

SOLVING THE MAXIMUM INDEPENDANT SET PROBLEM WITH EINSUM NETWORKS

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ABSTRACT

Solving the maximum independent set size problem by mapping the problem to an einsum network. We obtain the maximum independent set size, the independence polynomial and the optimal configuration by using different tensor elements to a commutative semiring.

1 COMPUTING INDEPENDENCE POLYNOMIAL

The independence polynomial (Ferrin, 2014; Harvey et al., 2017) of a graph G is defined as

$$I(G, x) = \sum_{k=1}^{\alpha(G)} a_k x^k, \quad (1)$$

where a_k is the number of independent sets of order k in G , and $\alpha(G)$ is the maximum independent set size of graph G . If we map a graph into an einsum network, as shown in Fig. 1, by putting a vertex tensor of size 2 at each vertex

$$W(x) = \begin{pmatrix} 1 \\ x \end{pmatrix}, \quad (2)$$

and a matrix of size 2×2 at each edge

$$B = \begin{pmatrix} 1 & 1 \\ 1 & 0 \end{pmatrix}. \quad (3)$$

where x is a variable. The contraction of this einsum network is the independence polynomial

$$I(G, x) = \sum_{s_1, s_2, \dots, s_n=0}^1 \prod_{i=1}^n W(x)_{s_i} \prod_{(i,j) \in E(G)} B_{s_i s_j}. \quad (4)$$

Here, the einsum enumerates all possible configurations and accumulates the product of tensor elements to the output. The product over vertices tensors provides a factor x^k , where $k = \sum_k s_k$ is the set size. While the product over edge tensors provides a factor 0 if the configuration is not an independent set, otherwise 1. The benefit of mapping the problem to einsum is, some existing algorithms can be used to decrease the complexity of the problem. Especially those techniques developed in recent quantum circuit simulations for finding the optimal exactly contraction order of a tensor network (Gray & Kourtis, 2021; Pan & Zhang, 2021). With these algorithms, one can easily find a good contraction order that reduces the complexity to a quantity related to the minimum treewidth of the graph. (Markov & Shi, 2008) Such mapping also makes it possible to utilize faster BLAS functions.

$$\begin{pmatrix} 1 & r\omega & r^2\omega^2 & \dots & r^m\omega^m \\ 1 & r\omega^2 & r^2\omega^4 & \dots & r^m\omega^{2m} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & r\omega^m & r^2\omega^{2m} & \dots & r^m\omega^{m^2} \end{pmatrix} \begin{pmatrix} a_0 \\ a_1 \\ \vdots \\ a_m \end{pmatrix} = \begin{pmatrix} y_0 \\ y_1 \\ \vdots \\ y_m \end{pmatrix} \quad (8)$$

When $r = 1$, the left side is a DFT matrix. We can obtain the factors using the relation $\vec{d} = \text{FFT}^{-1}(\omega) \cdot \vec{y}$. In the special case that $\omega = e^{-2\pi i/(m+1)}$, it is directly solvable with inverse fast fourier transformation algorithm in package FFTW.

$$\text{FFT}(\omega) \cdot \vec{d}_r = \vec{y} \quad (9)$$

where $(\vec{d}_r)_k = a_k r^k$, by choosing different r , we can obtain better precision in low independent set size region ($\omega < 1$) and high independent set size region ($\omega > 1$).

1.4 FINITE FIELD ALGEBRA

It is possible to compute the independence polynomials exactly using basic integer types only, even when the degeneracy is bigger than that can be represented by 64 bit integers. What we need is computing with the finite field algebra $GF(p)$

$$\begin{aligned} x \pmod{p} \oplus y \pmod{p} &= x + y \pmod{p} \\ x \pmod{p} \odot y \pmod{p} &= xy \pmod{p} \\ \mathbb{0} &= 0 \\ \mathbb{1} &= 1 \end{aligned} \quad (10)$$

It is a field because the multiplicative inverse can be defined on non-zero values, which can be computed with the extended Euclidean algorithm. In a finite field algebra, we have the following observations

1. One can still use Gaussian elimination (Golub & Van Loan, 2013) to solve the linear equations in a finite field, since the multiplicative inverse is properly defined in a finite field,
2. Let remainders of an integer x over a set of coprime integers $\{p_1, p_2, \dots, p_n\}$ being $\{x \pmod{p_1}, x \pmod{p_2}, \dots, x \pmod{p_n}\}$, then $\{x \pmod{p_1 \times p_2 \times \dots \times p_n}\}$ can be obtained using the chinese remainder theorem.

With these observations, we use Algorithm 1 to compute degeneracies. In the algorithm, except the computation of chinese remainder theorem, all computations are done with integers with fixed width.

Algorithm 1: Compute independence polynomial exactly without integer overflow

```

1 Let  $P = 1$ ,  $X = 0, 1, 2, \dots, m$  and  $\hat{X}_{ij} = X_{ij}^j$ , where  $i, j = \{0, 1, \dots, m\}$ ;
2 while true do
3   compute the largest prime 64 bit integer  $p$  that  $\text{gcd}(p, P) = 1$ ;
4   compute the tensor network contraction on  $GF(p)$  and obtain  $Y \pmod{p} = \{y_0 \pmod{p}, y_1 \pmod{p}, \dots, y_m \pmod{p}\}$ ;
5    $A \pmod{p} = \{a_0 \pmod{p}, a_1 \pmod{p}, \dots, a_m \pmod{p}\} = \text{gaussian\_elimination}(\hat{X}, Y \pmod{p})$ ;
6    $A \pmod{P \times p} = \text{chinese\_remainder}(A \pmod{P}, A \pmod{p})$ ;
7   if  $A \pmod{P} = A \pmod{P \times p}$  then
8     return  $A \pmod{P}$ ;
9   end
10   $P = P \times p$ ;
11 end
```

2 COMPUTING MAXIMUM INDEPENDENT SET SIZE AND ITS DEGENERACY

Obtaining the maximum independent set size and its degeneracy is much easier. Let $x = \infty$, then the independence polynomial becomes

$$I(G, \infty) = a_k \infty^{\alpha(G)}, \quad (11)$$

where the lower orders terms disappear automatically. We can define a new algebra as

$$\begin{aligned} a_x \infty^x \oplus a_y \infty^y &= \begin{cases} (a_x + a_y) \infty^{\max(x,y)}, & x = y \\ a_y \infty^{\max(x,y)}, & x < y \\ a_x \infty^{\max(x,y)}, & x > y \end{cases} \\ a_x \infty^x \odot a_y \infty^y &= a_x a_y \infty^{x+y} \\ \mathbf{0} &= -\infty \\ \mathbf{1} &= 0 \end{aligned} \quad (12)$$

In the program, we only store the power x and the corresponding factor a_x that initialized to 1. This is exactly the tropical tensor network (Liu et al., 2021) used in solving spin glass ground states. Here if we are only interested in obtaining $\alpha(G)$, we can only consider the powers, which corresponds to the tropical algebra (Maclagan & Sturmfels, 2015).

One may want to obtain all ground state configurations, it can be achieved replacing the factors a_x with a set of bit strings s_x initialized to a vertex dependent value $s_i = \{\delta_{ij=1, \dots, n}\}$. The bit strings has a redefined algebra

$$\begin{aligned} s_x \oplus s_y &= s_x \cup s_y \\ s_x \odot s_y &= \{\sigma \vee \tau | \sigma \in s_x, \tau \in s_y\} \end{aligned} \quad (13)$$

The \vee is the bit-wise or.

2.1 UTILIZING SPARSITY

So far, we have used the language of einsum networks for contraction. When using sparse tropical tensor networks to find the maximum independent set, we can introduce a new rule to compress the tensor by removing elements that are not helpful. As show in Fig. 2, after we contract the tensors in a subregion $R \subseteq G$ of a graph G , and obtain a resulting tensor A of rank $|C|$, where C is the set of vertex tensors at the cut. Each tensor entry A_σ defines a local maximum independent set size with a fixed boundary configuration $\sigma \in \{0, 1\}^{|C|}$ by marginalizing the inner degrees of freedom. We call this properly the *local maximum rule*.

In addition, suppose we have two entries with the same local maximum independent set size corresponding to local configurations shown in (a) and (b), it is safe to claim configuration (a) being better than configuration (b) and remove the entry A_{σ_b} . This is because the boundary of (a) is less restrictive to the rest of the graph while having the equally good local maximum independent set size, i.e. any exterior configuration $\bar{\sigma}$ making $\bar{\sigma} \cup \sigma_b$ a global maximum independent set also makes $\bar{\sigma} \cup \sigma_a$ a maximum independent set. Hence an entry A_{σ_a} is “better” than A_{σ_b} can be defined as

$$(\sigma_a \wedge \sigma_b = \sigma_a) \wedge (A_{\sigma_a} \geq A_{\sigma_b}), \quad (14)$$

where \wedge is a bitwise and operations. The first term means that whenever a bit in σ_a has boolean value 1, the corresponding bit in σ_b is also 1. While the second term means the maximum independent set size with boundary configuration fixed to σ_a is not less than that fixed to σ_b . The word “better” means the best solution with boundary configuration σ_a is never worse than that with σ_b . When Eq. 14 holds, It is easy to see that if $\sigma_b \cup \bar{\sigma}$ is one of the solutions for maximum independent sets in G , $\sigma_a \cup \bar{\sigma}$ is also a solution. With this observation, it is safe to set A_{σ_b} to tropical zero. We call this property the *least restrictive rule*, that is, among boundary configurations with equal local maximum independent sizes, we only retain those least restrictive (less ones at the boundary) to exterior configurations.

2.2 THE TENSOR NETWORK COMPRESSION DETECTS BRANCHING RULES AUTOMATICALLY

In the following, we are going to show *local maximum rule* and *least restrictive rule* can automatically discover branching rules and use it for compressing tensor elements, i.e. sparse tropical tensor network can not be worse than branching in truncating the search space.

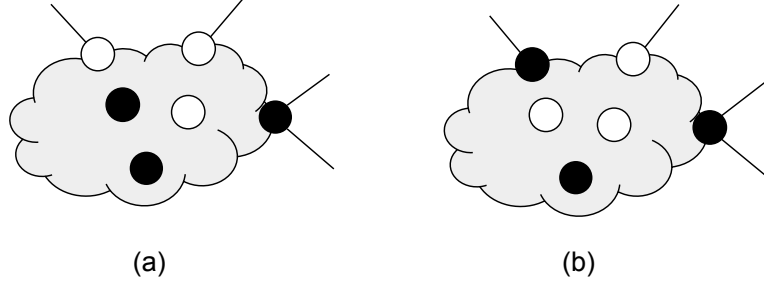


Figure 2: Two configurations with same local independent size $v_a = v_b = 3$ and different boundary configurations (a) 001 and (b) 101, where black nodes are 1s (in the independent set) and white nodes are 0s (not in the independent set).

We are going to verify the Lemmas used for branching in book (Fomin & Kaski, 2013).

Branching Rule 1. *If a vertex v is in an independent set I , then none of its neighbors can be in I . On the other hand, if I is a maximum (and thus maximal) independent set, and thus if v is not in I then at least one of its neighbors is in I .*

Contract $N[v]$ and the resulting tensor A has a rank $|N(v)|$. Each tensor entry A_σ corresponds to a locally maximized independent set size with fixed boundary configuration $\sigma \in \{0, 1\}^{|N(v)|}$. If the boundary configuration is a bit string of 0s, σ_v will take value 1 to maximize the local independent set size.

Branching Rule 2. *Let $G = (V, E)$ be a graph, let v and w be adjacent vertices of G such that $N[v] \subseteq N[w]$. Then*

$$\alpha(G) = \alpha(G \setminus w). \quad (15)$$

Contract $N[w]$, both $\{v, w\}$ disappear from the tensor indices because they are inner degrees of freedom. If w is one, then $N[v]$ are all zeros, the resulting tensor element can not be larger than setting $v = 1$ and $w = 0$. By the maximization rule, the local tensor does not change if we remove w .

Branching Rule 3. *Let $G = (V, E)$ be a graph and let v be a vertex of G . If no maximum independent set of G contains v then every maximum independent set of G contains at least two vertices of $N(v)$.*

Contract $N[v]$, the minimum tensor element is 2, otherwise letting v be in the independent set is one of the solutions. This is again captured by the *local maximum rule*.

Branching Rule 4. *Let $G = (V, E)$ be a graph and v a vertex of G . Then*

$$\alpha(G) = \max(1 + \alpha(G \setminus N[v]), \alpha(G \setminus (M(v) \cup \{v\}))). \quad (16)$$

Here, $M(v)$ is the set of mirrors of v in G . A vertex $w \in N^2(v)$ is called a mirror of v if $N(v) \setminus N(w)$ is a clique. This rule states that if v is not in M , there exists an MIS I that $M(v) \not\subseteq I$. otherwise, there must be one of $N(v)$ in the MIS (*local maximum rule*). If w is in I , then none of $N(v) \cap N(w)$ is in I , then there must be one of node in the clique $N(v) \setminus N(w)$ in I (*local maximum rule*), since clique has at most one node in the MIS, the tensor compression will eliminate this solution by moving the occupied node to the interior. Hence, the *least restrictive rule* captures the mirror rule.

Branching Rule 5. *Let $G = (V, E)$ be a graph and v be a vertex of G such that $N[v]$ is a clique. Then*

$$\alpha(G) = 1 + \alpha(G \setminus N[v]). \quad (17)$$

Contract $N[v]$, and this rule can be captured by the *local maximum rule*.

Branching Rule 6. *Let G be a graph, let S be a separator of G and let $I(S)$ be the set of all subsets of S being an independent set of G . Then*

$$\alpha(G) = \max_{A \in I(S)} |A| + \alpha(G \setminus (S \cup N[A])). \quad (18)$$

This rule corresponds to first contract tensors in subgraph S , then $N[S]$, and then contract the rest parts. These branching rule can be captured by the *local maximum rule*.

Branching Rule 7. Let $G = (V, E)$ be a disconnected graph and $C \subseteq V$ a connected component of G . Then

$$\alpha(G) = \alpha(G[C]) + \alpha(G \setminus C). \quad (19)$$

Contract by the disconnected parts, and this rule can be captured by the *local maximum rule*.

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A TECHNICAL GUIDE

OMEinsum a package for einsum,

OMEinsumContractionOrders a package for finding the optimal contraction order for einsum
<https://github.com/Happy-Diode/OMEinsumContractionOrders.jl>,

TropicalGEMM a package for efficient tropical matrix multiplication (compatible with OMEinsum),

TropicalNumbers a package providing tropical number types and tropical algebra, one of the dependency of TropicalGEMM,

LightGraphs a package providing graph utilities, like random regular graph generator,

Polynomials a package providing polynomial algebra and polynomial fitting,

Mods and Primes packages providing finite field algebra and prime number generators.

One can install these packages by opening a julia REPL, type `]` to enter the `pkg>` mode and type, e.g.

```
pkg> add OMEinsum
```

An example of computing the maximum independent set size

```
julia> using OMEinsum, LightGraphs, TropicalNumbers

julia> n, k = 6, 3
(6, 3)

julia> g = LightGraphs.random_regular_graph(n, k)
{6, 9} undirected simple Int64 graph

julia> ix = [minmax(e.src,e.dst) for e in LightGraphs.edges(g)]
9-element Vector{Tuple{Int64, Int64}}:
 (1, 2)
 (1, 4)
 (1, 5)
 (2, 3)
 (2, 6)
 (3, 5)
 (3, 6)
 (4, 5)
 (4, 6)

julia> code = EinCode((ixs..., [(i,) for i in LightGraphs.vertices(g)]...), ())
1∘2, 1∘4, 1∘5, 2∘3, 2∘6, 3∘5, 3∘6, 4∘5, 4∘6, 1, 2, 3, 4, 5, 6 ->

julia> optimized_code = optimize_greedy(code, Dict([i=>2 for i=1:6]))
4∘6, 4∘6 ->
├─ 4∘6, 6 -> 4∘6
│   └─ 6
│       └─ 4∘6
└─ 2∘5∘4, 5∘2∘6 -> 4∘6
    └─ 2∘5∘6, 6∘2 -> 5∘2∘6
        └─ 2∘6, 2 -> 6∘2
            └─ 2
                └─ 2∘6
            └─ 2∘3, 5∘6∘3 -> 2∘5∘6
                └─ 3∘5, 6∘3 -> 5∘6∘3
                    └─ 3∘6, 3 -> 6∘3
                        └─ ⋮
                            └─ 3∘5, 5 -> 3∘5
                                └─ ⋮
                                    └─ 2∘3
                    └─ 2∘4∘1, 1∘4∘5 -> 2∘5∘4
                        └─ 1∘5, 5∘4 -> 1∘4∘5
                            └─ 4∘5, 4 -> 5∘4
                                └─ 4
                                    └─ 4∘5
                                └─ 1∘5
                            └─ 2∘1, 1∘4 -> 2∘4∘1
                                └─ 1∘4
                                    └─ 1∘2, 1 -> 2∘1
                                        └─ 1
                                            └─ 1∘2

julia> function mis_contract(code::OMEinsum.NestedEinsum, x::T) where T
    tensors = map(OMEinsum.getixs(Iterators.flatten(code))) do ix
        @assert length(ix) == 1 || length(ix) == 2
        length(ix) == 1 ? [one(T), x] : [one(T) one(T); one(T) zero(T)]
    end
    code(tensors...)
end

mis_contract (generic function with 1 method)

julia> result = mis_contract(optimized_code, Tropical(1.0))
0-dimensional Array{TropicalF64, 0}:
2.0,
```

For larger graph, we hight recommend using two additional packages, `TropicalGEMM` for BLAS speed tropical matrix multiplication and `OMEinsumContractionOrders` for hyper optimized contraction order (the KaHyPar ([Schlag, 2020](#)) approach) ([Kourtis et al., 2019](#); [Gray & Kourtis, 2021](#); [Pan & Zhang, 2021](#)).