Mecánica Cuántica

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1 Problema 4

Problem 4

Use Equations 1, 2, and 3, to construct Y_0^0 and Y_2^1 . Check that they are normalized and orthogonal.

$$P_l^m(x) = (-1)^m (1 - x^2)^{m/2} \left(\frac{d}{dx}\right)^m P_l(x)$$
(1)

$$P_l(x) = \frac{1}{2^l l!} \left(\frac{d}{dx}\right)^l (x^2 - 1)^l \tag{2}$$

$$Y_{l}^{m}(\theta,\phi) = \sqrt{\frac{(2l+1)}{4\pi} \frac{(l-m)!}{(l+m)!}} e^{im\phi} P_{l}^{m}(\cos\theta)$$
(3)

Solución: Usando la formula $Y_l^m(\theta,\phi) = \sqrt{\frac{(2l+1)}{4\pi}\frac{(l-m)!}{(l+m)!}}e^{im\phi}P_l^m(cos\theta)$

$$Y_0^0(\theta,\phi) = \sqrt{\frac{(2\cdot 0+1)}{4\pi} \frac{(0-0)!}{(0+0)!}} e^{i\cdot 0\phi} P_0^0(\cos\theta)$$

$$P_0^0(x) = (-1)^0 (1-x^2)^{0/2} (\frac{d}{dx})^0 P_0(x)$$

$$P_0(x) = \frac{1}{2^0 \cdot 0!} (\frac{d}{dx})^0 (x^2 - 1)^0 = 1 \Rightarrow P_0^0 = 1 \cdot 1 \cdot 1 = 1$$

$$Y_0^0(\theta,\phi) = \sqrt{\frac{1}{4\pi} \cdot 1} \cdot 1 \cdot 1 = \sqrt{\frac{1}{4\pi}}$$

Verificamos que esté normalizada

$$\int_{0}^{\pi} \int_{0}^{2\pi} (\sqrt{\frac{1}{4\pi}} \sqrt{\frac{1}{4\pi}} sen\theta) d\theta d\phi = \frac{1}{4\pi} \int_{0}^{\pi} \int_{0}^{2\pi} sen\theta d\theta d\phi = \frac{4\pi}{4\pi} = 1$$

Para el siguiente usamos la misma formula

$$Y_2^1(\theta,\phi) = \sqrt{\frac{(2\cdot 2+1)}{4\pi} \frac{(2-1)!}{(2+1)!}} e^{i\phi} P_2^1(\cos\theta)$$

$$P_2^1(x) = (-1)^1 (1-x^2)^{1/2} (\frac{d}{dx}) P_2(x)$$

$$P_2(x) = \frac{1}{2^2 \cdot 2!} (\frac{d}{dx})^2 (x^2 - 1)^2 = \frac{1}{2} (3x^2 - 1) \Rightarrow P_2^1(x) = -(1-x^2)^{1/2} \frac{d}{dx} \frac{1}{2} (3x^2 - 1) = -3x\sqrt{1-x^2}$$

$$Y_2^1(\theta,\phi) = \sqrt{\frac{5}{4\pi}} \frac{1}{6} e^{i\phi} (-3\cos\theta\sqrt{1-\cos^2\theta}) = -\sqrt{\frac{15}{8\pi}} e^{i\phi}\cos\theta \sin\theta$$

Verificamos que este normalizada

$$\int_{0}^{\pi} \int_{0}^{2\pi} (-\sqrt{\frac{15}{8\pi}} e^{-i\phi} cos\theta sen\theta) (-\sqrt{\frac{15}{8\pi}} e^{i\phi} cos\theta sen\theta) sen\theta d\theta d\phi$$

$$\frac{15}{8\pi} \int_{0}^{\pi} \int_{0}^{2\pi} sen^{3}\theta cos^{2}\theta d\phi d\theta = \frac{15}{8\pi} 2\pi \int_{0}^{\pi} sen^{3}\theta cos^{2}\theta d\theta$$

$$\frac{15}{4} \int_{0}^{\pi} sen^{3}\theta (1 - sen^{2}\theta) d\theta = \frac{15}{4} (\int_{0}^{\pi} sen^{3}\theta d\theta - \int_{0}^{\pi} sen^{5}\theta d\theta)$$

$$\frac{15}{4} (\frac{cos^{3}\theta}{3} - cos\theta + \frac{5cos^{5}\theta}{5} - \frac{2cos^{3}\theta}{3} + cos\theta)|_{0}^{\pi} = \frac{15}{5} \frac{4}{15} = 1$$

Ahora verificamos que sean ortogonales entre si

$$-\int_{0}^{\pi} \int_{0}^{2\pi} (\sqrt{\frac{1}{4\pi}} \sqrt{\frac{15}{8\pi}} e^{i\phi} cos\theta sen\theta sen\theta) d\theta d\phi = -\sqrt{\frac{15}{32\pi^{2}}} \int_{0}^{2\pi} e^{i\phi} d\phi \int_{0}^{\pi} sen^{2}\theta cos\theta d\theta - \sqrt{\frac{15}{32\pi^{2}}} (-ie^{i\phi})|_{0}^{2\pi} (\frac{sen^{3}\theta}{3})|_{0}^{\pi} = -\sqrt{\frac{15}{32\pi^{2}}} \cdot 0 \cdot 0 = 0$$

Lo que demuestra que son ortogonales entre si

2 Problema 5

Problem 5

Show that

$$\Theta(\theta) = Aln[tan(\theta/2)]$$

satisfies the θ equation (Equation 4), for . This is the unacceptable "second solution"—what's wrong with it?

$$\sin\theta \frac{d}{d\theta} \left(\sin\theta \frac{d\Theta}{d\theta} \right) + \left[l(l+1)\sin^2\theta - m^2 \right] \Theta = 0 \tag{4}$$

Solución: La ecuación 4 para l=m=0 es

$$\begin{split} sen\theta\frac{d}{d\theta}(sen\theta\frac{d\Theta}{d\theta}) + (0\cdot(0+1)sen^2\theta - 0^2)\Theta &= sen\theta\frac{d}{d\theta}(sen\theta\frac{d\Theta}{d\theta}) \\ \frac{d\Theta}{d\theta} &= A\frac{sec^2(\theta/2)}{2tan(\theta/2)} = \frac{A}{2}\frac{1}{cos^2(\theta/2)}\frac{cos(\theta/2)}{sen(\theta/2)} \\ &= \frac{A}{2}\frac{1}{cos(\theta/2)sen(\theta/2)} = \frac{A}{2\frac{sen\theta}{2}} = \frac{A}{sen\theta} \\ sen\theta\frac{d}{d\theta}(sen\theta\frac{d\Theta}{d\theta}) &= sen\theta\frac{d}{d\theta}(sen\theta\frac{A}{sen\theta}) = sen\theta\frac{dA}{d\theta} = 0 \end{split}$$

Lo que cumple la ecuación 4, sin embargo podemos ver que Θ diverge tanto en 0 como en π , por lo que no es una solución aceptable

Problem 6

Using Equation 5 and footnote 5, show that

$$Y_l^{-m}(\theta,\phi) = (-1)^m (Y_l^m)^*$$

$$Y_{l}^{m}(\theta,\phi) = \sqrt{\frac{(2l+1)(l-m)!}{4\pi}} e^{im\phi} P_{l}^{m}(cos\theta)$$
 (5)

$$P_l^m(x) = (-1)^m (1 - x^2)^{m/2} \left(\frac{d}{dx}\right)^m P_l(x)$$
 (6)

Footnote 5: Some books (including earlier editions of this one) do not include the factor $(-1)^m$ in the definition of P_l^m . Equation 6 assumes that $m \ge 0$; for negative values we define

$$P_l^{-m}(x) = (-1)^m \frac{(l-m)!}{(l+m)!} P_l^m(x)$$

Solución:

La ecuación (5) nos dice que:

$$Y_{l}^{m}(\theta,\phi) = \sqrt{\frac{(2l+1)}{4\pi} \frac{(l-m)!}{(l+m)!}} e^{im\phi} P_{l}^{m}(\cos\theta),$$

y el pie de nota 5 nos dice que para valores negativos (-m) tenemos:

$$P_l^{-m}(x) = (-1)^m \frac{(l-m)!}{(l+m)!} P_l^m(x). \tag{7}$$

Veamos la ecuación (5) para valores (-m):

$$Y_{l}^{-m}(\theta,\phi) = \sqrt{\frac{(2l+1)}{4\pi} \frac{(l-(-m)!}{(l+(-m)!)}} e^{i(-m)\phi} P_{l}^{-m}(\cos\theta),$$

pero de acuerdo a la ecuación (7) tenemos que:

$$P_l^{-m}(\cos\theta) = (-1)^m \frac{(l-m)!}{(l+m)!} P_l^m(\cos\theta).$$

Por lo que:

$$Y_l^{-m}(\theta,\phi) = \sqrt{\frac{(2l+1)}{4\pi} \frac{(l+m)!}{(l-m)!}} e^{i(-m)\phi} (-1)^m \frac{(l-m)!}{(l+m)!} P_l^m(\cos\theta),$$

$$\Rightarrow Y_{l}^{-m}(\theta,\phi) = (-1)^{m} \sqrt{\frac{(2l+1)}{4\pi} \frac{(l+m)!}{(l-m)!}} e^{-im\phi} \sqrt{\frac{((l-m)!)^{2}}{((l+m)!)^{2}}} P_{l}^{m}(\cos\theta)$$

$$\Rightarrow Y_{l}^{-m}(\theta,\phi) = (-1)^{m} \sqrt{\frac{(2l+1)}{4\pi} \frac{(l+m)!}{(l-m)!} \frac{((l-m)!)^{2}}{((l+m)!)^{2}}} e^{-im\phi} P_{l}^{m}(\cos\theta)$$

$$\Rightarrow Y_{l}^{-m}(\theta,\phi) = (-1)^{m} \sqrt{\frac{(2l+1)}{4\pi} \frac{(l+m)!}{(l-m)!} \frac{(l-m)!}{(l+m)!}} e^{-im\phi} P_{l}^{m}(\cos\theta)$$

$$\Rightarrow Y_{l}^{-m}(\theta,\phi) = (-1)^{m} \sqrt{\frac{(2l+1)}{4\pi} \frac{(l-m)!}{(l+m)!}} e^{-im\phi} P_{l}^{m}(\cos\theta), \tag{8}$$

Ahora notemos lo siguiente:

$$(e^{im\phi})^* = e^{-im\phi}$$

por otra parte, $P_l^m(cos\theta)$ es una función real por lo que:

$$(P_l^m(cos\theta))^* = P_l^m(cos\theta),$$

y como sabemos que l es un entero no negativo entonces tenemos que $\frac{(2l+1)}{4\pi}$ es un número real y $\frac{(l-m)!}{(l+m)!}$ es el cociente de dos factoriales, los cuales son siempre números reales. Por lo que si:

$$Y_{l}^{m} = \sqrt{\frac{(2l+1)}{4\pi} \frac{(l-m)!}{(l+m)!}} e^{im\phi} P_{l}^{m}(\cos\theta) \Longrightarrow (Y_{l}^{m})^{*} = \sqrt{\frac{(2l+1)}{4\pi} \frac{(l-m)!}{(l+m)!}} e^{-im\phi} P_{l}^{m}(\cos\theta)$$

Sustituimos esto en la ecuación (8):

$$Y_l^{-m}(\theta,\phi) = (-1)^m \sqrt{\frac{(2l+1)(l-m)!}{4\pi}} e^{-im\phi} P_l^m(\cos\theta) = (-1)^m (Y_l^m)^*$$

por lo tanto concluimos que:

$$Y_l^{-m}(\theta,\phi) = (-1)^m (Y_l^m)^*$$

Problem 7

Using Equation 9, find $Y_l^l(\theta, \phi)$ and $Y_3^2(\theta, \phi)$. (You can take P_3^2 from Table 4.2, but you'll have to work P_l^l out from Equations 10 and 11.) Check that they satisfy the angular equation (Equation 12), for the appropriate values of l and m.

$$Y_{l}^{m}(\theta,\phi) = \sqrt{\frac{(2l+1)(l-m)!}{4\pi}} e^{im\phi} P_{l}^{m}(\cos\theta)$$
 (9)

$$P_l^m(x) = (-1)^m (1 - x^2)^{m/2} \left(\frac{d}{dx}\right)^m P_l(x)$$
 (10)

$$P_l(x) = \frac{1}{2^l l!} \left(\frac{d}{dx}\right)^l (x^2 - 1)^l \tag{11}$$

$$sin\theta \frac{\partial}{\partial \theta} \left(sin\theta \frac{\partial Y}{\partial \theta} \right) + \frac{\partial^2 Y}{\partial \phi^2} = -l(l+1)sin^2 \theta Y \tag{12}$$

Solución:

Usando la ecuación 9

$$Y_{l}^{l}(\theta,\phi) = \sqrt{\frac{(2l+1)}{4\pi} \frac{(l-l)!}{(l+l)!}} e^{il\phi} P_{l}^{l}(\cos\theta) = \sqrt{\frac{(2l+1)}{4\pi(2l)!}} e^{il\phi} P_{l}^{l}(\cos\theta)$$

En donde al usar la Ecuación 10

$$P_l^l(x) = (-1)^l (1 - x^2)^{l/2} \left(\frac{d}{dx}\right)^l P_l(x) = (-1)^l (1 - x^2)^{l/2} \left(\frac{d}{dx}\right)^l \left[\frac{1}{2^l l!} \left(\frac{d}{dx}\right)^l (x^2 - 1)^l\right]$$

$$P_l^l(x) = (-1)^l (1 - x^2)^{l/2} \frac{1}{2^l l!} \left(\frac{d}{dx}\right)^{2l} \left[(x^2 - 1)^l\right]$$

Veamos como se comporta esta derivada

$$l = 0 ; \left(\frac{d}{dx}\right)^{0} (x^{2} - 1)^{0} = (2(0))!$$

$$l = 1 ; \frac{d^{2}}{dx^{2}} (x^{2} - 1) = 2 = (2(1))!$$

$$l = 2 ; \frac{d^{4}}{dx^{4}} (x^{2} - 1)^{2} = 24 = (2(2))!$$

$$l = 3 ; \frac{d^{6}}{dx^{6}} (x^{2} - 1)^{3} = 720 = (2(3))!$$

Así

$$\frac{d^{2l}}{dx^{2l}}(x^2-1)^l = (2l)!$$

Por lo tanto

$$P_l^l(x) = (-1)^l (1 - x^2)^{l/2} \frac{1}{2^l l!} (2l)!$$

Hemos calculado que

$$Y_{l}^{l}(\theta,\phi) = \sqrt{\frac{(2l+1)}{4\pi(2l)!}}e^{il\phi}(-1)^{l}(1-\cos^{2}\theta)^{l/2}\frac{(2l)!}{2^{l}l!} = \frac{(-1)^{l}}{2^{l}l!}\sqrt{\frac{(2l+1)(2l)!}{4\pi}}e^{il\phi}(sen^{2}\theta)^{l/2}$$

$$Y_{l}^{l}(\theta,\phi) = \frac{(-1)^{l}}{2^{l}l!}\sqrt{\frac{(2l+1)!}{4\pi}}e^{il\phi}sen^{l}\theta = Be^{il\phi}sen^{l}\theta$$

En donde

$$B = \frac{(-1)^l}{2^l l!} \sqrt{\frac{(2l+1)!}{4\pi}}$$

Veamos que cumple la ecuación angular, calcularemos las derivadas:

$$\begin{split} \frac{\partial}{\partial \phi} Y_l^l &= Bsen^l \theta(il) e^{il\phi} \implies \frac{\partial^2}{\partial \phi^2} Y_l^l = (il)^2 Y_l^l \\ sin\theta \frac{\partial Y_l^l}{\partial \theta} &= lcos\theta sin\theta B e^{il\phi} sen^{l-1}\theta = cos\theta l B e^{il\phi} sen^l \theta = lcos\theta Y_l^l \\ \implies sin\theta \frac{\partial}{\partial \theta} \left(sin\theta \frac{\partial Y_l^l}{\partial \theta} \right) &= sin\theta \left[lcos\theta \frac{\partial Y_l^l}{\partial \theta} - lY_1^1 sin\theta \right] = \left[lcos\theta sin\theta \frac{\partial Y_l^l}{\partial \theta} - lY_1^1 sin^2\theta \right] \\ \implies sin\theta \frac{\partial}{\partial \theta} \left(sin\theta \frac{\partial Y_1^l}{\partial \theta} \right) &= \left[l^2 cos^2\theta - lsin^2\theta \right] Y_1^1 \end{split}$$

Sustituimos los valores en la ecuación:

$$\left\lceil l^2 cos^2 \theta - lsin^2 \theta \right\rceil Y_l^l - l^2 Y_l^l = Y_l^l \left\lceil l^2 (cos^2 \theta - 1) - lsin^2 \theta \right\rceil$$

$$Y_{1}^{1}\bigg[-l^{2}sin^{2}\theta-lsin^{2}\theta\bigg] = -Y_{1}^{1}sin^{2}\theta(l^{2}+l) = -l(l+1)sin^{2}\theta Y_{1}^{1}$$

Por lo tanto, vemos que sí es solución de la ecuación angular

$$sin\theta \frac{\partial}{\partial \theta} \left(sin\theta \frac{\partial Y_1^1}{\partial \theta} \right) + \frac{\partial^2 Y_1^1}{\partial \phi^2} = -l(l+1)sin^2\theta Y_1^1$$

Ahora, vamos a calcular $Y_3^2(\theta,\phi)$ y veremos que se satisfaga la ecuación angular.

$$Y_3^2(\theta,\phi) = \sqrt{\frac{(2(3)+1)(3-2)!}{4\pi}} e^{i2\phi} P_3^2(\cos\theta) = \sqrt{\frac{7}{4\pi}} \frac{1}{5!} e^{i2\phi} P_3^2(\cos\theta)$$
$$Y_3^2(\theta,\phi) = \sqrt{\frac{7}{480\pi}} e^{i2\phi} P_3^2(\cos\theta)$$

Calculamos $P_3^2(\cos\theta)$

$$P_3^2(\cos\theta) = (-1)^2 (1 - \cos^2\theta)^{2/2} \frac{d^2}{d\theta^2} P_3(\cos\theta) = (1 - \cos^2\theta) \frac{d^2}{d\theta^2} [P_3(\cos\theta)]$$

En donde $P_3(cos\theta)$ se calcula como sigue:

$$P_{3}(\cos\theta) = \frac{1}{2^{3}3!} \frac{d^{3}}{d\theta^{3}} (\cos^{2}\theta - 1)^{3} = \frac{1}{2^{3}3!} \frac{d^{3}}{d\theta^{3}} ((1 - \sin^{2}\theta) - 1)^{3} = -\frac{1}{2^{3}3!} \frac{d^{3}}{d\theta^{3}} \sin^{6}\theta$$

$$\implies P_{3}(\cos\theta) = -\frac{1}{2^{3}3!} \frac{d^{2}}{d\theta^{2}} (6 \sin^{5}\theta \cos\theta) = -\frac{6}{2^{3}3!} \frac{d}{d\theta} (-\sin^{6}\theta + 5 \sin^{4}\theta \cos^{2}\theta)$$

$$\implies P_{3}(\cos\theta) = -\frac{6}{2^{3}3!} [-6 \sin^{5}\theta \cos\theta + 5(4 \sin^{3}\theta \cos^{3}\theta - 2 \cos\theta \sin^{5}\theta)]$$

$$\implies P_{3}(\cos\theta) = -\frac{6}{2^{3}3!} [-16 \sin^{5}\theta \cos\theta + 20 \sin^{3}\theta \cos^{3}\theta]$$

$$\implies P_{3}(\cos\theta) = -\frac{6}{8(6)} [-16 \sin^{5}\theta \cos\theta + 20 \sin^{3}\theta \cos^{3}\theta] = 2 \sin^{5}\theta \cos\theta - \frac{10}{4} \sin^{3}\theta \cos^{3}\theta$$

Así, sustituimos en la siguiente ecuación:

$$P_3^2(\cos\theta) = (1 - \cos^2\theta) \frac{d^2}{d\theta^2} [P_3(\cos\theta)] = \sin^2\theta \frac{d^2}{d\theta^2} [2\sin^5\theta \cos\theta - \frac{10}{4}\sin^3\theta \cos^3\theta]$$

$$\implies P_3^2(\cos\theta) = \sin^2\theta \frac{d}{d\theta} [2(5\sin^4\theta \cos^2\theta - \sin^6\theta) - \frac{10}{4}(3\sin^2\theta \cos^4\theta - 3\sin^4\theta \cos^2\theta)]$$

$$\implies P_3^2(\cos\theta) = \sin^2\theta \frac{d}{d\theta} [10\sin^4\theta \cos^2\theta - 2\sin^6\theta - \frac{30}{4}\sin^2\theta \cos^4\theta + \frac{30}{4}\sin^4\theta \cos^2\theta]$$

$$\implies P_3^2(\cos\theta) = \sin^2\theta \frac{d}{d\theta} \left[\frac{70}{4}\sin^4\theta \cos^2\theta - 2\sin^6\theta - \frac{30}{4}\sin^2\theta \cos^4\theta \right]$$

$$\implies P_3^2(cos\theta) = sen^2\theta \left[\frac{70}{4} (4sen^3cos^3\theta - 2cos\theta sen^5\theta) + 12sen^5\theta cos\theta - \frac{30}{4} (4cos^3\theta sen^3\theta - 2sen\theta cos^5\theta) \right]$$

$$P_3^2(cos\theta) = sen^2\theta \left[40sen^3\theta cos^3\theta - 23cos\theta sen^5\theta + 15sen\theta cos^5\theta \right]$$

$$P_3^2(cos\theta) = 40sen^5\theta cos^3\theta - 23sen^7\theta cos\theta + 15sen^3\theta cos^5\theta$$

Así finalmente, el armónico esferico resulta ser:

$$\begin{split} Y_3^2(\theta,\phi) &= \sqrt{\frac{7}{480\pi}} e^{i2\phi} P_3^2(\cos\theta) \\ \Longrightarrow & Y_3^2(\theta,\phi) = \sqrt{\frac{7}{480\pi}} e^{i2\phi} [40sen^5\theta cos^3\theta - 23sen^7\theta cos\theta + 15sen^3\theta cos^5\theta] \\ & Y_3^2(\theta,\phi) = B[40sen^5\theta cos^3\theta - 23sen^7\theta cos\theta + 15sen^3\theta cos^5\theta] \end{split}$$

En donde

$$B = \sqrt{\frac{7}{480\pi}}e^{i2\phi}$$

Ahora, vamos a corroborar que satisface la ecuación de ángulo.

$$\frac{\partial}{\partial \phi} Y_3^2 = 2iY_3^2$$

$$\implies \frac{\partial^2}{\partial \phi^2} Y_3^2 = 2i\frac{\partial}{\partial \phi} Y_3^2 = 2i(2iY_3^2) = -4Y_3^2$$

Calculamos la derivada con respecto a θ

$$\frac{\partial}{\partial \theta} Y_3^2 = B[40(5sen^4\theta cos^4\theta - 3cos^2\theta sen^6\theta) - 23(7sen^6\theta cos^2\theta - sen^8\theta) \\ + 15(3sen^2\theta cos^6\theta - 5cos^4\theta sen^4\theta)]$$

$$\implies \frac{\partial}{\partial \theta} Y_3^2 = B[125sen^4\theta cos^4\theta - 281cos^2\theta sen^6\theta + 23sen^8\theta + 45sen^2\theta cos^6\theta]$$

$$\implies \frac{\partial}{\partial \theta} Y_3^2 = B[125sen^4\theta cos^4\theta - 281cos^2\theta sen^6\theta + 23sen^8\theta + 45sen^2\theta cos^6\theta]$$
Usando que $sen^2\theta + cos^2\theta = 1 \implies cos^2\theta = 1 - sen^2\theta$

$$\implies \frac{\partial}{\partial \theta} Y_3^2 = B[125sen^4\theta(\cos^2\theta)^2 - 281cos^2\theta sen^6\theta + 23sen^8\theta + 45sen^2\theta(\cos^2\theta)^3]$$

$$\implies \frac{\partial}{\partial \theta} Y_3^2 = B[125sen^4\theta(1 - sen^2\theta)^2 - 281(1 - sen^2\theta)sen^6\theta + 23sen^8\theta$$

$$+45sen^2\theta(1-sen^2\theta)^3$$
]

$$\Rightarrow \frac{\partial}{\partial \theta} Y_3^2 = B[125sen^4\theta(1 - 2sen^2\theta + sen^4\theta) - 281(1 - sen^2\theta)sen^6\theta + 23sen^8\theta + 45sen^2\theta(1 - 3sen^2\theta + 3sen^4\theta - sen^6\theta)]$$

$$\Rightarrow \frac{\partial}{\partial \theta} Y_3^2 = B\left(45\sin^2\theta - 10\sin^4\theta - 396\sin^6\theta + 384\sin^8\theta\right)$$

$$\Rightarrow sen\theta\frac{\partial}{\partial \theta} Y_3^2 = B\left(45\sin^3\theta - 10\sin^5\theta - 396\sin^7\theta + 384\sin^9\theta\right)$$

$$\Rightarrow \frac{\partial}{\partial \theta} \left(sen\theta\frac{\partial}{\partial \theta} Y_3^2\right) = B\left(135\sin^2\theta - 50\sin^4\theta - 2583\sin^6\theta + 3456\sin^8\theta\right)$$

$$\Rightarrow sen\theta\frac{\partial}{\partial \theta} \left(sen\theta\frac{\partial}{\partial \theta} Y_3^2\right) = B\left(135\sin^3\theta - 50\sin^5\theta - 2583\sin^7\theta + 3456\sin^9\theta\right)$$

$$\Rightarrow sen\theta\frac{\partial}{\partial \theta} \left(sen\theta\frac{\partial}{\partial \theta} Y_3^2\right) = B\left(135\sin^3\theta - 50\sin^5\theta - 2583\sin^7\theta + 3456\sin^9\theta\right)$$

Con todo esto, podemos sustituir

$$sen\theta \frac{\partial}{\partial \theta} \left(sen\theta \frac{\partial}{\partial \theta} Y_3^2 \right) + \frac{\partial^2}{\partial \phi^2} Y_3^2 = sen\theta \frac{\partial}{\partial \theta} \left(sen\theta \frac{\partial}{\partial \theta} Y_3^2 \right) - 4Y_3^2$$

Ahora, como:

$$Y_3^2 = B[40sen^5\theta cos^3\theta - 23sen^7\theta cos\theta + 15sen^3\theta cos^5\theta]$$

$$\implies Y_3^3 = [40sen^5\theta cos^3\theta - 23sen^7\theta cos\theta + 15sen^3\theta cos^5\theta]$$

5 Problema 8

Problem 8

Starting from the Rodrigues formula, derive the orthonormality condition for Legendre polynomials:

$$\int_{-1}^{1} P_l(x) P_{l'}(x) dx = \left(\frac{2}{2l+1}\right) \delta_{ll'}$$

Hint: Use integration by parts

Solución:

Problem 9

- a) From the definition (Equation 13), construct $n_1(x)$ and $n_2(x)$.
- b) Expand the sines and cosines to obtain approximate formulas for $n_1(x)$ and $n_2(x)$, valid when $x \ll 1$. Confirm that they blow up at the origin.

$$j_l(x) = (-x)^l \left(\frac{1}{x} \frac{d}{dx}\right) \frac{\sin x}{x} \; ; \; n_l(x) = -(-x)^l \left(\frac{1}{x} \frac{d}{dx}\right)^l \frac{\cos x}{x}$$
 (13)

Solución:

Inciso a)

Usando la ecuación 13 calculamos $n_1(x)$

$$n_1(x) = -(-x)^1 \left(\frac{1}{x} \frac{d}{dx}\right)^1 \frac{\cos x}{x}$$

$$n_1(x) = x \frac{1}{x} \frac{d}{dx} \left(\frac{\cos x}{x}\right) = \frac{d}{dx} \left(\frac{\cos x}{x}\right)$$

$$n_1(x) = \frac{-x \sin x - \cos x}{x^2}$$

Ahora, calculamos $n_2(x)$

$$n_2(x) = -(-x)^2 \left(\frac{1}{x} \frac{d}{dx}\right)^2 \frac{\cos x}{x}$$

$$n_2(x) = -x^2 \frac{1}{x^2} \frac{d^2}{dx^2} \left(\frac{\cos x}{x}\right) = -\frac{d}{dx} \left[\frac{d}{dx} \left(\frac{\cos x}{x}\right)\right]$$

$$n_2(x) = -\frac{d}{dx} \left[\frac{-x \operatorname{sen} x - \cos x}{x^2}\right]$$

$$n_2(x) = -\frac{(-x \operatorname{cos} x)x^2 - 2x(-x \operatorname{sen} x - \cos x)}{x^4}$$

$$n_2(x) = \frac{x^3 \operatorname{cos} x - 2x^2 \operatorname{sen} x - 2x \operatorname{cos} x}{x^4}$$

$$n_2(x) = \frac{x^2 \operatorname{cos} x - 2x \operatorname{sen} x - 2\cos x}{x^3}$$

Inciso b)

Sabemos que la expansión del coseno, resulta ser

$$\cos(x) = \sum_{k=0}^{\infty} \frac{(-1)^k x^{2k}}{(2k)!}$$

Así

$$\frac{\cos(x)}{x} = \sum_{k=0}^{\infty} \frac{(-1)^k x^{2k-1}}{(2k)!}$$

Con esto, podemos calcular $n_1(x)$ y $n_2(x)$

$$n_1(x) = -(-x)^1 \left(\frac{1}{x} \frac{d}{dx}\right)^1 \frac{\cos x}{x} = \frac{d}{dx} \left(\frac{\cos x}{x}\right) = \frac{d}{dx} \left(\sum_{k=0}^{\infty} \frac{(-1)^k x^{2k-1}}{(2k)!}\right)$$

$$n_1(x) = \sum_{k=0}^{\infty} (2k-1) \frac{(-1)^k x^{2k-2}}{(2k)!}$$

$$n_2(x) = -(-x)^2 \left(\frac{1}{x} \frac{d}{dx}\right)^2 \frac{\cos x}{x} = -\frac{d^2}{dx^2} \left(\frac{\cos x}{x}\right) = -\frac{d^2}{dx^2} \left(\sum_{k=0}^{\infty} \frac{(-1)^k x^{2k-1}}{(2k)!}\right)$$

$$n_2(x) = \sum_{k=0}^{\infty} (1-2k)(2k-2) \frac{(-1)^k x^{2k-3}}{(2k)!}$$

Expandiendo la serie, tenemos lo siguiente:

$$n_1(x) \approx (-1)\frac{x^{-2}}{0!} + (1)\frac{(-1)^1 x^0}{2!} + (3)\frac{(-1)^2 x^2}{4!} + (5)\frac{(-1)^3 x^4}{6!} + (7)\frac{(-1)^4 x^6}{8!}$$
$$n_1(x) \approx -\frac{1}{x^2} - \frac{1}{2} + \frac{3x^2}{24} - \frac{5x^4}{720} + \frac{7x^6}{40320}$$

$$n_2(x) \approx (1)(-2)\frac{(-1)^0 x^{-3}}{0!} + (-3)(2)\frac{(-1)^2 x^1}{4!} + (-5)(4)\frac{(-1)^3 x^3}{6!} + (-7)(6)\frac{(-1)^4 x^5}{8!}$$
$$n_2(x) \approx -\frac{2}{x^3} - \frac{6x}{24} + \frac{20x^3}{720} - \frac{42x^5}{40320}$$

Puesto que $x \ll 1$, los términos que van a dominar en la aproximación resultan ser:

$$n_1(x) \approx -\frac{1}{x^2} - \frac{1}{2} \implies \lim_{x \to 0} n_1(x) = -\infty$$

 $n_2(x) \approx -\frac{2}{x^3} \implies \lim_{x \to 0} n_2(x) = -\infty$

Problem 10

- a) Check that $Arj_1(kr)$ satisfies the radial equation with V(r) = 0 and l = 1.
- b) Determine graphically the allowed energies for the infinite spherical well, when l=1. Show that for large N, $E_{N1}\approx (\hbar^2\pi^2/2ma^2)(N+1/2)^2$. Hint: First show that $j_1(x)=0 \implies x=tanx$. Plot x and tanx on the same graph, and locate the points of intersection.

Solución:

Inciso a)

Tenemos que la función para l=1

$$u_1 = Arj_1(x) = Ar(-x)^1 \left(\frac{1}{x}\frac{d}{dx}\right)^1 \frac{sen(x)}{x} = (-Ar)\frac{cos(x)x - sen(x)}{(x)^2} = Ar\frac{sen(x) - xcos(x)}{x^2}$$

Tomamos x = kr

$$u_1 = Ar \frac{sen(kr) - krcos(kr)}{(kr)^2} = A \left[\frac{sen(kr)}{k^2r} - \frac{cos(kr)}{k} \right] = \frac{A}{kr} \left[\frac{sen(kr)}{k} - cos(kr)r \right]$$

La ecuación radial para V(r) = 0 es

$$\frac{d^2u_1}{dr^2} + \left[k^2 - \frac{l(l+1)}{r^2}\right]u_1 = 0$$

Calculamos las derivadas

$$\frac{du_1}{dr} = -\frac{A}{kr^2} \left[\frac{sen(kr)}{k} - cos(kr)r \right] + A \left[sen(kr) \right]$$

$$\frac{d^2u_1}{dr^2} = -\frac{A}{kr^2} \left[sen(kr)kr \right] + \frac{2A}{kr^3} \left[\frac{sen(kr)}{k} - cos(kr)r \right] + Acos(kr)k$$

$$\frac{d^2u_1}{dr^2} = \frac{A}{kr} \left[-sen(kr)k + k^2rcos(kr) \right] + \frac{2A}{kr^3} \left[\frac{sen(kr)}{k} - cos(kr)r \right]$$

$$\left[k^2 - \frac{l(l+1)}{r^2} \right] u_1 = \frac{Ak}{r} \left[\frac{sen(kr)}{k} - cos(kr)r \right] - \frac{l(l+1)}{r^2} \frac{A}{kr} \left[\frac{sen(kr)}{k} - cos(kr)r \right]$$

Como estamos en el caso en que l=1

$$\left[k^2 - \frac{l(l+1)}{r^2}\right]u_1 = \frac{A}{kr}\left[sen(kr)k - cos(kr)k^2r\right] - \frac{2A}{kr^3}\left[\frac{sen(kr)}{k} - cos(kr)r\right]$$

Así, sustituyendo las derivadas en la ecuación

$$\frac{d^2u_1}{dr^2} + \left[k^2 - \frac{l(l+1)}{r^2}\right]u_1 = \frac{A}{kr}\left[-sen(kr)k + k^2rcos(kr)\right] + \frac{2A}{kr^3}\left[\frac{sen(kr)}{k} - cos(kr)r\right] + \frac{A}{kr}\left[sen(kr)k - cos(kr)k^2r\right] - \frac{2A}{kr^3}\left[\frac{sen(kr)}{k} - cos(kr)r\right] = 0$$

Por lo tanto se comprueba que la función es solución a la ecuación radial con V(r)=0 y l=1

Inciso b)

Sabemos que:
$$Aaj_1(ka) = 0 \implies j_1(ka) = 0$$

$$\implies \frac{sen(ka) - (ka)cos(ka)}{(ka)^2} = 0$$

$$\implies tan(ka) - ka = 0$$

$$\implies tan(ka) = ka$$

o bien, como x = ka. Llegamos a la ecuación trascendental

$$tan(x) = x$$

La cual graficamos en la siguiente Figura ??:

Notamos que para valores grandes de N, los valores en donde la función tan(x) y la función x son iguales, estos se aproximan a $x = \frac{\pi}{2} + N\pi$ que es donde la función tan(x) diverge.

Ahora como x = ka, para que la ecuación trascendental sea válida, se tiene que

$$ka = \frac{\pi}{2} + N\pi$$

$$a\sqrt{\frac{2mE_N}{\hbar^2}} = \frac{\pi}{2} + N\pi$$

$$\frac{2mE_N}{\hbar^2} = a^{-2}\pi^2(N+1/2)^2$$

$$E_N = \frac{\hbar^2\pi^2}{2ma^2}(N+1/2)^2$$

Problem 11

A particle of mass m is placed in a finite spherical well:

$$V(r) = \begin{cases} -V_0. & r \leq a; \\ 0. & r > a \end{cases}$$
 Find the ground state, by solving the radial equation with

l=0. Show that there is no bound state if $V_0a^2 < \pi^2\hbar^2/8m$.

Solución: