

COMPUTING PROPERTIES OF INDEPENDENT SETS BY GENERIC PROGRAMMING TENSOR NETWORKS *

XXX[†] AND YYY[‡]

Abstract. We introduce a method that uses generic programming tensor network to compute various properties of a class of hard graph problems that can be quantified locally. Let the independent set problem be an example, we show how to compute, for example, the maximum several independent set sizes, the counting of independent sets of a give size, the enumeration of independent sets of a given size. By making use of generic programming, our algorithms are very simple to implement and one can in addition directly utilize recent advances in tensor network contraction techniques such as near-optimal contraction order finding and slicing to achieve high performance. The algorithmic complexity of this approach is $2^{\text{tw}(G)}$, where $\text{tw}(G)$ is the treewidth of the problem graph. Other hard graph problems solvable in our framework includes the cutting problem, the K-coloring problem, the maximal clique problem, the dominating set problem, and satisfiability problem, among others. To demonstrate the versatility of this tool, we apply it to a few examples including the calculations of the entropy constant for some hardcore lattice gases on 2D square lattices and the study of the overlap gap property on three regular graphs. We show how our method helps people understand these hard problems better.

Key words. independent set, tensor network, maximum independent set, independence polynomial, generic programming

AMS subject classifications. 05C31, 14N07

[[ML: a comment](#)] and [[ML: a resolved comment](#)]

1. Introduction. In graph theory and combinatorial optimization, there is an important class of problems that can be quantified over a vertex and its neighborhood, including the independent set problem, the cutting problem, the K-coloring problem, the satisfiability problem, the maximal clique problem, and the vertex cover problem et al [43]. There is a wide-spread interest of these problem because they have wide range of applications in scheduling, logistics, wireless networks and telecommunication, and computer vision, etc. [13, 52]. [[JG: need more diversified applications not limited to IS](#)]

The most studied version of the above problems is finding a maximum/minimum set satisfying the constraints and the corresponding set size, while many of them belong to the hardness category NP-hard [30], i.e. it is unlikely to solve them in a polynomial time. In this paper, we focus on the less studied *property* computation which can be harder but is crucial towards understanding a hard problem. Here, the *property* refers to a class of quantities that not only include the maximum/minimum set size, but also include the number of sets at a give size, enumeration of all sets at a give size or direct sampling of such sets when they are not enumerable. For the weighted version of the above problems, it also includes finding largest/smallest k sets and their sizes.

We propose to view the computation of all these properties from an abstract algebra perspective and use a unified framework, generic tensor-network, to solve these seeming unrelated properties all. Tensor network is also known as the sum product network or einsum in different fields, where its tensor element types are typically restricted to standard number types such as real numbers and complex numbers. We extend this concept to generic tensor network by generalizing the tensor element type to any type being a commutative semiring. Given a good contraction order, the space required to contract a tensor network is 2 power the tree width of its line graph [42]. Recent progress in simulating quantum circuits with

*

Funding: ...

[†]XXX (email, website).

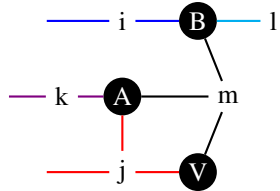
[‡]yyyyy (yyyy, email).

random tensor networks [28, 45, 34] enables us to find such a good contraction order for a tensor network with thousands of tensors in a reasonable time. When using tensor network for solving one of the best sets with the maximum size, this time complexity is just on par with the dynamic programming algorithms [16, 23], however, tensor network has richer algebraic structure, and this structure is crucial for property computation.

To avoid the discussion being unnecessarily diversified, we mainly focus on the independent set problem, and discuss other combinatorial optimization problems in Appendix C. The paper is organized as follows. We first introduce the basic concepts of tensor networks and generic programming in Sec. 2 and Sec. 3. Then we show how to reduce an independent set problem to a tensor network contraction problem Sec. 4 and show how to engineer the element types in Sec. 6, Sec. 7 and Sec. 8 to count independent sets at a given size, enumerate/sample independent sets at a given size and find largest set sizes in weighted graphs. Lastly, we benchmark our algorithms in Sec. 9 and provide a few examples in Sec. 10 to demonstrate the versatility of our tool. We show how our method can help computing the entropy constant for some hardcore lattice gases on 2D square lattices and analyzing the presence of overlap gap property on three regular graphs.

2. Tensor networks. Tensor network [15, 44] is also known as einsum, factor graph or sum-product network [10] in different contexts; it can be viewed as a generalization of binary matrix multiplication to n-ary tensor contraction. Einstein’s notation is often used to represent a tensor network, e.g. it represents the matrix multiplication between two matrices A and B as $C_{ik} = A_{ij}B_{jk}$, where we use a label in the subscript to represent a degree of freedom. One can enumerate these degrees of freedom and accumulate the product of tensor elements to the output tensor. In the standard Einstein’s notation for tensor networks in physics, each index appears precisely twice. Hence a tensor network can be represented as a simple graph, where a tensor is mapped to a vertex and a label shared by two tensors is mapped to an edge. In this work, we do not impose such a restriction, so an index can appear an arbitrary number of times. The graphical representation of a generalized tensor network is a hypergraph, in which an edge (a label in subscript) can be shared by an arbitrary number of vertices (tensors). These two notations are equivalent in representation power because one can easily translate a generalized tensor network to the standard notation by adding δ tensors, a high dimensional equivalence of the identity matrix. However, introducing δ tensors may significantly increase the complexity of the tensor contraction. We illustrate this subtle point in Appendix B.

Example 1. $C_{ijk} = A_{jkm}B_{mil}V_{jm}$ is a tensor network that can be evaluated as $C_{ijk} = \sum_{ml} A_{jkm}B_{mil}V_{jm}$. Its hypergraph representation is shown below, where we use different colors to represent different hyperedges.



3. Generic programming tensor contractions. In previous works relating tensor networks and combinatoric problems [35, 9], the elements in the tensor networks are limited to standard number types such as floating point numbers and integers. Owing to the development of modern compiling technology, we no longer need to re our imagination to standard number types. One of the key concepts that push the technology forward is called generic programming:

DEFINITION 3.1 (Generic programming [50]). *Generic programming is an approach to programming that focuses on designing algorithms and data structures so that they work in the most general setting without loss of efficiency.*

This definition of generic programming contains two major aspects: a single program works in the most general setting and efficiency. To understand the first aspect on generality, suppose we want to write a function that raises an element to a power, $f(x, n) := x^n$. One can easily write a function for standard number types that computes the power of x in $O(\log(n))$ steps using the multiply and square trick. Generic programming does not require x to be a standard number type, instead it treats x as an element with an associative multiplication operation \odot and a multiplicative identity $\mathbb{1}$. In such a way, when the program takes a matrix as an input, it computes the matrix power without extra efforts. The second aspect is about performance. For dynamically typed languages such as Python, one can easily write very general codes, but the efficiency is not guaranteed; for example, the speed of computing the matrix multiplication between two numpy arrays with python objects as elements is much slower than statically typed languages such as C++ and Julia [8]. C++ uses templates for generic programming while Julia takes advantage of just-in-time compilation and multiple dispatch. When these languages “see” a new input type, the compiler can recompile the generic program for the new type. A myriad of optimizations can be done during the compilation, such as inlining immutable elements with fixed sizes in an array to decrease the cache miss rate when accessing data. In Julia, these inlined arrays can even be compiled to GPU devices for faster computation [7].

This motivates us to think about what is the most general element type allowed in a tensor network contraction program. We find that as long as the algebra of tensor elements forms a commutative semiring, the tensor network contraction result will be valid and be independent of the contraction order. A commutative semiring is a ring with its multiplication operation being commutative and without an additive inverse. To define a commutative semiring with the addition operation \oplus and the multiplication operation \odot on a set R , the following relations must hold for any arbitrary three elements $a, b, c \in R$.

$$(a \oplus b) \oplus c = a \oplus (b \oplus c) \quad \triangleright \text{commutative monoid } \oplus \text{ with identity } \mathbb{0}$$

$$a \oplus \mathbb{0} = \mathbb{0} \oplus a = a$$

$$a \oplus b = b \oplus a$$

$$(a \odot b) \odot c = a \odot (b \odot c) \quad \triangleright \text{commutative monoid } \odot \text{ with identity } \mathbb{1}$$

$$a \odot \mathbb{1} = \mathbb{1} \odot a = a$$

$$a \odot b = b \odot a$$

$$a \odot (b \oplus c) = a \odot b \oplus a \odot c \quad \triangleright \text{left and right distributive}$$

$$(a \oplus b) \odot c = a \odot c \oplus b \odot c$$

$$a \odot \mathbb{0} = \mathbb{0} \odot a = \mathbb{0}$$

We require the algebra being multiplicative commutative because the contraction order optimizer will change the multiplication order of elements, while we wish this contraction order does not change the contraction results. In the following sections, we show how to compute a number of properties of independent sets by designing tensor element types as specific commutative semirings without changing the tensor network contraction

Element type	Property to compute
\mathbb{R}	Counting of all independent sets
Polynomial (Eq. (5.2: PN))	Independence polynomial
Tropical (Eq. (6.3: T))	Maximum independent set size
Extended tropical of order k (Eq. (8.1: Tk))	Largest k independent set sizes
Polynomial truncated to k -th order (Eq. (6.2: P1) and Eq. (6.5: P2))	k largest independent sizes and their counting
Set (Eq. (7.1: SN))	Enumeration of independent sets
Sum Product Tree (Eq. (7.6: SPT))	Sampling of independent sets
Polynomial truncated to highest order combined with Bit string (Eq. (7.5: S1))	Maximum independent set size and one of such configurations
Polynomial truncated to k -th order combined with Set (Eq. (7.3: P1+SN))	k largest independent set sizes and their enumeration

Table 1: Tensor element types and the independent set properties that can be computed using them.

program [50]. The Table 1 summarizes those properties that can be computed by various tensor element types.

4. Tensor network representation of independent sets. This section is about reducing the independent set problem to a tensor network contraction problem. For people with physics background, we suggest using an alternative approach described in Appendix A to understand this reduction from an energy model. Let $G = (V, E)$ be a weighted undirected graph with each vertex $i \in V$ being associated with a weight w_i . An independent set $I \subseteq V$ is a set of vertices that for any vertex pair $u, v \in I$, $(u, v) \notin E$, and we denote the maximum independent set (MIS) size $\alpha(G) \equiv \max_I \sum_{v \in I} w_v$. We map a vertex $v \in V$ to a label $s_v \in \{0, 1\}$ of dimension 2 in a tensor network, where we use 0 (1) to denote a vertex is absent (present) in the set. For each label s_v , we defined a parametrized rank-one vertex tensor $W(x_v, w_v)$ as

$$(4.1) \quad W(x_v, w_v) = \begin{pmatrix} 1_v \\ x_v^{w_v} \end{pmatrix}.$$

where we use subscripts to annotate different vertices, these annotation are only used in configuration enumeration and sampling. One can index tensor elements by subscripting tensors, e.g. $W(x_i, w_i)_0 = 1$ is the first element associated with $s_i = 0$ and $W(x_i, w_i)_1 = x_i^{w_i}$ is the second element associated with $s_i = 1$. Similarly, on each edge (u, v) , we define a matrix B indexed by s_u and s_v as

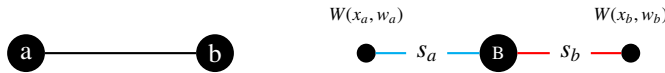
$$(4.2) \quad B = \begin{pmatrix} 1 & 1 \\ 1 & 0 \end{pmatrix}.$$

The corresponding tensor network contraction gives

$$(4.3) \quad P(G, \{x_1^{w_1}, x_2^{w_2}, \dots, x_{|V|}^{w_{|V|}}\}) = \sum_{s_1, s_2, \dots, s_{|V|}=0}^1 \prod_{i=1}^{|V|} W(x_i, w_i)_{s_i} \prod_{(i,j) \in E(G)} B_{s_i s_j},$$

where the summation runs over all vertex configurations $\{s_1, s_2, \dots, s_{|V|}\}$ and accumulates the product of tensor elements to the output P (see Example 3 for a concrete example). The edge tensor element $B_{s_i=1, s_j=1} = 0$ encodes the independent set constraint, meaning vertex i and j cannot be both in the independent set if they are connected by an edge (i, j) .

Example 2. In the following, we show a minimum example of mapping the independent problem of a 2 vertex complete graph K2 (left) to a tensor network (right).



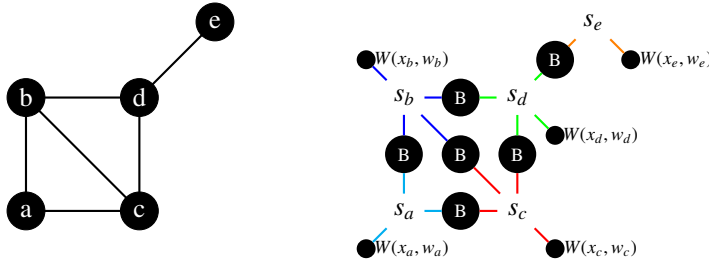
In the graphical representation of tensor network on the right panel, we use a circle to represent a tensor, a hyperedge in cyan color to represent the degree of freedom s_a , and a hyperedge in red color to represent a degree of freedom s_b . Tensors sharing a same degree of freedom are connected by that hyperedge. The contraction of this tensor network can be evaluated manually:

$$(4.4) \quad P(K2, \{x_a^{w_a}, x_b^{w_b}\}) = \begin{pmatrix} 1 & x_a^{w_a} \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ x_b^{w_b} \end{pmatrix} = 1 + x_a^{w_a} + x_b^{w_b}$$

The resulting polynomial represents 3 different independent sets $\{\}, \{a\}$ and $\{b\}$, with weights being 0, w_a and w_b respectively.

It is in general not computational efficient to evaluate Eq. (4.3) by directly summing up the $2^{|V|}$ products. The standard approach to evaluate a tensor network is to contract two tensors at a time with a certain order utilizing the associativity and commutativity of tensor contraction. A fully optimized contraction order can reduce the time to $O(2^{\text{tw}(G)})$ [42], while requiring a space at the same order of magnitude to store intermediate contraction results. The pairwise tensor contraction also makes it possible to utilize basic linear algebra subprograms (BLAS) functions to speed up the computation for certain tensor element types.

Example 3. In the following, we mapping a 5 vertex graph (left) to a tensor network (right) and show how the contraction order simplifies the computation.



The contraction of this tensor network can be done in a pairwise order utilizing the asso-

ciativity, additive commutativity, and multiplicative commutativity of tensor elements:

$$\begin{aligned}
& \sum_{s_a, s_b, s_c, s_d, s_e} W(x_a, w_a)_{s_a} W(x_b, w_b)_{s_b} W(x_c, w_c)_{s_c} W(x_d, w_d)_{s_d} W(x_e, w_e)_{s_e} B_{s_a s_b} B_{s_b s_d} B_{s_c s_d} B_{s_a s_c} B_{s_b s_c} B_{s_d s_e} \\
&= \sum_{s_b, s_c} \left(\sum_{s_d} \left(\left(\left(\sum_{s_e} B_{s_d s_e} W(x_e, w_e)_{s_e} \right) W(x_d, w_d)_{s_d} \right) (B_{s_b s_d} W(x_b, w_b)_{s_b}) \right) (B_{s_c s_d} W(x_c, w_c)_{s_c}) \right) \\
& \quad \left(B_{s_b s_c} \left(\sum_{s_a} B_{s_a s_b} (B_{s_a s_c} W(x_a, w_a)_{s_a}) \right) \right) \\
&= 1 + x_a^{w_a} + x_b^{w_b} + x_c^{w_c} + x_d^{w_d} + x_e^{w_e} + x_a^{w_a} x_d^{w_d} + x_a^{w_a} x_e^{w_e} + x_c^{w_c} x_e^{w_e} + x_b^{w_b} x_e^{w_e}
\end{aligned}$$

One can easily check the maximum tensor during computing is rank 2, and the total number of multiplication operations is reduced significantly.

5. Independence polynomial. Let $x_i = x$ and $w_i = 1$, the evaluation of Eq. (4.3) corresponds to the independence polynomial [29, 21], which is a useful graph characteristic related to, for example, the partition functions [36, 53] and Euler characteristics of the independence complex [11, 38]. The independence polynomial is an important graph polynomial that contains the counting information of independent sets. It is defined as

$$(5.1) \quad I(G, x) = \sum_{k=0}^{\alpha(G)} a_k x^k,$$

where a_k is the number of independent sets of size k in graph $G = (V, E)$. The total number of independent sets is thus equal to $I(G, 1)$. Its connection to Eq. (4.3) can be understood as follows: the product over vertex tensor elements produces a factor x^k , where $k = \sum_i s_i$ counts the set size, and the product over edge tensor elements gives a factor 1 for a configuration being in an independent set and 0 otherwise. The summation counts the number of independent sets of size k . One can feed the standard computer numbers types into this tensor network and evaluate the independence polynomial by contracting it numerically. However, we are more interested in finding the coefficients of this polynomial, because this quantity tells us the counting of independent sets at each size. To achieve this, let us first elevate the tensor elements 0s and 1s in tensors $W(x, 1)$ and B from integers and floating point numbers to the additive identity, $\mathbb{0}$, and multiplicative identity, $\mathbb{1}$, of a commutative semiring as discussed in Sec. 3. Let us create a polynomial type and represent a polynomial $a_0 + a_1 x + \dots + a_k x^k$ as a coefficient vector $a = (a_0, a_1, \dots, a_k) \in \mathbb{R}^k$, so, e.g., x is represented as $(0, 1)$. We define the algebra between the polynomials a of order k_a and b of order k_b as

$$\begin{aligned}
(5.2: \text{PN}) \quad & a \oplus b = (a_0 + b_0, a_1 + b_1, \dots, a_{\max(k_a, k_b)} + b_{\max(k_a, k_b)}), \\
& a \odot b = (a_0 + b_0, a_1 b_0 + a_0 b_1, a_2 b_0 + a_1 b_1 + a_0 b_2, \dots, a_{k_a} b_{k_b}), \\
& \mathbb{0} = (), \\
& \mathbb{1} = (1).
\end{aligned}$$

Here, the multiplication operation can be evaluated efficiently using the convolution theorem [49]. These operations are standard addition and multiplication operations of polynomials, and the polynomial type forms a commutative ring. The tensors W and B can thus be written as

$$(5.3) \quad W^{\text{PN}} = \begin{pmatrix} \mathbb{1} \\ (0, 1) \end{pmatrix}, \quad B^{\text{PN}} = \begin{pmatrix} \mathbb{1} & \mathbb{1} \\ \mathbb{1} & \mathbb{0} \end{pmatrix}.$$

By contracting the tensor network with the polynomial type, we have the exact representation of the independence polynomial. In practise, using the polynomial type suffers a space overhead proportional to $\alpha(G)$ because each polynomial requires a vector of such size to store the coefficients. In Appendix E, we provide an accurate and memory efficient method to find the polynomial by fitting it in a finite field. Alternatively, one can use the method in Appendix F to fit this polynomial more time efficiently with complex number at the cost of precision.

6. Maximum independent sets and its counting.

6.1. Tropical algebra for finding the MIS size and counting MISs. The MIS size $\alpha(G)$, or the independence number is an important graph characteristics. The method we use to compute this quantity is based on the following observations. Let $x = \infty$, the independence polynomial in the previous section becomes

$$(6.1) \quad I(G, \infty) = a_{\alpha(G)} \infty^{\alpha(G)},$$

where the lower-order terms vanish. We can thus replace the polynomial type $a = (a_0, a_1, \dots, a_k)$ with a new type with two fields: the largest exponent k and its coefficient a_k . From this, we can define a new algebra as

$$(6.2: P1) \quad \begin{aligned} a_x \infty^x \oplus a_y \infty^y &= \begin{cases} (a_x + a_y) \infty^{\max(x,y)}, & x = y \\ a_y \infty^{\max(x,y)}, & x < y \\ a_x \infty^{\max(x,y)}, & x > y \end{cases} \\ a_x \infty^x \odot a_y \infty^y &= a_x a_y \infty^{x+y} \\ \mathbb{0} &= 0 \infty^{-\infty} \\ \mathbb{1} &= 1 \infty^0. \end{aligned}$$

To implement this algebra programmatically, we create a data type with two fields (x, a_x) to store the MIS size and its counting, and define the above operations and constants correspondingly. If one is only interested in finding the MIS size, one can drop the counting field. The algebra of the exponents becomes the max-plus tropical algebra [40, 43]:

$$(6.3: T) \quad \begin{aligned} x \oplus y &= \max(x, y) \\ x \odot y &= x + y \\ \mathbb{0} &= -\infty \\ \mathbb{1} &= 0. \end{aligned}$$

Algebra Eq. (6.3: T) and Eq. (6.2: P1) are the same as those used in Liu et al. [39] to calculate and count spin glass ground states. For independent set calculations here, the vertex tensor and edge tensor becomes:

$$(6.4) \quad W^T = \begin{pmatrix} \mathbb{1} \\ \infty \end{pmatrix}, \quad B^T = \begin{pmatrix} \mathbb{1} & \mathbb{1} \\ \mathbb{1} & \mathbb{0} \end{pmatrix}.$$

6.2. Truncated polynomial algebra for counting independent sets of large size. Instead of counting just the MISs, one may also be interested in counting the independent sets with largest several sizes. For example, if one is interested in counting only $a_{\alpha(G)}$ and $a_{\alpha(G)-1}$, we can define a truncated polynomial algebra by keeping only the largest two coefficients in

the polynomial in Eq. (5.2: PN) as:

$$\begin{aligned}
(6.5: P2) \quad & a \oplus b = (a_{\max(k_a, k_b)-1} + b_{\max(k_a, k_b)-1}, a_{\max(k_a, k_b)} + b_{\max(k_a, k_b)}), \\
& a \odot b = (a_{k_a-1}b_{k_b} + a_{k_a}b_{k_b-1}, a_{k_a}b_{k_b}), \\
& \mathbb{0} = (), \\
& \mathbb{1} = (1).
\end{aligned}$$

In the program, we thus need a data structure that contains three fields, the largest order k , and the coefficients for the two largest orders a_k and a_{k-1} . This approach can clearly be extended to calculate more independence polynomial coefficients and is more efficient than calculating the entire independence polynomial. Similarly, one can also truncate the polynomial and keep only its smallest several orders, which can be used for e.g. counting the minimum maximal independent sets. As will be shown below, this algebra can also be extended to enumerate those large-size independent sets.

7. Enumerating and sampling independent sets.

7.1. Set algebra for configuration enumeration. The configuration enumeration of independent sets includes, for example, the enumeration of all independent sets, the enumeration of all MISs and the enumeration of independent sets with maximum several sizes. The most studied enumeration of maximal independent sets [12, 19, 33] requires using a different tensor network structure, which is discussed in Appendix C.1. To enumerate all independent sets, we designed an algebra defined on sets of bit strings:

$$\begin{aligned}
(7.1: SN) \quad & s \oplus t = s \cup t \\
& s \odot t = \{\sigma \vee^\circ \tau \mid \sigma \in s, \tau \in t\} \\
& \mathbb{0} = \{\} \\
& \mathbb{1} = \{0^{\otimes |V|}\}.
\end{aligned}$$

where s and t are each a set of $|V|$ -bit strings and \vee° is the bitwise OR operation over two bit strings.

Example 4. For elements being bit strings of length 5, we have the following set algebra

$$\begin{aligned}
& \{00001\} \oplus \{01110, 01000\} = \{01110, 01000\} \oplus \{00001\} = \{00001, 01110, 01000\} \\
& \{00001\} \oplus \{\} = \{00001\} \\
& \{00001\} \odot \{01110, 01000\} = \{01110, 01000\} \odot \{00001\} = \{01111, 01001\} \\
& \{00001\} \odot \{\} = \{\} \\
& \{00001\} \odot \{00000\} = \{00001\}.
\end{aligned}$$

To enumerate all independent sets, we initialize variable x_i in the vertex tensor to $x_i = \{e_i\}$, where e_i is a basis bit string of size $|V|$ that has only one non-zero value at location i . The vertex and edge tensors are thus

$$(7.2) \quad W^{\text{SN}}(\{e_i\}) = \begin{pmatrix} \mathbb{1} \\ \{e_i\} \end{pmatrix}, \quad B^{\text{SN}} = \begin{pmatrix} \mathbb{1} & \mathbb{1} \\ \mathbb{1} & \mathbb{0} \end{pmatrix}.$$

This set algebra can serve as the coefficients in Eq. (5.2: PN) to enumerate independent sets of all different sizes, Eq. (6.2: P1) to enumerate all MISs, or Eq. (6.5: P2) to enumerate

all independent sets of size $\alpha(G)$ and $\alpha(G) - 1$. As long as the coefficients in a truncated polynomial forms a commutative semiring, the polynomial itself is still a commutative semiring. For example, to enumerate only the MISs, with the tropical algebra, we can define $s_k \infty^k$, where the coefficients follow the algebra in Eq. (7.1: SN) and the orders (exponents) follow the max-plus tropical algebra. The combined operations become:

$$(7.3: \text{P1+SN}) \quad \begin{aligned} s_x \infty^x \oplus s_y \infty^y &= \begin{cases} (s_x \cup s_y) \infty^{\max(x,y)}, & x = y \\ s_y \infty^{\max(x,y)}, & x < y, \\ s_x \infty^{\max(x,y)}, & x > y \end{cases} \\ s_x \infty^x \odot s_y \infty^y &= \{\sigma \vee^\circ \tau \mid \sigma \in s_x, \tau \in s_y\} \infty^{x+y}, \\ \mathbb{0} &= \{\} \infty^{-\infty}, \\ \mathbb{1} &= \{0^{\otimes |V|}\} \infty^0. \end{aligned}$$

Clearly, the vertex tensor and edge tensor become

$$(7.4) \quad W^{\text{P1+SN}}(\{e_i \infty^1\}) = \begin{pmatrix} \mathbb{1} \\ \{e_i\} \infty^1 \end{pmatrix}, \quad B^{\text{P1+SN}} = \begin{pmatrix} \mathbb{1} & \mathbb{1} \\ \mathbb{1} & \mathbb{0} \end{pmatrix}.$$

The contraction of the corresponding tensor network yields an enumeration of all MIS configurations. Direct implementing this algebra for configuration MISs might incur significant space overheads for keeping too many intermediate states not contributing to the final maximum independent sets, one can avoid this issue by using the bounding technique in Appendix D.

If one is interested in obtaining only a single MIS configuration, one can just keep a single configuration in the intermediate computations to save the computational effort. Here is a new algebra defined on the bit strings, replacing the sets of bit strings in Eq. (7.1: SN),

$$(7.5: \text{S1}) \quad \begin{aligned} \sigma \oplus \tau &= \text{select}(\sigma, \tau), \\ \sigma \odot \tau &= (\sigma \vee^\circ \tau), \\ \mathbb{0} &= 1^{\otimes |V|}, \\ \mathbb{1} &= 0^{\otimes |V|}, \end{aligned}$$

where the `select` function picks one of σ and τ by some criteria to make the algebra commutative and associative, e.g. by picking the one with a smaller integer value.

7.2. Sampling extremely large configuration space. When the graph size goes large, the configurations might be impossible to fit into any type of storage. Then instead of using the set algebra, one can use a binary sum-product tree as a compact representation of the configurations. We use a quadruple (*type*, *data*, *left*, *right*) to represent a node in the tree, where *type* is one of LEAF, ZERO, SUM and PROD, *data* is a bit string as the content in a LEAF node, *left* and *right* are left and right operands of SUM and PROD nodes.

$$(7.6: \text{SPT}) \quad \begin{aligned} s \oplus t &= (\text{SUM}, /, s, t) \\ s \odot t &= (\text{PROD}, /, s, t) \\ \mathbb{0} &= (\text{ZERO}, /, /, /) \\ \mathbb{1} &= (\text{LEAF}, 0^{\otimes |V|}, /, /). \end{aligned}$$

The vertex tensor and edge tensor become

$$(7.7) \quad W^{\text{SPT}}((\text{LEAF}, e_i, /, /)) = \begin{pmatrix} \mathbb{1} \\ (\text{LEAF}, e_i, /, /) \end{pmatrix}, \quad B^{\text{SPT}} = \begin{pmatrix} \mathbb{1} & \mathbb{1} \\ \mathbb{1} & 0 \end{pmatrix}.$$

This algebra is a commutative semiring because we define the equivalence of two sum-product trees by comparing their expanded (using Eq. (7.1: SN)) forms, rather than their storages. Although one is probably unable to enumerate the configurations represented by this tree due to the extremely large configuration space size, but he can directly sample configurations from this tree. This is because once we have this sum-product representation of the configuration space, one can compute the number of configurations in each branch rigorously and efficiently. To generate samples, one just descends this sum product tree to left or right sibling with probability decided by the size of each branch while doing the computation specified by the first field.

8. Weighted graphs. In the previous discussion, we have limited ourselves to unweighted graphs. For integer weighted graphs, all results still holds, while for general weighted graphs, the independence polynomial is not computable anymore. For general weighted graphs, we are more interested to know the k most weighted sets and their sizes. Hence we define the extended tropical numbers as the following

$$(8.1: \text{Tk}) \quad \begin{aligned} s \oplus t &= \text{largest}(s \cup t, k) \\ s \odot t &= \text{largest}(\{a + b | a \in s, b \in t\}, k) \\ 0 &= -\infty^{\otimes k} \\ 1 &= -\infty^{\otimes k-1} \otimes 0 \end{aligned}$$

where $\text{largest}(s, k)$ means truncating the set by only keeping k largest values in the set s . The computation of $s \odot t$ is a maximum sum combination problem which can be done in time $\sim k \log(k)$. The vertex tensor and edge tensor become

$$(8.2) \quad W^{\text{Tk}}(-\infty^{\otimes k-1} \otimes 1, w_v) = \begin{pmatrix} \mathbb{1} \\ -\infty^{\otimes k-1} \otimes w_v \end{pmatrix}, \quad B^{\text{Tk}} = \begin{pmatrix} \mathbb{1} & \mathbb{1} \\ \mathbb{1} & 0 \end{pmatrix},$$

where we have used the relation $(-\infty^{\otimes k-1} \otimes 1)^{w_v} = -\infty^{\otimes k-1} \otimes w_v$. Similar to the tropical numbers, we can combine it with a bit string sampler (Eq. (7.5: S1)) for finding the configuration of each size. Since the \odot operation of configuration sampler is not used, the resulting configurations are deterministic and complete.

9. Performance benchmarks. We run a single thread benchmark on CPU Intel(R) Xeon(R) CPU E5-2686 v4 @ 2.30GHz, and its CUDA version on a GPU Tesla V100-SXM2 16G. The results are summarized in Figure 1. The graphs that we use in benchmarks are random three regular graphs, a typical type of sparse graphs that has a small treewidth that asymptotically smaller than $|V|/6$ [22]. In this benchmark, we do not include traditional algorithms for finding the MIS size like the branching algorithms [51, 48] and dynamic programming based algorithms [16, 23], because to the best of our knowledge, most of the properties mentioned in this paper are not computable within these frameworks. The main goal of this section is to show the relative timing for computing different properties.

Figure 1(a) shows the time and space complexity of tensor network contraction for different graph sizes, where the space complexity is the same as the treewidth of the problem graph. The contraction order is obtained using the local search algorithm in Ref. [34]. In

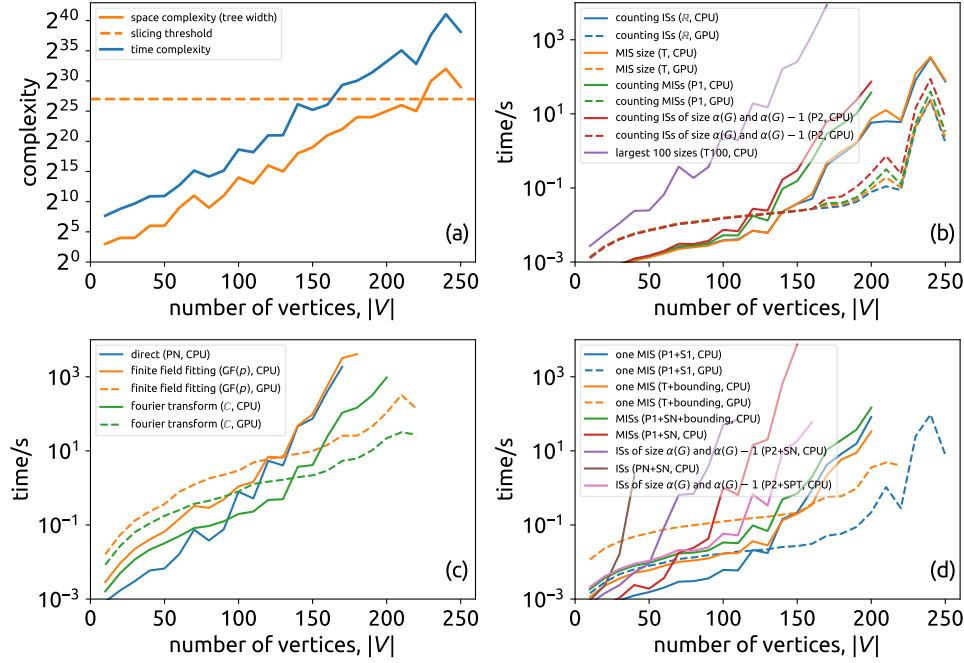


Figure 1: Benchmark results for computing different properties of independent sets of a random three regular graph with different tensor element types. The time in these plots only includes tensor network contraction, without taking into account the contraction order finding and just-in-time compilation time. Legends are properties, algebra and devices that we used in the computation; one can find the corresponding computed property in Table 1. (a) time and space complexity versus the number of vertices for the benchmarked graphs. (b) The computing time for calculating the MIS size and for counting the number of all independent sets (ISs), the number of MISs, the number of independent sets having size $\alpha(G)$ and $\alpha(G) - 1$, and finding 100 largest sets sizes. (c) The computing time for calculating the independence polynomials with different approaches. (d) The computing time for configuration enumeration, including the enumeration of all independent set configurations, a single MIS configuration, all MIS configurations, all independent set configurations having size $\alpha(G)$ and $\alpha(G) - 1$, and all independent set configurations having size $\alpha(G)$ and $\alpha(G) - 1$ in the sum product tree format, with or without bounding the enumeration space.

practice, slicing technique is used for graphs with treewidth greater than 27 to fit the computation into a 16GB memory. One can see that all the computing times in figure 1(b), (c) and (d) have a strong correlation with the treewidth. Among these benchmarks, computational tasks with data types T (CPU), \mathbb{R} (CPU), \mathbb{R} (GPU), \mathbb{C} (CPU), \mathbb{C} (GPU) and T+bounding (CPU) can utilize fast BLAS functions to evaluate tensor contractions, hence are much faster comparing to non-BLAS element types in the same category. GPU computes much faster than CPU in all cases when the problem scale is large enough such that the actual computing time is comparable or larger than the launching overhead of CUDA kernels. Most algebras can be computed on a GPU, except those requiring dynamic sized structures, i.e. PN, SN, T100 and SPT. In figure 1(c), one can see the Fourier transformation based method is the fastest in computing the independence polynomial, but it

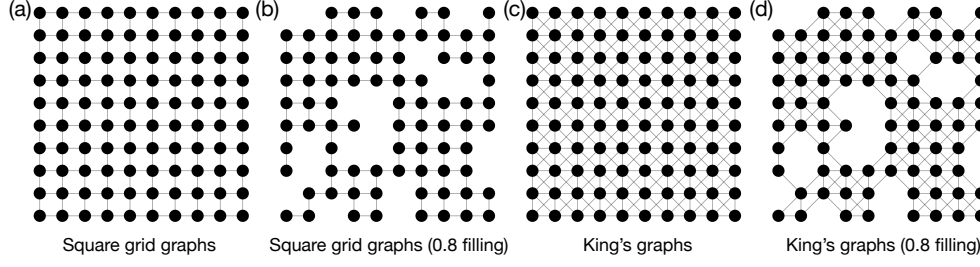


Figure 2: The types of graphs used in the case study in section 10.1. The lattice dimensions are $L \times L$. (a) Square grid graphs. (b) Square grid graphs with a filling factor $p = 0.8$. (c) King's graphs. (d) King's graphs with a filling factor $p = 0.8$.

may suffer from round-off errors. The finite field ($\text{GF}(p)$) approach is the only method that does not have round-off errors and can be run on a GPU. In figure 1(d), one can see the technique to bound the enumeration space improves the performance for more than one order of magnitude in enumerating the MISs. Bounding can also reduce the memory usage significantly, without which the largest computable graph size is only ~ 150 on a device with 32GB main memory.

10. Example case studies. In this section, we give a few examples where the different properties of independence sets are used.

10.1. Number of independent sets and entropy constant for some hardcore lattice gases. We compute the counting of all independent sets for graphs shown in Figure 2, where vertices are all placed on square lattices with lattice dimensions $L \times L$. The types of graphs include: the square grid graphs (Figure 2(a)); the square grid graphs with a filling factor $p = 0.8$, which means $\lfloor pL^2 \rfloor$ square grids are occupied with vertices (Figure 2(b)); the King's graphs (Figure 2(c)); the King's graphs with a filling factor $p = 0.8$ (Figure 2(d)). The number of independent sets for square grid graphs of size $L \times L$ form a well-known integer sequence (OEIS A006506), which is thought as a two-dimensional generalization of the Fibonacci numbers. We computed the integer sequence for $L = 38$ and $L = 39$, which, to the best of our knowledge, is not known before. In the computation, we used the finite-field algebra for contracting arbitrarily high precision integer tensor networks.

A useful number that can be computed using the number of independent sets is the entropy constant for the hardcore lattice gases on these graphs. For the square grid graphs, it is called the *hard square entropy constant* (OEIS A085850), which is defined as $\lim_{L \rightarrow \infty} F(L, L)^{1/L^2}$, where $F(L, L)$ is the number of independent sets of a given lattice dimensions $L \times L$. This quantity arises in statistical mechanics of hard-square lattice gases [6, 46] and is used to understand phase transitions for these systems. This entropy constant is not known to have an exact representation, but it is accurately known in many digits. Similarly, we can define entropy constants for other lattice gases. In Fig. 3, we look at how $F(L, L)^{1/L^2}$ scales as a function of the grid size L for all types of graphs shown in Figure 2. Our results match the known results for the non-disordered square grid and King's graphs. For disordered square grid and King's graphs with a filling factor $p = 0.8$, we randomly sample 1000 graph instances. To our knowledge, the entropy constants for these disordered graphs have not been studied before. Interestingly, the variations due to different random instances are negligible for this entropy quantity.

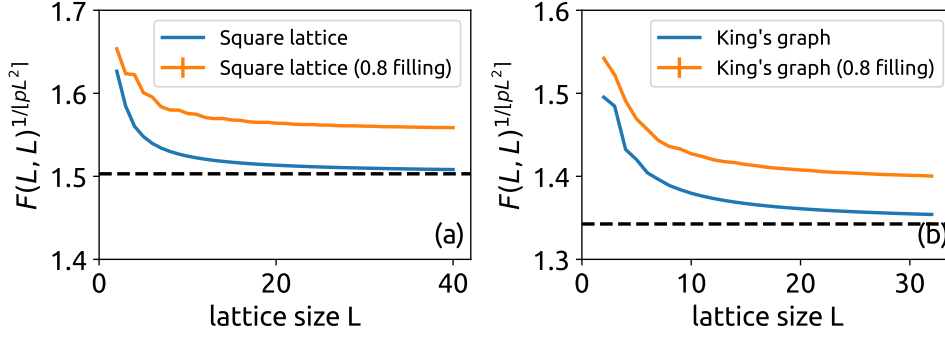


Figure 3: Mean entropy for lattice gases on graphs defined in Fig. 2. We sampled 1000 instances for $p = 0.8$ lattices and the error bar is too small to be visible. The horizontal black dashed lines are for $\lim_{L \rightarrow \infty} F(L, L)^{1/L^2}$ for the corresponding non-disordered square grid and King’s graphs.

10.2. The overlap gap property. With the tool to enumerate configurations, one can try to understand the structure of the independent set configuration space, such as the optimization landscape for finding the MISs and the geometry of the solution space. One of the known barriers for finding the MIS is the so-called overlap gap property [25, 24]. Intuitively, if the overlap gap property is present, it means every two large independent sets either have a significant intersection or very small intersection; it implies that large independent sets are clustered together. This clustering property has been used to prove the limitations of local algorithms in finding the MISs [25, 24]. Here, we use the computed configurations of large independent sets to investigate the presence or absence of the overlap gap property for some given graphs. In particular, in Fig. 4, we enumerated all ISs of size $\geq \alpha(G) - 1$ for two random three regular graph instances of size 100 and show the pairwise Hamming distance statistics for the enumerated configurations. In Fig. 4(a), we observed the multiple peak structure for the pairwise Hamming distance for one of the graphs. This indicates the presence of the overlap gap property and disconnected clusters exist in the configuration space for large independent sets. We expect our numerical tool can be used to understand this phenomenon better and to further investigate the graph properties and the geometry of the configuration spaces for a variety of graph instances.

11. Discussion and conclusion. In this paper, we introduce a method to use generic programming tensor networks to compute a number of different properties of a hard problem, including the largest several independent sets sizes, the counting of independent sets of a given size and the enumeration of independent sets of a given size. For each property, we design an algebra being a commutative semiring and let it be the tensor element type. With different algebras, the contraction of the same tensor network gives us the different properties. The data types introduced in the main text is summarized in the diagram in Fig. 5. Other than the case with real numbers, we have polynomials and truncated polynomials. We combine them with the bit string algebra, the set algebra, and sum product tree algebra for finding, enumeration, and sampling of independent sets at a given size. These algebras are of generic use: they can be utilized to compute properties of maximal independent sets and a variety of other combinatorial problems such as the matching problem, the k-coloring problem, the max-cut problem, and the set packing problem, as

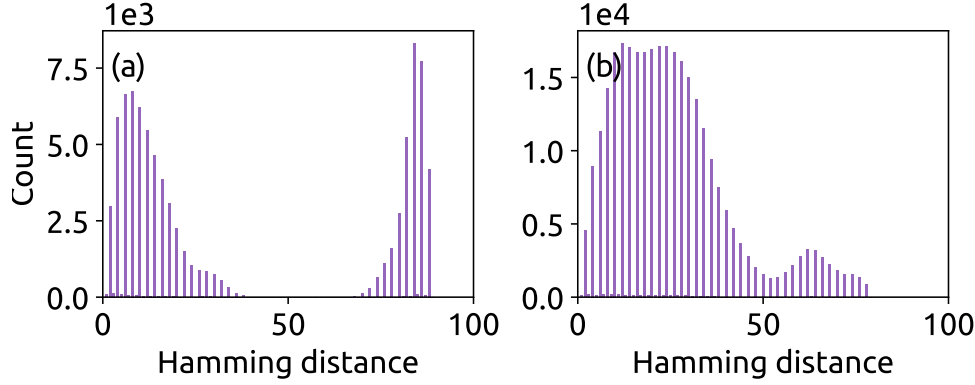


Figure 4: The distribution of pairwise Hamming distances between every two independent sets of size $\geq \alpha(G) - 1$ for two random three regular graph instances of size 100.

detailed in Appendix C. Moreover, since the independence polynomial is closely related to the matching polynomial [37], the clique polynomial [31], and the vertex cover polynomial [4], our algorithm to compute the independence polynomial can also be used to compute these graph polynomials. By adapting to the style of generic programming, we can implement all the algorithms without much effort. We show some of the Julia language implementations in Appendix H. A complete implementation can be found in our Github repository [1]. We expect our tool can be used to understand and study many interesting applications of independent sets and beyond. We also hope this approach of generic tensor networks can inspire future works on tensor network computations.

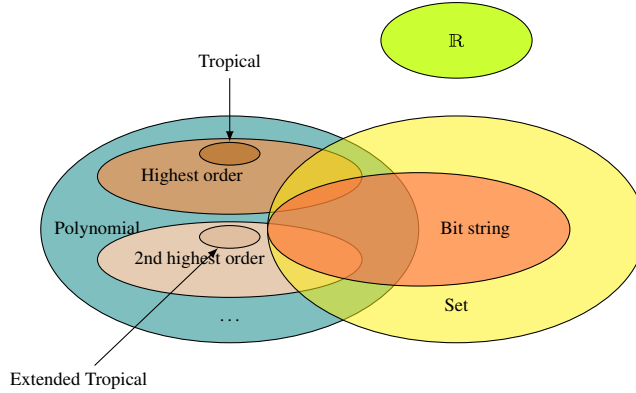


Figure 5: The the tensor network element types used in this work, while the purpose of these element types can be found in Table 1. The overlap between two algebra indicates that a new algebra can be created by combining those two types of algebra.

Acknowledgments. We would like to thank Pan Zhang for sharing his python code for optimizing contraction orders of a tensor network. We acknowledge Sepehr Ebadi, Maddie Cain, and Leo Zhou for coming up with many interesting questions about independent sets

and their questions strongly motivated the development of this project. We thank Chris Elord for helping us to write the fastest matrix multiplication library for tropical numbers, TropicalGEMM.jl. We would also like to thank a number of open-source software developers, including Roger Luo, Time Besard, and Katharine Hyatt, for actively maintaining their packages and resolving related issues voluntarily. [JG: funding information]

REFERENCES

- [1] <https://github.com/Happy-Diode/GraphTensorNetworks.jl>.
- [2] <https://github.com/JuliaGraphs/Graphs.jl>.
- [3] <https://github.com/TensorBFS/TropicalGEMM.jl>.
- [4] S. AKBARI AND M. R. OBOUDI, *On the edge cover polynomial of a graph*, European Journal of Combinatorics, 34 (2013), pp. 297–321.
- [5] S. ALIKHANI AND Y. HOCK PENG, *Introduction to domination polynomial of a graph*, 2009, <https://arxiv.org/abs/0905.2251>.
- [6] R. J. BAXTER, I. G. ENTING, AND S. K. TSANG, *Hard-square lattice gas*, Journal of Statistical Physics, 22 (1980), pp. 465–489, <https://doi.org/10.1007/BF01012867>, <https://doi.org/10.1007/BF01012867>.
- [7] T. BESARD, C. FOKET, AND B. D. SUTTER, *Effective extensible programming: Unleashing julia on gpus*, CoRR, abs/1712.03112 (2017), <http://arxiv.org/abs/1712.03112>, <https://arxiv.org/abs/1712.03112>.
- [8] J. BEZANSON, S. KARPINSKI, V. B. SHAH, AND A. EDELMAN, *Julia: A fast dynamic language for technical computing*, 2012, <https://arxiv.org/abs/1209.5145>, <https://arxiv.org/abs/1209.5145>.
- [9] J. BIAMONTE AND V. BERGHOLM, *Tensor networks in a nutshell*, 2017, <https://arxiv.org/abs/1708.00006>.
- [10] C. M. BISHOP, *Pattern Recognition and Machine Learning*, Springer, 2006.
- [11] M. BOUSQUET-MÉLOU, S. LINUSSON, AND E. NEVO, *On the independence complex of square grids*, Journal of Algebraic combinatorics, 27 (2008), pp. 423–450.
- [12] C. BRON AND J. KERBOSCH, *Algorithm 457: finding all cliques of an undirected graph*, Communications of the ACM, 16 (1973), pp. 575–577.
- [13] S. BUTENKO AND P. M. PARDALOS, *Maximum Independent Set and Related Problems, with Applications*, PhD thesis, USA, 2003. AAI3120100.
- [14] P. BUTERA AND M. PERNICI, *Sums of permanent minors using grassmann algebra*, 2014, <https://arxiv.org/abs/1406.5337>.
- [15] I. CIRAC, D. PEREZ-GARCIA, N. SCHUCH, AND F. VERSTRAETE, *Matrix Product States and Projected Entangled Pair States: Concepts, Symmetries, and Theorems*, arXiv e-prints, (2020), arXiv:2011.12127, p. arXiv:2011.12127, <https://arxiv.org/abs/2011.12127>.
- [16] B. COURCELLE, *The monadic second-order logic of graphs. i. recognizable sets of finite graphs*, Information and computation, 85 (1990), pp. 12–75.
- [17] J. C. DYRE, *Simple liquids’ quasiuniversality and the hard-sphere paradigm*, Journal of Physics: Condensed Matter, 28 (2016), p. 323001.
- [18] S. EBADI, A. KEESLING, M. CAIN, T. T. WANG, H. LEVINE, D. BLUVSTEIN, G. SEMEGHINI, A. OMRAN, J. LIU, R. SAMAJDAR, X.-Z. LUO, B. NASH, X. GAO, B. BARAK, E. FARHI, S. SACHDEV, N. GEMELKE, L. ZHOU, S. CHOI, H. PICHLER, S. WANG, M. GREINER, V. VULETIC, AND M. D. LUKIN, *Quantum optimization of maximum independent set using rydberg atom arrays*, 2022, <https://arxiv.org/abs/2202.09372>.
- [19] D. EPPSTEIN, M. LÖFFLER, AND D. STRASH, *Listing all maximal cliques in sparse graphs in near-optimal time*, in Algorithms and Computation, O. Cheong, K.-Y. Chwa, and K. Park, eds., Berlin, Heidelberg, 2010, Springer Berlin Heidelberg, pp. 403–414.
- [20] H. C. M. FERNANDES, J. J. ARENZON, AND Y. LEVIN, *Monte carlo simulations of two-dimensional hard core lattice gases*, The Journal of Chemical Physics, 126 (2007), p. 114508, <https://doi.org/10.1063/1.2539141>, <https://doi.org/10.1063/1.2539141>, <https://arxiv.org/abs/https://doi.org/10.1063/1.2539141>.
- [21] G. M. FERRIN, *Independence polynomials*, (2014).
- [22] F. V. FOMIN AND K. HØIE, *Pathwidth of cubic graphs and exact algorithms*, Information Processing Letters, 97 (2006), pp. 191–196.
- [23] F. V. FOMIN AND P. KASKI, *Exact exponential algorithms*, Communications of the ACM, 56 (2013), pp. 80–88.
- [24] D. GAMARNIK AND A. JAGANNATH, *The overlap gap property and approximate message passing algorithms for p-spin models*, 2019, <https://arxiv.org/abs/1911.06943>.
- [25] D. GAMARNIK AND M. SUDAN, *Limits of local algorithms over sparse random graphs*, 2013, <https://arxiv.org/abs/1304.1831>.
- [26] S. GASPERS, D. KRATSCHE, AND M. LIEDLOFF, *On independent sets and bicliques in graphs*, Algorithmica, 62 (2012), pp. 637–658, <https://doi.org/10.1007/s00453-010-9474-1>, <https://doi.org/10.1007/s00453-010-9474-1>.

- [27] G. H. GOLUB AND C. F. VAN LOAN, *Matrix computations*, vol. 3, JHU press, 2013.
- [28] J. GRAY AND S. KOURTIS, *Hyper-optimized tensor network contraction*, *Quantum*, 5 (2021), p. 410, <https://doi.org/10.22331/q-2021-03-15-410>, <http://dx.doi.org/10.22331/q-2021-03-15-410>.
- [29] N. J. HARVEY, P. SRIVASTAVA, AND J. VONDRÁK, *Computing the independence polynomial: from the tree threshold down to the roots*, in *Proceedings of the Twenty-Ninth Annual ACM-SIAM Symposium on Discrete Algorithms*, SIAM, 2018, pp. 1557–1576.
- [30] J. HASTAD, *Clique is hard to approximate within $n^{1-\epsilon}$* , in *Proceedings of 37th Conference on Foundations of Computer Science*, IEEE, 1996, pp. 627–636.
- [31] C. HOEDE AND X. LI, *Clique polynomials and independent set polynomials of graphs*, *Discrete Mathematics*, 125 (1994), pp. 219–228.
- [32] H. HU, T. MANSOUR, AND C. SONG, *On the maximal independence polynomial of certain graph configurations*, *Rocky Mountain Journal of Mathematics*, 47 (2017), pp. 2219 – 2253, <https://doi.org/10.1216/RMJ-2017-47-7-2219>, <https://doi.org/10.1216/RMJ-2017-47-7-2219>.
- [33] D. S. JOHNSON, M. YANNAKAKIS, AND C. H. PAPADIMITRIOU, *On generating all maximal independent sets*, *Information Processing Letters*, 27 (1988), pp. 119–123, [https://doi.org/https://doi.org/10.1016/0020-0190\(88\)90065-8](https://doi.org/https://doi.org/10.1016/0020-0190(88)90065-8), <https://www.sciencedirect.com/science/article/pii/0020019088900658>.
- [34] G. KALACHEV, P. PANTELEEV, AND M.-H. YUNG, *Recursive multi-tensor contraction for xeb verification of quantum circuits*, 2021, <https://arxiv.org/abs/2108.05665>.
- [35] S. KOURTIS, C. CHAMON, E. MUCCILO, AND A. RUCKENSTEIN, *Fast counting with tensor networks*, *SciPost Physics*, 7 (2019), <https://doi.org/10.21468/scipostphys.7.5.060>, <http://dx.doi.org/10.21468/SciPostPhys.7.5.060>.
- [36] T.-D. LEE AND C.-N. YANG, *Statistical theory of equations of state and phase transitions. ii. lattice gas and ising model*, *Physical Review*, 87 (1952), p. 410.
- [37] V. E. LEVIT AND E. MANDRESCU, *The independence polynomial of a graph-a survey*, in *Proceedings of the 1st International Conference on Algebraic Informatics*, vol. 233254, Aristotle Univ. Thessaloniki Thessaloniki, 2005, pp. 231–252.
- [38] V. E. LEVIT AND E. MANDRESCU, *The independence polynomial of a graph at -1*, 2009, <https://arxiv.org/abs/0904.4819>.
- [39] J.-G. LIU, L. WANG, AND P. ZHANG, *Tropical tensor network for ground states of spin glasses*, *Physical Review Letters*, 126 (2021), <https://doi.org/10.1103/physrevlett.126.090506>, <http://dx.doi.org/10.1103/PhysRevLett.126.090506>.
- [40] D. MACLAGAN AND B. STURMFELS, *Introduction to tropical geometry*, vol. 161, American Mathematical Soc., 2015, <http://www.cs.technion.ac.il/~janos/COURSES/238900-13/Tropical/MaclaganSturmfels.pdf>.
- [41] F. MANNE AND S. SHARMIN, *Efficient counting of maximal independent sets in sparse graphs*, in *International Symposium on Experimental Algorithms*, Springer, 2013, pp. 103–114.
- [42] I. L. MARKOV AND Y. SHI, *Simulating quantum computation by contracting tensor networks*, *SIAM Journal on Computing*, 38 (2008), p. 963–981, <https://doi.org/10.1137/050644756>, <http://dx.doi.org/10.1137/050644756>.
- [43] C. MOORE AND S. MERTENS, *The nature of computation*, OUP Oxford, 2011.
- [44] R. ORÚS, *A practical introduction to tensor networks: Matrix product states and projected entangled pair states*, *Annals of Physics*, 349 (2014), pp. 117–158.
- [45] F. PAN AND P. ZHANG, *Simulating the sycamore quantum supremacy circuits*, 2021, <https://arxiv.org/abs/2103.03074>.
- [46] P. A. PEARCE AND K. A. SEATON, *A classical theory of hard squares*, *Journal of Statistical Physics*, 53 (1988), pp. 1061–1072, <https://doi.org/10.1007/BF01023857>, <https://doi.org/10.1007/BF01023857>.
- [47] H. PICHLER, S.-T. WANG, L. ZHOU, S. CHOI, AND M. D. LUKIN, *Computational complexity of the rydberg blockade in two dimensions*, *arXiv preprint arXiv:1809.04954*, (2018).
- [48] J. M. ROBSON, *Algorithms for maximum independent sets*, *Journal of Algorithms*, 7 (1986), pp. 425–440.
- [49] A. SCHÖNHAGE AND V. STRASSEN, *Schnelle multiplikation grosser zahlen*, *Computing*, 7 (1971), pp. 281–292.
- [50] A. A. STEPANOV AND D. E. ROSE, *From mathematics to generic programming*, Pearson Education, 2014.
- [51] R. E. TARJAN AND A. E. TROJANOWSKI, *Finding a maximum independent set*, *SIAM Journal on Computing*, 6 (1977), pp. 537–546.
- [52] Q. WU AND J.-K. HAO, *A review on algorithms for maximum clique problems*, *European Journal of Operational Research*, 242 (2015), pp. 693–709, <https://doi.org/https://doi.org/10.1016/j.ejor.2014.09.064>, <https://www.sciencedirect.com/science/article/pii/S0377221714008030>.
- [53] C.-N. YANG AND T.-D. LEE, *Statistical theory of equations of state and phase transitions. i. theory of condensation*, *Physical Review*, 87 (1952), p. 404.

Appendix A. An alternative way to construct the tensor network.

Let us characterize the independent set problem on graph $G = (V, E)$ as an energy model

$$(A.1) \quad \mathcal{E}(G, s) = - \sum_{i \in V(G)} w_i s_i + \infty \sum_{\langle i, j \rangle \in E(G)} s_i s_j$$

where s_i is a spin on vertex $i \in V$ and w_i is an onsite energy term associated with it. The first term is the energy that corresponds to the inverse of the independent set size and the second term describes the independence constraint. In physical systems, this constraint term usually corresponds to the Rydberg blockade [47, 18] in cold atom arrays or the repulsive force in hard core lattice model [17, 20]. Its partition function is defined as

$$(A.2) \quad \begin{aligned} Z(G, \beta) &= \sum_s e^{-\beta \mathcal{E}(G, s)} = \sum_{s \in \mathcal{I}(G)} e^{\beta \sum w_i s_i} \\ &= \sum_{k=0}^{\alpha(G)} a(k) e^{\beta k} \quad (k = \sum w_i s_i) \end{aligned}$$

where $\mathcal{I}(G)$ is the set of independent sets of graph G . $\alpha(G)$ is the absolute value of the minimum energy (maximum independent set size). $a(k)$ is the number of spin configurations with energy $-k$ (independent sets of size k). The partition function can be evaluated as a tensor network contraction, where the tensor network is constructed by placing a vertex tensor on each spin i

$$(A.3) \quad W(\beta, w_i) = \begin{pmatrix} 1 \\ e^{\beta w_i} \end{pmatrix},$$

and an edge tensor on each bond

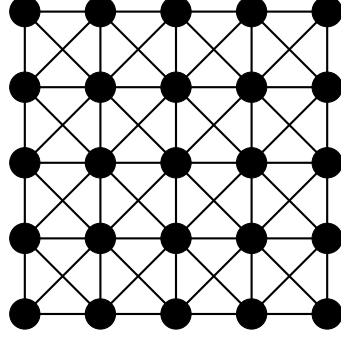
$$(A.4) \quad B = \begin{pmatrix} 1 & 1 \\ 1 & 0 \end{pmatrix},$$

where the 0 in the bond tensor coming from $e^{-\beta \infty}$ in the second term of Eq. (A.1). By letting $x = e^\beta$, we get the tensor network for computing the independence polynomial as described by Eq. (4.1) and Eq. (4.2), and by letting $w_i = 1$, the second line of Eq. (A.2) is the independence polynomial.

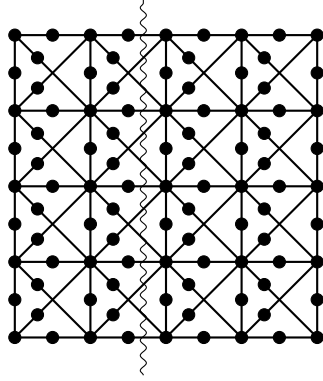
Appendix B. An example of contraction complexity increase due to introducing δ tensors. As we have mentioned in the main text, a standard tensor network notation is equivalent to the generalized tensor network by introducing δ tensors, where a δ tensor of rank d is defined as

$$(B.1) \quad \delta_{i_1, i_2, \dots, i_d} = \begin{cases} 1, & i_1 = i_2 = \dots = i_d, \\ 0, & \text{otherwise.} \end{cases}$$

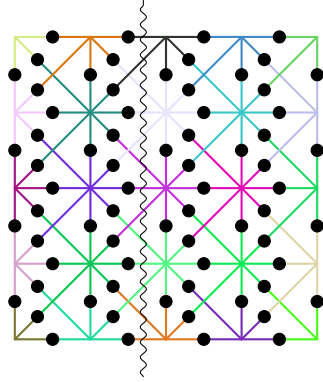
In the following, we are going to show introducing δ tensors might increase the contraction complexity of a tensor network. Let us consider the following King's graph.



By mapping the independent set problem to a standard tensor network with δ tensors, we have the following graphical representation.



In this diagram, the circle on each vertex in the original graph is a δ tensor of rank 8 or less. If we contract this tensor network in a naive column-wise order, the maximum intermediate tensor has rank $\sim 3L$, requiring a storage of size $\approx 2^{3L}$. If we relax the restriction that each label must appear exactly twice, we have the following hypergraph representation of a generalized tensor network.



Here, we use different colors to distinguish different hyperedges. A vertex tensor always has rank 1 and is not shown here since it does not change the contraction complexity. If we contract this tensor network in the column-wise order, the maximum intermediate tensor rank is $\sim L$, which can be seen by counting the number of colors at the cut.

Appendix C. Hard problems and tensor networks.

C.1. Maximal independent sets and maximal cliques. Finding maximal independent sets of a graph is equivalent to finding the maximal cliques of its complement graph, so in

the following we mainly discuss how to find maximal independent sets. Let us denote the neighborhood of a vertex v as $N(v)$ and denote $N[v] = N(v) \cup \{v\}$. A maximal independent set I_m is an independent set where there exists no vertex $v \in V$ such that $I_m \cap N[v] = \emptyset$. Similar to the independence polynomial, the maximal independence polynomial counts the number of maximal independent sets of various sizes [32], which can be helpful for understanding why the program for solving MIS can be trapped in a local minimum. Concretely, it is defined as

$$(C.1) \quad I_{\max}(G, x) = \sum_{k=0}^{\alpha(G)} b_k x^k,$$

where b_k is the number of maximal independent sets of size k in graph $G = (V, E)$. Comparing with the independence polynomial in Eq. (5.1), we have $b_k \leq a_k$ and $b_{\alpha(G)} = a_{\alpha(G)}$. $I_{\max}(G, 1)$ counts the total number of maximal independent sets [26, 41], where, to our knowledge, the fastest algorithm currently has a runtime of $O(1.3642^{|V|})$ [26]. If we want to find an MIS, b_k counts the number of local optimum at size $k < \alpha(G)$, and can, in some cases, provide hints on the difficulty of finding the MIS using local algorithms [18]. The uni-modality, log-concavity, and real-rootness properties of the maximal independence polynomial for special classes of graphs have also been studied [32].

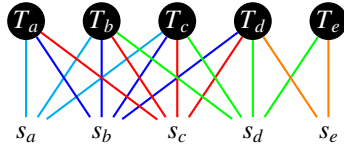
We can modify the tensor network for computing the independence polynomial to include this restriction. Instead of defining the restriction on vertices and edges, it is more natural to define it on $N[v]$:

$$(C.2) \quad T(x_v, w_v)_{s_1, s_2, \dots, s_{|N(v)|}, s_v} = \begin{cases} s_v x_v^{w_v} & s_1 = s_2 = \dots = s_{|N(v)|} = 0, \\ 1 - s_v & \text{otherwise.} \end{cases}$$

Intuitively, it means if all the neighbourhood vertices are not in I_m , i.e., $s_1 = s_2 = \dots = s_{|N(v)|} = 0$, then v should be in I_m and contribute a factor x_v , otherwise, if any of the neighbourhood vertices is in I_m , then v cannot be in I_m . As an example, for a vertex of degree 2, the resulting rank-3 tensor is

$$(C.3) \quad T(x_v, w_v) = \begin{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 1 \end{pmatrix} \\ x_v^{w_v} \\ 0 \end{pmatrix}.$$

By contracting this tensor network with generic element types, we can compute the maximal independent set properties such as the maximal independence polynomial and the enumeration of maximal independent sets. Let us consider the example in Sec. 3: its corresponding tensor network structure for computing the maximal independent polynomial becomes



One can see that the average degree of a tensor is increased. The computational complexity of this new tensor network contraction is often greater than the one for computing the independence polynomial. However, for most sparse graphs, this tensor network contraction approach is still much faster than enumerating all the maximal cliques on its complement graph using the Bron-Kerbosch algorithm [12], which is the standard algorithm that we are aware of to compute the maximal independence polynomial. We show

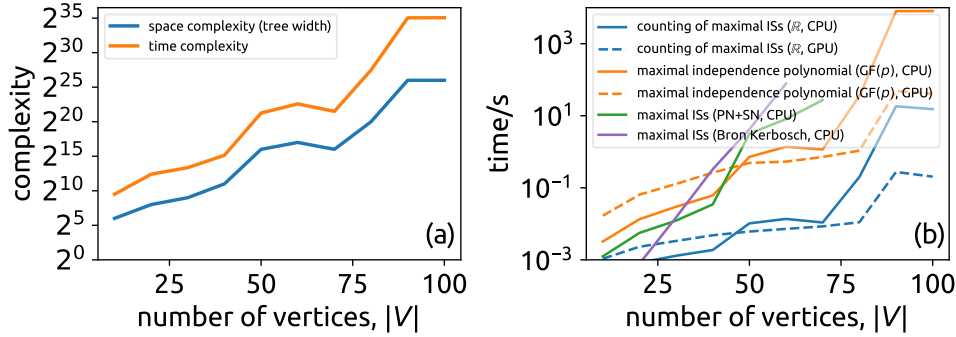


Figure 6: Benchmark results for computing different properties of maximal independent sets on a random three regular graph with different tensor element types. (a) treewidth versus the number of vertices for the benchmarked graphs. (b) The computing time for calculating the number of independent sets and enumerate all MISs.

the benchmark of computing the maximal independent set properties in Fig. 6, including a comparison to the Bron-Kerbosch algorithm from Julia package Graphs [2]. The treewidth of this tensor network is significantly larger, so benchmark with only a smaller graph size is feasible. The time for the tensor network approach and the Bron-Kerbosch approach to enumerate all maximal independent sets are comparable, while the tensor network does counting much more efficiently. Due to the memory limit, this Bron-Kerbosch algorithm is only feasible up to a graph size around 70.

C.2. Matching problem. A matching polynomial of a graph G is defined as

$$(C.4) \quad M(G, x) = \sum_{k=1}^{|V|/2} c_k x^k,$$

where k is the size of a matching, and coefficients c_k are the corresponding counting of the number of matchings of the size k . We map an edge $(u, v) \in E$ to a label $\langle u, v \rangle \in \{0, 1\}$ in a tensor network, where 1 means two vertices of an edge are matched and 0 otherwise. Then we define a tensor of rank $d(v) = |N(v)|$ on vertex v such that,

$$(C.5) \quad W_{\langle v, n_1 \rangle, \langle v, n_2 \rangle, \dots, \langle v, n_{d(v)} \rangle} = \begin{cases} 1, & \sum_{i=1}^{d(v)} \langle v, n_i \rangle \leq 1, \\ 0, & \text{otherwise,} \end{cases}$$

and a tensor of rank 1 on the bond

$$(C.6) \quad B(w_{\langle u, v \rangle})_{\langle u, v \rangle} = \begin{cases} 1, & \langle u, v \rangle = 0 \\ x_{\langle u, v \rangle}^{w_{\langle u, v \rangle}}, & \langle u, v \rangle = 1, \end{cases}$$

where label $\langle v, u \rangle$ is equivalent to $\langle u, v \rangle$. Here, a vertex tensor specifies the restriction that a vertex cannot be matched by more than one edge, while an edge tensor contributes the variable in the polynomial.

C.3. k -Coloring. Let us use the 3-coloring problem defined on vertices as an example. For a vertex v , we define the degrees of freedom $c_v \in \{0, 1, 2\}$ and a vertex tensor labelled by

it as

$$(C.7) \quad W(v) = \begin{pmatrix} r_v \\ g_v \\ b_v \end{pmatrix},$$

where r_v , g_v and b_v are colors used to label the configurations. For an edge (u, v) , we define an edge tensor as a matrix labelled by (c_u, c_v) to specify the constraint

$$(C.8) \quad B(x, w_{uv}) = \begin{pmatrix} 0 & x^{w_{uv}} & x^{w_{uv}} \\ x^{w_{uv}} & 0 & x^{w_{uv}} \\ x^{w_{uv}} & x^{w_{uv}} & 0 \end{pmatrix},$$

where weights are associated with edges. The number of possible colorings can be obtained by contracting this tensor network by setting r_v , g_v and b_v to 1. By designing generic types as tensor elements, one can get other properties. Similarly, one can also define the k -coloring problem on edges by switching the roles of edges and vertices. Let $x^{w_{uv}} = x$, we have a graph polynomial, where its k th coefficient is the number of coloring with k bonds are satisfied. If a graph is colorable, the maximum order of this polynomial should be equal to the number of edges in this graph.

C.4. Max-Cut problem. The Max-Cut problem is also known as the boolean spin glass problem. For a vertex $v \in V$, we define a boolean degree of freedom $s_v \in \{0, 1\}$. Then the Max-Cut problem can be encoded to tensor networks by mapping an edge $(i, j) \in E$ to an edge matrix labelled by $s_i s_j$

$$(C.9) \quad B(x, w_{(i,j)}) = \begin{pmatrix} 1 & x_j^{w_{(i,j)}} \\ x_i^{w_{(i,j)}} & 1 \end{pmatrix},$$

where variables x_i and x_j represents a cut on edge (i, j) or a domain wall of an Ising spin glass. Let $x_i^{w_{(i,j)}} = x_j^{w_{(i,j)}} = x$, we have a graph polynomial similar to previous ones, where its k th coefficient is two times the number of configurations of cut size k .

C.5. Dominating Set. Let $G = (V, E)$ be the target graph that we want to solve. The tensor network representation map a vertex $v \in V$ to a boolean degree of freedom $s_v \in \{0, 1\}$. We defined the restriction on a vertex and its neighbouring vertices $N(v)$:

$$(C.10) \quad T(x, w_v)_{s_1, s_2, \dots, s_{|N(v)|}, s_v} = \begin{cases} 0_v & s_1 = s_2 = \dots = s_{|N(v)|} = s_v = 0, \\ 1_v & s_v = 0, \\ x^{w_v} & \text{otherwise,} \end{cases}$$

where w_v is the weight of vertex v . This tensor means if both v and its neighbouring vertices are not in D , i.e., $s_1 = s_2 = \dots = s_{|N(v)|} = s_v = 0$, this configuration is forbidden because v is not adjacent to any member in the set. otherwise, if v is in D , it has a contribution x^{w_v} to the final result. The graph polynomial for the dominating set problem is known as the domination polynomial [5], which is defined as

$$(C.11) \quad D(G, x) = \sum_{k=0}^{\gamma(G)} d_k x^k,$$

where d_k is the number of dominating sets of size k in graph $G = (V, E)$.

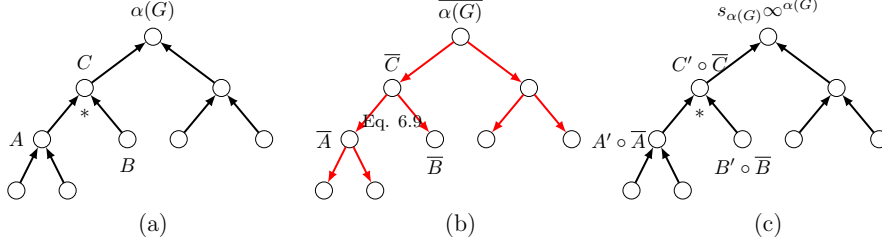


Figure 7: Bounded enumeration of maximum independent sets. In these graphs, a circle is a tensor, an arrow specifies execution direction of a function and \circ is the Hadamard (element-wise) multiplication. \overline{A} means the boolean mask of A . (a) is the forward pass with algebra Eq. (6.3: T) for computing $\alpha(G)$. (b) is the backward pass for computing boolean gradients as masks. (c) is the masked forward pass with algebra Eq. (7.3: P1+SN) for enumerating configurations.

C.6. Satisfiability Problem. We can construct a satisfiability problem to solve the above problem. To generate a tensor network, we map a boolean variable a and its negation $\neg a$ to a degree of freedom (label) $s_a \in \{0, 1\}$, where 0 stands for variable a having value false while 1 stands for having value true. Then we map a clause to a tensor. For example, the k -th clause $\neg a \vee b \vee \neg c$ can be mapped to a tensor labeled by (s_a, s_b, s_c) .

$$(C.12) \quad C_k = \begin{pmatrix} x^{w_k} & x_b^{w_k} \\ x_a^{w_k} & x_{ab}^{w_k} \\ x_c^{w_k} & x_{bc}^{w_k} \\ 1 & x_{abc}^{w_k} \end{pmatrix},$$

where a weight is associated with a clause. There is only one entry $(s_x, s_y, s_z) = (1, 0, 1)$ that makes this clause unsatisfied. If we contract this tensor network, we will get a multiplicative factor x whenever there is a clause satisfied. Let $x_{ab\dots} = x$, then we can get a polynomial that its k -th coefficients gives us the number of assignments with k clauses satisfied.

C.7. Set packing. Set packing is the hypergraph generalization of the maximum independent set problem, where a set corresponds to a vertex and an element corresponds to a hyperedge. To solve the set packing problem, we can just remove the rank-2 restriction of the edge tensor in Eq. (4.2)

$$(C.13) \quad B_{v,w,\dots,z} = \begin{cases} 1, & v + w + \dots + z \leq 1, \\ 0, & \text{otherwise.} \end{cases}$$

Appendix D. Bounding the MIS enumeration space.

When we use the algebra in Eq. (7.3: P1+SN) to enumerate all MIS configurations, we find that the program stores significantly more intermediate configurations than necessary and thus incur significant overheads in space. To speed up the computation and reduce space overhead, we bound the searching space using the information from the computation of the MIS size $\alpha(G)$. As shown in Fig. 7, (a) we first compute the value of $\alpha(G)$ with tropical algebra and cache all intermediate tensors. (b) Then, we compute a boolean mask for each cached tensor, where we use a boolean true to represent a tensor element having a contribution to the MIS (i.e. with a non-zero gradient) and boolean false otherwise. (c) Finally, we perform masked tensor network contraction (i.e. discarding the unnecessary

intermediate configurations) using the element type with the algebra in Eq. (7.3: P1+SN) to obtain all MIS configurations. Note that these masks in fact correspond to tensor elements with non-zero gradients with respect to the MIS size; we compute these masks by back-propagating the gradients. To derive the back-propagation rule for tropical tensor contraction, we first reduce the problem to finding the back-propagation rule of a tropical matrix multiplication $C = AB$. Since $C_{ik} = \bigoplus_j A_{ij} \odot B_{jk} = \max_j A_{ij} \odot B_{jk}$ with tropical algebra, we have the following inequality

$$(D.1) \quad A_{ij} \odot B_{jk} \leq C_{ik}.$$

Here \leq on tropical numbers are the same as the real-number algebra. The equality holds for some j' , which means $A_{ij'}$ and $B_{j'k}$ have contributions to C_{ik} . Intuitively, one can use this relation to identify elements with nonzero gradients in A and B , but if doing this directly, one loses the advantage of using BLAS libraries [3] for high performance. Since $A_{ij} \odot B_{jk} = A_{ij} + B_{jk}$, one can move B_{jk} to the right hand side of the inequality:

$$(D.2) \quad A_{ij} \leq C_{ik} \odot B_{jk}^{\circ-1}$$

where \circ^{-1} is the element-wise multiplicative inverse on tropical algebra (which is the additive inverse on real numbers). The inequality still holds if we take the minimum over k :

$$(D.3) \quad A_{ij} \leq \min_k (C_{ik} \odot B_{jk}^{\circ-1}) = \left(\max_k (C_{ik}^{\circ-1} \odot B_{jk}) \right)^{\circ-1} = \left(\bigoplus_k (C_{ik}^{\circ-1} \odot B_{jk}) \right)^{\circ-1} = (C^{\circ-1} B^T)_{ij}^{\circ-1}.$$

On the right hand side, we transform the operation into a tropical matrix multiplication so that we can utilize the fast tropical BLAS routines [3]. Again, the equality holds if and only if the element A_{ij} has a contribution to C (i.e. having a non-zero gradient). Let the gradient mask for C be \bar{C} ; the back-propagation rule for gradient masks reads

$$(D.4) \quad \bar{A}_{ij} = \delta \left(A_{ij}, \left((C^{\circ-1} \circ \bar{C}) B^T \right)_{ij}^{\circ-1} \right),$$

where δ is the Dirac delta function that returns one if two arguments have the same value and zero otherwise, \circ is the element-wise product, boolean false is treated as the tropical number $\mathbb{0}$, and boolean true is treated as the tropical number $\mathbb{1}$. This rule defined on matrix multiplication can be easily generalized to tensor contraction by replacing the matrix multiplication between $C^{\circ-1} \circ \bar{C}$ and B^T by a tensor contraction. With the above method, one can significantly reduce the space needed to store the intermediate configurations by setting the tensor elements masked false to zero during contraction.

Appendix E. The fitting approach to computing the independence polynomial.

In this section, we propose to find the independence polynomial by fitting $\alpha(G) + 1$ random pairs of x_i and $y_i = I(G, x_i)$. One can then compute the independence polynomial coefficients a_i by solving the linear equation:

$$(E.1) \quad \begin{pmatrix} 1 & x_0 & x_0^2 & \dots & x_0^{\alpha(G)} \\ 1 & x_1 & x_1^2 & \dots & x_1^{\alpha(G)} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & x_{\alpha(G)} & x_{\alpha(G)}^2 & \dots & x_{\alpha(G)}^{\alpha(G)} \end{pmatrix} \begin{pmatrix} a_0 \\ a_1 \\ \vdots \\ a_{\alpha(G)} \end{pmatrix} = \begin{pmatrix} y_0 \\ y_1 \\ \vdots \\ y_{\alpha(G)} \end{pmatrix}.$$

With this approach, we do not incur the linear overhead in space. However, because the independence polynomial coefficients can have a huge order-of-magnitude range, if we use

floating point numbers in the computation, the round-off errors can be significant for the counting of large-size independent sets. In addition, the number could easily overflow if we use fixed-width integer types. The big integer type is also not a good option because big integers with varying width can be very slow and is incompatible with GPU devices. These problems can be solved by introducing a finite-field algebra $\text{GF}(p)$:

$$\begin{aligned} x \oplus y &= x + y \pmod{p}, \\ x \odot y &= xy \pmod{p}, \\ \mathbb{0} &= 0, \\ \mathbb{1} &= 1. \end{aligned} \tag{E.2: GF}(p)$$

With a finite-field algebra, we have the following observations:

1. One can use Gaussian elimination [27] to solve the linear equation Eq. (E.1) since it is a generic algorithm that works for any elements with field algebra. The multiplicative inverse of a finite-field algebra can be computed with the extended Euclidean algorithm.
2. Given the remainders of a larger unknown integer x over a set of co-prime integers $\{p_1, p_2, \dots, p_n\}$, $x \pmod{p_1 \times p_2 \times \dots \times p_n}$ can be computed using the Chinese remainder theorem. With this, one can infer big integers from small integers.

With these observations, we develop Algorithm E.1 to compute the independence polynomial exactly without introducing space overheads. The algorithm iterates over a sequence of large prime numbers until convergence. In each iteration, we choose a large prime number p , and contract the tensor networks to evaluate the polynomial for each variable $\chi = (x_0, x_1, \dots, x_{\alpha(G)})$ on $\text{GF}(p)$ and denote the outputs as $(y_0, y_1, \dots, y_{\alpha(G)}) \pmod{p}$. Then we solve Eq. (E.1) using Gaussian elimination on $\text{GF}(p)$ to find the coefficient modulo p , $A_p \equiv (a_0, a_1, \dots, a_{\alpha(G)}) \pmod{p}$. As the last step of each iteration, we apply the Chinese remainder theorem to update $A \pmod{P}$ to $A \pmod{P \times p}$, where P is a product of all prime numbers chosen in previous iterations. If this number does not change compared with the previous iteration, it indicates the convergence of the result and the program terminates. All computations are done with integers of fixed width W except the last step of applying the Chinese remainder theorem, where we use arbitrary precision integers to represent the counting.

Algorithm E.1 Computing the independence polynomial exactly without integer overflow

Let $P = 1$, W be the integer width, vector $\chi = (0, 1, 2, \dots, \alpha(G))$, matrix $X_{ij} = (\chi_i)^j$, where $i, j = 0, 1, \dots, \alpha(G)$

```

while true do
  compute the largest prime  $p$  that  $\gcd(p, P) = 1$  and  $p < 2^W$ 
  for  $i = 0 \dots \alpha(G)$  do
     $y_i \pmod{p} = \text{contract\_tensor\_network}(\chi_i \pmod{p})$  ; // on  $\text{GF}(p)$ 
  end
   $A_p = (a_0, a_1, \dots, a_{\alpha(G)}) \pmod{p} = \text{gaussian\_elimination}(X, (y_0, y_1, \dots, y_{\alpha(G)}) \pmod{p})$ 
   $A_{P \times p} = \text{chinese\_remainder}(A_p, A_p)$ 
  if  $A_p = A_{P \times p}$  then
    return  $A_p$  ; // converged
  end
   $P = P \times p$ 
end

```

Appendix F. The discrete Fourier transform approach to computing the independence polynomial.

In section 5, we show that the independence polynomial can be obtained by solving the linear equation Eq. (E.1). Since the coefficients of the independence polynomial can range many orders of magnitude, the round-off errors in fitting can be significant if we use random floating point numbers for x_i . In the main text, we propose to use a finite field $\text{GF}(p)$ to circumvent integer overflow and round-off errors. One drawback of using finite field algebra is its matrix multiplication is less computational efficient compared with floating point matrix multiplication. Here, we give an alternative method based on discrete Fourier transform with controllable round off errors. Instead of choosing x_i as random numbers, we can choose them such that they form a geometric sequence in the complex domain $x_j = r\omega^j$, where $r \in \mathbb{R}$ and $\omega = e^{-2\pi i/(\alpha(G)+1)}$. The linear equation thus becomes

$$(F.1) \quad \begin{pmatrix} 1 & r & r^2 & \dots & r^{\alpha(G)} \\ 1 & r\omega & r^2\omega^2 & \dots & r^{\alpha(G)}\omega^{\alpha(G)} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & r\omega^{\alpha(G)} & r^2\omega^{2\alpha(G)} & \dots & r^{\alpha(G)}\omega^{\alpha(G)^2} \end{pmatrix} \begin{pmatrix} a_0 \\ a_1 \\ \vdots \\ a_{\alpha(G)} \end{pmatrix} = \begin{pmatrix} y_0 \\ y_1 \\ \vdots \\ y_{\alpha(G)} \end{pmatrix}.$$

Let us rearrange the coefficients r^j to a_j , the matrix on the left side becomes the discrete Fourier transform matrix. Thus, we can obtain the coefficients by inverse Fourier transform $\vec{a}_r = \text{FFT}^{-1}(\omega) \cdot \vec{y}$, where $(\vec{a}_r)_j = a_j r^j$. By choosing different r , one can obtain better precision in low independent set size region by choosing $r < 1$ or high independent set size region by choosing $r > 1$.

Appendix G. Integer sequence formed by the number of independent sets. We computed the number of independent sets on square lattices and King’s graphs with our generic tensor network contraction algorithm on GPUs. The tensor element type is finite-field algebra so that we can reach arbitrary precision. We also computed the independence polynomial exactly for these lattices which can be found in our [Github repo](#).

Table 2: The number of independent sets for square grid graphs of size $L \times L$. This forms the integer sequence [OEIS A006506](#). Here we only include two updated entries for $L = 38, 39$, which, to our knowledge, has not been computed before [14].

L	square grid graphs
38	616 412 251 028 728 207 385 738 562 656 236 093 713 609 747 387 533 907 560 081 990 229 746 115 948 572 583 817 557 035 128 726 922 565 913 748 716 778 414 190 432 479 964 245 067 083 441 583 742 870 993 696 157 129 887 194 203 643 048 435 362 875 885 498 554 979 326 352 127 528 330 481 118 313 702 375 541 902 300 956 879 563 063 343 972 979
39	29 855 612 447 544 274 159 031 389 813 027 239 335 497 014 990 491 494 036 487 199 167 155 042 005 286 230 480 609 472 592 158 583 920 411 213 748 368 073 011 775 053 878 033 685 239 323 444 700 725 664 632 236 525 923 258 394 737 964 155 747 730 125 966 370 906 864 022 395 459 136 352 378 231 301 643 917 282 836 792 261 715 266 731 741 625 623 207 330 411 607

Appendix H. Technical guide. This project depends on multiple open source packages in the Julia ecosystem. We list the Julia packages that play important roles in our code base as follows:

OMEinsum and **OMEinsumContractionOrders** are packages providing the support for Einstein’s (or tensor network) notation and contraction order optimizations. **OMEinsumContractionOrders** implements state-of-the-art algorithms for finding

the optimal contraction order for a tensor network, including the KaHypar+Greedy [28, 45] and local transformation based approaches [34], **TropicalNumbers** and **TropicalGEMM** are packages providing tropical number and efficient tropical matrix multiplication, **Graphs** is a package providing graph utilities, like random regular graph generator, **Polynomials** is a package providing polynomial algebra and polynomial fitting, **Mods** and **Primes** are packages providing finite field algebra and prime number generations. One can install these packages by opening a Julia REPL, type `]` to enter the `pkg>` mode and type, e.g.

```
pkg> add OMEinsum Graphs Mods Primes Polynomials TropicalNumbers OMEinsumContractionOrders
```

The reader may find it surprising that the Julia implementation of the algorithms introduced in this work is so short that except the bounding algorithm, all are contained in this appendix. This is due to the power of generic programming. After installing the required packages, one can open a Julia REPL and copy the following codes into it.

```
using OMEinsum, OMEinsumContractionOrders
using Graphs
using Random

# generate a random regular graph of size 100, degree 3
graph = (Random.seed!(2); Graphs.random_regular_graph(50, 3))

# generate einsum code, i.e. the labels of tensors
code = EinCode((minmax(e.src,e.dst) for e in Graphs.edges(graph))..., # labels for edge tensors
               [(i,) for i in Graphs.vertices(graph)]..., ()) # labels for vertex
               tensors

# an einsum contraction without a contraction order specified is called `EinCode`,
# an einsum contraction having a contraction order (specified as a tree structure) is called `
#   NestedEinsum`.
# assign each label a dimension-2, it will be used in the contraction order optimization
# `uniquelabels` function extracts the tensor labels into a vector.
size_dict = Dict{String{<int>}, Int{<int>}}{s->2 for s in uniquelabels(code)}
# optimize the contraction order using the `TreeSA` method; the target space complexity is 2^17
optimized_code = optimize_code(code, size_dict, TreeSA())
println("time/space complexity is $(OMEinsum.timespace_complexity(optimized_code, size_dict))")

# a function for computing the independence polynomial
function independence_polynomial(x::T, code) where {T}
    xs = map(getixsv(code)) do ix
        # if the tensor rank is 1, create a vertex tensor.
        # otherwise the tensor rank must be 2, create a bond tensor.
        length(ix)==1 ? [one(T), x] : [one(T) one(T); one(T) zero(T)]
    end
    # both `EinCode` and `NestedEinsum` are callable, inputs are tensors.
    code(xs...)
end

##### COMPUTING THE MAXIMUM INDEPENDENT SET SIZE AND ITS COUNTING/DEGENERACY #####

# using Tropical numbers to compute the MIS size and the MIS degeneracy.
using TropicalNumbers
mis_size(code) = independence_polynomial(TropicalF64(1.0), code)[]
println("the maximum independent set size is $(mis_size(optimized_code).n)")

# A `CountingTropical` object has two fields, tropical field `n` and counting field `c`.
mis_count(code) = independence_polynomial(CountingTropical{Float64,Float64}(1.0, 1.0), code)[]
println("the degeneracy of maximum independent sets is $(mis_count(optimized_code).c)")

##### COMPUTING THE INDEPENDENCE POLYNOMIAL #####

# using Polynomial numbers to compute the polynomial directly
```

```

using Polynomials
println("the independence polynomial is $(independence_polynomial(Polynomial([0.0, 1.0]),
    optimized_code[]))")

##### FINDING MIS CONFIGURATIONS #####

# define the set algebra
struct ConfigEnumerator{N}
    # NOTE: BitVector is dynamic and it can be very slow; check our repo for the static version
    data::Vector{BitVector}
end
function Base.:+(x::ConfigEnumerator{N}, y::ConfigEnumerator{N}) where {N}
    res = ConfigEnumerator{N}(vcat(x.data, y.data))
    return res
end
function Base.:*(x::ConfigEnumerator{L}, y::ConfigEnumerator{L}) where {L}
    M, N = length(x.data), length(y.data)
    z = Vector{BitVector}(undef, M*N)
    for j=1:N, i=1:M
        z[(j-1)*M+i] = x.data[i] .| y.data[j]
    end
    return ConfigEnumerator{L}(z)
end
Base.zero(::Type{ConfigEnumerator{N}}) where {N} = ConfigEnumerator{N}(BitVector[])
Base.one(::Type{ConfigEnumerator{N}}) where {N} = ConfigEnumerator{N}([falses(N)])

# the algebra sampling one of the configurations
struct ConfigSampler{N}
    data::BitVector
end
function Base.:+(x::ConfigSampler{N}, y::ConfigSampler{N}) where {N} # biased sampling: return
    `x`
    return x # randomly pick one
end
function Base.:*(x::ConfigSampler{L}, y::ConfigSampler{L}) where {L}
    ConfigSampler{L}(x.data .| y.data)
end

Base.zero(::Type{ConfigSampler{N}}) where {N} = ConfigSampler{N}(trues(N))
Base.one(::Type{ConfigSampler{N}}) where {N} = ConfigSampler{N}(falses(N))

# enumerate all configurations if `all` is true; compute one otherwise.
# a configuration is stored in the data type of `StaticBitVector`; it uses integers to represent
    bit strings.
# `ConfigTropical` is defined in `TropicalNumbers`. It has two fields: tropical number `n` and
    optimal configuration `config`.
# `CountingTropical{T,<:ConfigEnumerator}` stores configurations instead of simple counting.
function mis_config(code; all=false)
    # map a vertex label to an integer
    vertex_index = Dict{[s=>i for (i, s) in enumerate(unique_labels(code))]}
    N = length(vertex_index) # number of vertices
    xs = map(getixsv(code)) do ix
        T = all ? CountingTropical{Float64, ConfigEnumerator{N}} : CountingTropical{Float64,
            ConfigSampler{N}}
        if length(ix) == 2
            return [one(T) one(T); one(T) zero(T)]
        else
            s = falses(N)
            s[vertex_index[ix[1]]] = true # one hot vector
            if all
                [one(T), T(1.0, ConfigEnumerator{N}([s]))]
            else
                [one(T), T(1.0, ConfigSampler{N}(s))]
            end
        end
    end
    return code(xs...)
end

println("one of the optimal configurations is $(mis_config(optimized_code; all=false)[].c.data)"
)

# direct enumeration of configurations can be very slow; please check the bounding version in

```

```
our Github repo.  
println("all optimal configurations are $(mis_config(optimized_code; all=true)[].c)")
```

For better performance, we recommend checking our GitHub repository for the full featured version: <https://github.com/Happy-Diode/GraphTensorNetworks.jl>. It can be installed in a similar style to other packages.