

The New
$$\lambda_{\alpha\beta} = \frac{\partial \chi}{\partial u^{\alpha}} + \frac{\partial \chi}{\partial u^{\beta}} + \frac{\partial \chi}{\partial u^{\beta}} + \frac{\partial \chi}{\partial u^{\beta}}$$
The New
$$\lambda_{\alpha\beta} = \frac{\partial \chi}{\partial u^{\alpha}} + \frac{\partial \chi}{\partial u^{\beta}} + \frac{\partial \chi}{\partial$$

$$\dot{T} = kN = kgQ + k_n N \qquad Q = n \times T,$$

$$kg = \left(\frac{\sigma}{\sigma} \frac{du^{\alpha}}{ds} \frac{du^{\beta}}{ds} + \frac{d^2u^{\alpha}}{ds^2} \right) \chi_{\sigma}$$

ODE.
$$\int_{a\beta}^{b} \frac{du^{\beta}}{ds} \frac{du^{\beta}}{ds} + \frac{d^{2}u^{5}}{ds^{2}} = 0$$

$$u^{\alpha}(0) = u^{\alpha}_{o}$$
 — start point $\frac{du^{\alpha}}{ds}(0) = v^{\alpha}_{o}$ — start vector $v = v^{\alpha}_{o} \times v^{\alpha}_{o}$

ODEROUS 3! 和性的

V V € TpM. IIVII=S

进户点双前的方面给何是的测地线

股(S) 取为(S)上一点Q, S.t. 混神中电学Q
到P点的距离为S. 对及V的格

 $exp_p: T_pM \longrightarrow M.$ $V \longmapsto Q$

Q×肝在中的邻域内是 一个微为问服。

 $Q = exp_{P}(V)$

若在PUSULTD平面下M处选版为正形势?O;enery。 Hon = V=Voe1+Voe2

V= || V|| V0 = 5 V0 = 5 V0 e1 + 5 V0 e2 } SV0 ,5 V0]

相分级的物对文层的物品

由台正得到Q的对方部为 对的局子。(Normal)

MM. 在助面M上一点中的邻城的这件科等下

$$g_{11}(p) = g_{22}(p) = 1$$
. $g_{12}(p) = 0$ $\frac{\partial g_{\alpha \beta}}{\partial y^{\beta}}(0) > 0$

\$?o,e,e, ETPM

电为初始行星的测地段 9~ 多数曲线

 y^{α} . 35/2 $\frac{\partial x}{\partial y^{\alpha}} = e_{\alpha} \implies g_{ii}(p) = g_{ii}(p) = g_{ii}(p) = 1$

$$\int_{\beta r}^{\alpha} \frac{du^{\beta}}{ds} \frac{du^{\gamma}}{ds} + \frac{d^{2}u^{\alpha}}{ds^{2}} = 0$$

$$\int_{\beta r}^{\alpha} \frac{dy^{\beta}}{ds} \frac{dy^{\beta}}{ds} + \frac{d^{2}y^{\alpha}}{ds^{2}} = 0$$

$$\int_{\beta r}^{\alpha} \frac{dy^{\beta}}{ds} \frac{dy^{\gamma}}{ds} + \frac{d^{2}y^{\alpha}}{ds^{2}} = 0$$

$$\mathcal{V}^{\alpha}(\mathfrak{I}) = \mathcal{V}^{\alpha} \qquad \qquad \mathcal{Y}^{\alpha}(\mathfrak{I}) = 0$$

$$\frac{d\mathcal{V}^{\alpha}}{d\mathfrak{I}}(\mathfrak{I}) = \mathcal{V}^{\alpha} \qquad \qquad \mathcal{Y}^{\alpha}(\mathfrak{I}) = 0$$

$$\frac{dy^{\alpha}}{ds}(0) = y^{\alpha}$$

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$$\begin{cases} y' = \rho \cos \theta \\ y^2 = \rho \sin \theta \end{cases}$$

Thm 在关于户点为中心的洲地极生称多下。

极些桥下. 限点有奇险。

$$\begin{array}{c}
Pf. & (f, o) = (q, q).
\end{array}$$

$$ds^2 = \overline{g}_{11}(dp)^2 + 2\overline{g}_{12}dpdo + \overline{g}_{21}(dp)^2$$

$$\overline{g_{i}} = \chi_{\overline{y}_{i}} \chi_{\overline{y}_{i}} = 1$$

$$\int_{\alpha\beta}^{\sigma} \frac{du^{\alpha}}{ds} \frac{du^{\beta}}{ds} + \frac{d^{2}u^{\alpha}}{ds^{2}} = 0 \qquad \Rightarrow 0 + \overline{P}_{11}^{2} \quad 1.1 = 0$$

$$\frac{du^{\alpha}}{ds}(0) = u^{\alpha}_{0}$$

$$\frac{du^{\alpha}}{ds}(0) = \overline{V}_{0}^{\alpha}$$

$$\mathcal{N}_{\alpha}(\omega) = \mathcal{N}_{\alpha}^{\circ}$$

$$\frac{du^{\alpha}}{ds}(0) = V_{\alpha}^{\alpha}$$

$$\overline{\int_{1}^{2}} = \frac{1}{2} \overline{g}^{24} \left(\frac{\partial \overline{g}_{K1}}{\partial \overline{y}^{1}} + \frac{\partial \overline{g}_{10}}{\partial \overline{y}^{1}} - \frac{\partial \overline{g}_{11}}{\partial g^{2}} \right)$$

$$= \overline{9^{22}} \frac{3\overline{9_{21}}}{3\overline{9_{1}}} = 0$$

$$\frac{3\overline{9_{21}}}{3\overline{9}} = 0 \quad \overline{9_{21}} \quad 5\overline{9_{22}}$$

$$\Rightarrow \widehat{q}_{\mathcal{H}}(\alpha) = \lim_{\rho \to \alpha} \widehat{q}_{21} \qquad \overline{q}_{21} = \chi_{\theta} \cdot \chi_{\rho}.$$

$$\chi_0 = -\chi_{y'} \rho_{sin\theta} + \chi_{y} \rho_{\omega 10}$$

$$\chi_{\rho} = \chi_{y'} \omega_{10} + \chi_{y} sin\omega$$

$$= 310. \text{ Mp} = -911 \left(\sin \theta \cos \theta + 922 \right) \sin \theta \cos \theta$$

$$+ 911 \left(\sin^2 \theta - \cos^2 \theta \right)$$

$$P \rightarrow 0. \quad \chi_0 \cdot \chi_0 = 0 \quad \Rightarrow \quad \tilde{J}_{12}(0) = 0$$

$$\frac{3}{\left(\frac{g_{11}}{g_{21}},\frac{g_{12}}{g_{11}}\right)} = \left(\frac{1}{0},\frac{0}{g_{21}}\right) \det\left(\frac{g_{\alpha\beta}}{g_{\alpha\beta}}\right) = \frac{g_{22}}{\int \det(\frac{g_{\alpha\beta}}{g_{\alpha\beta}})} = \frac{1}{2}g_{22}$$

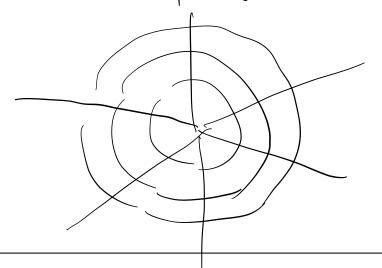
$$= \int det g_{ab} \frac{\partial (y,y)}{\partial (p,0)}$$

$$= \int det (g_{ab})$$

$$\lim_{p\to 0} \sqrt{g_{22}} = \lim_{p\to 0} \int \int det g_{ab} = 0$$

$$(\sqrt{g_{12}})_p = (\int det g_{ab})_p = \int det g_{ab} + \int (\int det g_{ab})_p$$

$$(\sqrt{det g_{ab}})_p = \frac{1}{2\sqrt{det g_{ab}}} \int_{\partial P} \int$$



$$2 = f(y)$$

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$$4 = \frac{f'f'}{4^2/4+(f')^3 x^2} = -\frac{1}{9^2}$$

$$= \frac{f'f''}{u^2(1+(f')^2)^2} = -\frac{1}{9^2}$$

$$\frac{1}{2} \mathcal{A}((f^{\prime})^{2})^{2} = -\frac{1}{2} \mathcal{A}(u^{2})^{2}$$

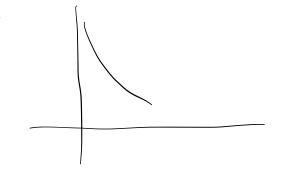
$$\frac{1}{1+(f')^2} = \frac{(u^{\eta})^2}{g^2} + C \qquad \text{In } C = 0$$

$$\int = \pm \frac{\int a^2 - (u^3)^2}{u^2}$$

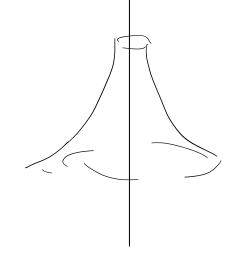
$$T_{N}'' - '' \qquad N^2 = a \cos \gamma$$

$$\Rightarrow f(\varphi) = \int_{\alpha} \frac{\sin^2 \varphi}{\cos \varphi} d\varphi$$

東份沒



份形的



正学的一种面

测地域 pyter $T_{9} = -\frac{dB}{dS} \cdot N = B \frac{dN}{dS} = \dot{N} \cdot (T_{1}N)$ =(N,T,N)* Rytor HRN= Rgat kin >> N=In $=) \quad T_{g} = (\dot{\eta}, \tau, \dot{\eta}) = (\eta_{\alpha} \frac{\alpha \dot{\eta}^{\alpha}}{\alpha \dot{\varsigma}}, \chi_{\beta} \frac{\alpha \dot{\eta}^{\beta}}{\alpha \dot{\varsigma}}, \dot{\eta})$ $= \frac{1}{\sqrt{\det 9}} \left((n_1 \times 2C_1)(x_1 \times x_2) \frac{(cln')^2}{(as')^2} \right)$

$$+\left((n_{1}\times\chi_{1})-(n_{1}\times\chi_{1})\right)(\chi_{1}\times\chi_{1})\left(\frac{\partial n_{1}}{\partial s}\frac{\partial n_{2}}{\partial s}\right)$$

$$=\left(n_{1}\chi_{1}\right)(\chi_{1}\chi_{1})\left(\chi_{1}\chi_{1}\right)\left(\chi_{1}\chi_{1}\right)\left(\chi_{1}\chi_{1}\right)\left(\chi_{1}\chi_{1}\right)\left(\chi_{1}\chi_{1}\right)$$

$$=-h_{11}g_{12}+h_{12}g_{11}$$

$$T_{9} = \frac{1}{\sqrt{3}} \left\{ \begin{array}{c} \frac{\partial u^{2}}{\partial s} \\ \frac{\partial u^{2}}{\partial s} \end{array} \right\} - \frac{\partial u^{3}}{\partial s} \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{2}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} - \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s} \right)^{2} \\ \frac{\partial u^{3}}{\partial s} \left(\frac{\partial u^{3}}{\partial s}$$