1.000.000

// Uniform in the domain
        double rho = gamma * ratio;
double alpha = 10;
         // Numerical Data
        double Dx = L / Nx, Dy = H / Ny;
double dt = 1e-4;
double t\_count = 0.0;
                                                                                                 // Time step
// Initial time count
// Convergence tolerance
// To choose the scheme
        double maxRes = 1e-6;
string method;
        // Start Timer
auto start = high_resolution_clock::now();
           / Compute Mesh
         MeshDefinition (H, L, W, Dx, Dy);
        // Compute Velocity Field VelocityField ();
          // Initial Map
        for (auto &row: phi) {
   for (auto &elem: row) {
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elem = phi_0;
                }
 }
 //Choose Interpolation Scheme cout << "Choose the method ('CDS' or 'UDS'): "; cin >> method;
 // Create an instance of the chosen method
InternalNodesCalculator* calculator = createCalculator(method);
  if (!calculator) {
                cerr << "Error: Discretization not valid!" << endl; return 1;
 // Compute Internal Nodes Coefficients calculator->computeCoefficients(rho, dt, gamma);
 // Compute Boundary Conditions
Boundary Conditions (alpha);
  // \ Loop \ Until \ Steady-State \ is \ Reached \\ Solver(maxRes, \ t\_count, \ dt, \ rho, \ gamma); 
       / Clear memory
  delete calculator;
  // File to check matrices
 // File to check matrices
ofstream TestFile ("Test.txt");
for (int i = 0; i < Ny+2; i++) {
    for (int j = 0; j < Nx+2; j++) {
        TestFile << xP[i][j] << "-";
                  TestFile << endl;
}
}
 // File to print data for Gnuplot 1
ofstream GnuplotDatal ("GnuplotDatal.txt");
for (int i = 0; i < Ny+2; i++) {
    for (int j = 0; j < Nx+2; j++) {
        GnuplotDatal << xP[i][j] << "-" << yP[i][j] << "-" << phi-1[i][j] << endl;
                 GnuplotData1 << "\n";
 }
  // File to print data for Gnuplot 2
ofstream GnuplotData2 ("GnuplotData2.txt");
\begin{array}{lll} & \text{ ofstream } & \text{ GnuplotData2 } \left( \text{ "GnuplotD} \right. \\ & \text{ double } & \text{ ref} \left[ 11 \right] \left[ 4 \right] & = \left\{ & \left\{ 0.0, 1.989, 2.0000, 2.000 \right\}, \\ & \left\{ 0.1, 1.402, 1.9990, 2.000 \right\}, \\ & \left\{ 0.2, 1.146, 1.9997, 2.000 \right\}, \\ & \left\{ 0.3, 0.946, 1.9850, 1.999 \right\}, \\ & \left\{ 0.4, 0.775, 1.8410, 1.964 \right\}, \\ & \left\{ 0.5, 0.621, 0.9510, 1.000 \right\}, \\ & \left\{ 0.6, 0.480, 0.1540, 0.036 \right\}, \\ & \left\{ 0.7, 0.349, 0.0010, 0.001 \right\}, \\ & \left\{ 0.8, 0.227, 0.0000, 0.000 \right\}, \\ & \left\{ 0.9, 0.111, 0.0000, 0.000 \right\}, \\ & \left\{ 1.0, 0.000, 0.0000, 0.000 \right\}, \\ & \left\{ 0.0, 0.000, 0.0000, 0.0000, 0.000 \right\}, \\ & \left\{ 0.0, 0.000, 0.0000, 0.0000, 0.0000, 0.0000 \right\}, \\ & \left\{ 0.0, 0.000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.
 \}; \\ // \ Prints \ automatically \ the \ correct \ column
 int index = 1;
if (ratio == 1000)
index = 2;
  else if (ratio = 1000000)
 index = 3;

for (int i = 0; i < 11; i++) {

    GnuplotData2 << ref[i][0] << "-" << ref[i][index] << std::endl;
 // File to print property distribution ofstream PhiDistribution ("PhiDistribution.txt"); for (int i=0;\ i< Ny+2;\ i++) { for (int j=0;\ j< Nx+2;\ j++) { PhiDistribution << phi-1[i][j] << "-";
                 PhiDistribution << endl;
 TestFile.close();
```