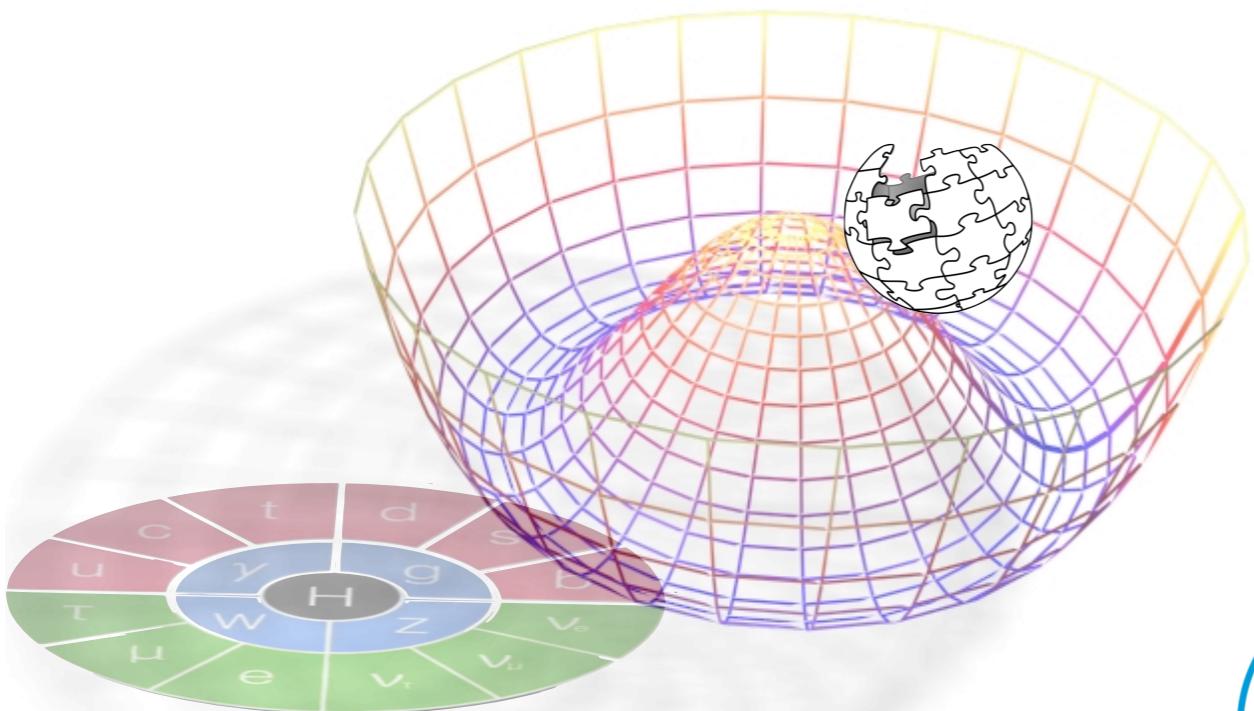


# The Standard Model of particle physics

*CERN summer student lectures 2023*

Lecture 3/5



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# Outline

## □ Monday: symmetry

- Lagrangians
- Lorentz symmetry - scalars, fermions, gauge bosons
- Gauge/local symmetry as dynamical principle - Example: U(1) electromagnetism

## □ Tuesday: SM symmetries

- Nuclear decay, Fermi theory and weak interactions: SU(2)
- Dimensional analysis: cross-sections and life-time computations made simple
- Strong interactions: SU(3)

## □ Wednesday: chirality of weak interactions

- Chirality of weak interactions
- Pion decay

## □ Thursday: Higgs mechanism

- Spontaneous symmetry breaking and Higgs mechanism
- Lepton and quark masses, quark mixings
- Neutrino masses

## □ Friday: quantum effects

- Running couplings
- Asymptotic freedom of QCD
- Anomalies cancelation



# Universality of Weak Interactions

$$\mu \rightarrow e \nu_\mu \bar{\nu}_e$$

$\tau_\mu \approx 10^{-6}$ s

$$n \rightarrow p e \bar{\nu}_e$$

$\tau_n \approx 900$ s

$$\mathcal{L} = G_F \psi^4$$

$$\Gamma_\mu = \frac{G_F^2 m_\mu^5}{192\pi^3} \sim 1/10^{-6}''$$

$$\Gamma_n = \frac{G_F^2 \Delta m^5}{192\pi^3} \sim 1/15'$$

[  
factor 192 not exactly correct  
because n and p are not elementary particles  
form factors are involved  
]

$$\mathcal{L} \stackrel{?}{=} G_F (\bar{n} p \bar{e} \nu_e + \bar{\mu} \nu_\mu \bar{e} \nu_e)$$

By analogy with electromagnetism, one can see the Fermi force as a current-current interaction (vector-vector interaction instead of scalar-scalar interaction)

$$\mathcal{L} = G_F J_\mu^* J^\mu \quad \text{with} \quad J^\mu \stackrel{?}{=} (\bar{n} \gamma^\mu p) + (\bar{e} \gamma^\mu \nu_e) + (\bar{\mu} \gamma^\mu \nu_\mu) + \dots$$

it can be shown (thanks to the transformation law of spin-1/2 field given before) that this Lagrangian is invariant under Lorentz transformation

The cross-terms generate both neutron decay and muon decay.

The life-times of the neutron and muon tell us that the relative factor between the e and the  $\mu$  in the current is of order one: the weak force has the **same strength for e and  $\mu$** .

# Pion decay(s)

What about  $\pi^\pm$  decay  $\tau_\pi \approx 10^{-8}$ s?

$$\pi^- \rightarrow \mu^-\bar{\nu}_\mu$$

$$\pi^- \rightarrow e^-\bar{\nu}_e$$

experimentally the pions decay dominantly into muons and not electrons.

Why  $\frac{\Gamma(\pi^- \rightarrow e^-\bar{\nu}_e)}{\Gamma(\pi^- \rightarrow \mu^-\bar{\nu}_\mu)} \sim 10^{-4}$  ? And not  $\frac{\Gamma(\pi^- \rightarrow e^-\bar{\nu}_e)}{\Gamma(\pi^- \rightarrow \mu^-\bar{\nu}_\mu)} \sim \frac{(m_\pi - m_e)^5}{(m_\pi - m_\mu)^5} \sim 500$  ?  
EXP TH

Does it mean that our way to compute decay rate is wrong?

Is pion decay mediated by another interaction?

The pion is a composite particle: does it mean that the form factors drastically change our estimates?

Is the weak interaction non universal, i.e. is the value of  $G_F$  processus dependent?

# Pathology at High Energy

What about weak scattering process, e.g.  $e\nu_e \rightarrow e\nu_e$ ?

$$\mathcal{L} = G_F J_\mu^* J^\mu \quad \text{with} \quad J^\mu = (\bar{n}\gamma^\mu p) + (\bar{e}\gamma^\mu \nu_e) + (\bar{\mu}\gamma^\mu \nu_\mu) + \dots$$

The same Fermi Lagrangian will thus also contain a term

$$G_F (\bar{e}\gamma^\mu \nu_e)(\bar{\nu}_e \gamma^\mu e)$$

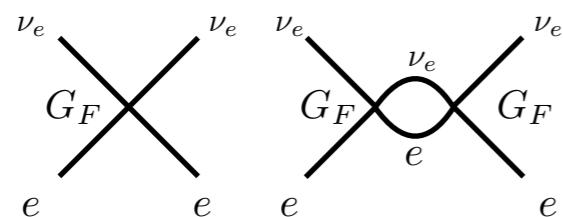
that will generate  $e$ - $\nu_e$  scattering whose cross-section can be guessed by dimensional arguments

$$\sigma \propto G_F^2 E^2$$

[mass]<sup>-2</sup>      [mass]<sup>-2×2</sup>      [mass]<sup>2</sup>

non conservation of probability  
(non-unitary theory)  
inconsistent at high energy

It means that, at high-energy, the quantum corrections to the classical contribution can be sizeable:



$$\sigma \propto G_F^2 E^2 + \frac{1}{16\pi^2} G_F^4 E^6 + \dots$$

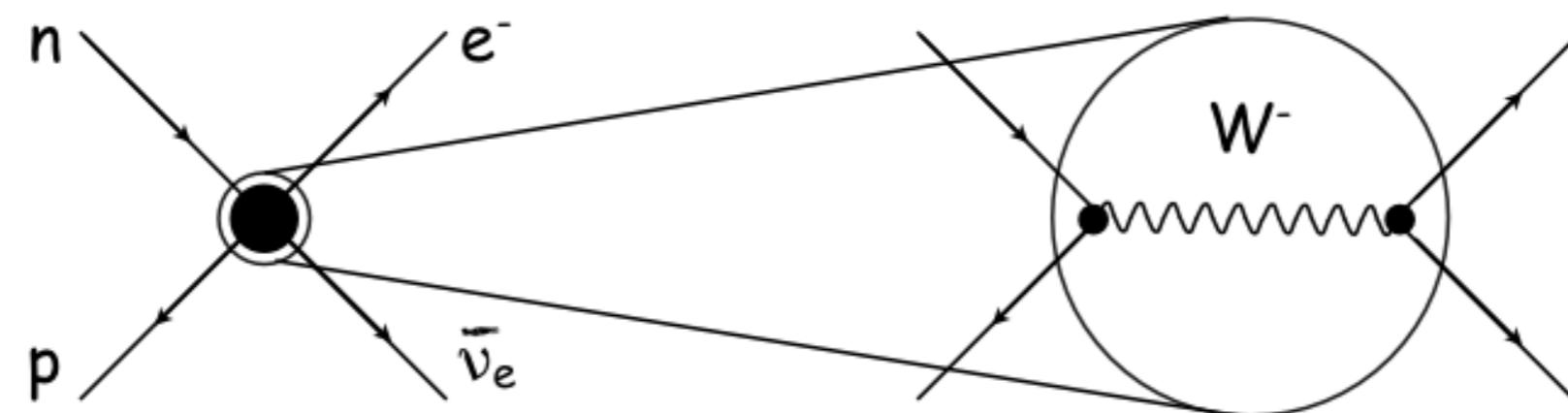
The theory becomes non-perturbative at an energy  $E_{\max} = \frac{2\sqrt{\pi}}{\sqrt{G_F}} \sim 100 \text{ GeV} - 1 \text{ TeV}$

unless new degrees of freedom appear before to change the behaviour of the scattering

# Electroweak Interactions

Low energy

High energy



$$\sigma \propto G_F^2 E^2$$

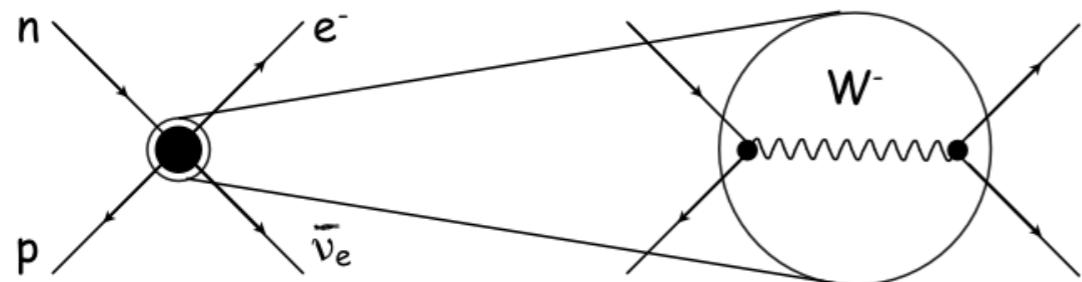
$$\sigma \propto g^4 \frac{E^2}{m_W^2 (E^2 + m_W^2)}$$

— matching —

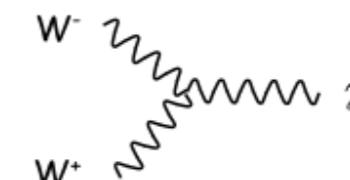
$$G_F \propto \frac{g^2}{m_W^2}$$

The Fermi interaction is not a fundamental interaction of Nature.  
It is a low energy effective interaction.

# Electroweak Interactions



charged  $W \Rightarrow$  must couple to photon:



$\Rightarrow$  non-abelian gauge symmetry  $[Q, T^\pm] = \pm T^\pm$

**1. No additional “force” (Georgi, Glashow '72)** mathematical consistency  $\Rightarrow$  **extra matter**

$SU(2)$

$$[T^a, T^b] = i\epsilon^{abc}T^c$$

$$[T^+, T^-] = Q \quad [Q, T^\pm] = + \pm T^\pm$$

$$T^\pm = \frac{1}{\sqrt{2}}(T^1 \pm iT^2)$$

$$\text{Tr}_{\text{irrep}} T^3 = 0 \quad \Rightarrow \text{extra matter}$$

$$\begin{pmatrix} X_L \\ \nu_L \\ e_L \end{pmatrix} \quad \begin{pmatrix} X_R \\ \nu_R \\ e_R \end{pmatrix}$$

$SU(1, 1)$

$$[T^+, T^-] = -Q$$

$$[Q, T^\pm] = + \pm T^\pm$$

non-compact  
unitary rep. has dim  $\infty$

$E_2$

2D Euclidean group

$$[T^+, T^-] = 0$$

$$[Q, T^\pm] = + \pm T^\pm$$

only one unitary rep.  
of finite dim. = trivial rep.

**2. No additional “matter” (Glashow '61, Weinberg '67, Salam '68):  $SU(2) \times U(1)$**

$\Rightarrow$  extra force

$$Q = T^3?$$

as Georgi-Glashow

$\Rightarrow$  extra matter

$$Q = Y?$$

$$Q(e_L) = Q(\nu_L)$$

$$Q = T^3 + Y!$$

Gell-Mann '56, Nishijima-Nakano '53

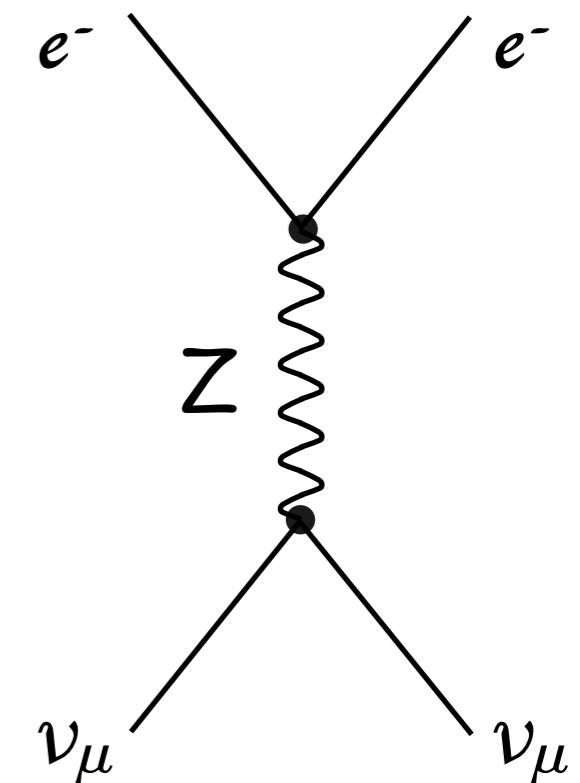
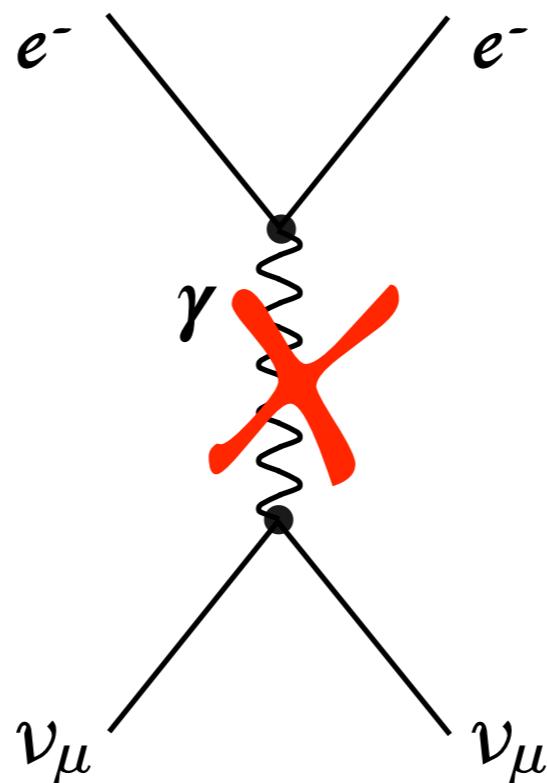
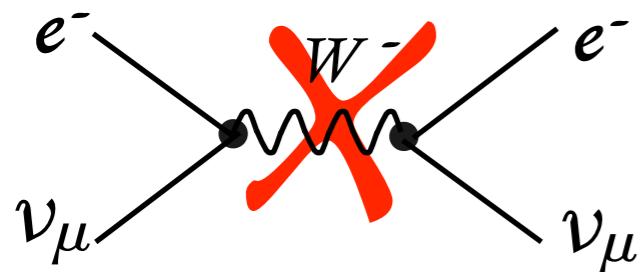
# Electroweak Interactions

**Gargamelle** experiment '73 first established the  $SU(2) \times U(1)$  structure

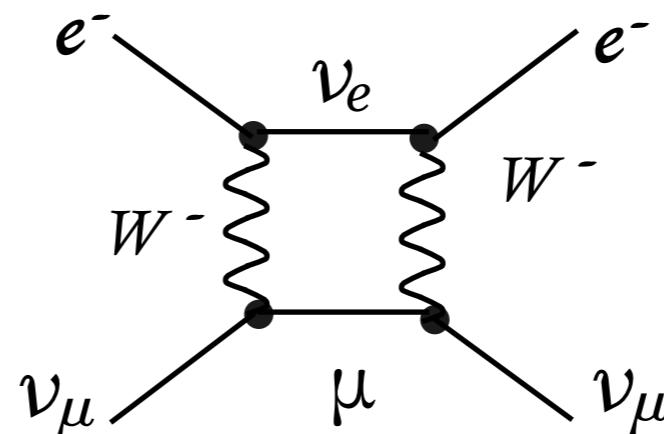
How?

rely on a particle that doesn't interact with photon to prove the existence a new neutral current process!

$$\nu_\mu e^- \rightarrow \nu_\mu e^-$$



loop-suppressed contribution from  $W$ :



# From Gauge Theory to Fermi Theory

We can derive the Fermi current-current contact interactions by “integrating out” the gauge bosons, i.e., by replacing in the Lagrangian the  $W$ 's by their equation of motion. Here is a simple derivation: (a better one should take into account the gauge kinetic term and the proper form of the fermionic current that we'll figure out tomorrow, for the moment, take it as a heuristic derivation)

$$\mathcal{L} = -m_W^2 W_\mu^+ W_\nu^- \eta^{\mu\nu} + g W_\mu^+ J_\nu^- \eta^{\mu\nu} + g W_\nu^- J_\mu^+ \eta^{\mu\nu}$$

$$J^{+\mu} = \bar{n}\gamma^\mu p + \bar{e}\gamma^\mu \nu_e + \bar{\mu}\gamma^\mu \nu_\mu + \dots \quad \text{and} \quad J^{-\mu} = (J^{+\mu})^*$$

The equation of motion for the gauge fields:  $\frac{\partial \mathcal{L}}{\partial W_\mu^+} = 0 \Rightarrow W_\mu^- = \frac{g}{m_W^2} J_\mu^-$

Plugging back in the original Lagrangian, we obtain an effective *Lagrangian* (valid below the mass of the gauge bosons):

$$\mathcal{L} = \frac{g^2}{m_W^2} J_\mu^+ J_\nu^- \eta^{\mu\nu}$$

which is the Fermi current-current interaction. The Fermi constant is given by  
(the correct expression involves a different normalisation factor)

$$G_F = \frac{g^2}{m_W^2}$$

But what is the origin of the  $W$  mass?

By the way, it is not invariant under  $SU(2)$  gauge transformation...

That's what the Higgs mechanism will take care of!

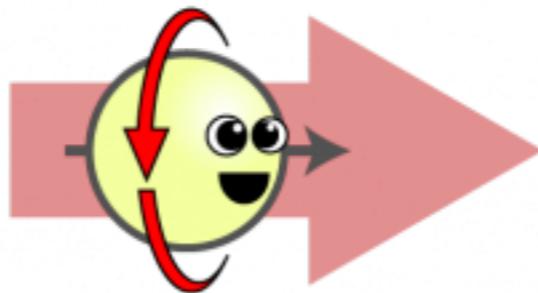
# Chirality & Masslessness

## Quantum Mechanics I.0.I

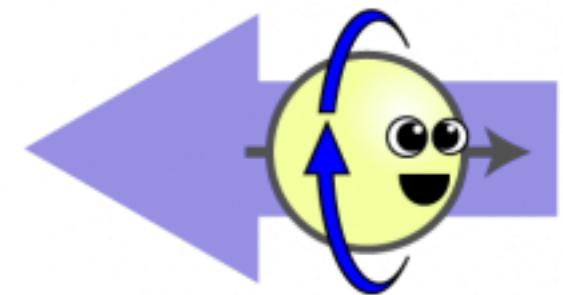
**Particle of spin  $s$  has  $2s+1$  polarisation states**

Particle spinning  
**anticlockwise** wrt its  
direction of motion

**electron**  
**has 2 polarisation**

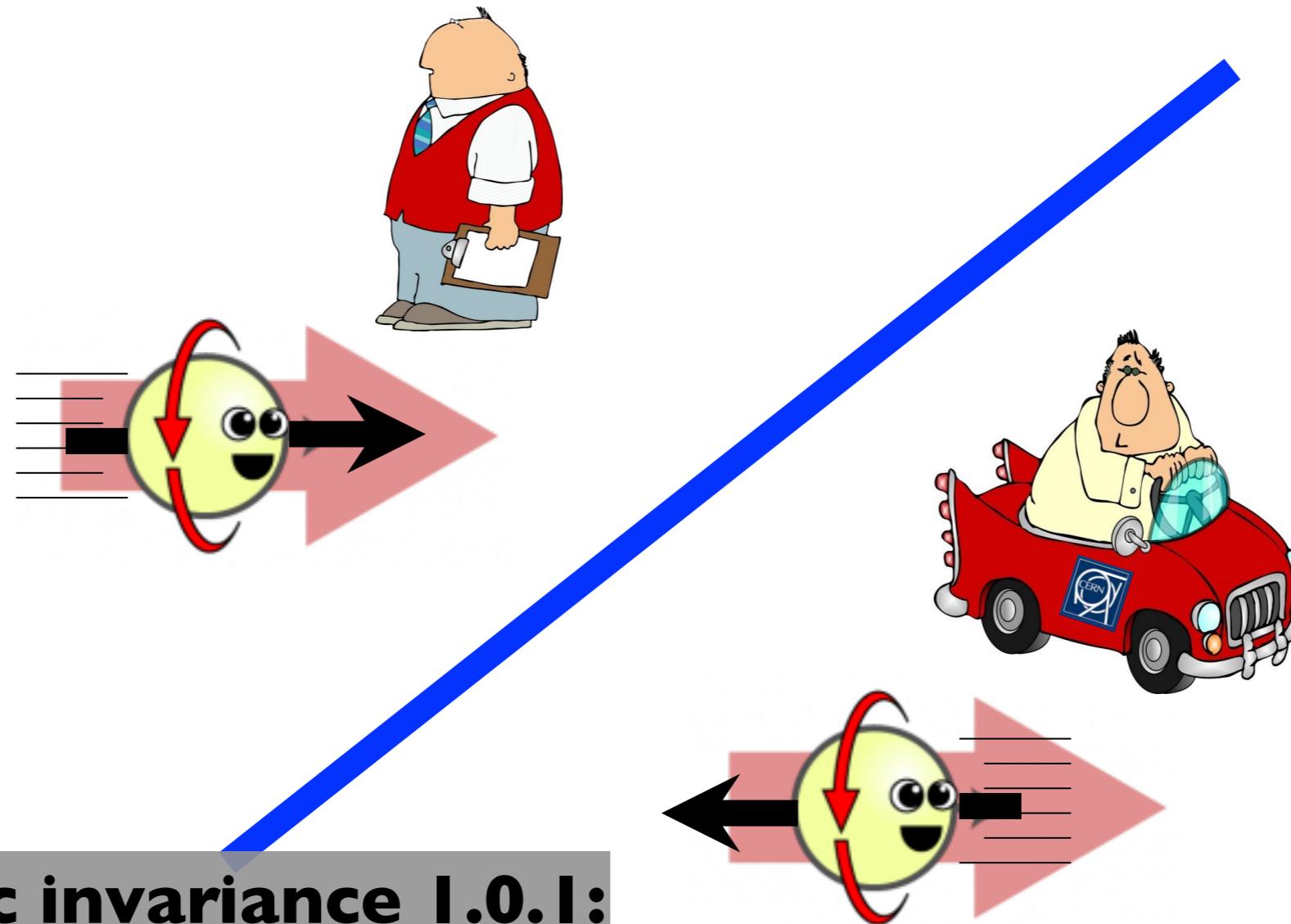


Particle spinning  
**clockwise** wrt its  
direction of motion



# Chirality & Masslessness

Picture courtesy to G. Giudice

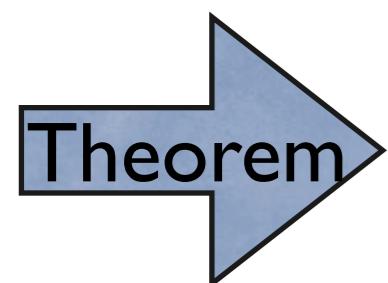


## Relativistic invariance I.0.1:

there must be no distinction for massive particles between particles spinning clockwise or anti-clockwise

[chirality operator doesn't commute with the Hamiltonian]

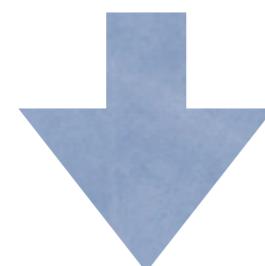
**If your theory sees a difference between  $e_L$  and  $e_R$ , either your theory is wrong or  $m_e=0$**



# Chirality of SM & Mass problem

**Weak interaction**  
(force responsible for neutron decay)  
is chiral!

[ $e_L$  and  $e_R$  are fundamentally two different particles  
Only an accident of the history of physics that they are both called electron]



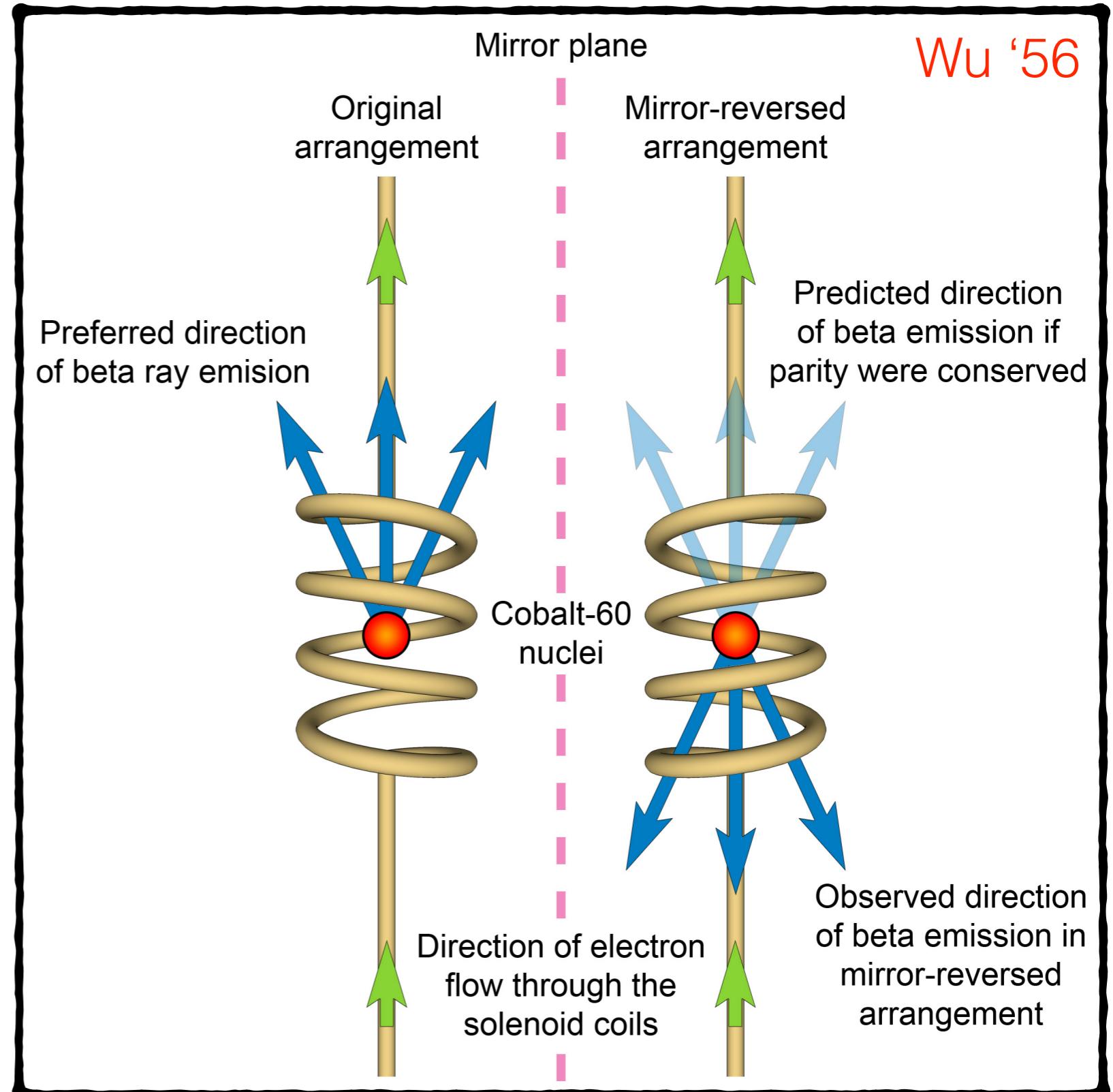
$$m_e = 0$$

but since we know it is not true, we

**need a new phenomena to generate mass:  
Higgs mechanism**

TH: Yang&Lee '56. EXP: Wu '57

Wu '56

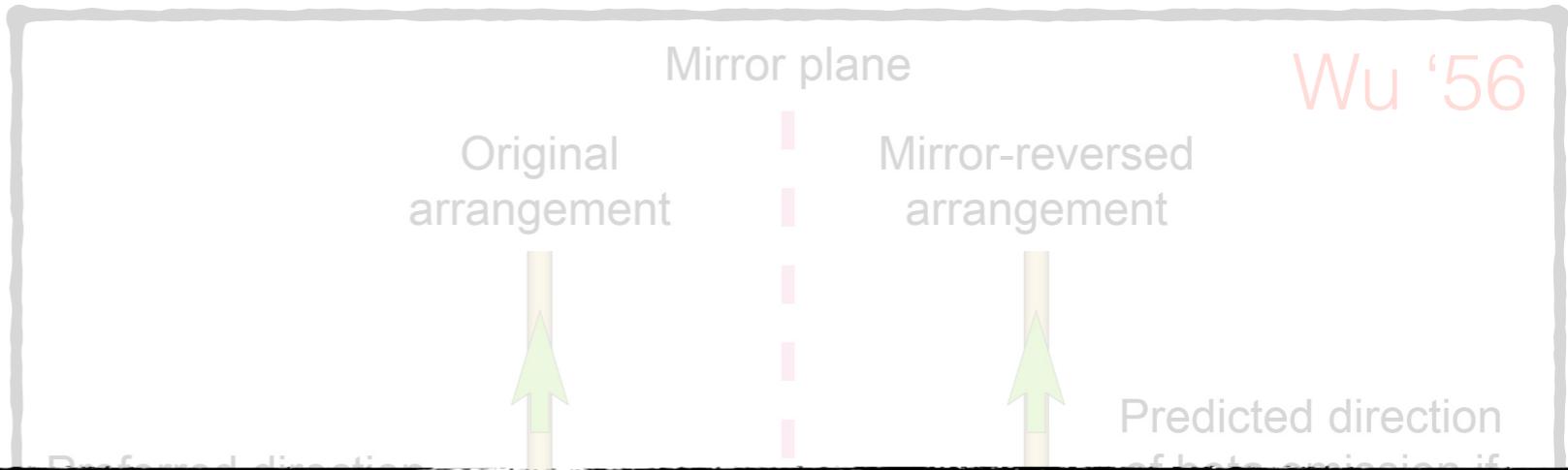


# Chirality of SM & Mass problem

TH: Yang&Lee '56. EXP: Wu '57

**Weak** interaction  
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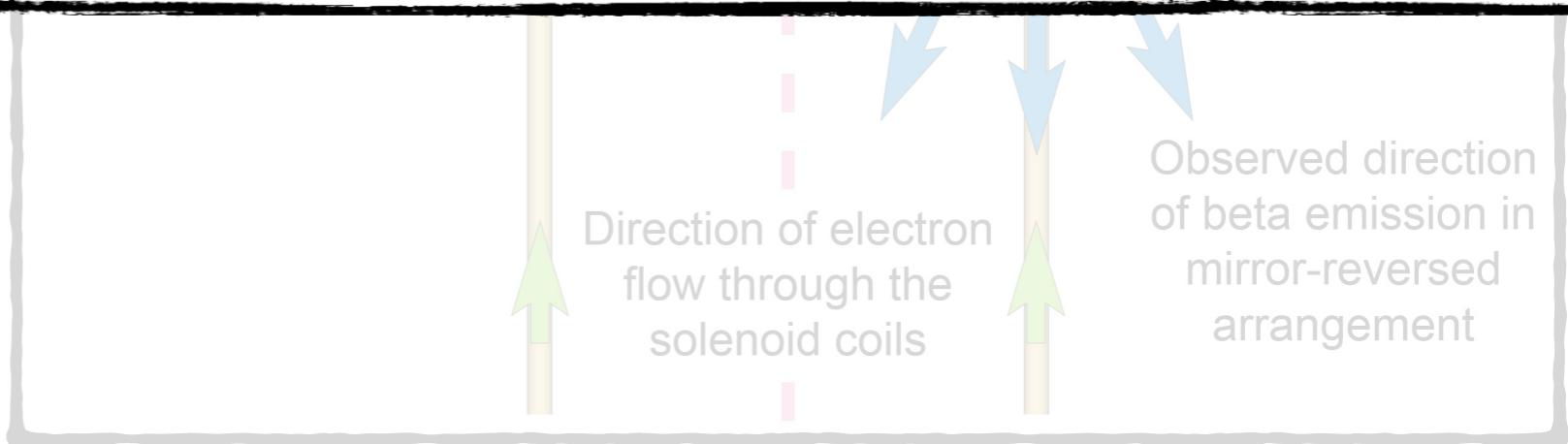
[ $e_L$  and  $e_R$  are fundamentally



Dextrorotation and Levorotation are essential for life to develop.  
To the best of our knowledge,  
in **molecular biology**, chirality seems an **emergent** property.  
At least, there is no clear evidence that it follows from chirality of the weak interactions.  
Are the chiral nature of the **weak** interactions **emergent** too?  
Some models of grand unification predict it. But we still don't know for sure.

but since we know it is not true, we

**need a new  
phenomena to  
generate mass:  
Higgs mechanism**



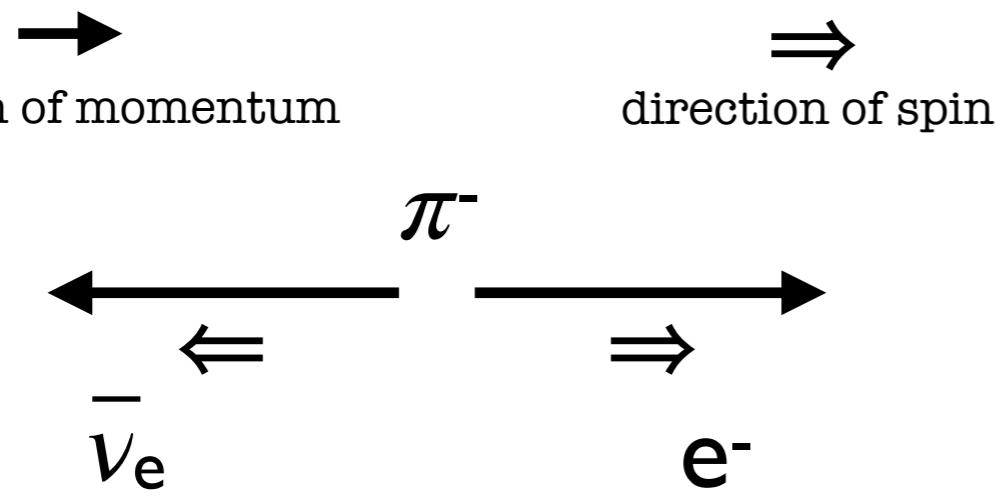
# SM is a Chiral Theory

Weak interactions maximally violates P



**Weak interactions act only on LH particles (and RH anti-particles)**

this property has an important consequence (aka selection rule) for pion decay



Conservation of momentum and spin  
imposes to have a RH  $e^-$

Weak decays proceed only w/ LH  $e^-$   
So the amplitude is prop. to  $m_e$

$$\mathcal{L}_{\text{Dirac}} = \bar{\psi}_L \gamma^\mu \partial_\mu \psi_L + \bar{\psi}_R \gamma^\mu \partial_\mu \psi_R + m (\bar{\psi}_L \psi_R + \bar{\psi}_R \psi_L)$$

Lorentz structure  
of fermion mass

$$\frac{\Gamma(\pi^- \rightarrow e^- \bar{\nu}_e)}{\Gamma(\pi^- \rightarrow \mu^- \bar{\nu}_\mu)} \propto \frac{m_e^2}{m_\mu^2} \sim 2 \times 10^{-5} \sim 10_{\text{obs}}^{-4}$$

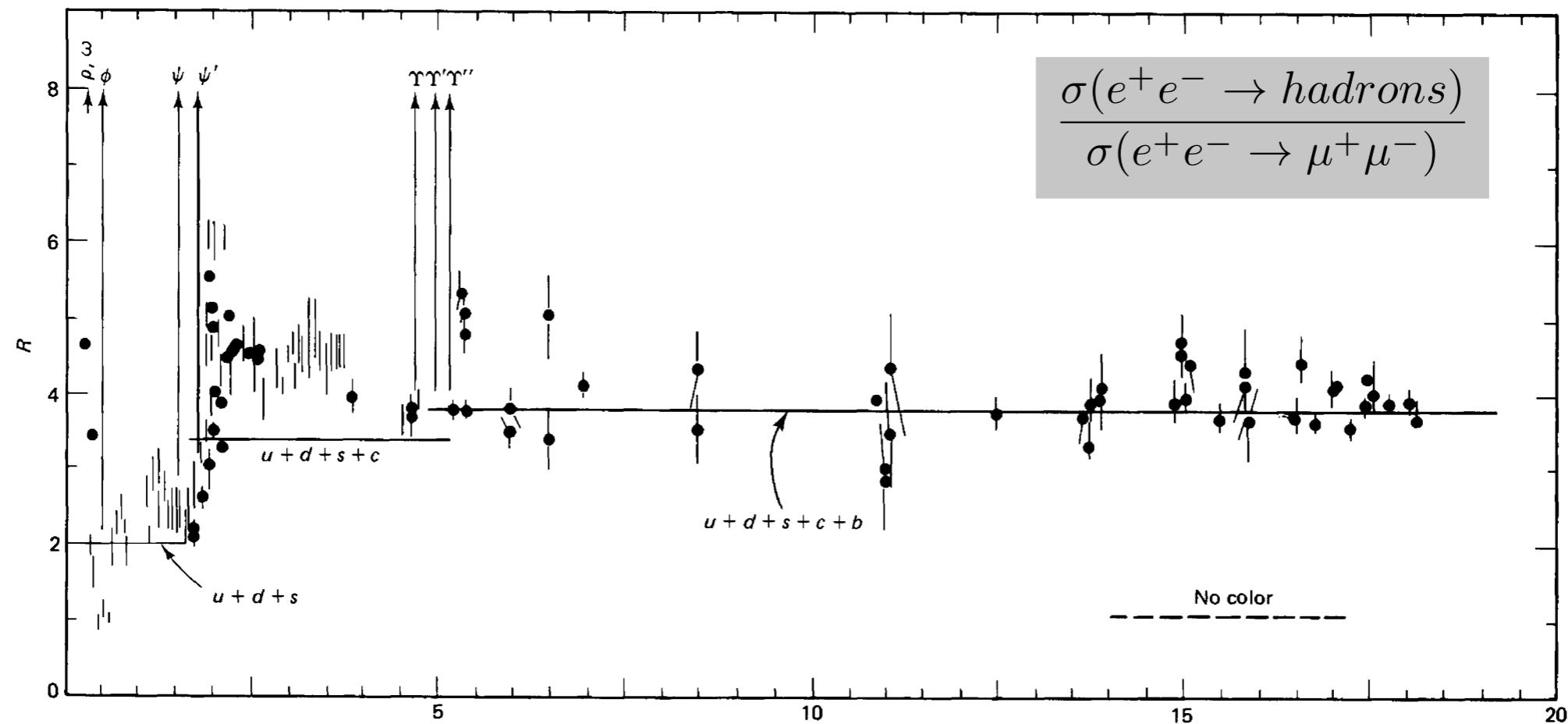
Extra phase-space factor

# SU(3) QCD

Experiments in the 60's revealed the internal structure of the neutrons and protons  
Gell-Mann and others proposed that they are made of "quarks"

Up quark (up, charm, top): spin-1/2,  $Q=2/3$   
Down quark (down, strange, bottom): spin-1/2,  $Q=-1/3$

SU(2) weak symmetry that changes neutrino into electron also changes up-quark into down-quark



counts the number of quarks and gives their electric charges  
another remarkable feature: at high energy, the quarks behaves like muons,  
i.e., not sensitive to strong interactions

**Asymptotic freedom of QCD!**

# SU(3) QCD

Deep inelastic experiments in the 60's revealed the internal structure of the neutrons and protons  
Gell-Mann and others proposed that they are made of “**quarks**”

Up quark: spin-1/2, Q=2/3  
Down quark: spin-1/2, Q=-1/3

**SU(2)** weak symmetry that changes neutrino into electron also changes up-quark into down-quark

**Quarks** carry yet another quantum number: “**colour**”

There are 3 possible colours and Nature is colour-blind, i.e., Lagrangian should remain the same when the colours of the quarks are changed, i.e., when we perform a rotation in the colour-space of quarks.

$$Q^a \rightarrow U^a{}_b Q^b \quad \begin{array}{l} \text{U: 3x3 matrix satisfying } U^\dagger U = 1_3 \\ \text{such that the quark kinetic term is invariant} \end{array} \quad \text{SU}(3)$$

**hadrons** (spin-1/2, #hadronic=1):  $p = uud$        $n = udd$

**mesons** (spin-0, #hadronic=0):  $\pi^0 = \frac{u\bar{u} + d\bar{d}}{\sqrt{2}}$        $\pi^+ = u\bar{d}$        $\pi^- = d\bar{u}$

(Each quark carries a baryon number =1/3)

There are (heavier) quarks and hence other baryons and mesons

**All the interactions of the SM preserve baryon and lepton numbers**

$$\mu \rightarrow e\nu_\mu\bar{\nu}_e \quad n \rightarrow p e\bar{\nu}_e \quad \pi^- \rightarrow \mu^-\bar{\nu}_\mu \quad \pi^0 \rightarrow \gamma\gamma \quad p \cancel{\rightarrow} \pi^0\bar{e}$$

# The Standard Model: Interactions

- $U(1)_Y$  electromagnetic interactions

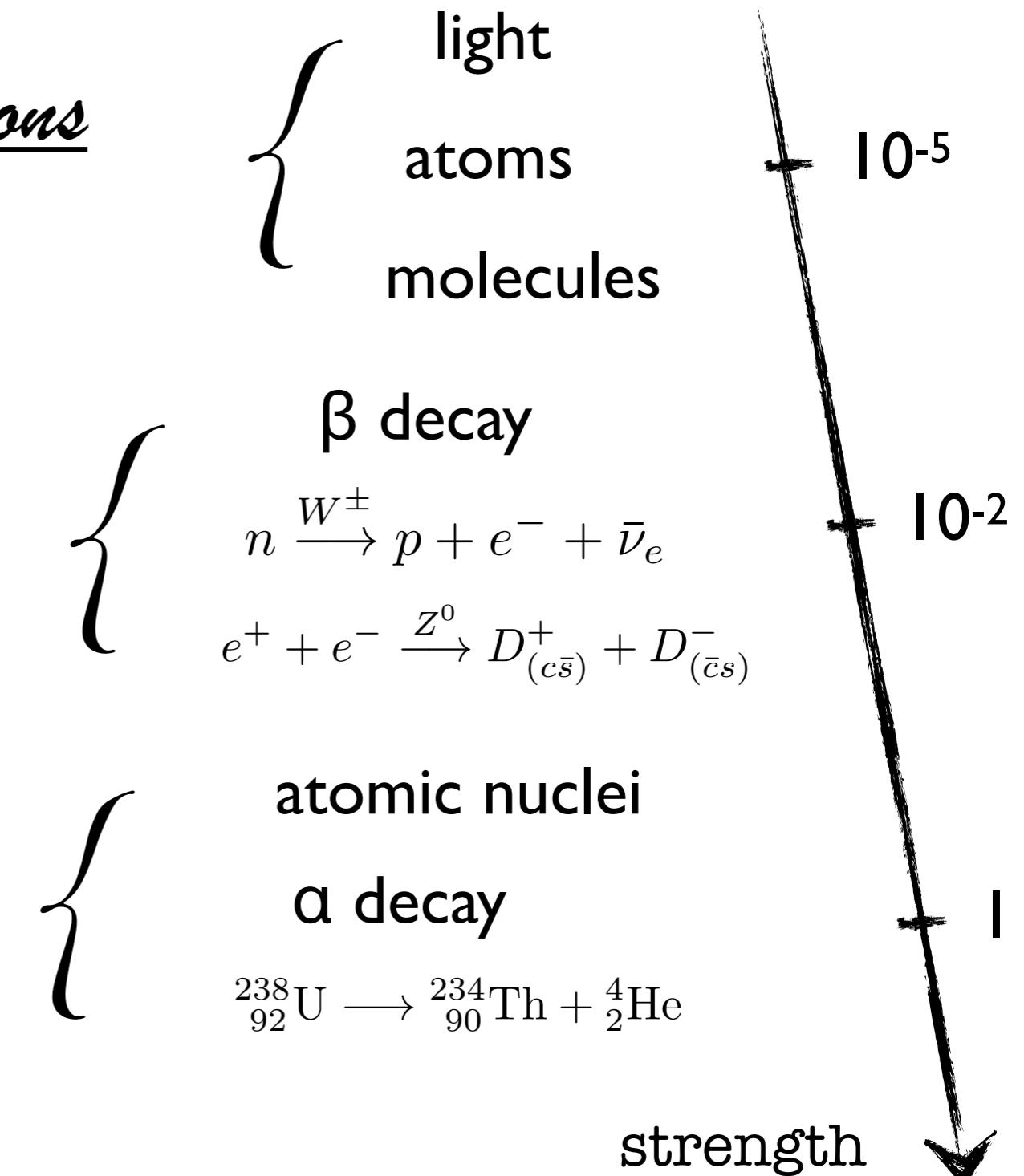
Photon  $\gamma$

- $SU(2)_L$  weak interactions

bosons  $W^\pm, Z^0$

- $SU(3)_c$  strong interactions

gluons  $g^a$



# Technical Details for Advanced Students

# Chirality

## Chirality matrix

$$\gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3$$

A few remarkable properties

$$(\gamma^5)^2 = 1_4$$

$$\{\gamma^5, \gamma^\mu\} = 0$$

$$\gamma^{5\dagger} = \gamma^5 = -\gamma^0\gamma^5\gamma^0$$

## Chiral/Weyl spinor

A **chiral/Weyl** spinor is an eigenvector of the chirality matrix  $\psi_{L,R} = \pm\gamma^5\psi_{L,R}$

From the Lorentz-transformation law of a spinor, it is obvious that the chirality condition is frame-independent

A Dirac spinor can also be written as a sum of two chiral spinors

$$\psi = \frac{1}{2} (1_4 + \gamma^5) \psi + \frac{1}{2} (1_4 - \gamma^5) \psi \equiv \psi_L + \psi_R$$

# Charge conjugation

In general,  $\psi$  and  $\psi^*$  do not transform in the same way under Lorentz transformations  
and the naive reality condition  $\psi = \psi^*$  is frame dependent

But it is possible to find a matrix  $C$ , called charge conjugation matrix, such that

$$\psi \quad \text{and} \quad \psi_C = C \psi^*$$

transform in the same way under Lorentz transformations

The matrix  $C$  needs to satisfy  $C\gamma^* = -\gamma^\mu C$

In the Dirac and Weyl representations,  $C = i\gamma^2$

In the Majorana representation,  $C = 1_4$

Basic properties of the charge conjugation matrix:  $C^2 = 1_4$ ,  $C^\dagger = C$ ,  $C^* = C$

The charge conjugated spinor,  $\psi_C$ , satisfies the same Dirac equation as  $\psi$ , with the same mass but opposite electric charge (when the spinor is minimally coupled to a U(1) gauge field)

A **Majorana** spinor satisfies the (Lorentz invariant!) condition  $\psi = \psi_C$

Note that in 4D, a spinor cannot be simultaneously chiral and Majorana

# Dirac and Majorana Masses

By construction, the following two mass terms in the Lagrangian are Lorentz-invariant

**Dirac mass:**  $\mathcal{L}_{\text{Dirac}} = m \bar{\psi} \psi$  (conserves fermion number)

**Majorana mass:**  $\mathcal{L}_{\text{Majorana}} = m \bar{\psi}_C \psi$  (changes fermion number by 2)

These two mass terms have different a chirality structure

$$\mathcal{L}_{\text{Dirac}} = m (\bar{\psi}_L \psi_R + \bar{\psi}_R \psi_L)$$

$$\mathcal{L}_{\text{Majorana}} = m (\bar{\psi}_{L_C} \psi_L + \bar{\psi}_{R_C} \psi_R)$$

A chiral fermion can have a Majorana mass  
A Dirac mass requires spinors of opposite chirality

Whether or not a Dirac or a Majorana mass can be included in the Lagrangian depends on transformation laws of the spinors under the gauge transformations

Within the SM (with the Higgs field), a Dirac mass can be written for the charged leptons and the quarks while a Majorana mass can be written for the neutrinos.

# Higgs Lifetime “Computation”

Using dimensional analysis arguments,  
compute the Higgs boson lifetime (or its inverse aka as the Higgs decay width)

— Hints —

Higgs couplings proportional are proportional to the mass of the particles it couples to.  
It will therefore decay predominantly decay into the heaviest particle that is lighter than  $m_H/2$

$$\Gamma \sim \frac{1}{8\pi} \left( \frac{m_b}{v} \right)^2 m_h \sim \frac{1}{10} \left( \frac{4 \text{ GeV}}{246 \text{ GeV}} \right)^2 125 \text{ GeV} \sim 1 \text{ MeV}$$

phase space      couplings to b-quark      dimensionally  $[\Gamma]_m=1$

$\rightarrow \quad \tau \sim 10^{-22} \text{ s}$

E	T	L
1eV	$10^{-16} \text{ s}$	$10^{-7} \text{ m}$

Putting all factors and considering the other decay modes, Higgs width = 4MeV in the SM

(for Z gauge boson:  $\Gamma_Z = \frac{7}{48\pi} g^2 m_Z \sim 2 \text{ GeV}$     i.e.     $\tau_Z \sim 10^{-25} \text{ s}$  )