

## Chapter 1

## Gauss Law

## **Gravitational Fields** 1.1

The divergence theorem for any field  $\bar{F}$  is denoted as

 $\iiint \nabla \cdot F \, dV = \oiint F \cdot \bar{n} \, ds, \text{ wherein } div[F] = \nabla \cdot F$  For a spherical mass  $m_1$ , the gravitational field generated by the mass  $m_1$  at any point:  $F_{field}^- = -\frac{Gm_1}{r^2}\hat{r}$ , wherein  $\hat{r}$  represents unit vector with direction from mass to object.

$$\hat{r} = \frac{1}{\sqrt{x^2 + y^2 + z^2}} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

$$\hat{r} = \begin{pmatrix} x(x^2 + y^2 + z^2)^{-\frac{1}{2}} \\ y(x^2 + y^2 + z^2)^{-\frac{1}{2}} \\ z(x^2 + y^2 + z^2)^{-\frac{1}{2}} \end{pmatrix}$$

Therefore, given that  $x \neq 0$ ,  $y \neq 0$ ,  $z \neq 0$ 

$$\bar{F}_f = -Gm_1 \begin{pmatrix} x(x^2 + y^2 + z^2)^{-\frac{3}{2}} \\ y(x^2 + y^2 + z^2)^{-\frac{3}{2}} \\ z(x^2 + y^2 + z^2)^{-\frac{3}{2}} \end{pmatrix}$$

$$\nabla \cdot \bar{F}_f = -Gm_1 \nabla \cdot \left[ \begin{pmatrix} x(x^2 + y^2 + z^2)^{-\frac{3}{2}} \\ y(x^2 + y^2 + z^2)^{-\frac{3}{2}} \\ z(x^2 + y^2 + z^2)^{-\frac{3}{2}} \end{pmatrix} \right]$$

Let 
$$k = \begin{pmatrix} f \\ g \\ h \end{pmatrix} = \begin{pmatrix} x(x^2 + y^2 + z^2)^{-\frac{3}{2}} \\ y(x^2 + y^2 + z^2)^{-\frac{3}{2}} \\ z(x^2 + y^2 + z^2)^{-\frac{3}{2}} \end{pmatrix}$$

$$\nabla \cdot k = f_x + g_y + h_z = \nabla \cdot \begin{pmatrix} x(x^2 + y^2 + z^2)^{-\frac{3}{2}} \\ y(x^2 + y^2 + z^2)^{-\frac{3}{2}} \\ z(x^2 + y^2 + z^2)^{-\frac{3}{2}} \end{pmatrix}$$

$$\frac{\partial f}{\partial x} = \frac{\partial}{\partial x} \left[ x(x^2 + y^2 + z^2)^{-\frac{3}{2}} \right]$$

$$\frac{\partial f}{\partial x} = -3x^2(x^2 + y^2 + z^2)^{-\frac{5}{2}} + (x^2 + y^2 + z^2)^{-\frac{3}{2}}$$

$$\frac{\partial g}{\partial y} = \frac{\partial}{\partial y} \left[ y(x^2 + y^2 + z^2)^{-\frac{3}{2}} \right]$$

$$\frac{\partial g}{\partial y} = -3y^2(x^2 + y^2 + z^2)^{-\frac{5}{2}} + (x^2 + y^2 + z^2)^{-\frac{3}{2}}$$

$$\frac{\partial f}{\partial z} = \frac{\partial}{\partial z} \left[ z(x^2 + y^2 + z^2)^{-\frac{5}{2}} + (x^2 + y^2 + z^2)^{-\frac{3}{2}} \right]$$

$$\nabla \cdot k = \frac{\partial f}{\partial x} + \frac{\partial g}{\partial y} + \frac{\partial h}{\partial z}$$

$$\nabla \cdot k = -3(x^2 + y^2 + z^2)(x^2 + y^2 + z^2)^{-\frac{5}{2}} + 3(x^2 + y^2 + z^2)^{-\frac{3}{2}}$$

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$$\nabla \cdot k = 0$$

$$\nabla \cdot F_f = -Gm_1 \times 0$$

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Therefore, for any point outside of the sphere, the divergence of the gravitational field is nonexistent. Consider a region R which is the space containing empty space and a mass sphere. Divergence in region R is the addition of divergence of free space and divergence of mass sphere. Since  $div(F_f) = 0$  outside of the mass sphere as shown in previous example, divergence in R is only divergence of mass sphere.

$$div(F_R) = div(F_{sphere}) + div(F_{freespace})$$

By previous example, 
$$div(F_{freespace}) = 0$$

$$div(F_R) = div(F_{sphere})$$

Consider a point mass of mass M of uniform density with radius R:

$$\bar{f}_f = \frac{GM}{r^2}\hat{r}$$

$$\iiint \nabla \cdot \bar{f}_f \, dV = \oiint \bar{f}_f \cdot \bar{n} \, ds$$

By considering a surface to be enveloping the mass sphere,

$$\iiint \nabla \cdot \bar{f}_f \, dV = 4\pi R^2 \bar{f}_f$$

For all points with radius r = R away from the center of the mass sphere,

$$\bar{f}_f = \frac{GM}{r^2}\hat{r}$$

$$\iiint \nabla \cdot \bar{f}_f \, dV = 4\pi R^2 \frac{GM}{R^2}$$

$$\iiint \nabla \cdot \bar{f}_f \, dV = 4\pi G M$$

By considering the predetermined property of this mass sphere to have uniform density:

$$(\nabla \cdot \bar{f}_f)V = 4\pi GM$$
 
$$(\nabla \cdot \bar{f}_f) \lim_{V \to 0} [V] = 4\pi G \lim_{M \to 0} [M]$$
 
$$(\nabla \cdot \bar{f}_f) dV = 4\pi G dM$$
 
$$\nabla \cdot \bar{f}_f = 4\pi G \frac{dM}{dV}$$
 wherein  $\rho$  is density,

$$\nabla \cdot \bar{f}_f = 4\pi G \rho$$

By considering the graviational field of any arbitrary object to be the summation of the gravitational field of small components of the object,

$$\bar{F} = \bar{f}_1 + \bar{f}_2 + \bar{f}_3 + \dots \bar{f}_i$$

$$\bar{F} = \sum_{n=1}^{i} \left[ \bar{f}_n \right]$$

 $\nabla \cdot$  could be considered as a linear transformation because  $\nabla \cdot (\bar{a} + \bar{b}) = (\nabla \cdot \bar{a}) + (\nabla \cdot \bar{b})$  and also  $\nabla \cdot (c\bar{a}) = c(\nabla \cdot \bar{a})$ . Therefore,

$$\nabla \cdot \bar{F} = \sum_{n=1}^{i} \left[ \nabla \cdot \bar{f}_{n} \right]$$

$$\nabla \cdot \bar{F} = \sum_{n=1}^{i} \left[ 4\pi G \rho_{n} \right]$$

$$\iiint_{R} \nabla \cdot \bar{F} \, dV = \lim_{\Delta V_{n} \to 0} \left[ \sum_{n=1}^{\infty} \left[ 4\pi G \rho_{n} \Delta V_{n} \right] \right]$$

$$\iiint_{R} \nabla \cdot \bar{F} \, dV = 4\pi G \lim_{\Delta V_{n} \to 0} \left[ \sum_{n=1}^{\infty} \left[ \rho_{n} \Delta V_{n} \right] \right]$$

$$\rho_{n} \Delta V_{n} = \left( \frac{dM_{n}}{dV_{n}} \right) dV_{n} = dM_{n}$$

$$\lim_{\Delta V_{n} \to 0} \left[ \sum_{n=1}^{\infty} \left[ \rho_{n} \Delta V_{n} \right] \right] = \int_{0}^{m_{e}} dM = m_{e}$$

$$\iiint_{R} \nabla \cdot \bar{F} \, dV = 4\pi G m_{e}$$

wherein  $m_e$  represents the mass enclosed by the surface. By reiterating The Divergence Theorem,

$$\iiint \nabla \cdot \bar{F} \, dV = \oiint \bar{F} \cdot \bar{n} \, ds$$
$$4\pi G m_e = \oiint \bar{F} \cdot \bar{n} \, ds$$

For the special case wherein  $\rho = f(r)$ , the gravitational field would be perpendicular to an imaginary spherical surface at radius r away from the center of the mass sphere. Therefore,

$$\bar{F}\cdot\bar{n}=|\bar{F}|$$
 
$$\oiint \bar{F}\cdot\bar{n}\,ds=|\bar{F}|\times 4\pi r^2$$

$$4\pi G m_e = |\bar{F}| \times 4\pi r^2$$
$$|\bar{F}| = \frac{G m_e}{r^2}$$

The  $|\bar{F}|$  represents gravitational field produced by a spherical mass with varying radial density,  $\rho = f(r)$ .  $m_e$  represents enclosed mass, which in this case is the entirety of the mass sphere. Since the force of gravity experienced by an object with mass  $m_2$  is  $F_{force} = F_{field}^- \times m_2$ ,

$$F_{force} = \frac{Gm_em_2}{r^2} = \frac{Gm_e}{r^2} \int_0^R f(r)dr$$

wherein R is the radius of the mass sphere of varying radial density.