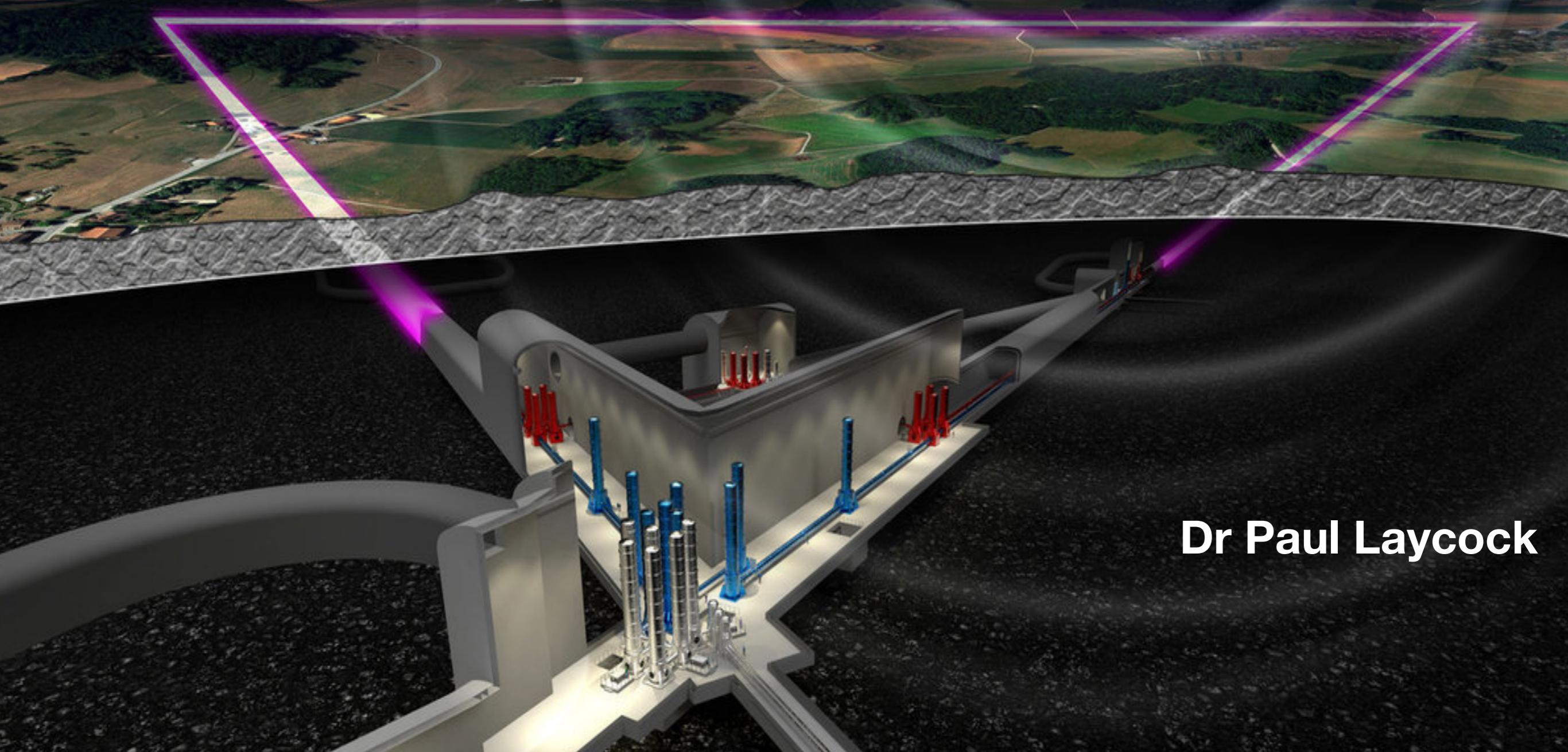
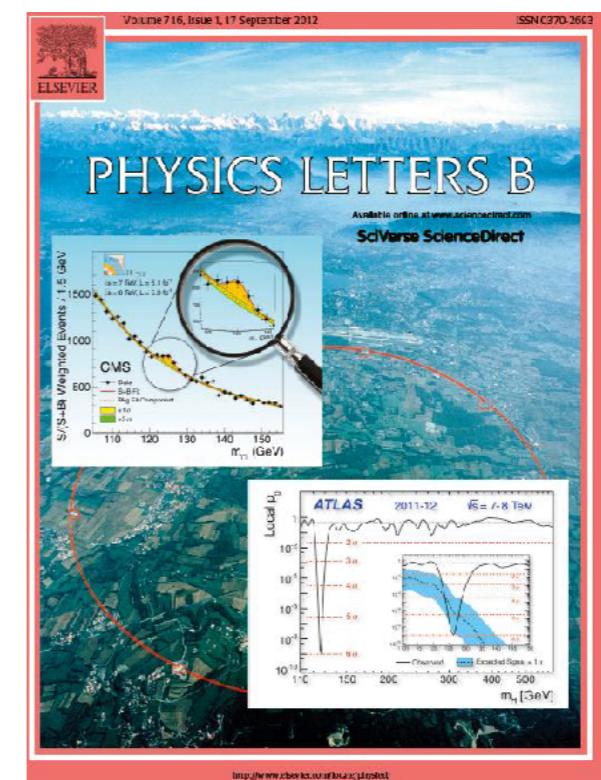
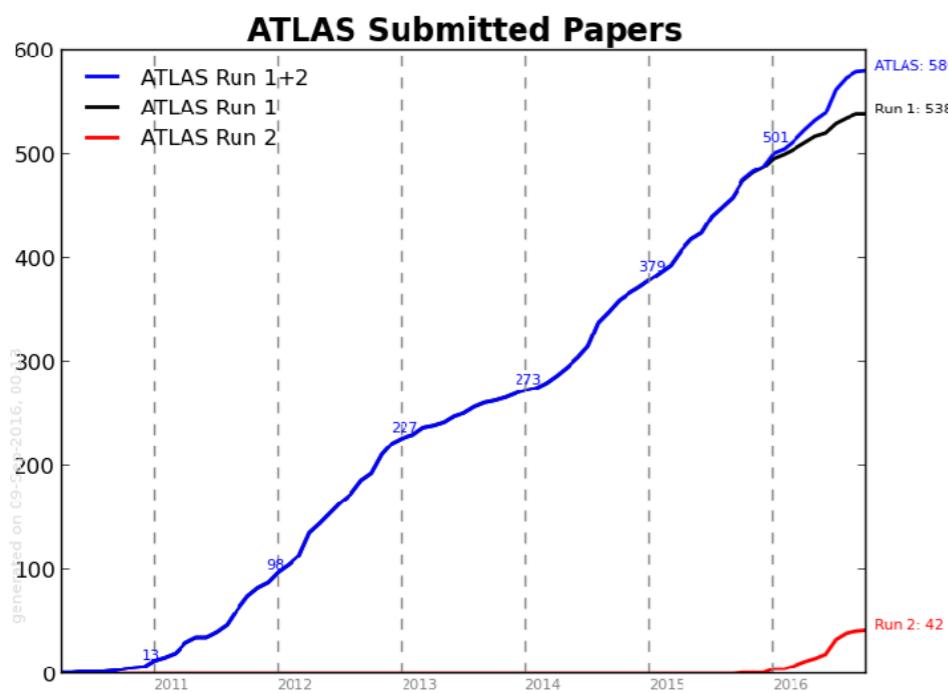
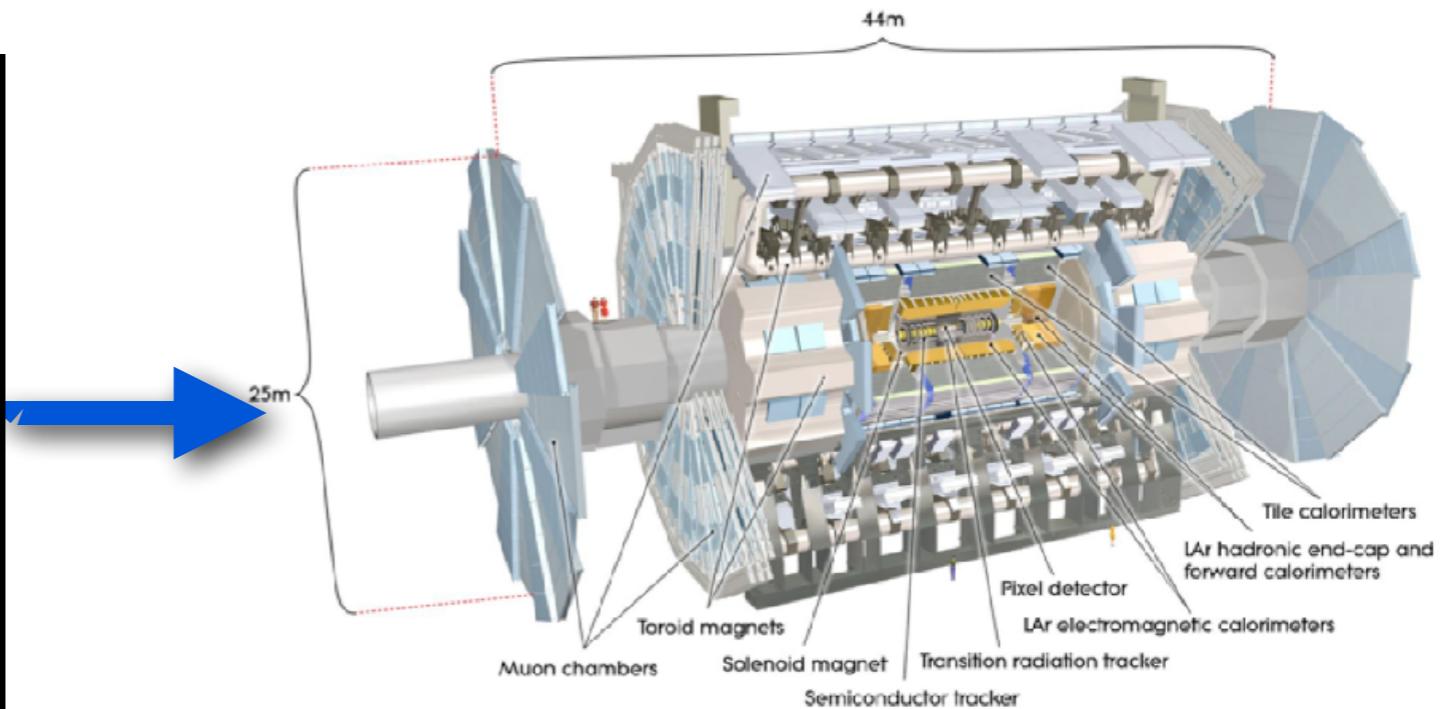
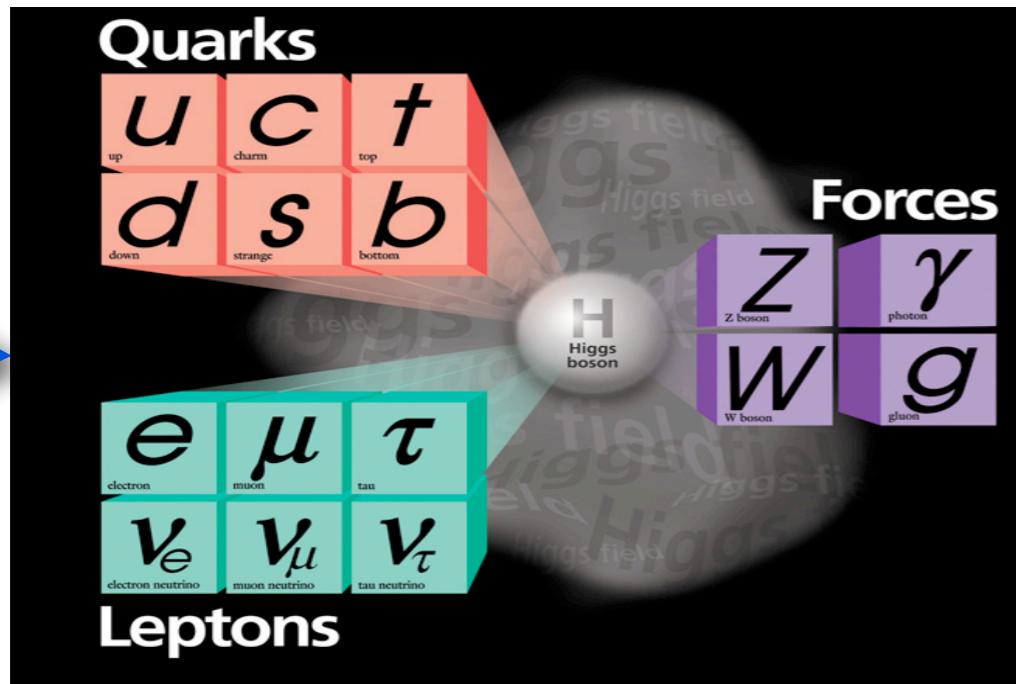


# From Raw data to Physics Results (3/3)



Dr Paul Laycock

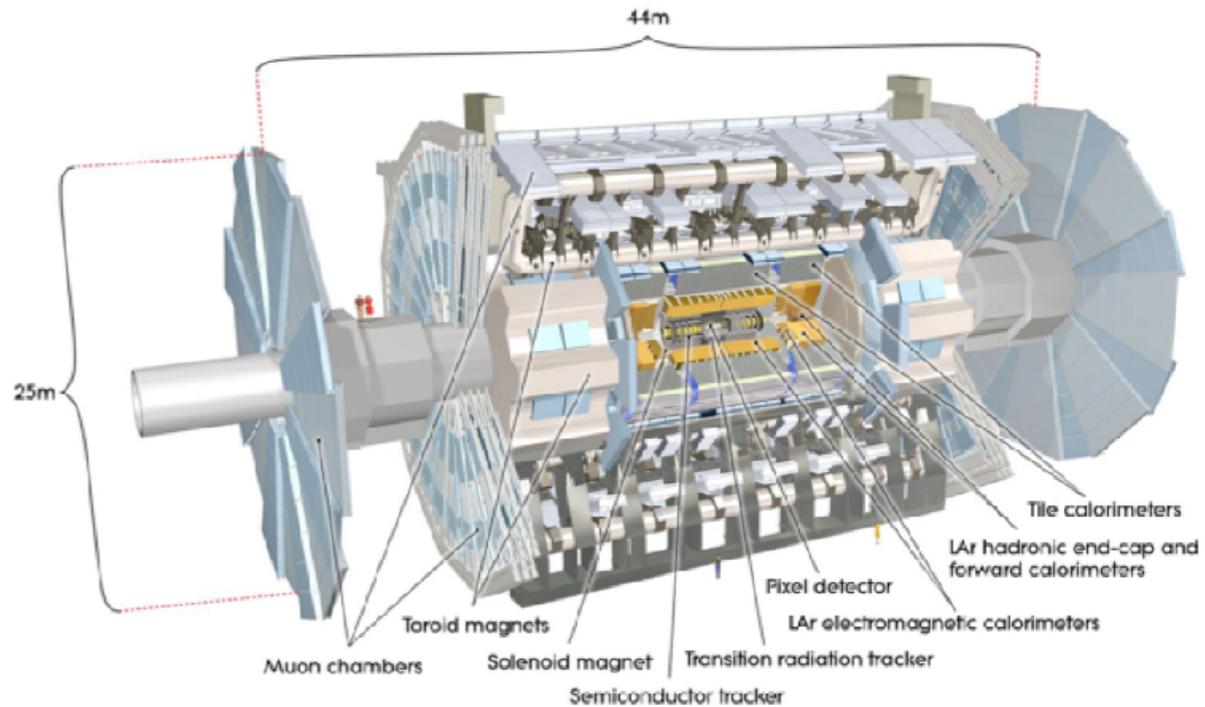
# The particle physics cycle



# Course outline

## • Lecture 1

- The journey of raw data from the detector to a publication

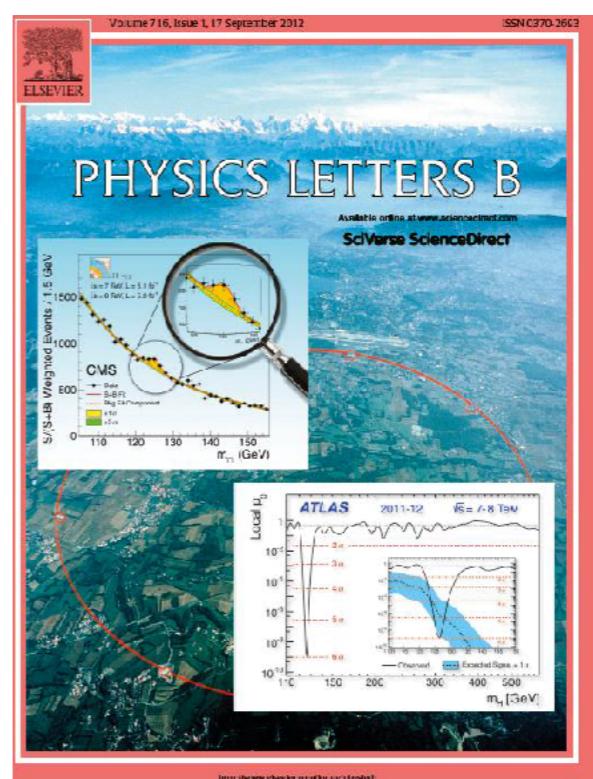


## • Lecture 2

- How we reconstruct fundamental physics processes from raw detector data

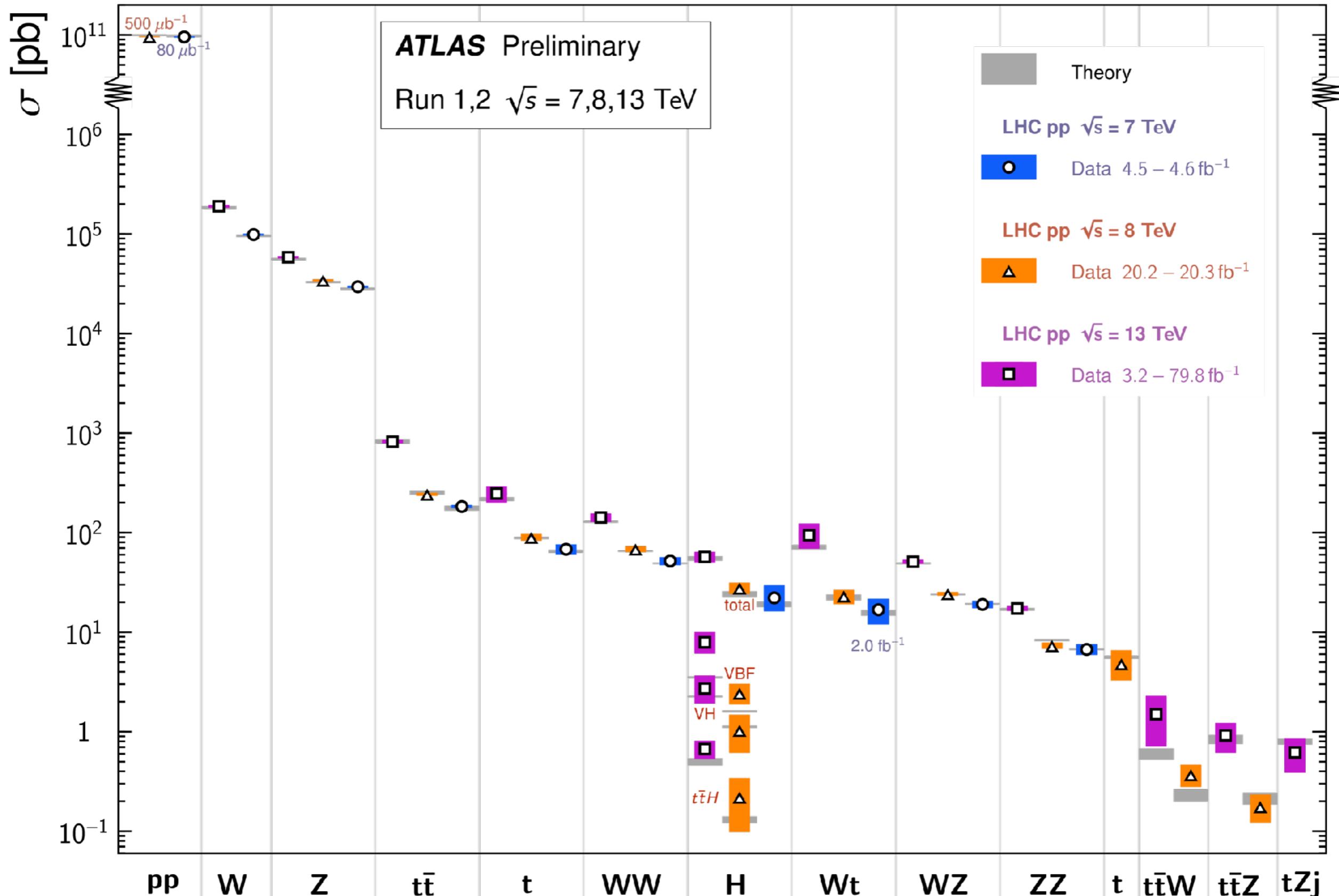
## • Lecture 3

- How we extract our signals from the mountain of data, finding needles in the haystack



# Standard Model Total Production Cross Section Measurements

Status: July 2018



pp    W    Z    t̄t    t    WW    H    Wt    WZ    ZZ    t    t̄W    t̄t̄Z    t̄Zj

t-chan

s-chan

# Measuring cross sections

$$\sigma = \frac{N}{L}$$

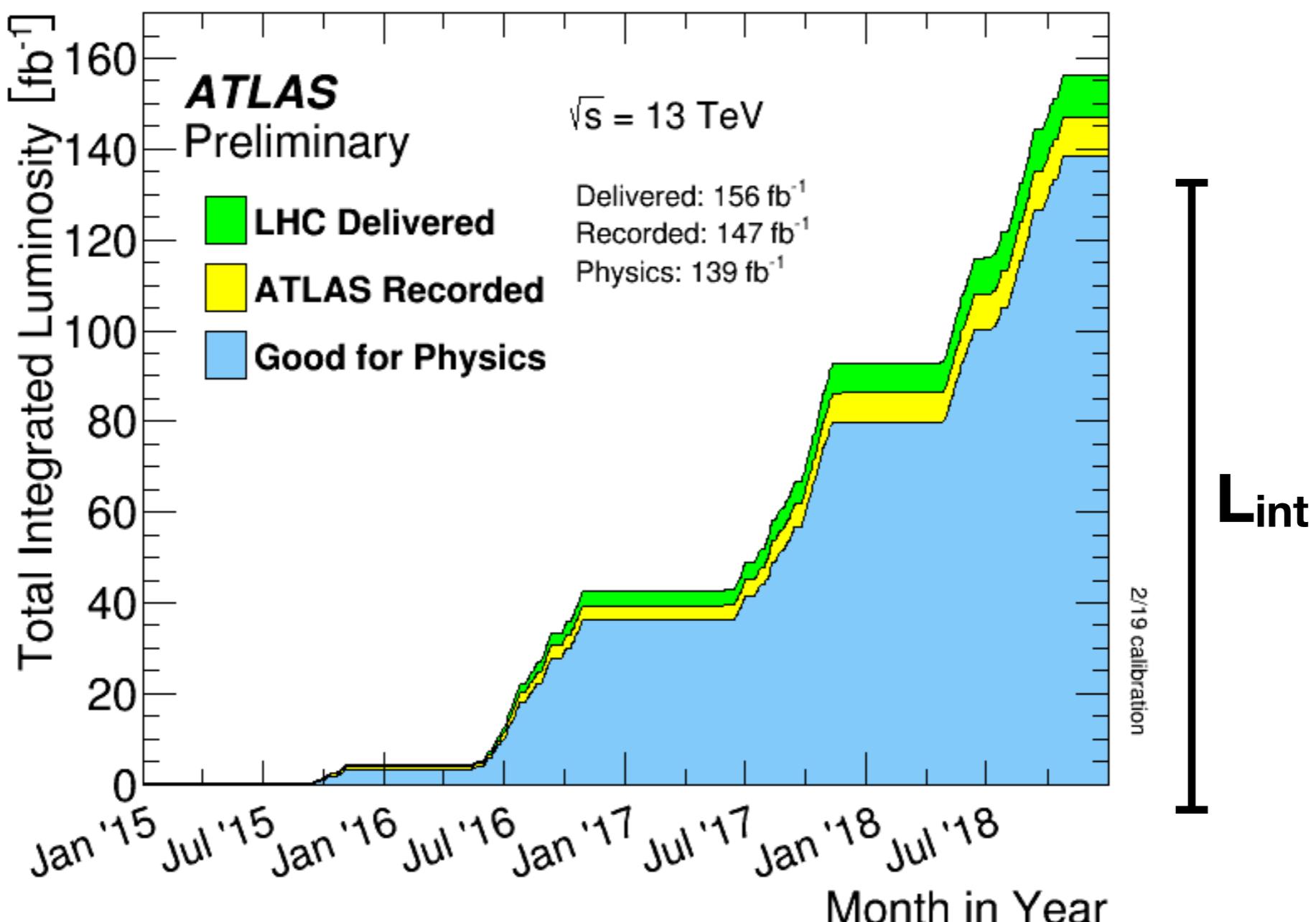
- The cross section for a process is defined as the number of events divided by luminosity

# Measuring cross sections

$$\sigma = \frac{N}{L_{int}}$$

- The cross section for a process is defined as the number of events divided by the integrated luminosity,  $L_{int}$ , which measures how much data we have collected

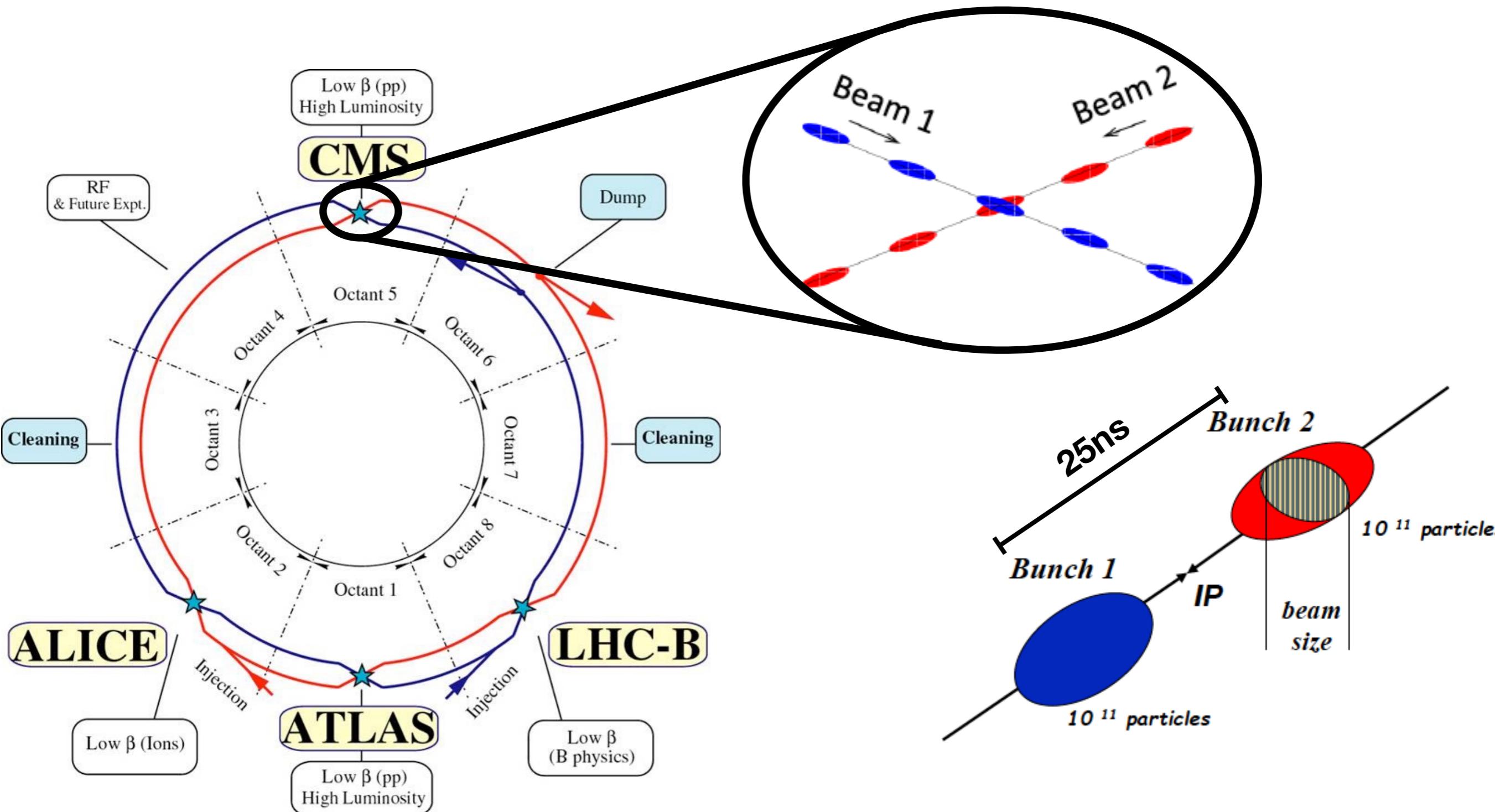
# ATLAS Luminosity



- **Question:** Why does ATLAS record less data than the LHC delivers?
- How do we know the *integrated luminosity delivered*?

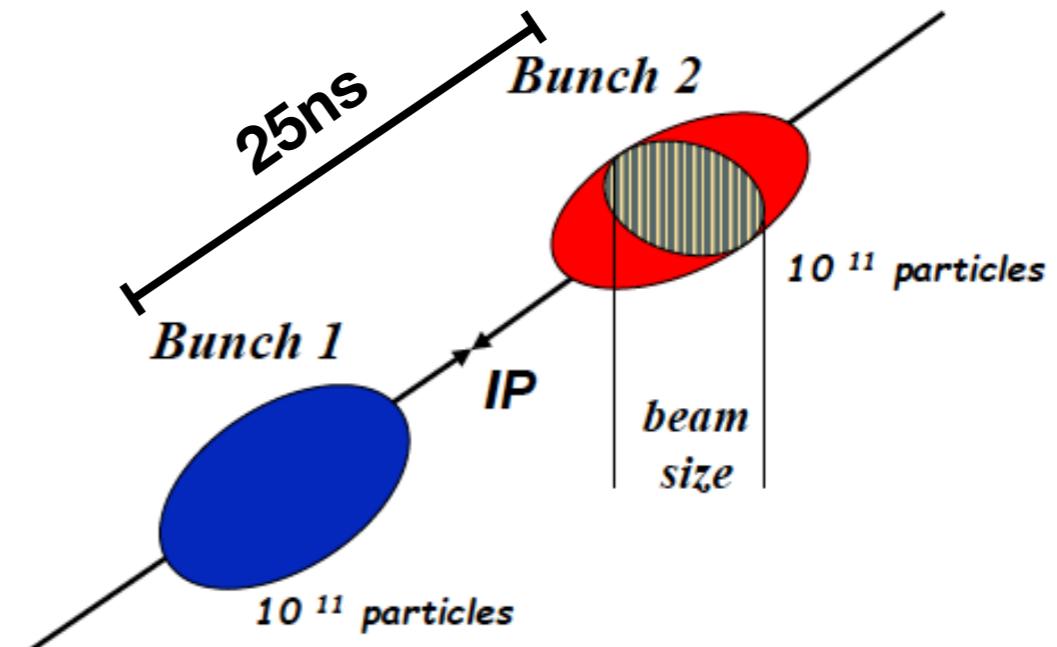
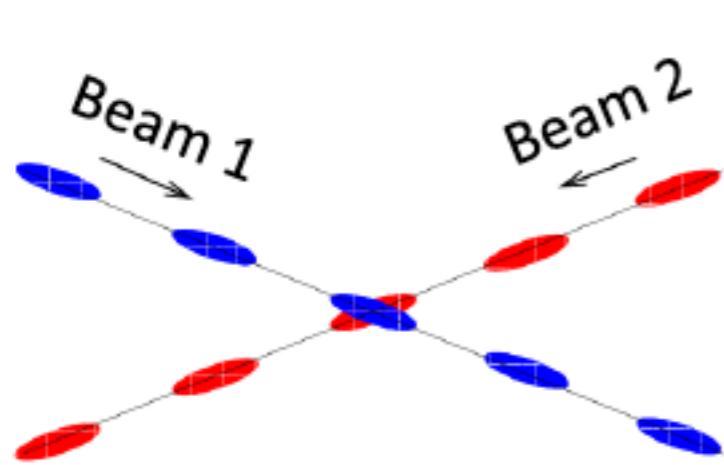
# LHC collisions

Figures adapted from Michaela Schaumann's third lecture (11/07/19) on "Particle Accelerators and Beam Dynamics"



- The LHC accelerates **bunches of  $10^{11}$  protons** separated by 25ns gaps

# Measuring Luminosity at the LHC



- Ingredients for a measurement of the luminosity
  - Measuring the **size** of the beams (for a certain LHC configuration)
    - This requires a dedicated measurement where we scan the beams across each other in the **horizontal** and **vertical** directions - a **van der Meer scan**
  - Measuring the beam **currents** in each bunch
    - This is done during collisions, integrating all of the bunch currents and knowing their size, we can calculate the luminosity
  - Make many **cross checks** because this is such a crucial measurement

# Measuring cross sections

$$\sigma = \frac{N}{L_{int}}$$

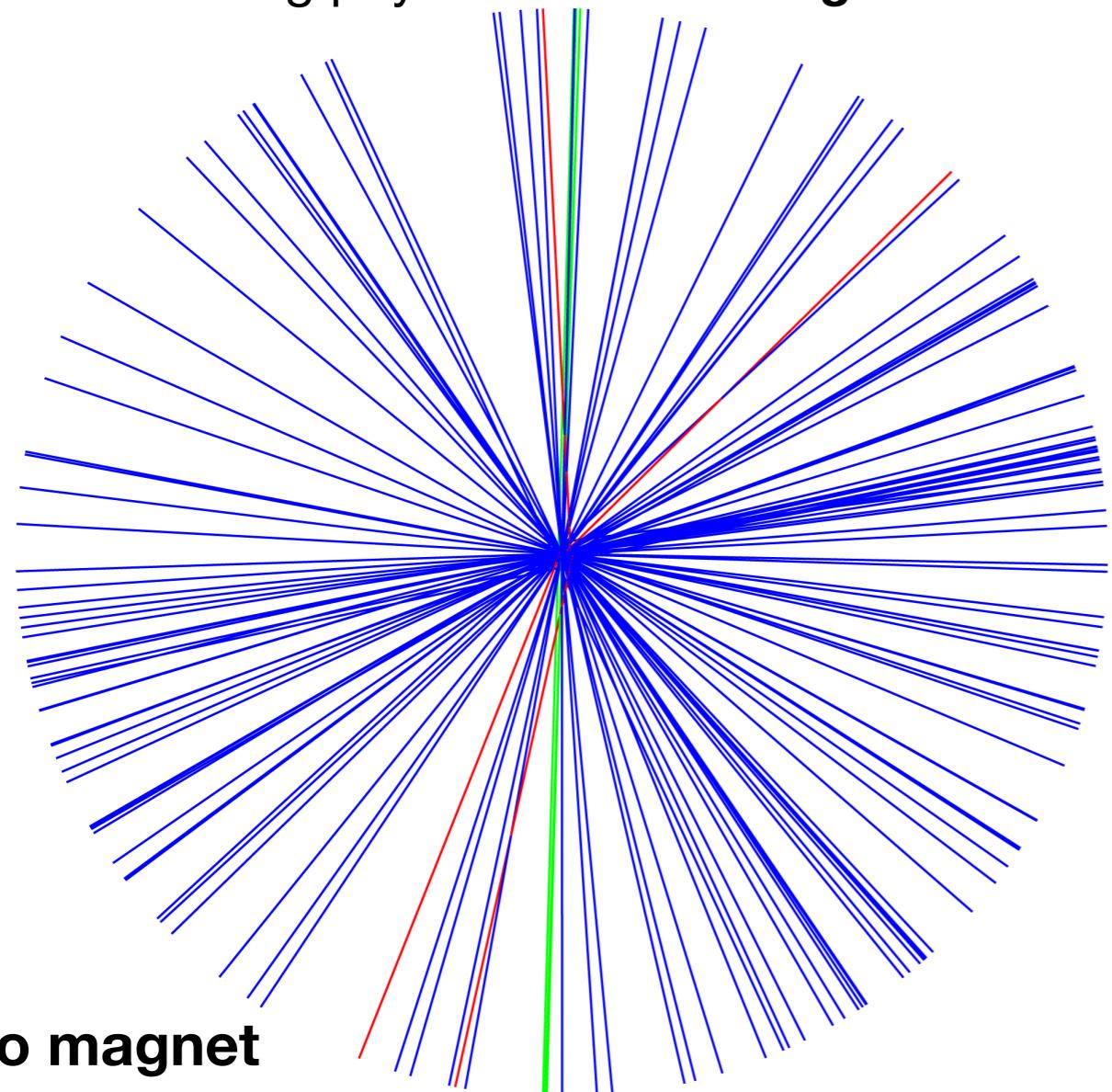
$$\sigma = \frac{N_{obs}}{A \cdot \epsilon \cdot L_{int}}$$

- The cross section for a process is defined as the number of events divided by the integrated luminosity,  $L_{int}$ , which measures how much data we have collected
- $N_{obs}$  in data needs to be corrected for the detector acceptance,  $A$ , for selecting those events. The reconstruction efficiency,  $\epsilon$ , is a product of all of the efficiencies that we need to measure and ensure that they are the same in our data and simulation

***Did I mention that simulation is important ?***

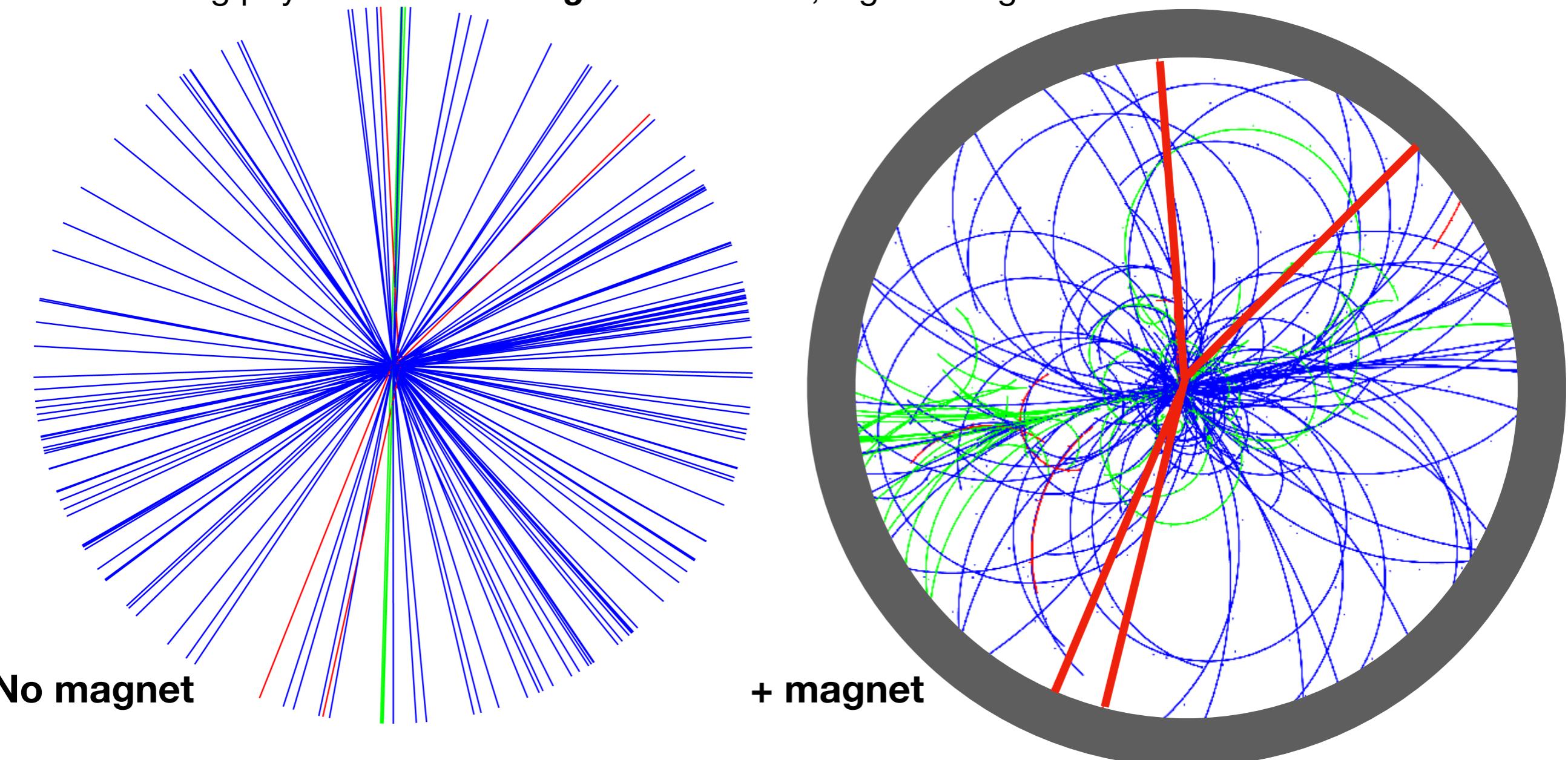
# Before the detector, came the simulation

- When designing detectors, we ***simulate detector response*** to physics of interest
- Adding a ***solenoid magnet*** makes it possible to measure momentum (and charge) in our tracker by measuring curvature in the transverse plane
- Interesting physics is often at ***high momentum***, e.g. four high momentum muon tracks here



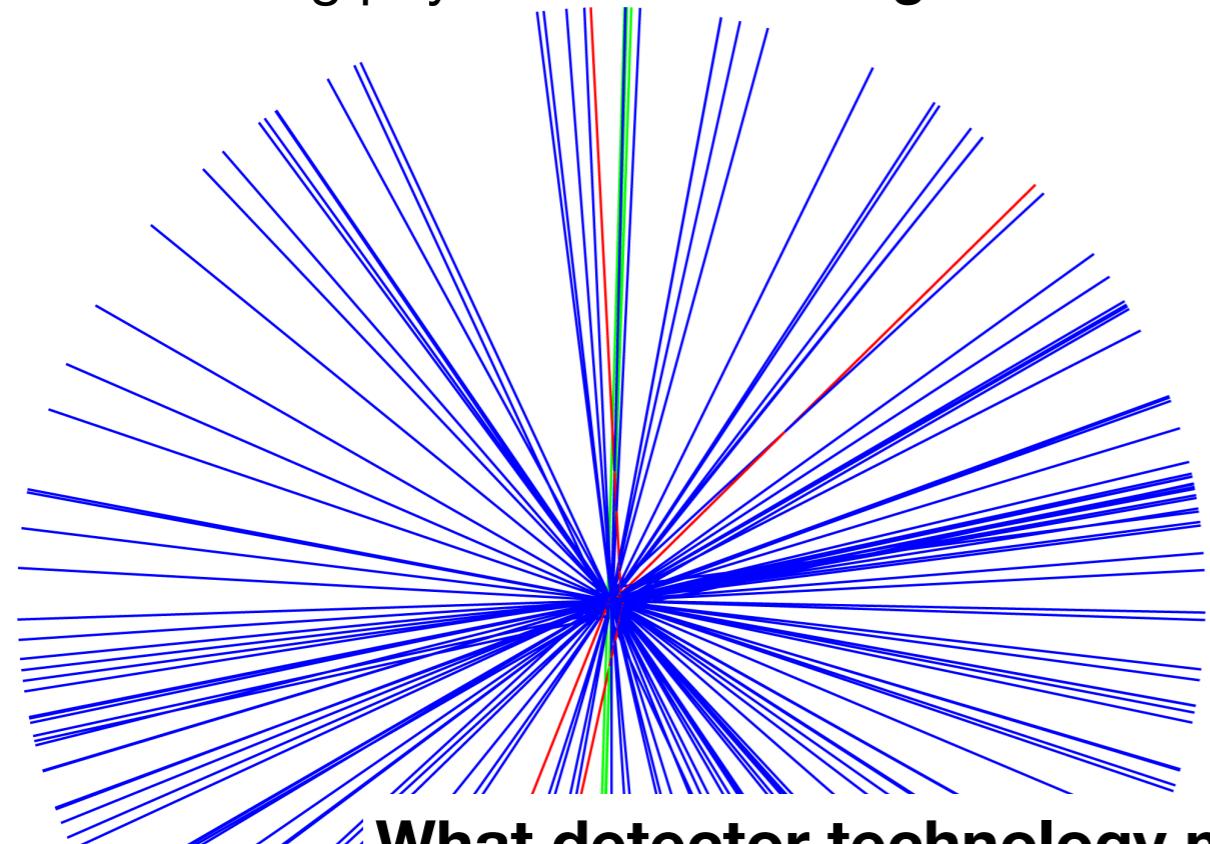
# Before the detector, came the simulation

- When designing detectors, we **simulate detector response** to physics of interest
- Adding a **solenoid magnet** makes it possible to measure momentum (and charge) in our tracker by measuring curvature in the transverse plane
- Interesting physics is often at **high momentum**, e.g. four high momentum muon tracks here



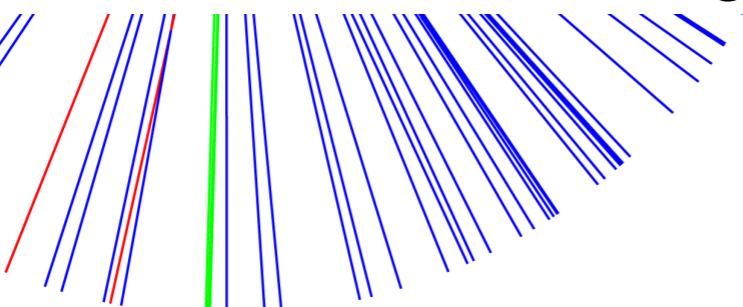
# Before the detector, came the simulation

- When designing detectors, we **simulate detector response** to physics of interest
- Adding a **solenoid magnet** makes it possible to measure momentum (and charge) in our tracker by measuring curvature in the transverse plane
- Interesting physics is often at **high momentum**, e.g. four high momentum muon tracks here

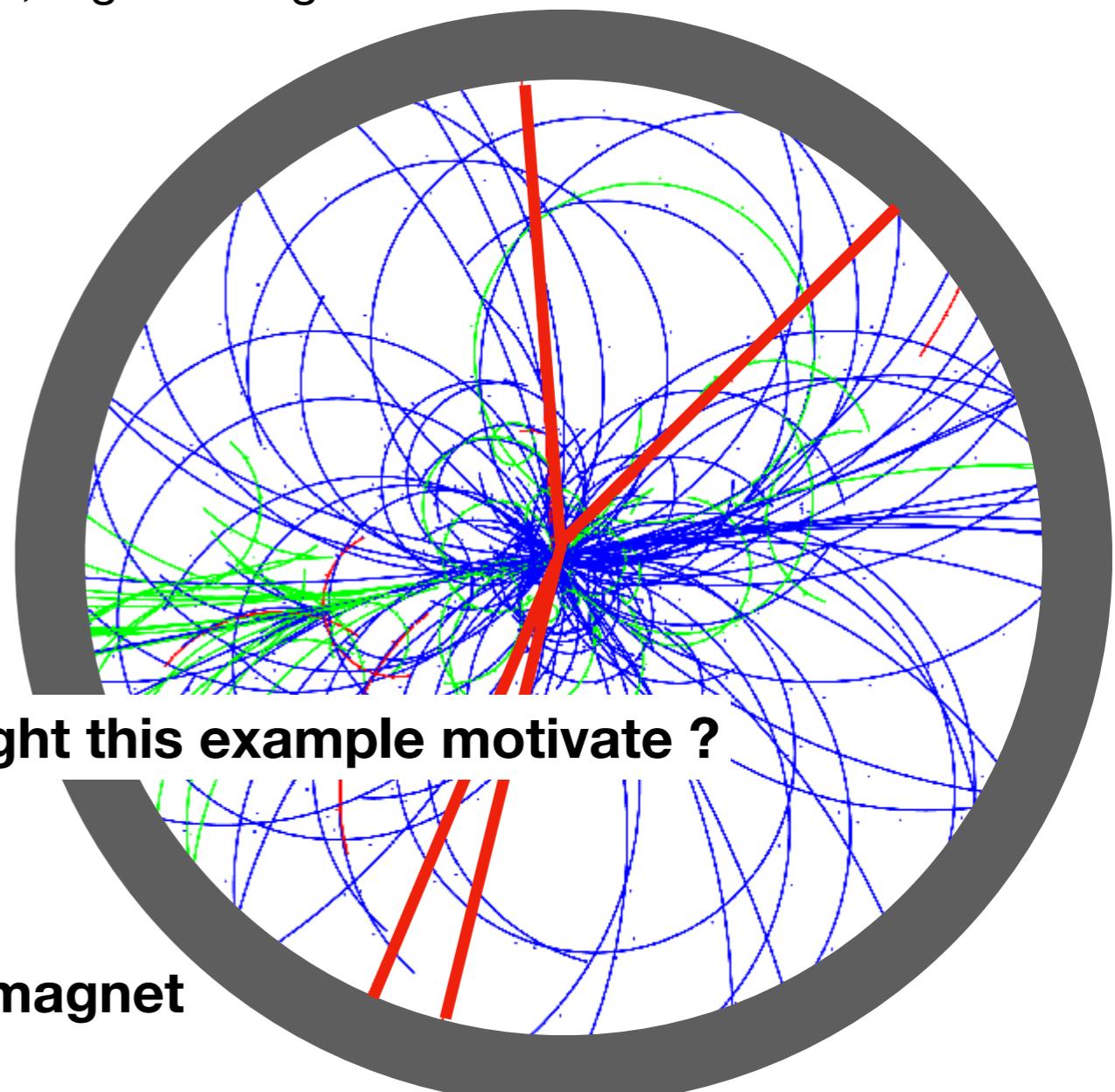


What detector technology might this example motivate ?

No magnet

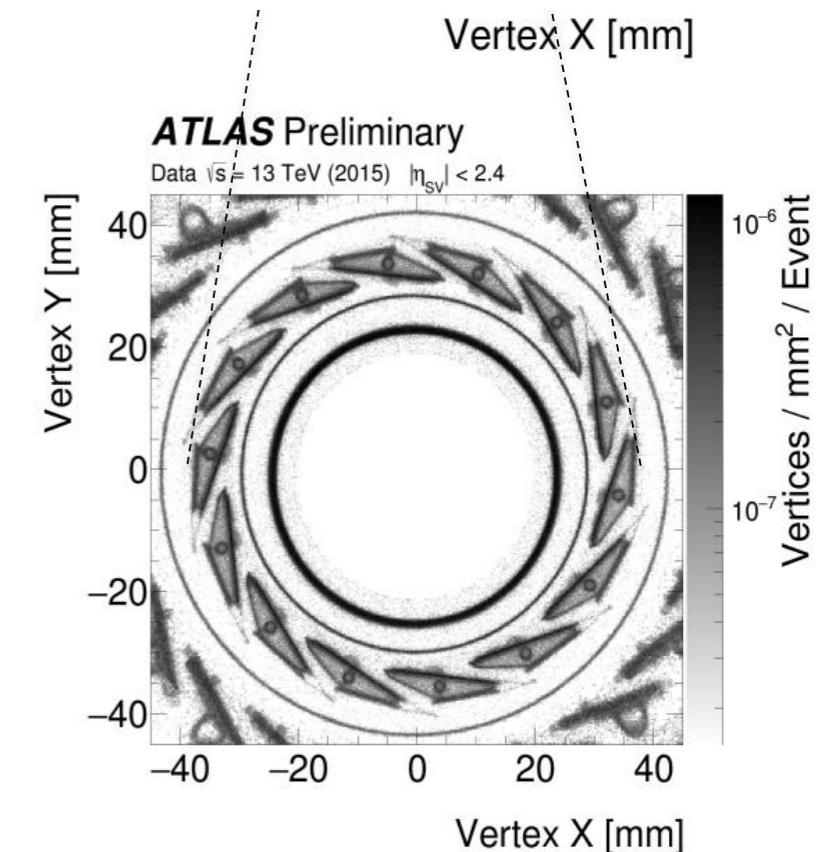
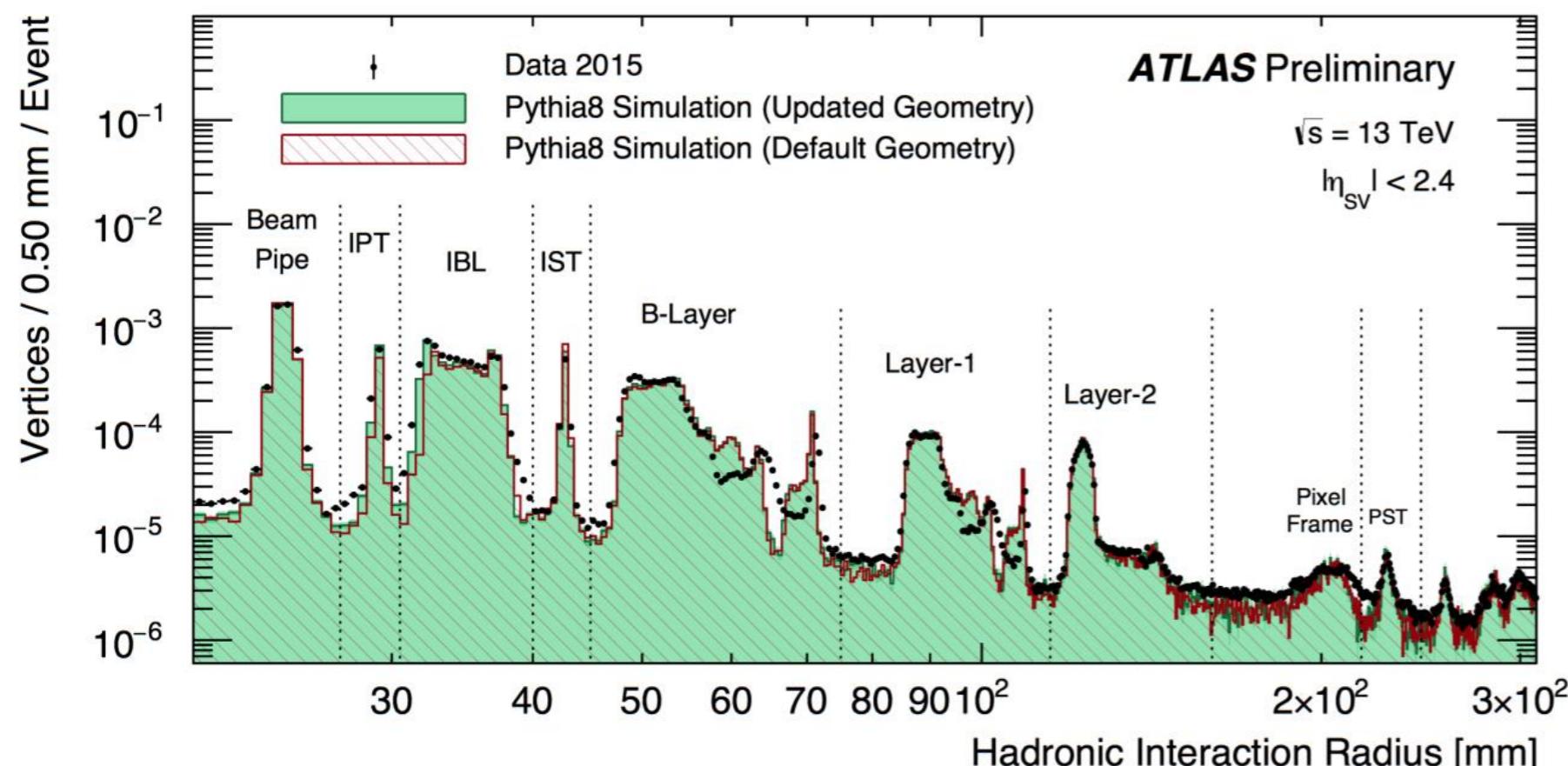


+ magnet

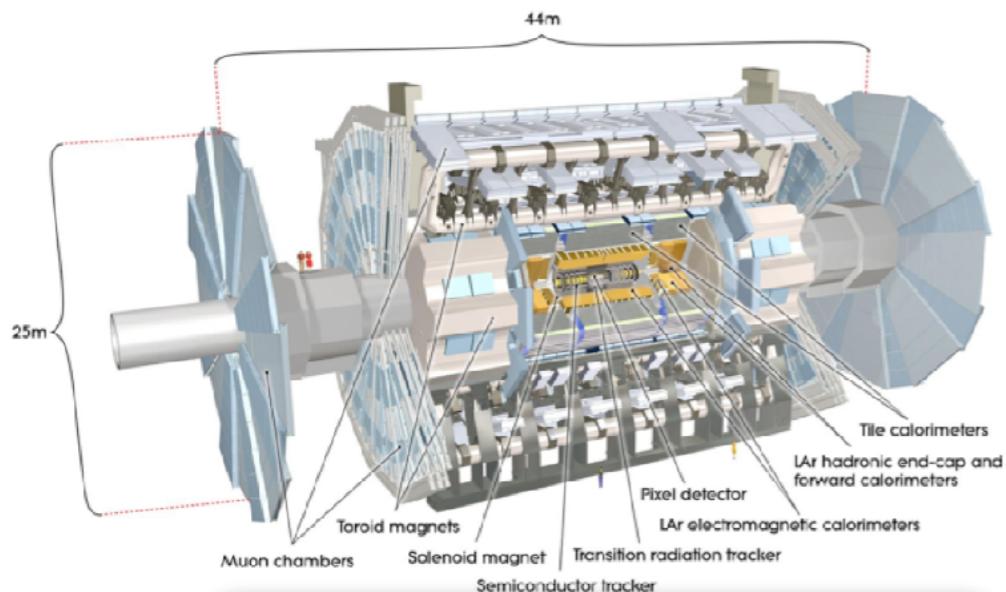


# Simulation and understanding detectors

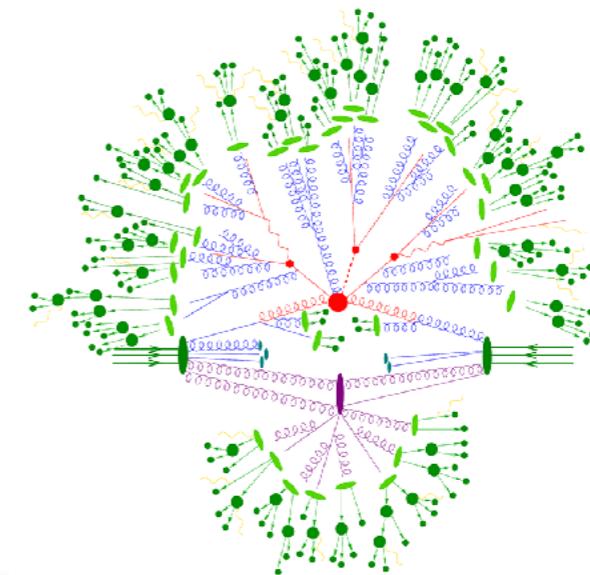
- We use **simulations** to model the detector as **accurately** and **precisely** as possible
- We then **test** that our simulations are accurate **using real data**
- We correct our simulations if necessary
- Once our simulation is an **accurate model** of our detector, we can use it to **correct the data for detector response**



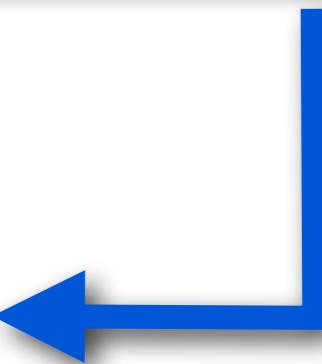
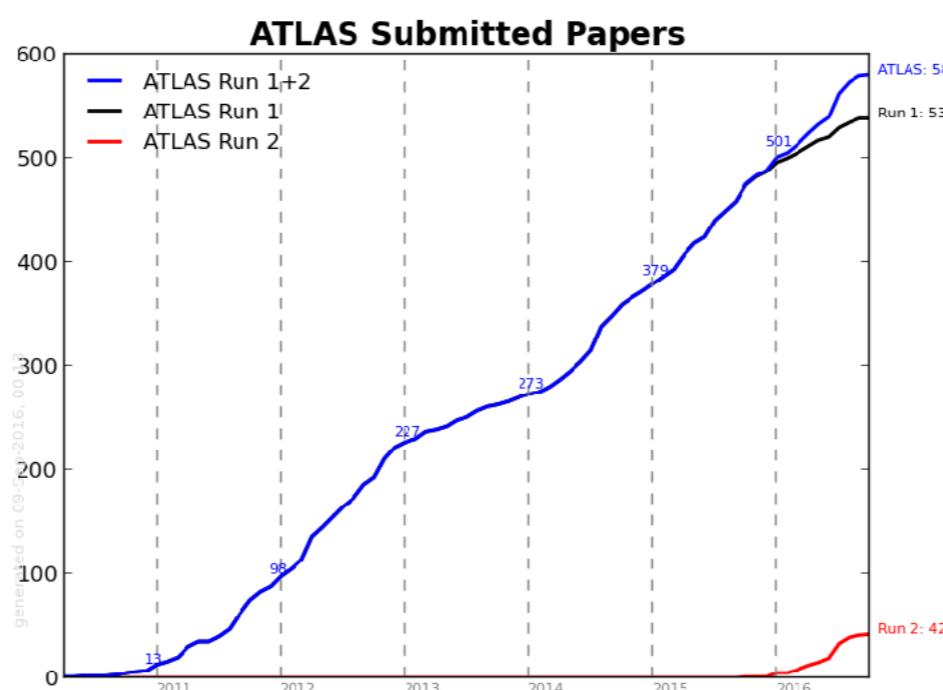
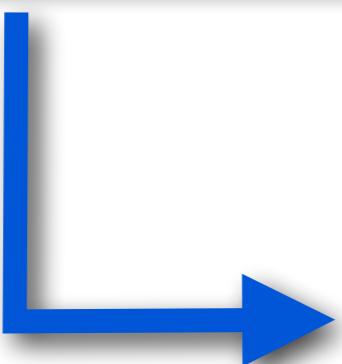
# Exabyte-scale physics analysis



Exabytes of Data

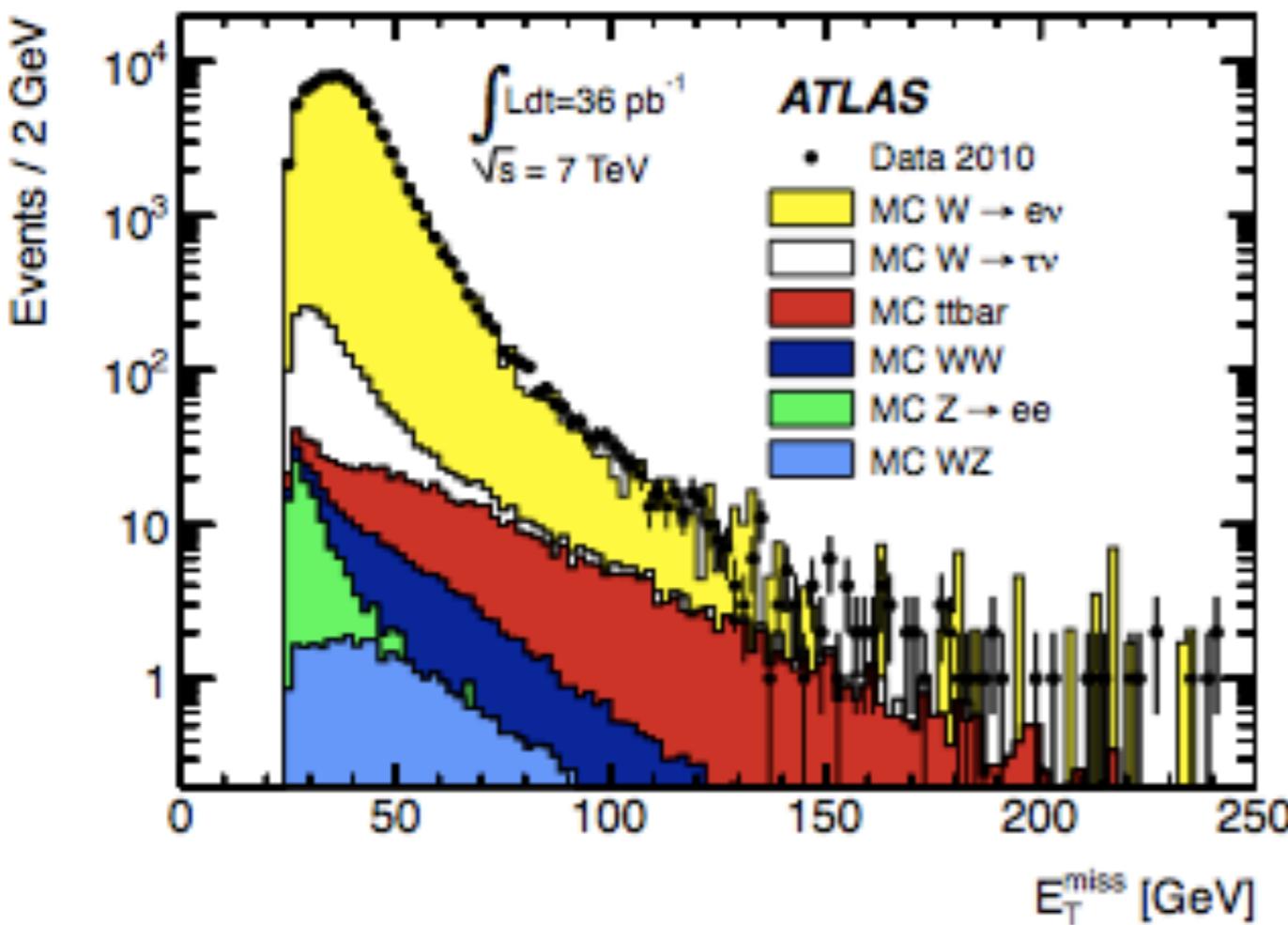


Exabytes of Simulation



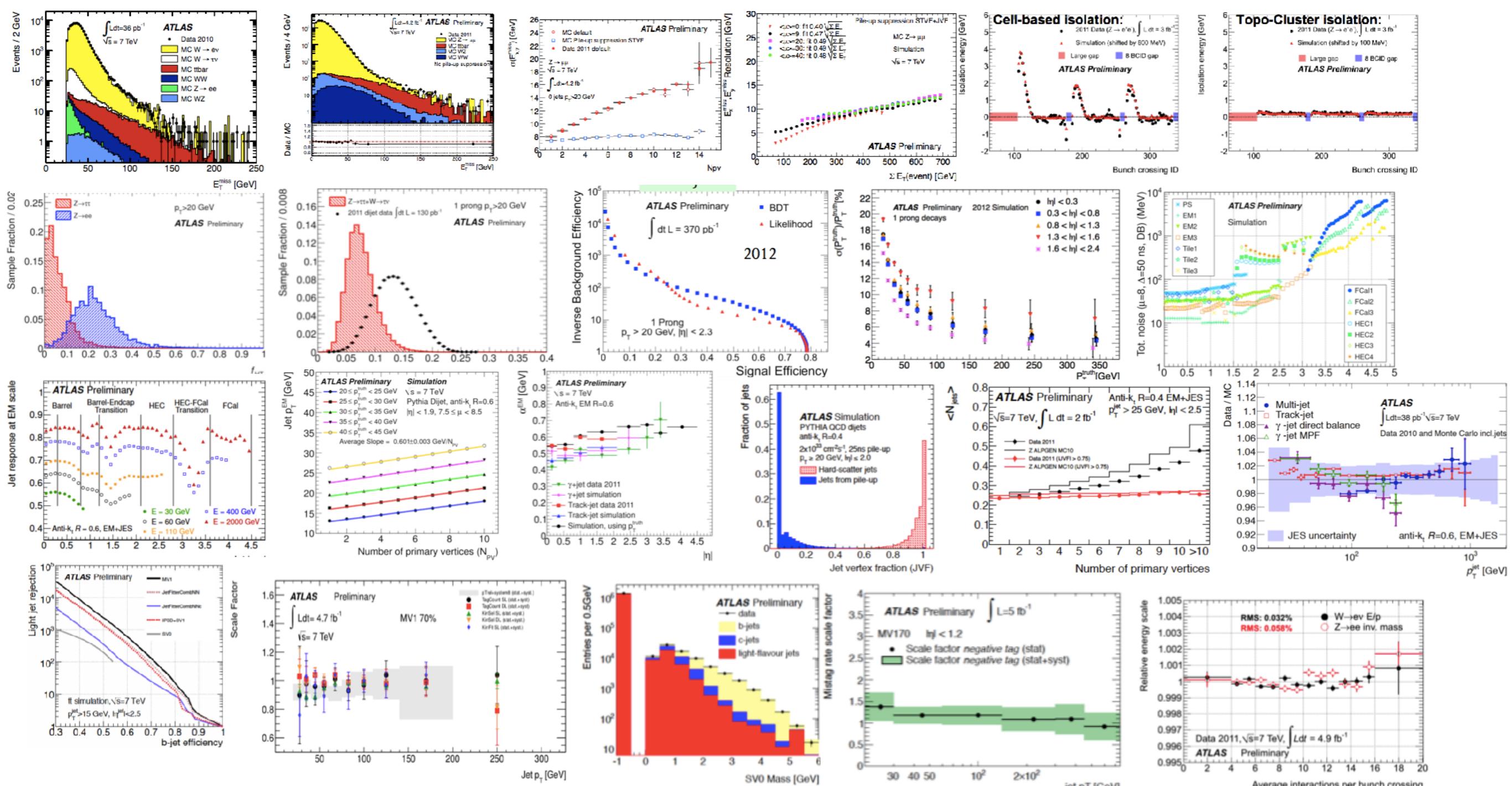
Publish!

# Ingredients to the ATLAS physics program



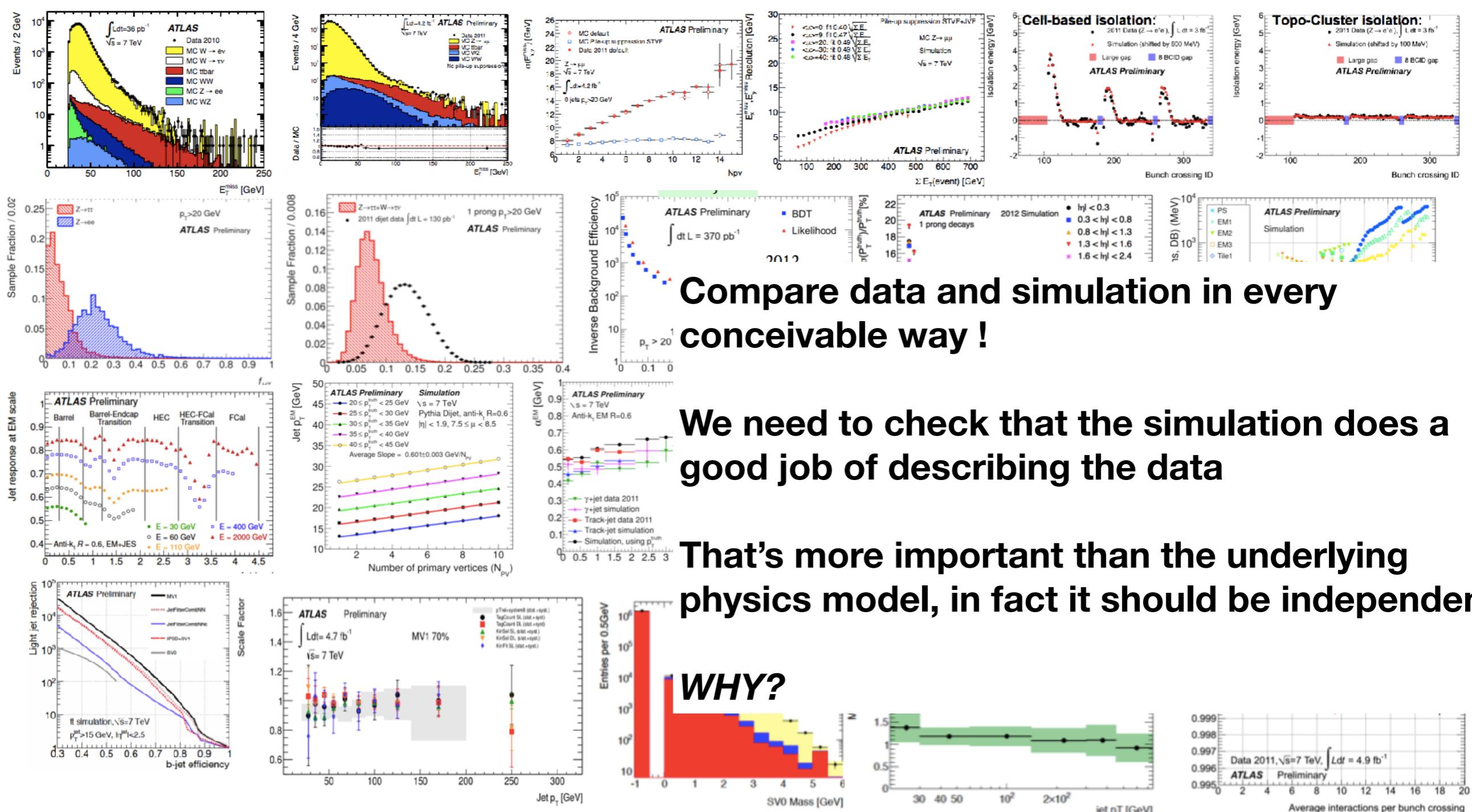
- We compare data with simulation

# Ingredients to the ATLAS physics program



- We make a LOT of comparisons of data and simulation

# Ingredients to the ATLAS physics program



Compare data and simulation in every conceivable way !

We need to check that the simulation does a good job of describing the data

That's more important than the underlying physics model, in fact it should be independent

# Measuring cross sections

$$\sigma = \frac{N}{L_{int}}$$

- The cross section for a process is defined as the number of events divided by the integrated luminosity,  $L_{int}$ , which measures how much data we have collected

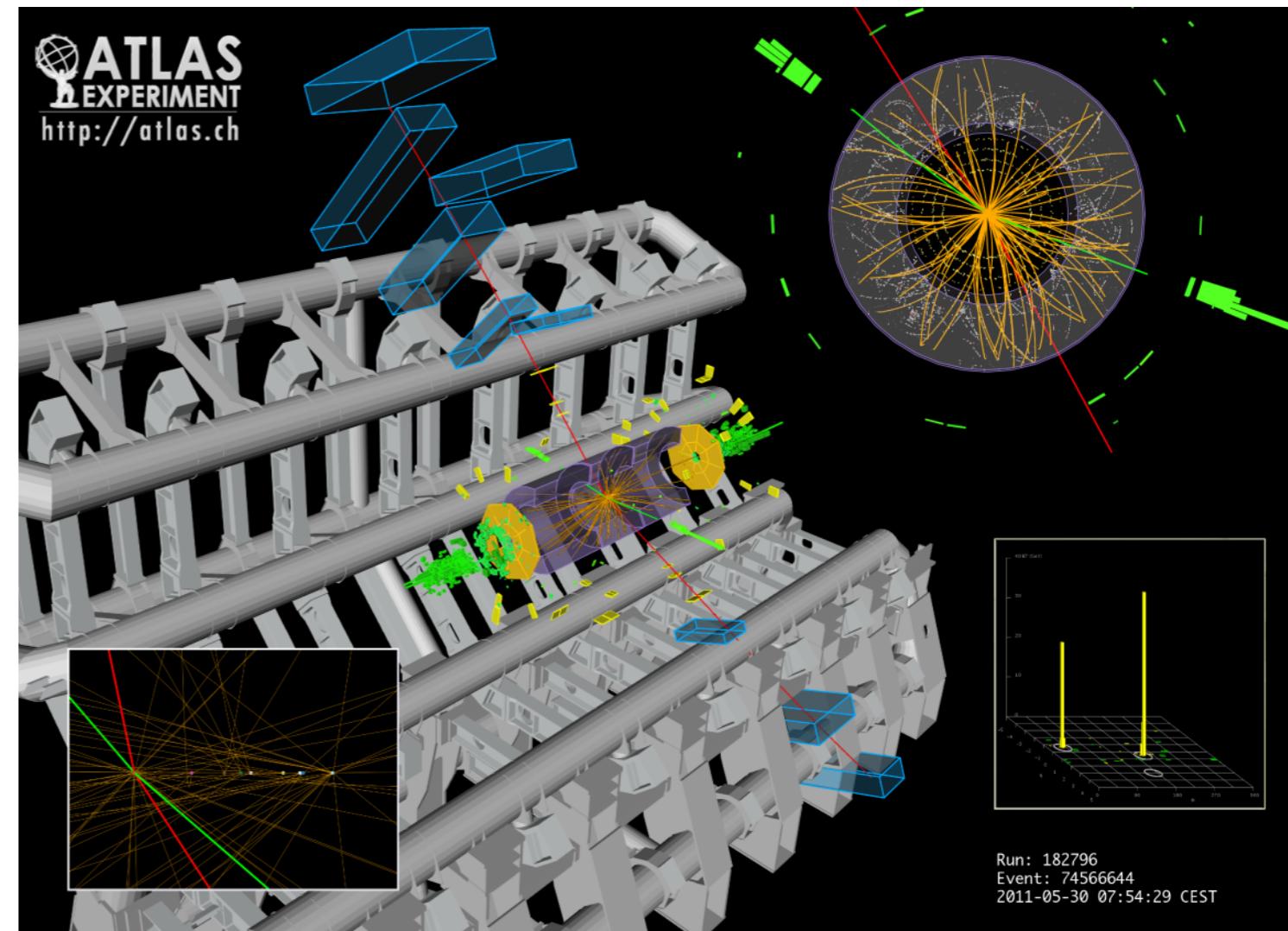
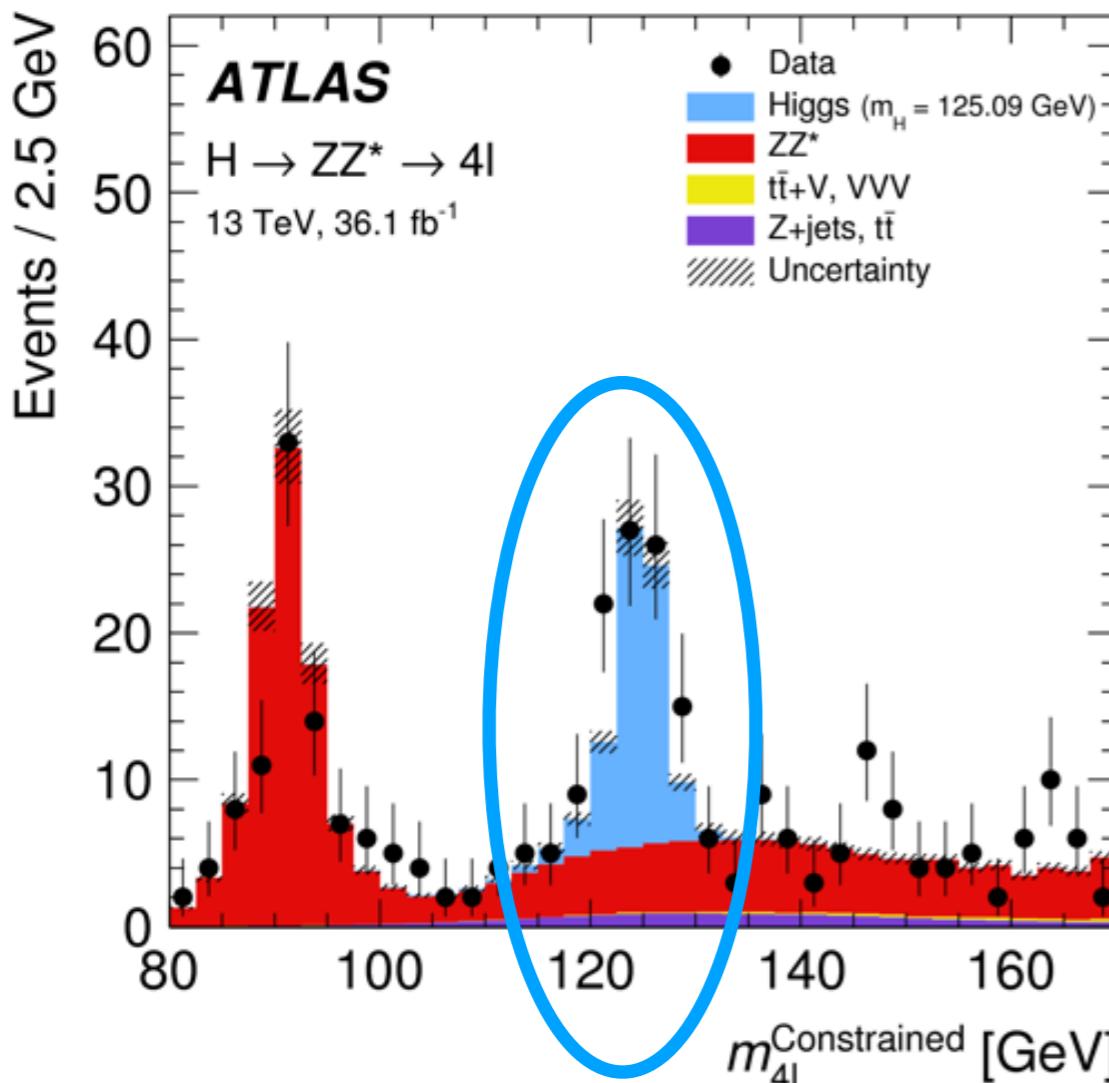
$$\sigma = \frac{N_{obs}}{A \cdot \epsilon \cdot L_{int}}$$

- $N_{obs}$  in data needs to be corrected for the detector acceptance,  $A$ , for selecting those events. The reconstruction efficiency,  $\epsilon$ , is a product of all of the efficiencies that we need to measure and ensure that they are the same in our data and simulation

$$\sigma = \frac{N_{obs} - N_{bkg}}{A \cdot \epsilon \cdot L_{int}}$$

- Finally, we need to measure and subtract background events that are not part of our signal process

# Discovering the Higgs Boson: $H \rightarrow ZZ \rightarrow 4l$



- We will (nearly) always have some irreducible background to the signal process that we are trying to measure

# Measuring cross sections

$$\sigma = \frac{N}{L_{int}}$$

- The cross section for a process is defined as the number of events divided by the integrated luminosity,  $L_{int}$ , which measures how much data we have collected

$$\sigma = \frac{N_{obs}}{A \cdot \epsilon \cdot L_{int}}$$

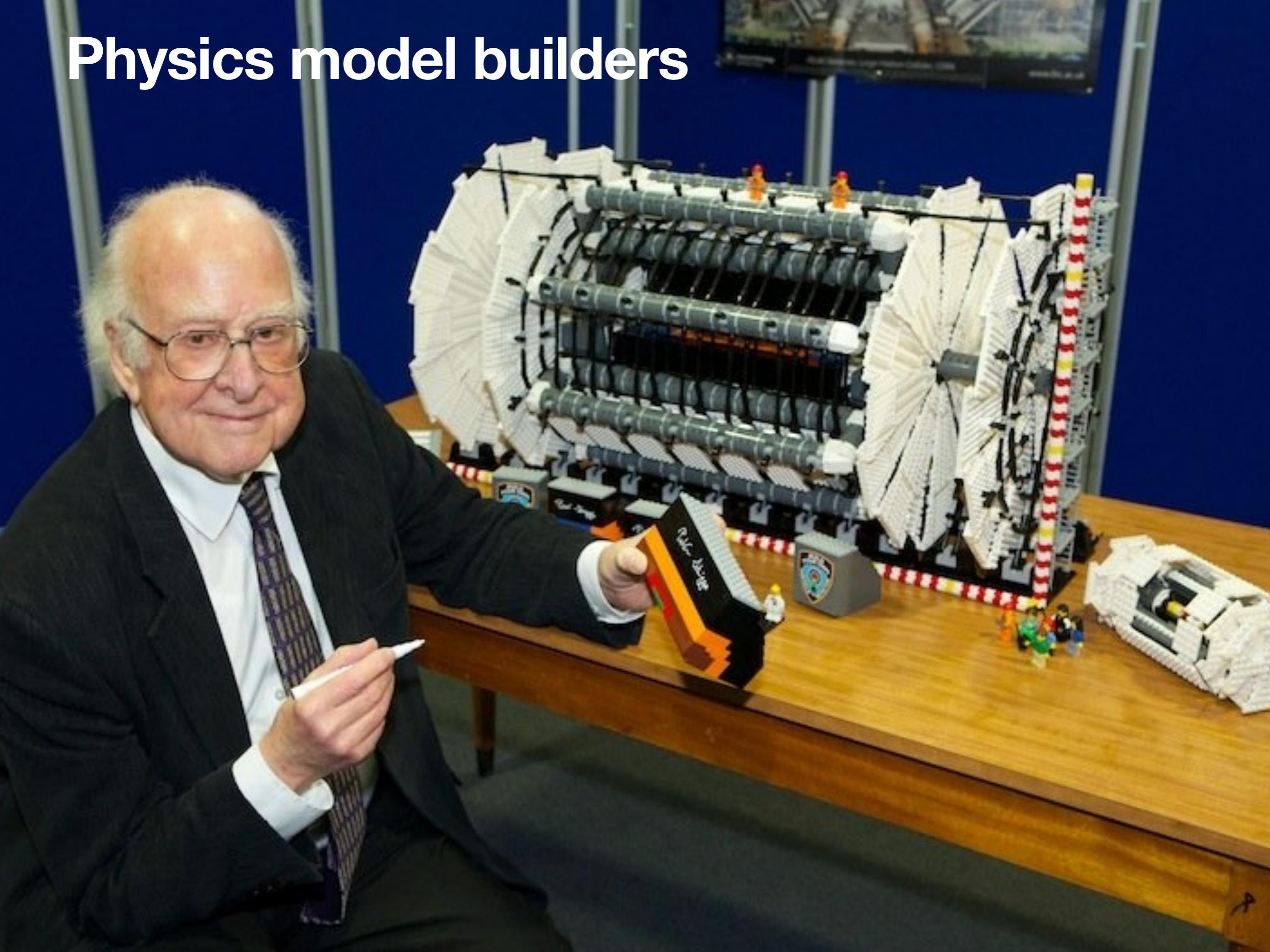
- $N_{obs}$  in data needs to be corrected for the detector acceptance,  $A$ , for selecting those events. The reconstruction efficiency,  $\epsilon$ , is a product of all of the efficiencies that we need to measure and ensure that they are the same in our data and simulation

$$\sigma = \frac{N_{obs} - N_{bkg}}{A \cdot \epsilon \cdot L_{int}}$$

- Finally, we need to measure and subtract background events that are not part of our signal process

**Now we can compare this to the theoretical cross section!**

# Physics model builders

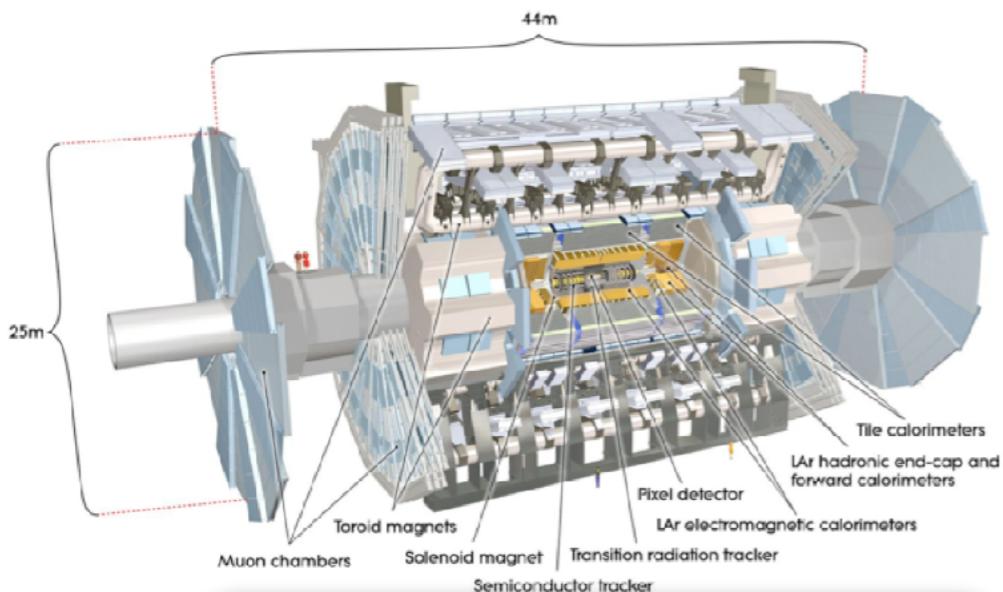


# Physics event generators

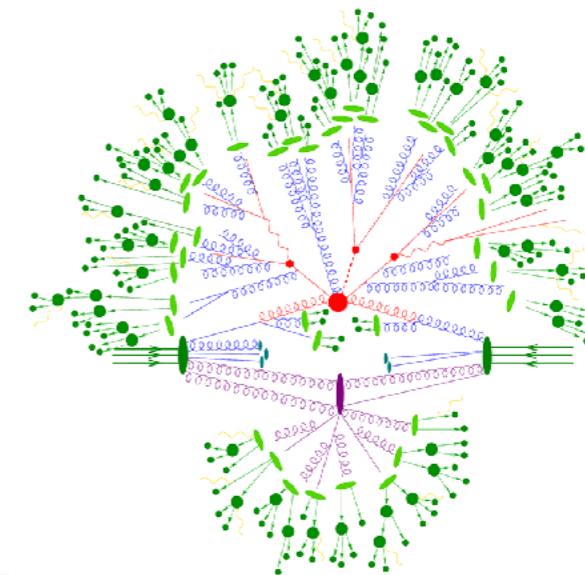


- There are lots of different physics models implemented in physics event generators, depending on the type of physics that you're interested in
- We want to see if reality looks like theory (and which one !)

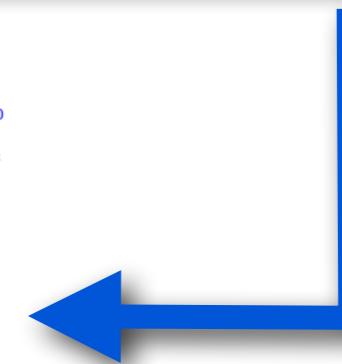
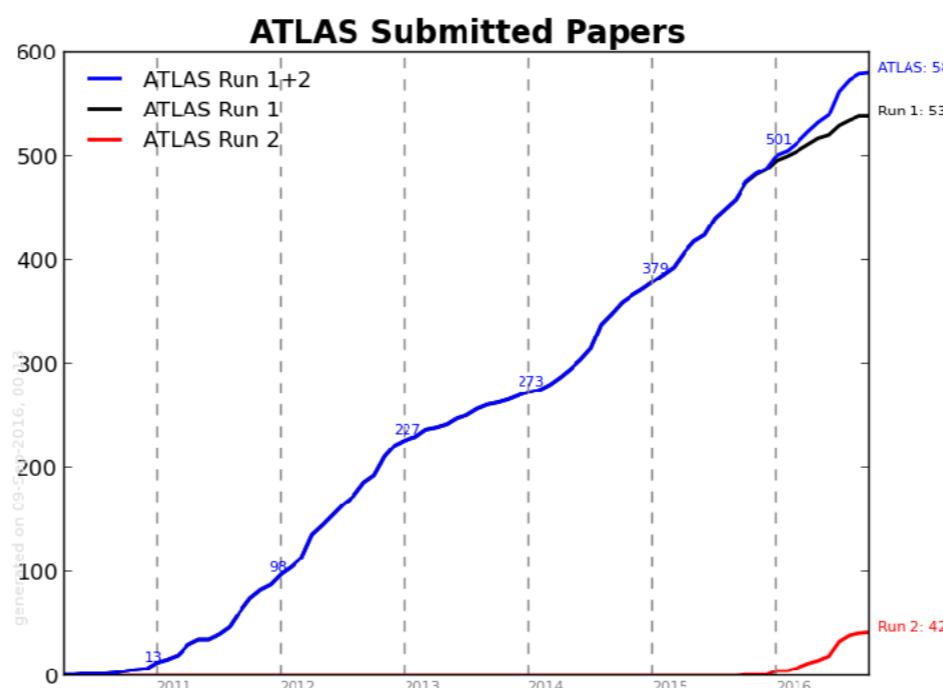
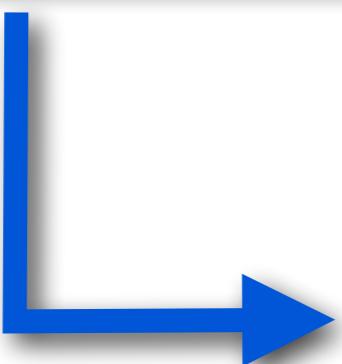
# Are we ready to do some exabyte-scale physics analysis?



Exabytes of Data

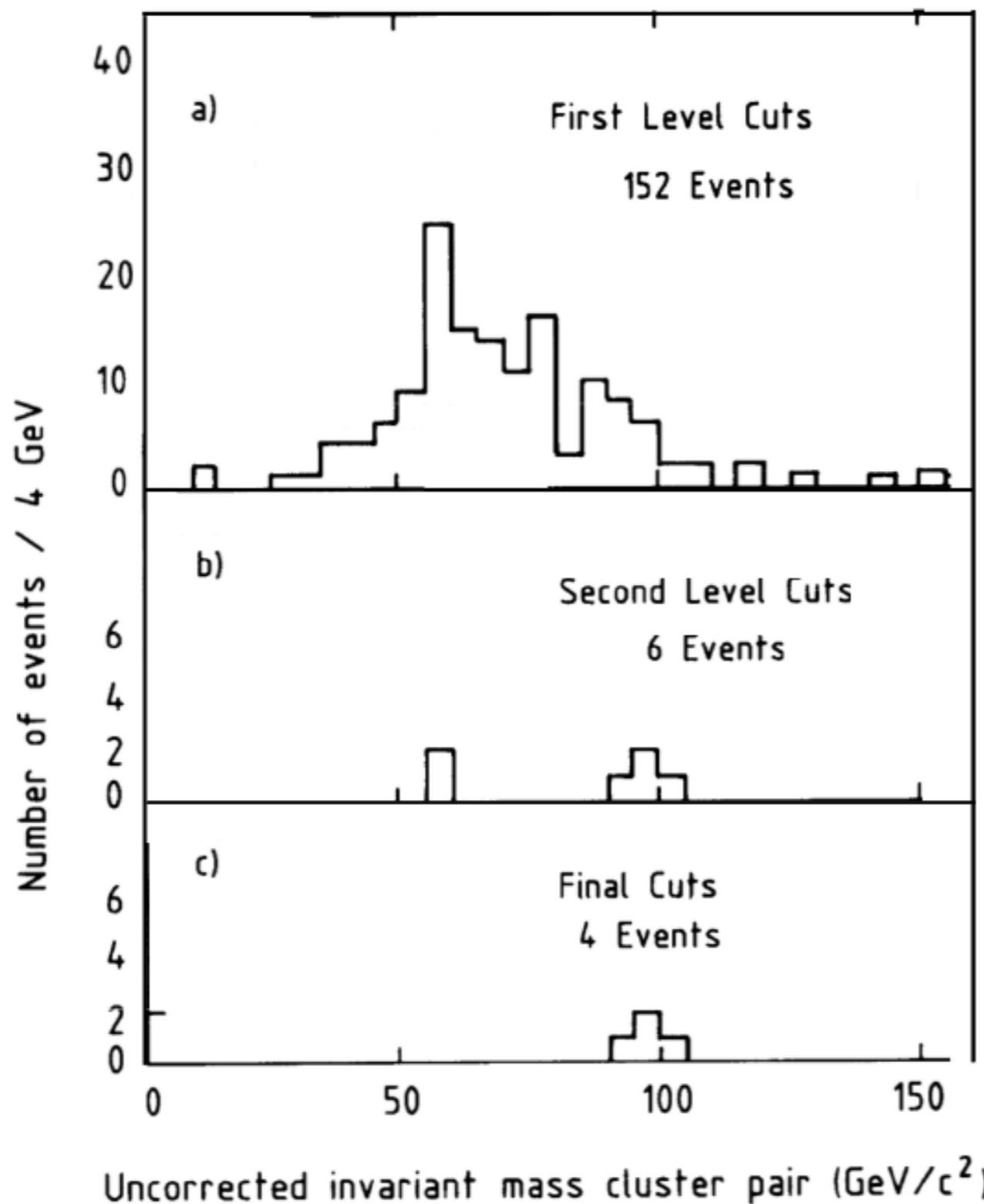


Exabytes of Simulation



Publish!

# First - measuring the Z boson



Z $\rightarrow$ ee in UA1

Two EM clusters with  $E_T > 25 \text{ GeV}$ .

As above plus a track with  $p_T > 7 \text{ GeV}$  pointing to the cluster. Hadronic and track isolation requirements applied.

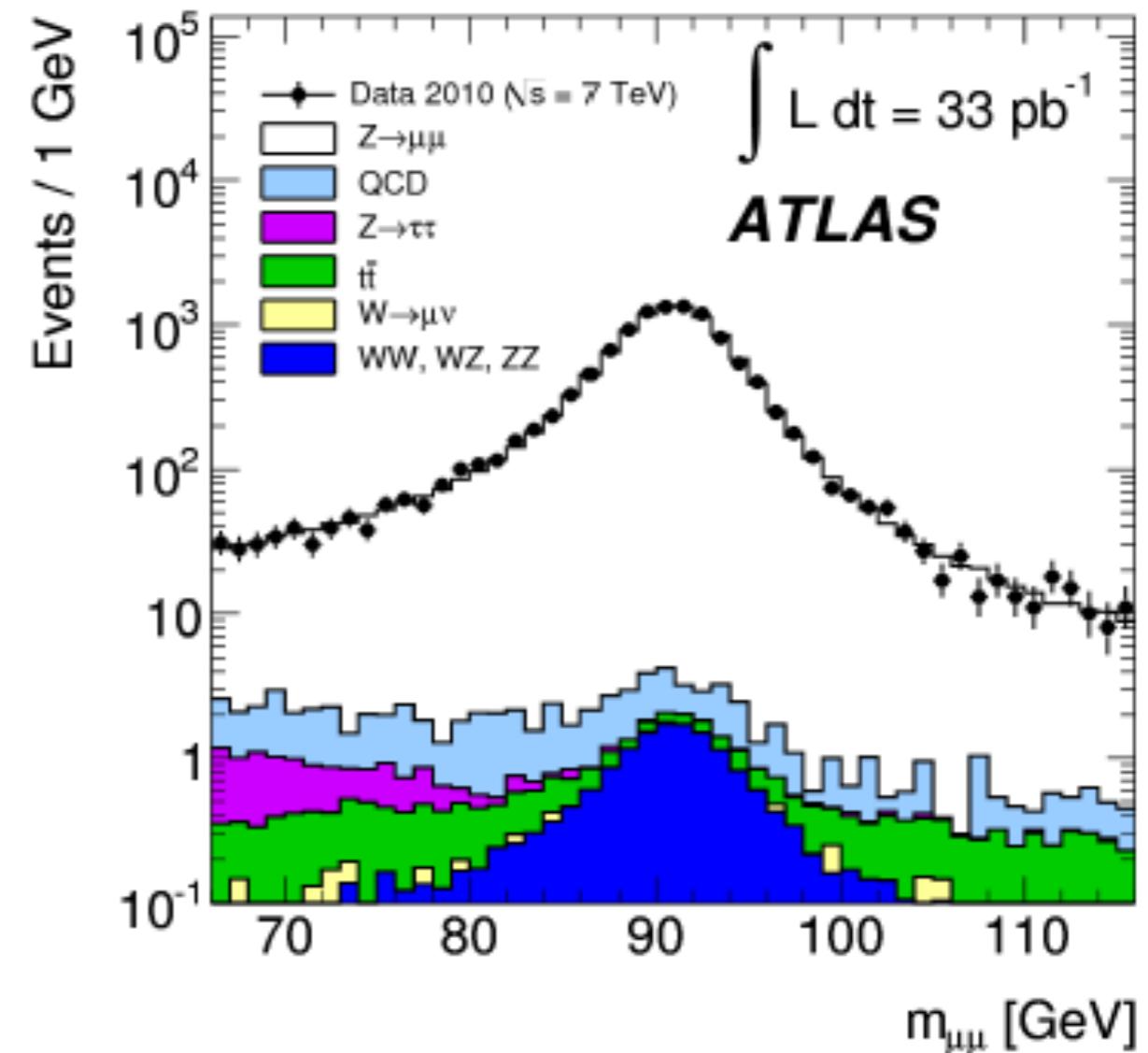
A second cluster has also an isolated track.



# Measuring the Z boson at ATLAS

$$\sigma = \frac{N_{obs} - N_{bkg}}{A \cdot \epsilon \cdot L_{int}}$$

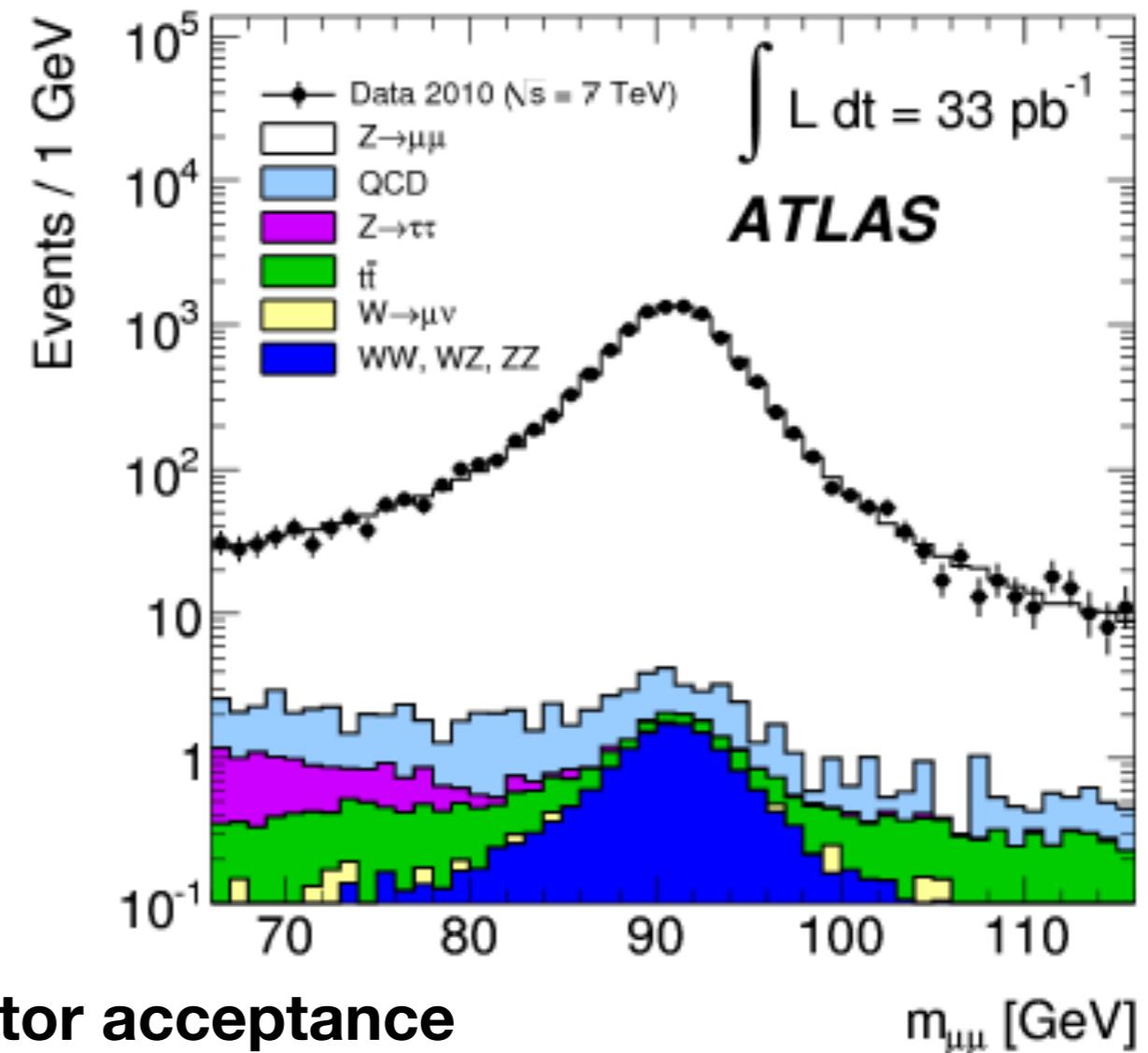
- Select events with (here) two muons
- **Question: what other selections can we apply to the muons?**
- Here I have only considered events with two muons
- **Question: is this the cross section for Z boson production?**



# Measuring the Z boson at ATLAS

$$\sigma = \frac{N_{obs} - N_{bkg}}{A \cdot \epsilon \cdot L_{int}}$$

- Backgrounds are small but still need to be measured and subtracted
- **We will quote a fiducial cross section corresponding to good detector acceptance**
- After making the event selection, applying the same selection to all of the simulations of background processes, and measuring my acceptance and efficiencies (and knowing the luminosity) - ***am I done?***



# Measuring the Z boson at ATLAS

$$\sigma = \frac{N_{obs} - N_{bkg}}{A \cdot \epsilon \cdot L_{int}}$$

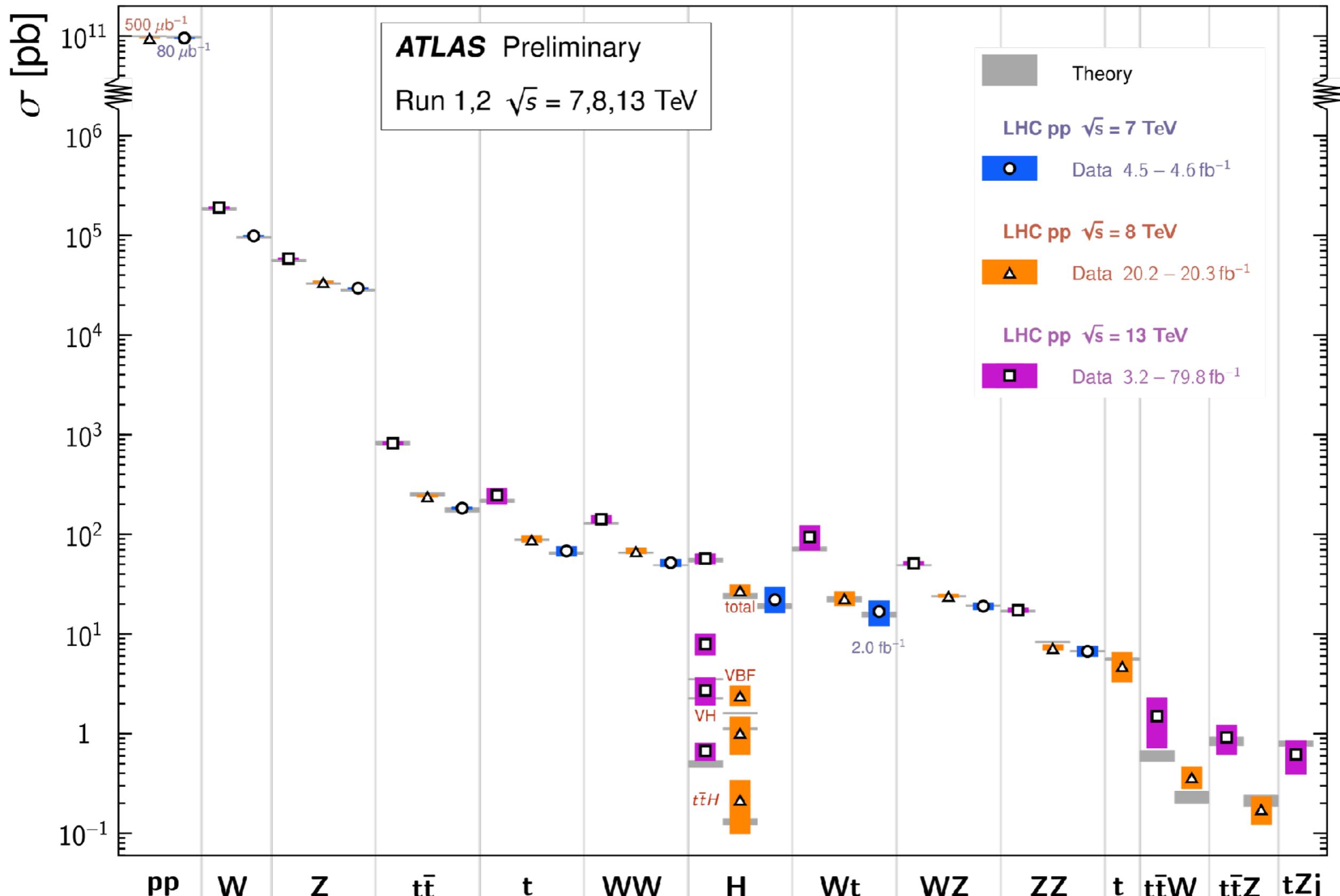
Table 5: Measured fiducial  $Z \rightarrow \ell^+ \ell^-$  differential and integrated cross sections for electron and muon channels.

$ y_{\ell\ell} ^{\min}$	$ y_{\ell\ell} ^{\max}$	$Z \rightarrow e^+ e^-$				$Z \rightarrow \mu^+ \mu^-$			
		$d\sigma/d y_{\ell\ell} $ [pb]	$\delta\sigma_{stat}$ [pb]	$\delta\sigma_{sys}$ [pb]	$\delta\sigma_{lumi}$ [pb]	$d\sigma/d y_{\ell\ell} $ [pb]	$\delta\sigma_{stat}$ [pb]	$\delta\sigma_{sys}$ [pb]	$\delta\sigma_{lumi}$ [pb]
0.0	0.5	99.9	2.5	1.6	1.9	105.2	2.4	1.1	2.0
0.5	1.0	100.3	2.7	1.6	1.9	101.9	2.3	1.0	1.9
1.0	1.5	89.2	2.7	1.4	1.7	89.8	2.1	0.8	1.7
1.5	2.0	59.6	2.4	1.2	1.1	61.0	1.8	0.6	1.1
2.0	2.5	19.6	1.3	0.7	0.4	20.3	1.2	0.2	0.4
0.0	2.5	369.0	5.3	4.7	6.9	377.9	4.4	3.4	7.1

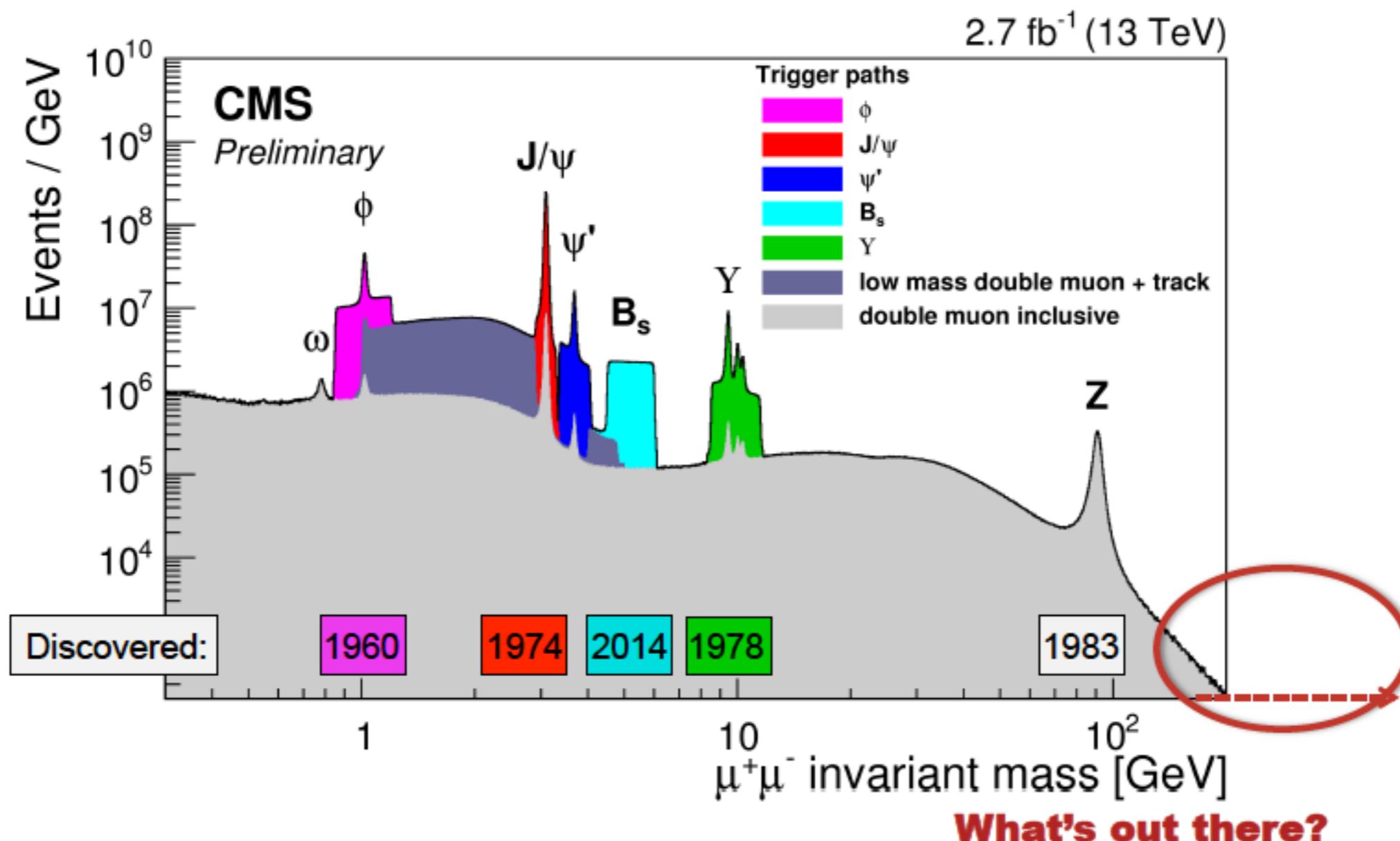
- No ! You would like to publish with the smallest ***uncertainties*** possible
- Every ingredient to the analysis comes with an uncertainty
- $N_{obs}$  has a ***statistical*** uncertainty
- $N_{bkg}$  is typically composed of several sources (different physics processes) with corresponding ***statistical*** and ***systematic*** contributions to the final uncertainty
- $A$  and particularly  $\epsilon$  have many ***systematic*** components stemming from each reconstruction algorithm that we used
- Finally,  $L_{int}$  also has an uncertainty that dictates how well we know the absolute scale of the measurement - a ***normalisation*** uncertainty

# Standard Model Total Production Cross Section Measurements

Status: July 2018

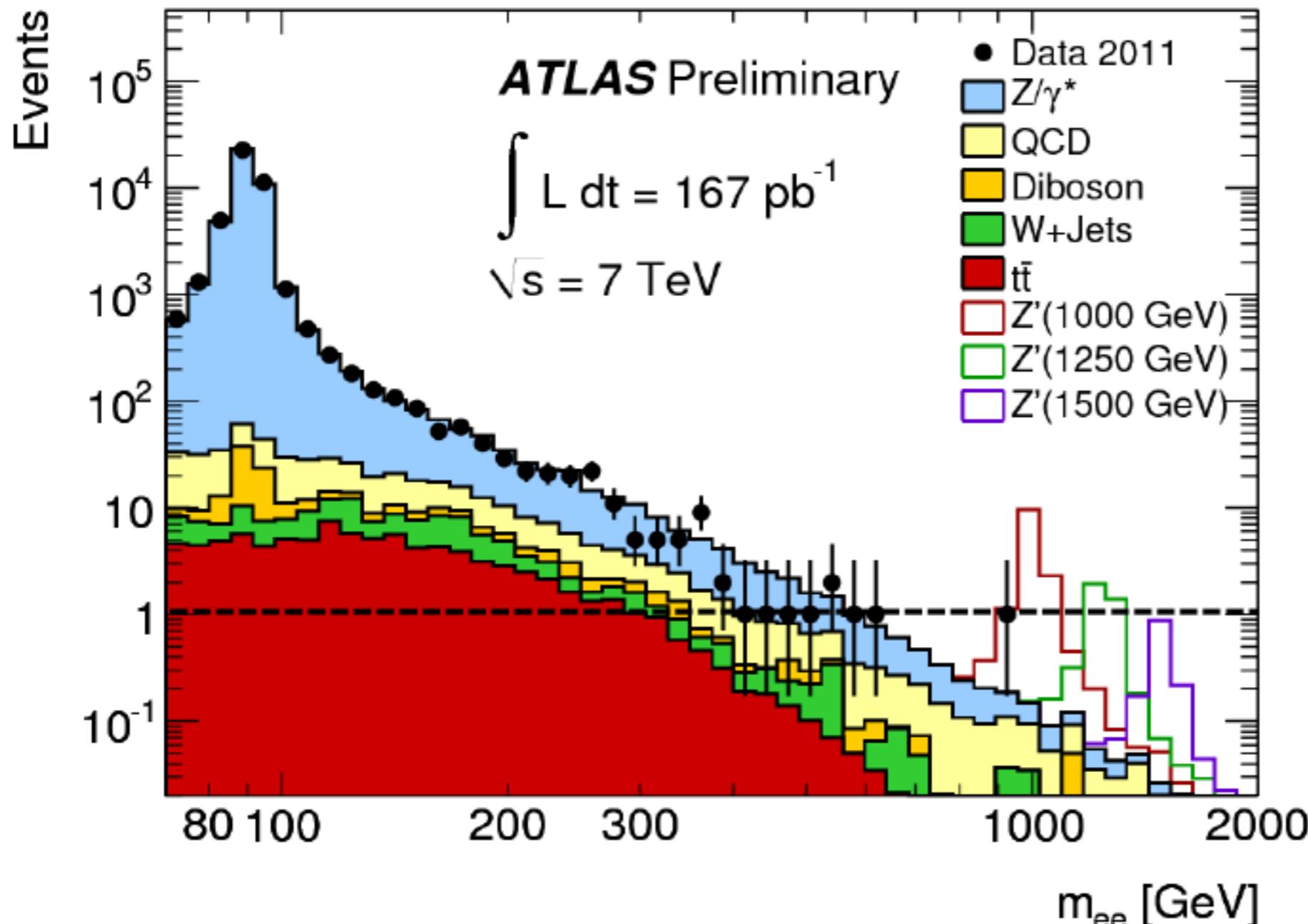


# Elements of a search



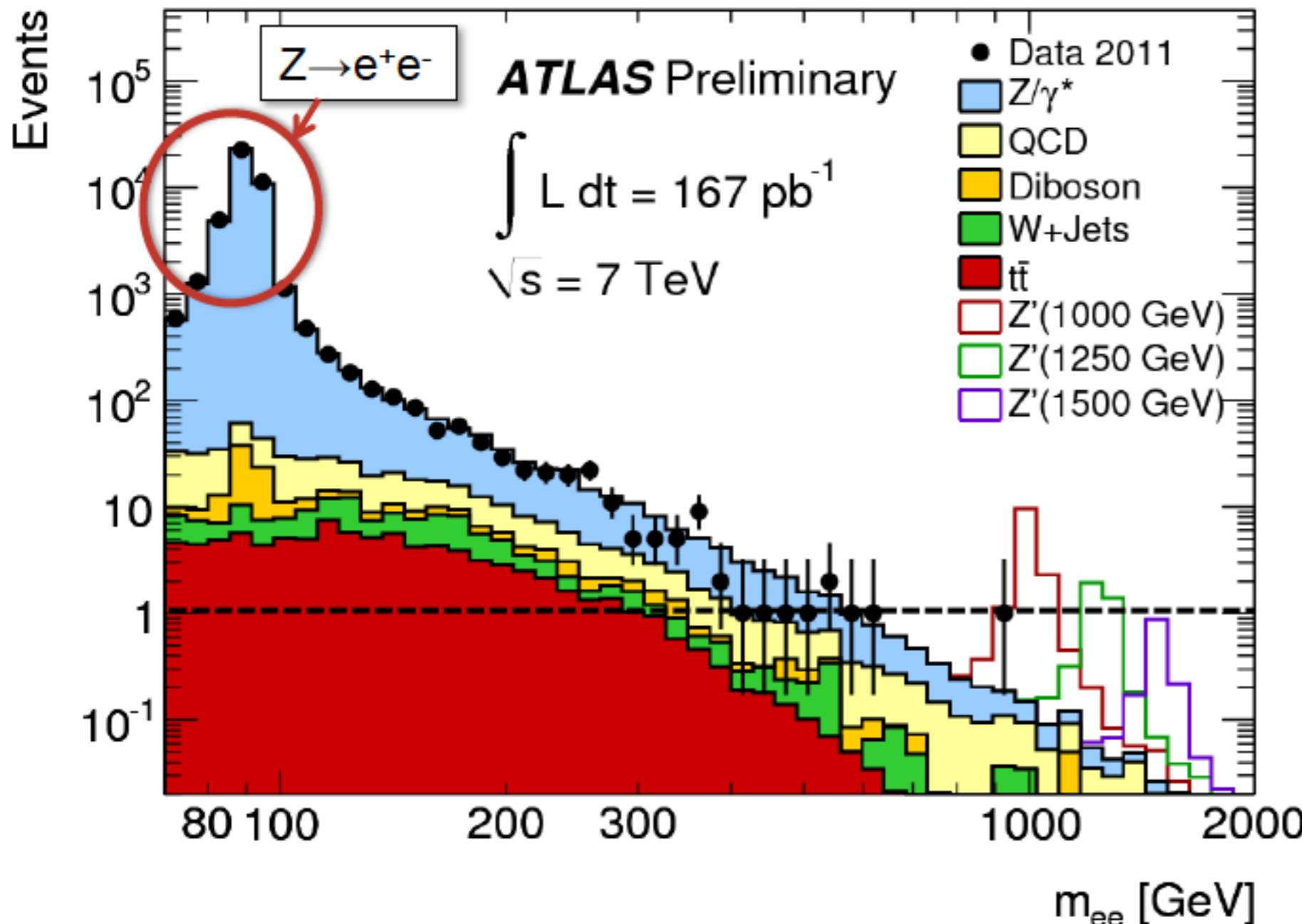
# Elements of a search

◎ Like Z->ee but at higher mass.



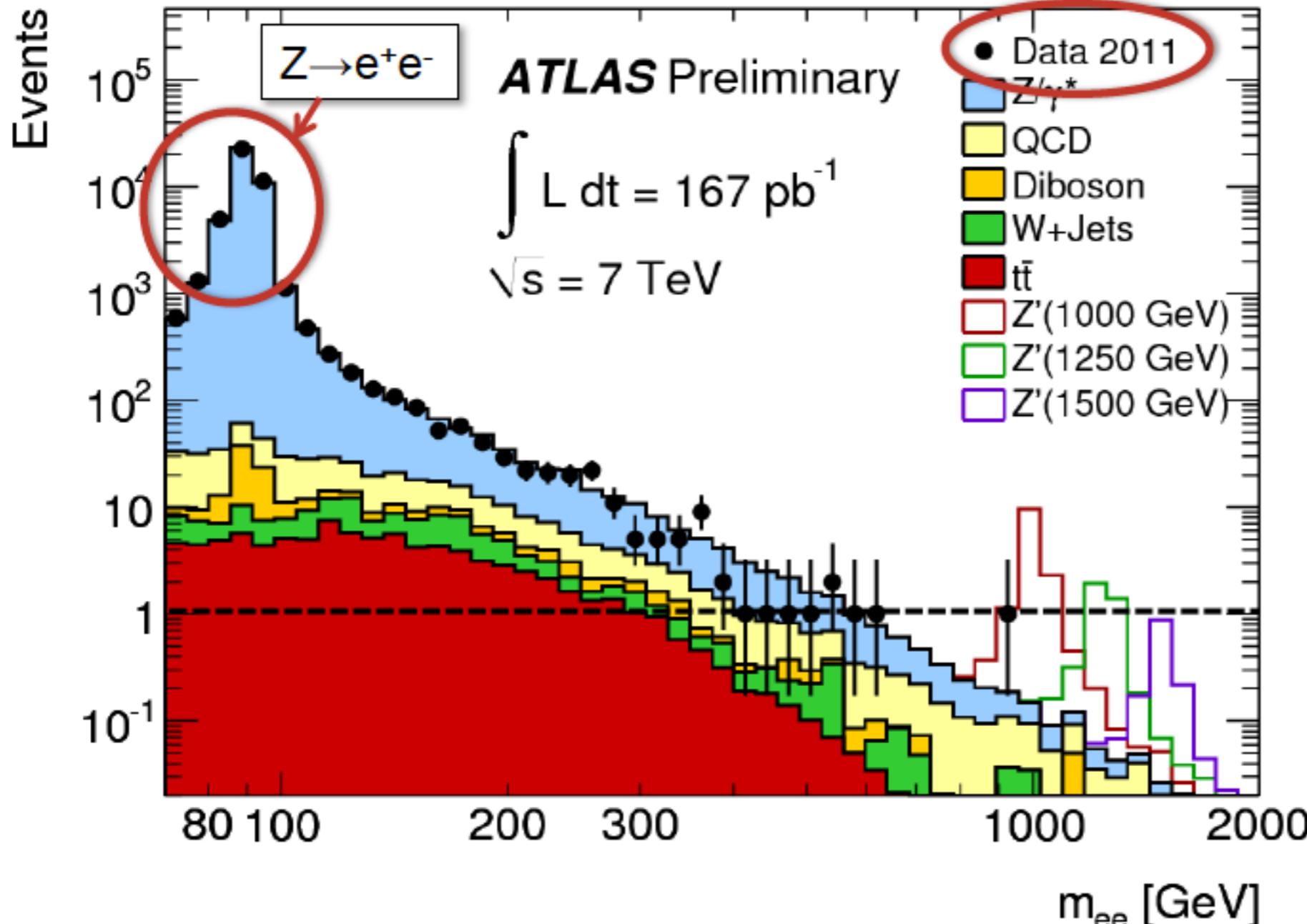
# Elements of a search

◎ Like  $Z \rightarrow ee$  but at higher mass.



# Elements of a search

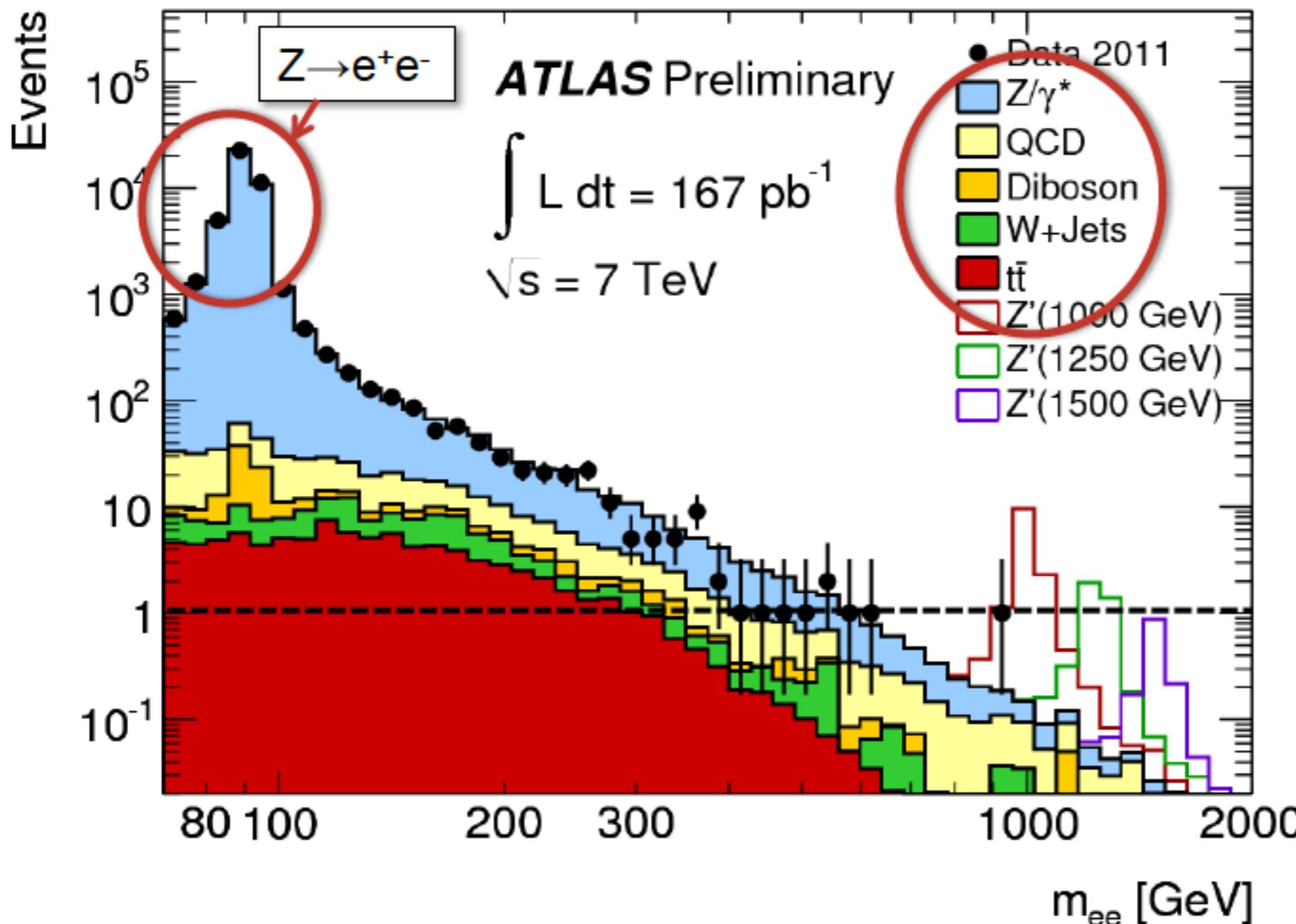
◎ Like  $Z \rightarrow ee$  but at higher mass.



Select 2 electron candidates and plot their invariant mass for:  
1. Data

# Elements of a search

◎ Like  $Z \rightarrow ee$  but at higher mass.

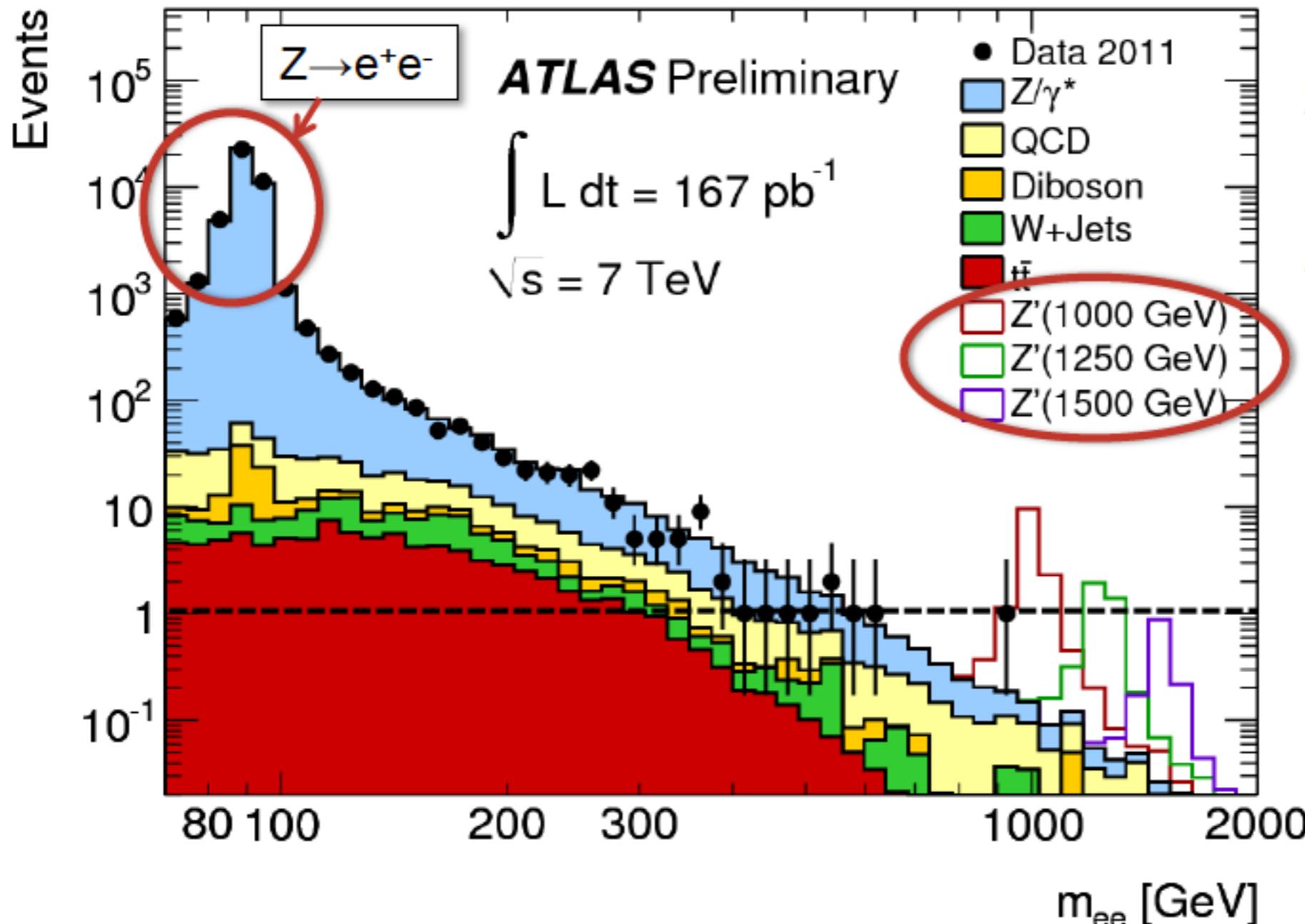


Select 2 electron candidates and plot their invariant mass for:

1. Data
2. Simulated background events

# Elements of a search

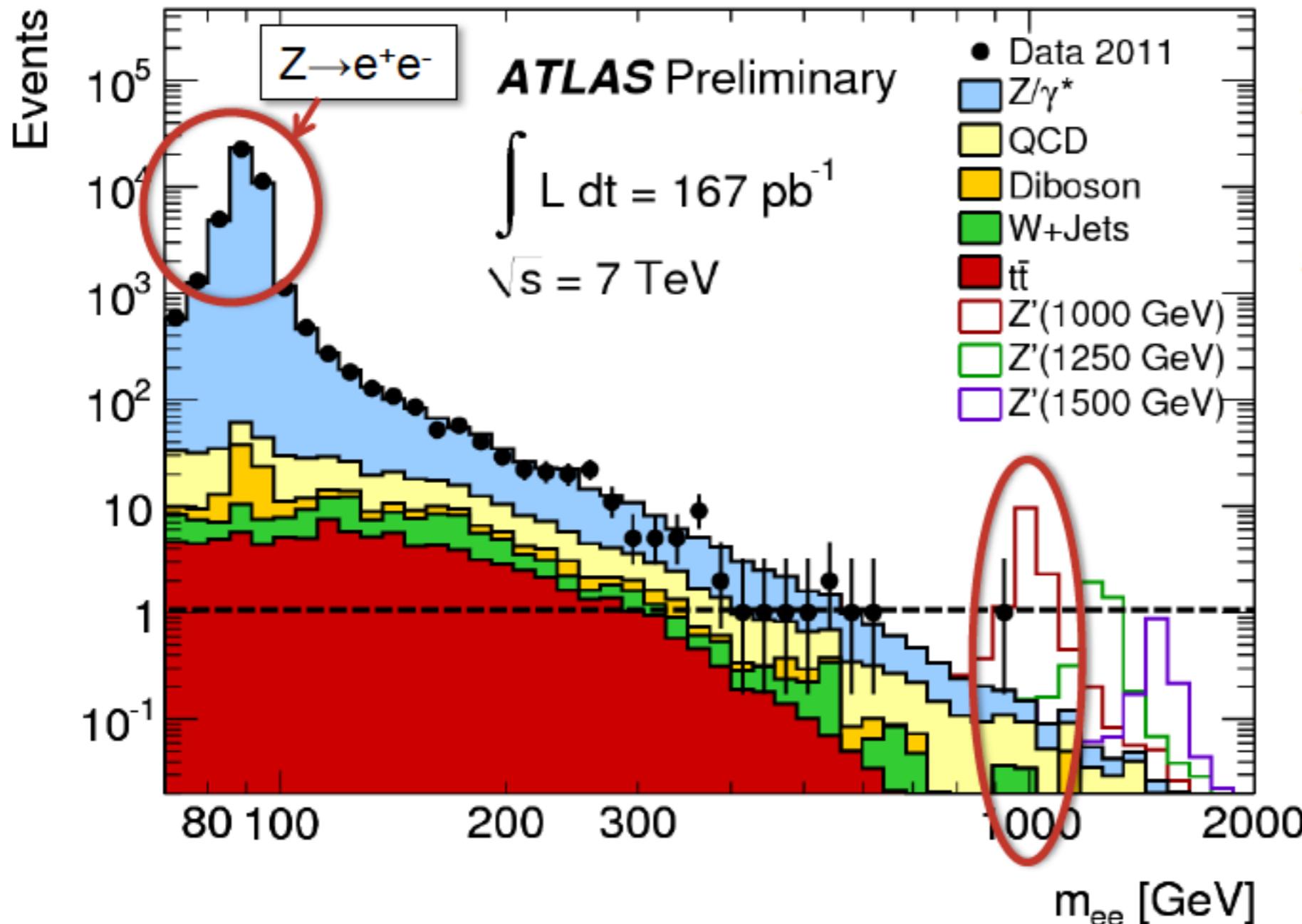
◎ Like  $Z \rightarrow ee$  but at higher mass.



- Select 2 electron candidates and plot their invariant mass for:
1. Data
  2. Simulated background events
  3. Simulated signal with different masses

# Elements of a search

◎ Like  $Z \rightarrow ee$  but at higher mass.



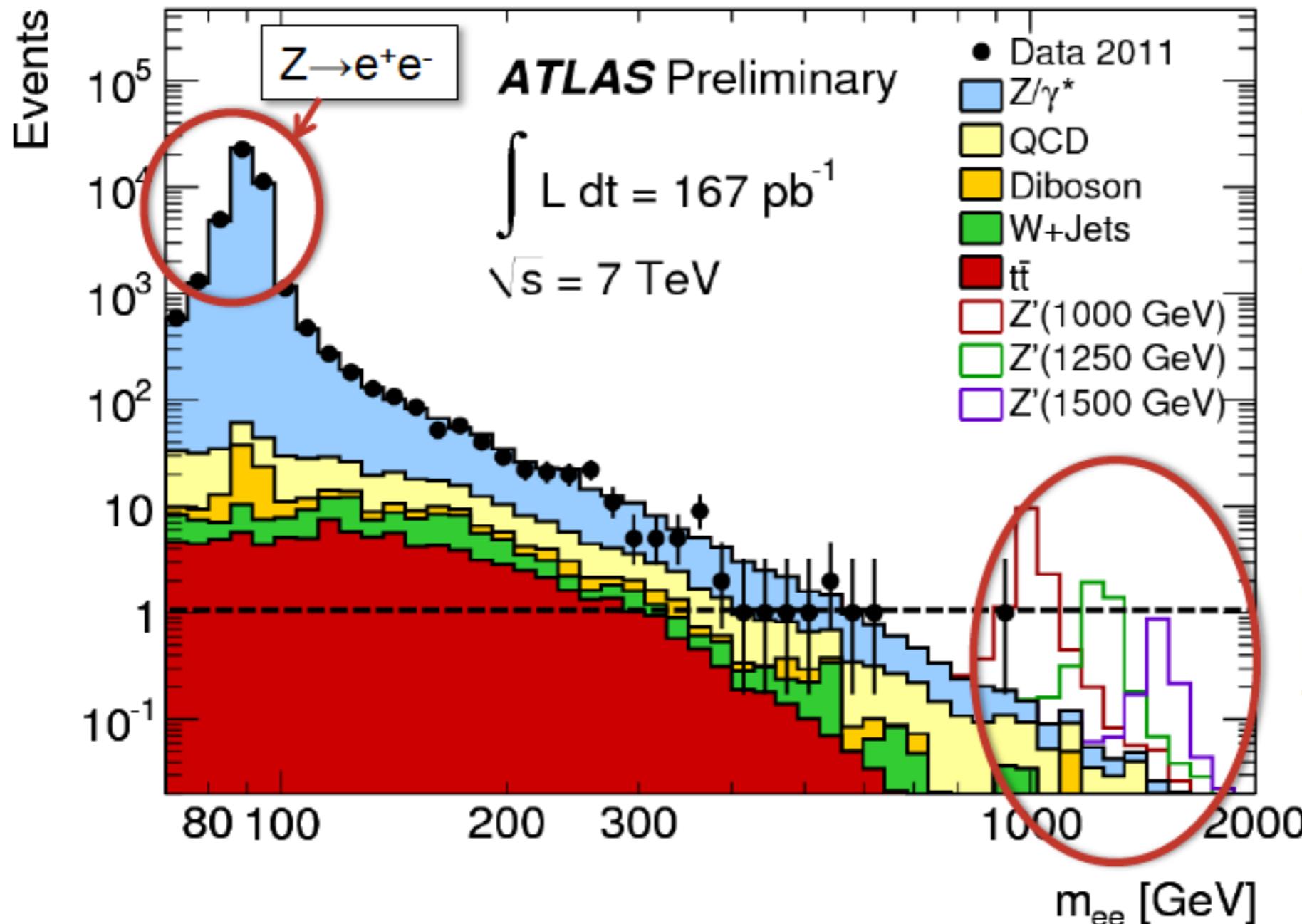
Select 2 electron candidates and plot their invariant mass for:

1. Data
2. Simulated background events
3. Simulated signal with different masses

Data inconsistent with a 1TeV  $Z'$

# Elements of a search

◎ Like  $Z \rightarrow ee$  but at higher mass.



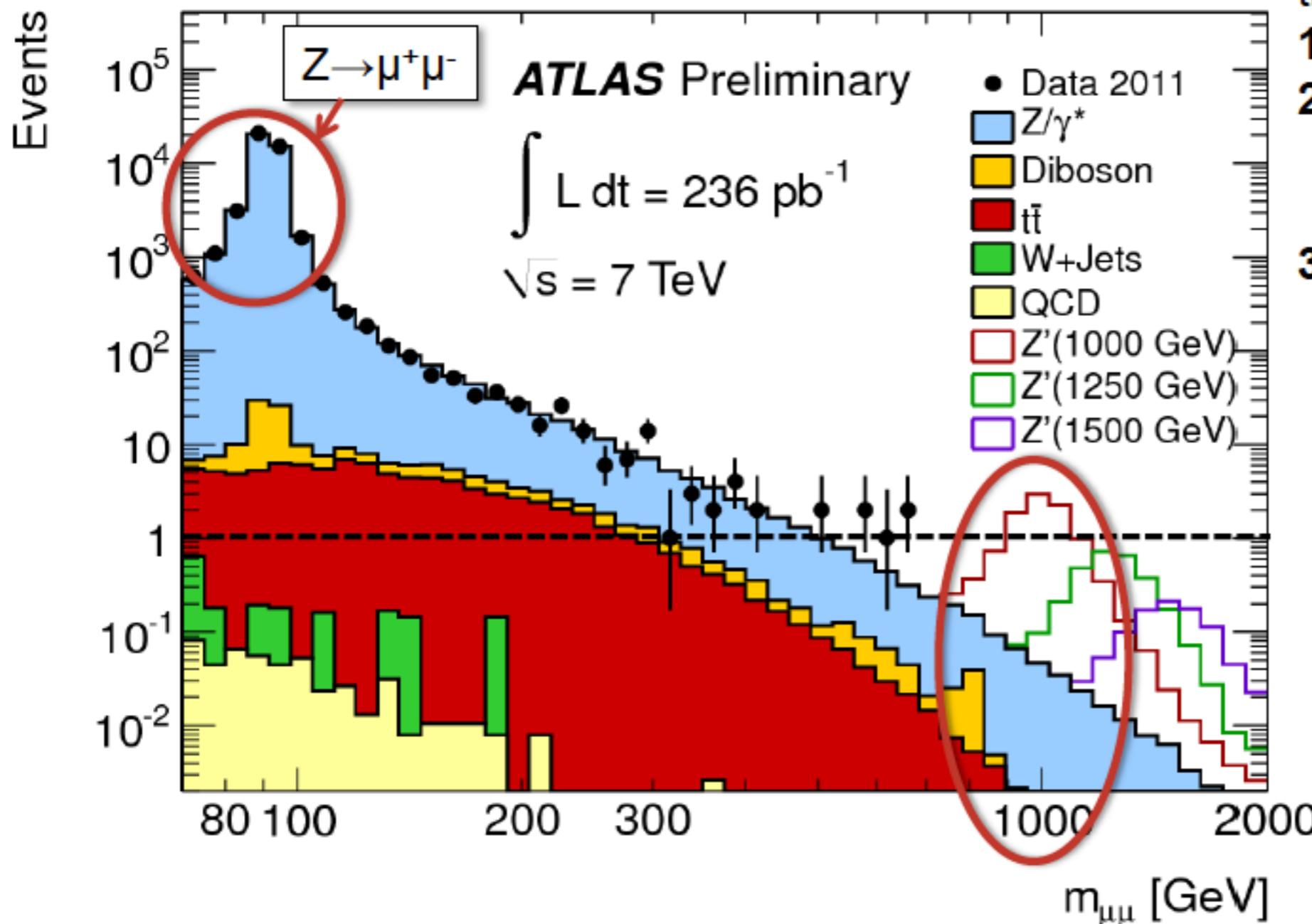
Select 2 electron candidates and plot their invariant mass for:

1. Data
2. Simulated background events
3. Simulated signal with different masses

Cross-section decreases with mass (higher the mass of the  $Z'$ , the more data needed to discover it)

# Elements of a search

## ◎ And similar for muons

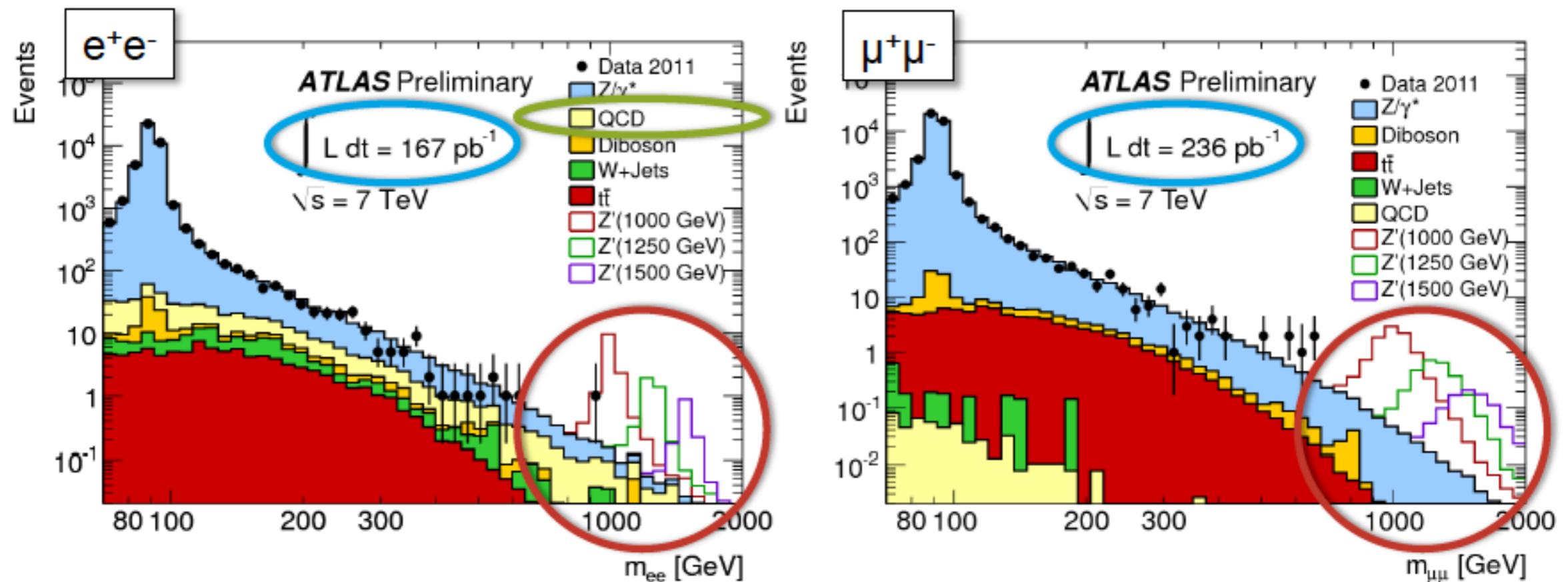


Select 2 muon candidates and plot their invariant mass for:

1. Data
2. Simulated background events
3. Simulated signal with different masses

Data inconsistent with a 1TeV  $Z'$

# Elements of a search

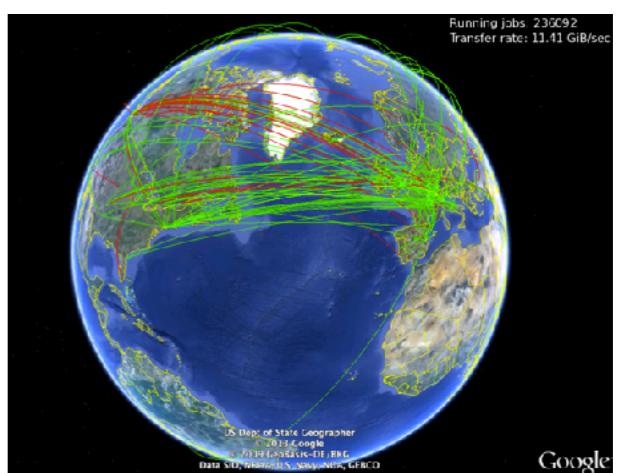
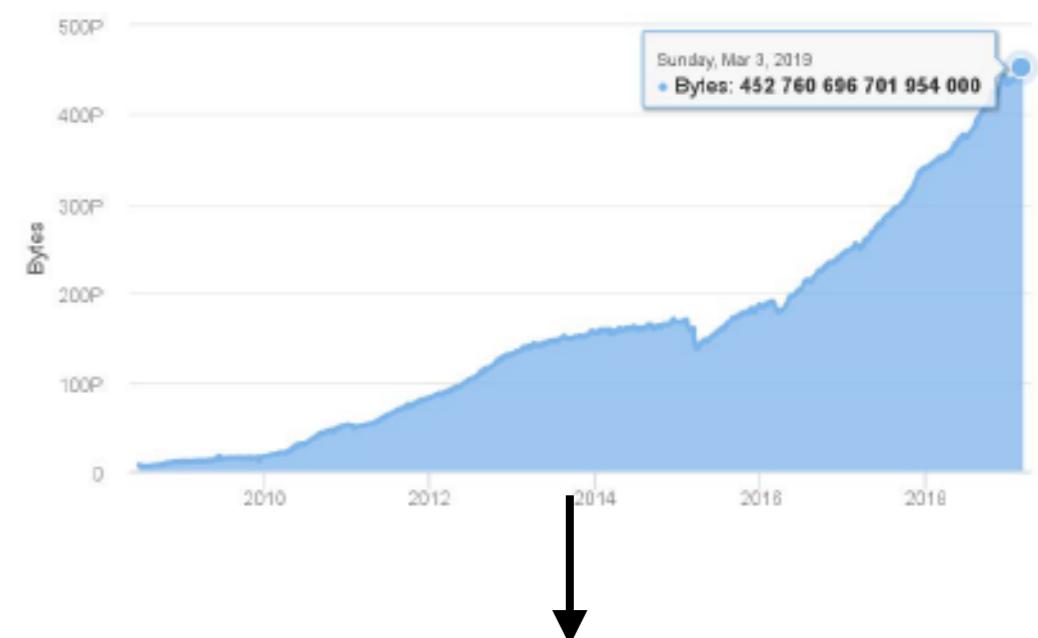
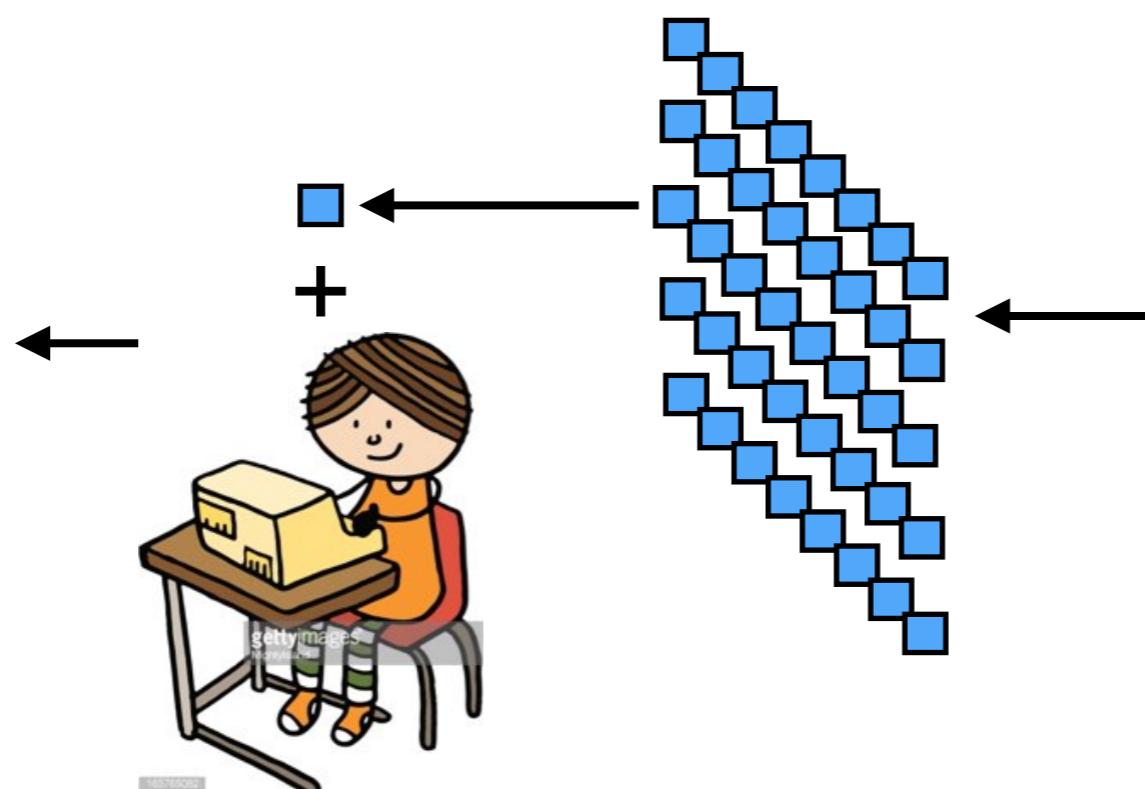
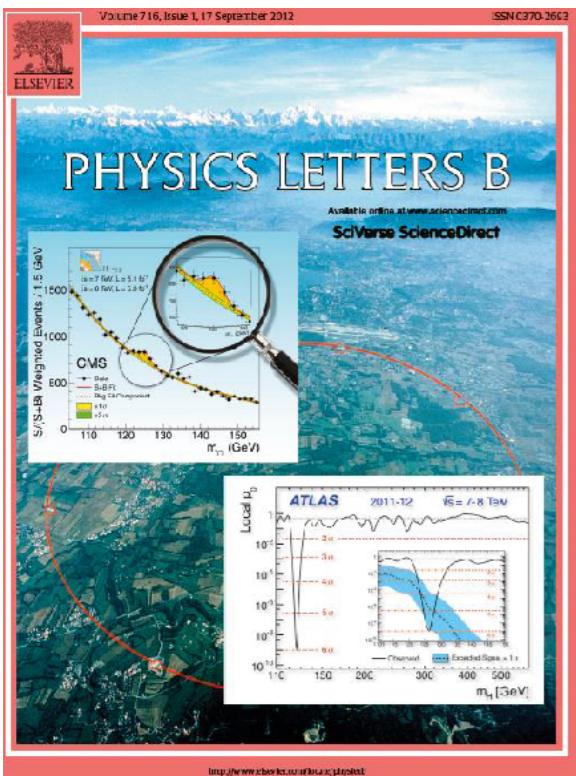


Differences in:

- ◎ Resolution
- ◎ Background composition
- ◎ Dataset

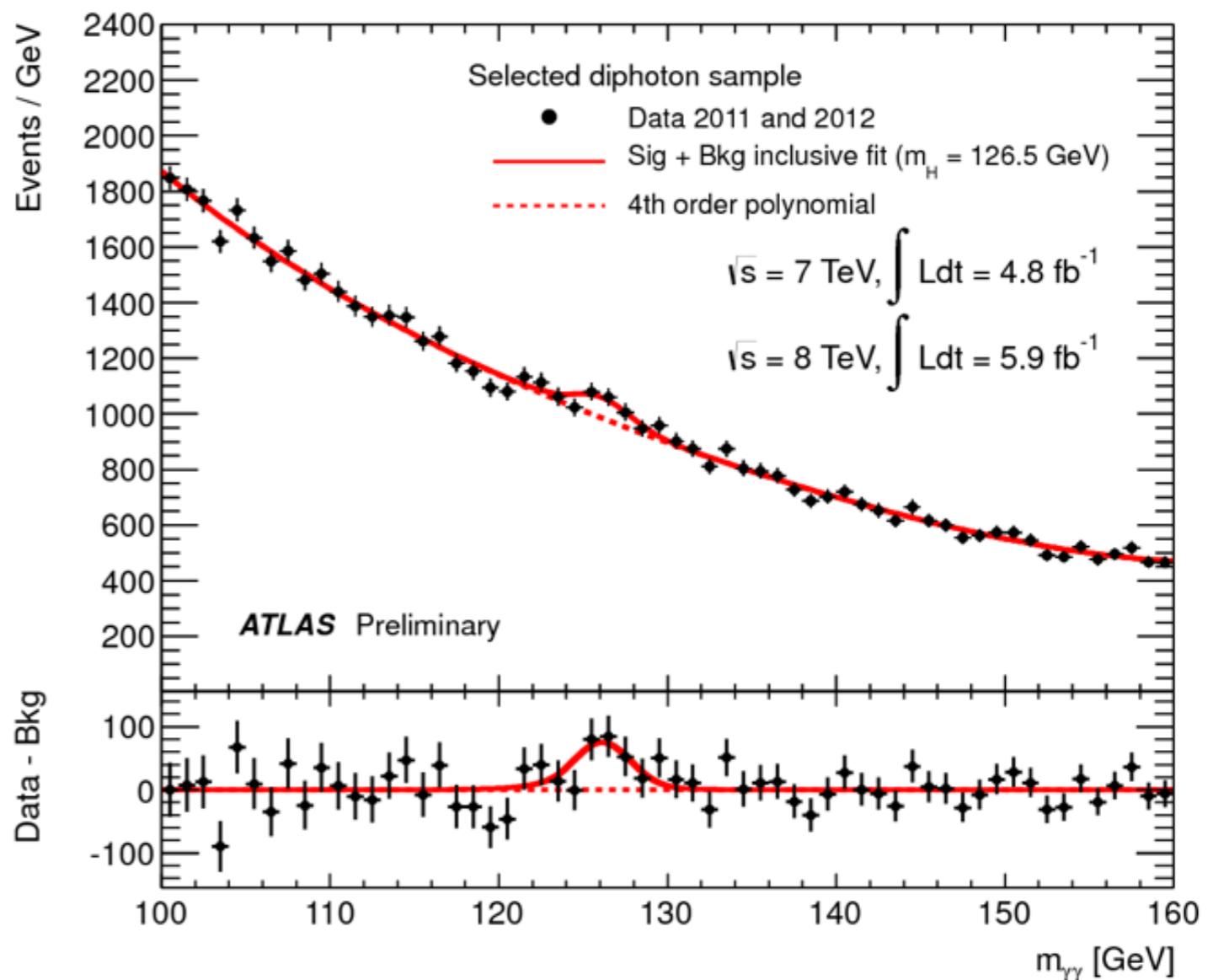
Why is the resolution worse in the muon channel?

# Data analysis

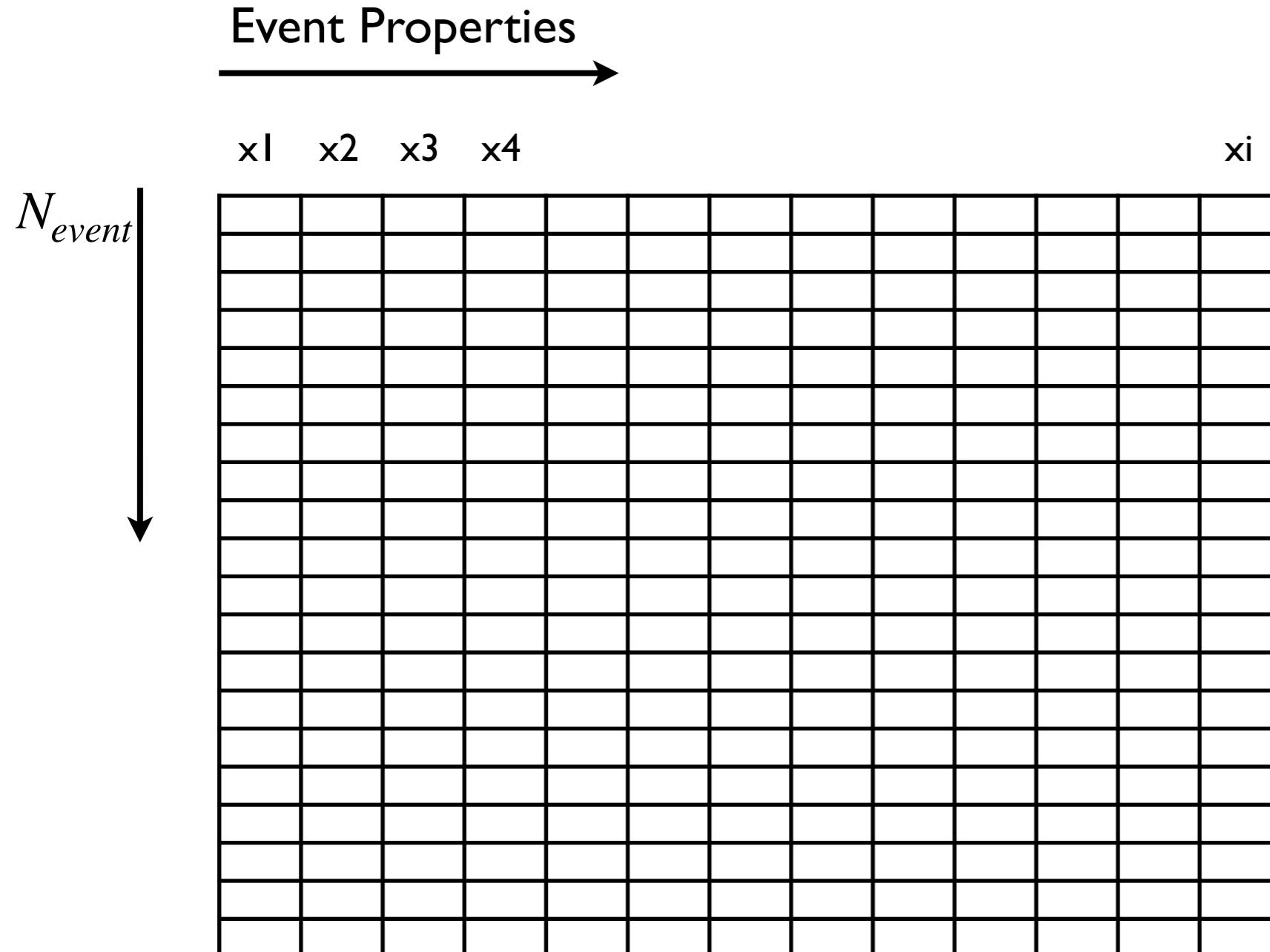


# Needles in haystacks

- We record billions of events
- The data are structured but each event is different - ***unique data science challenge***
- Data reduction proceeds via a two-pronged approach:
  - Select only the events that you are interested in
    - e.g. ***events with two photons***
  - Keep only the information you need
    - Throw away the rest !
  - Final statistical inference is only performed on the reduced data



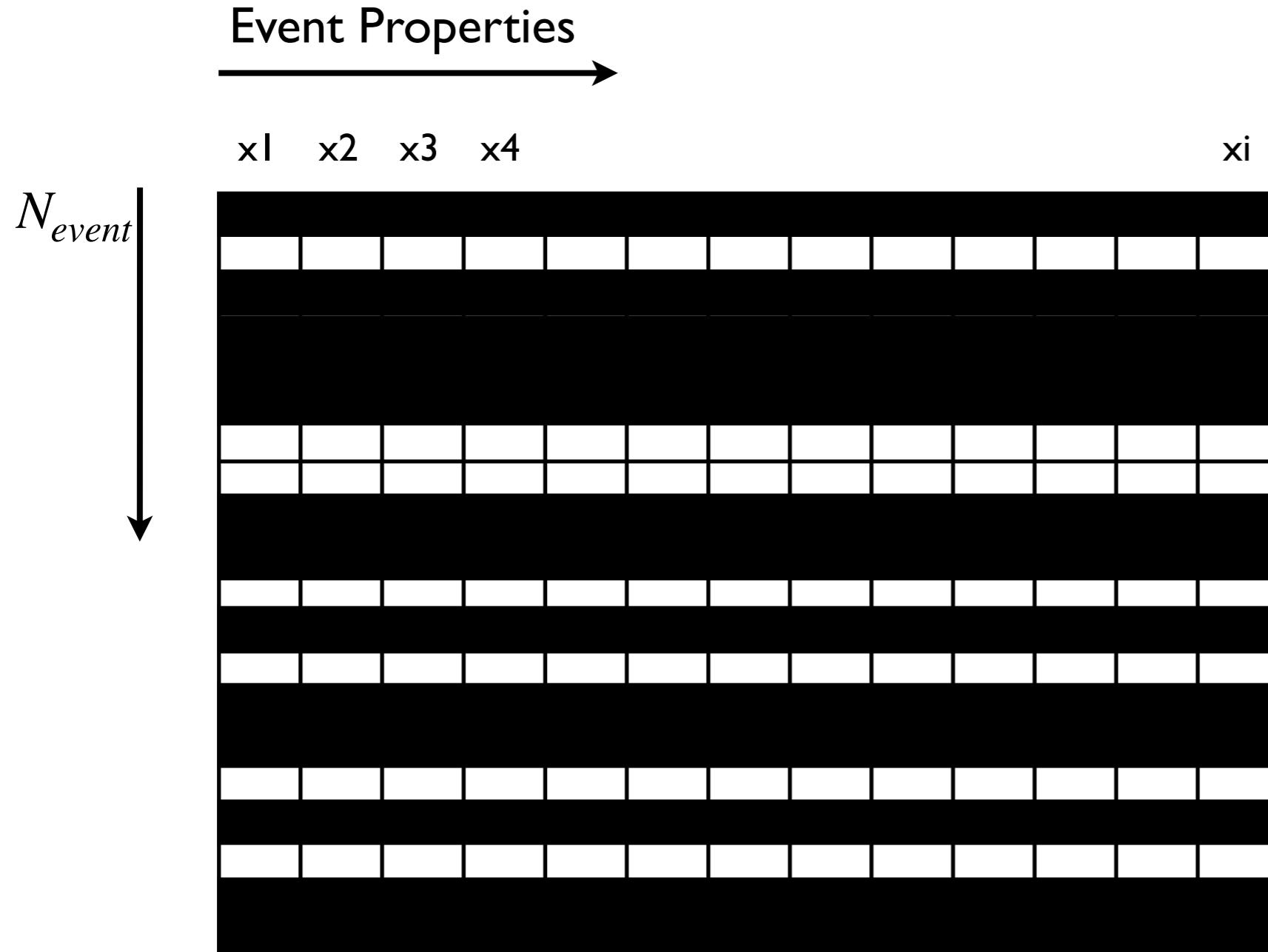
# Slice and dice - data reduction



Simplified picture, in reality the event properties depend on the event content found by event reconstruction

- There are two dimensions to our data challenge, one is the (billions) of individual events, the other is the properties of each event

# Slice and dice - data reduction



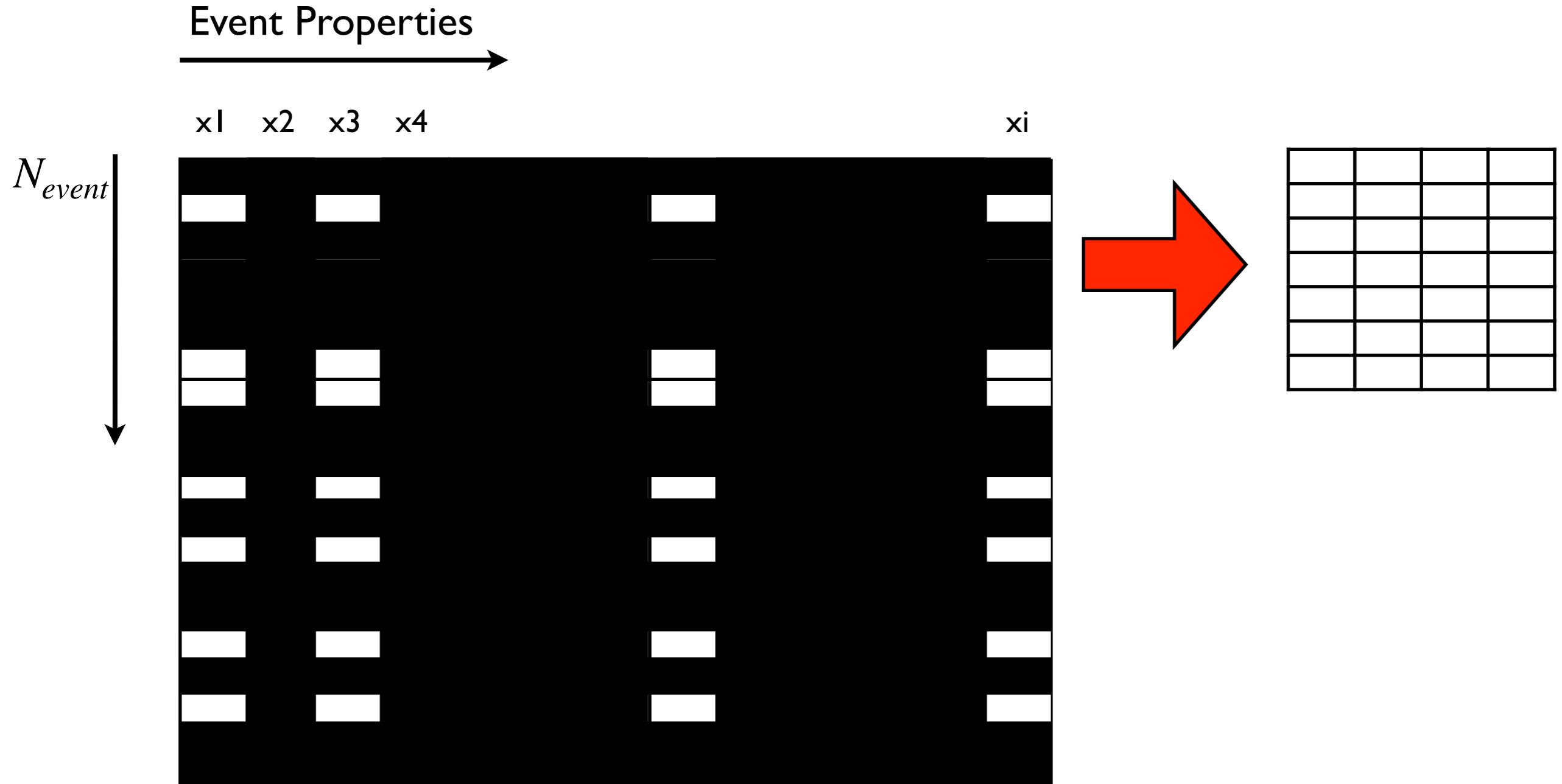
- We can reduce data by selecting only our interesting events

# Slice and dice - data reduction



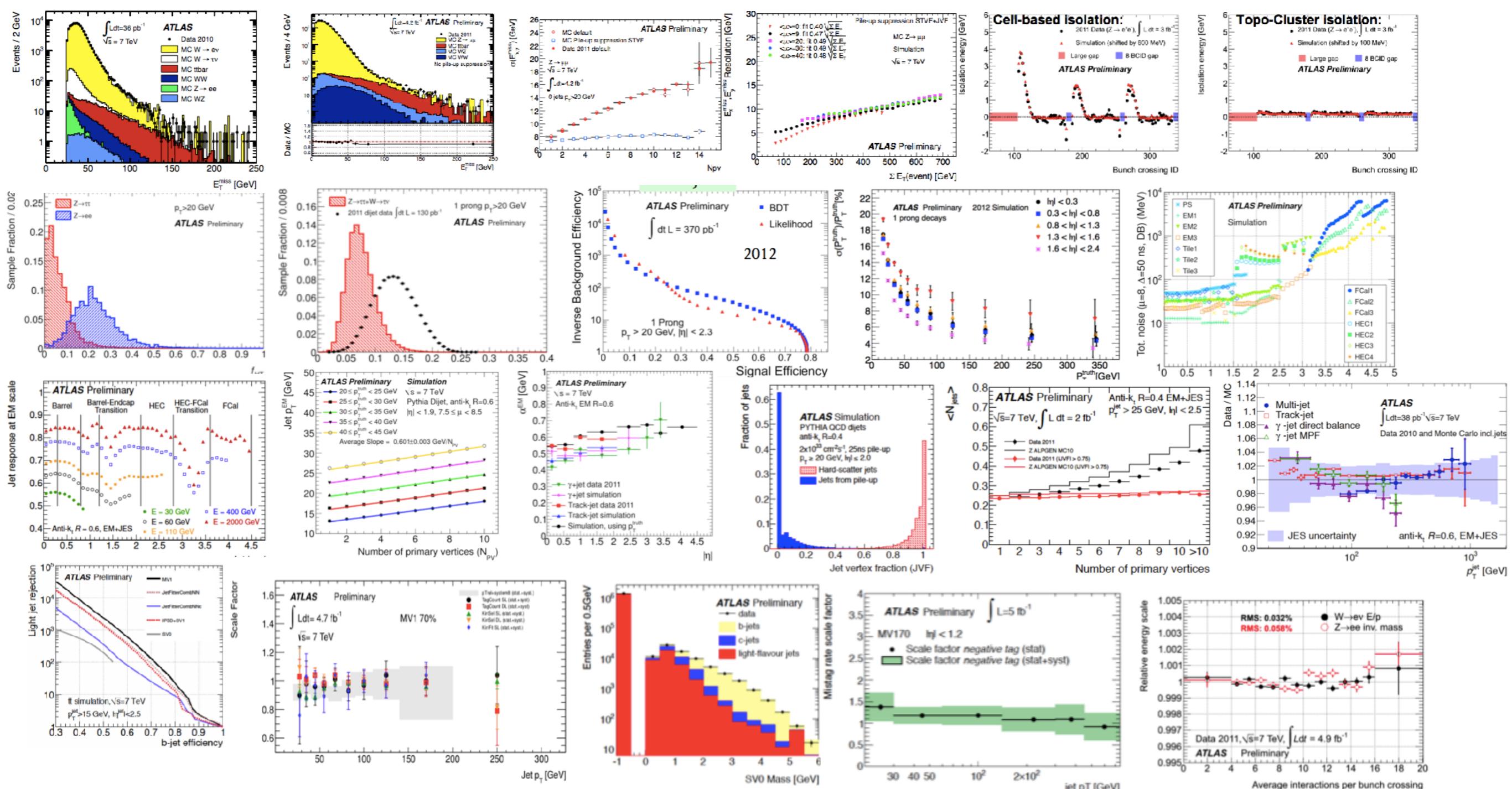
- And we can reduce data by selecting only the properties needed for our analysis

# Slice and dice - data reduction



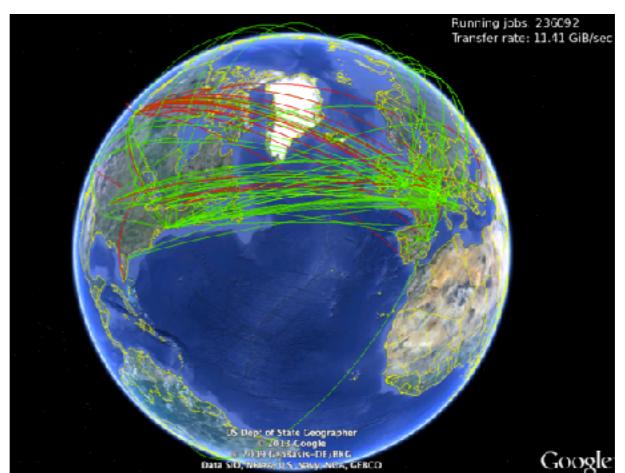
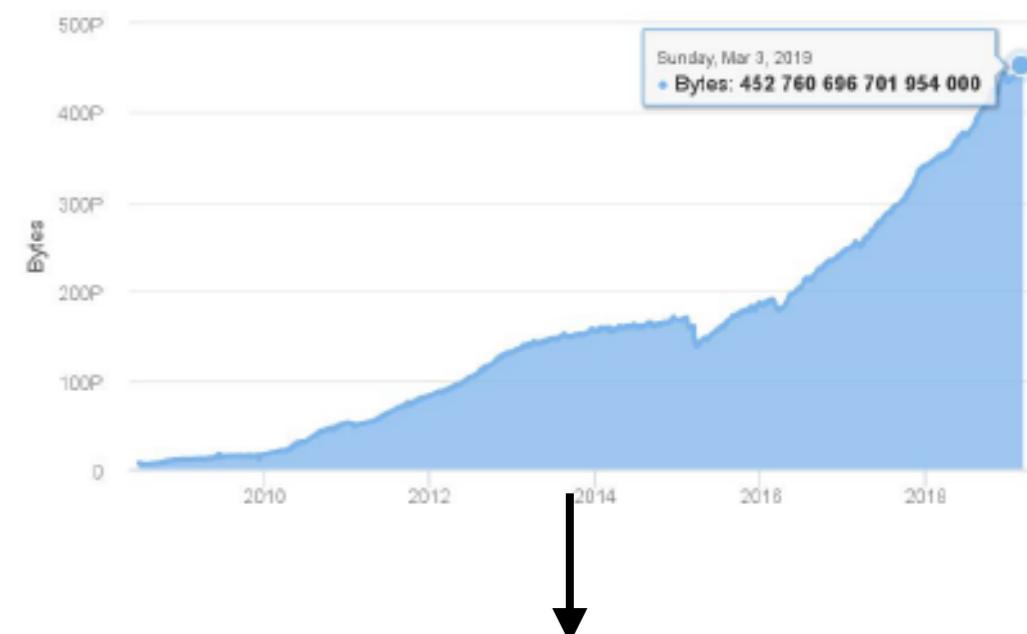
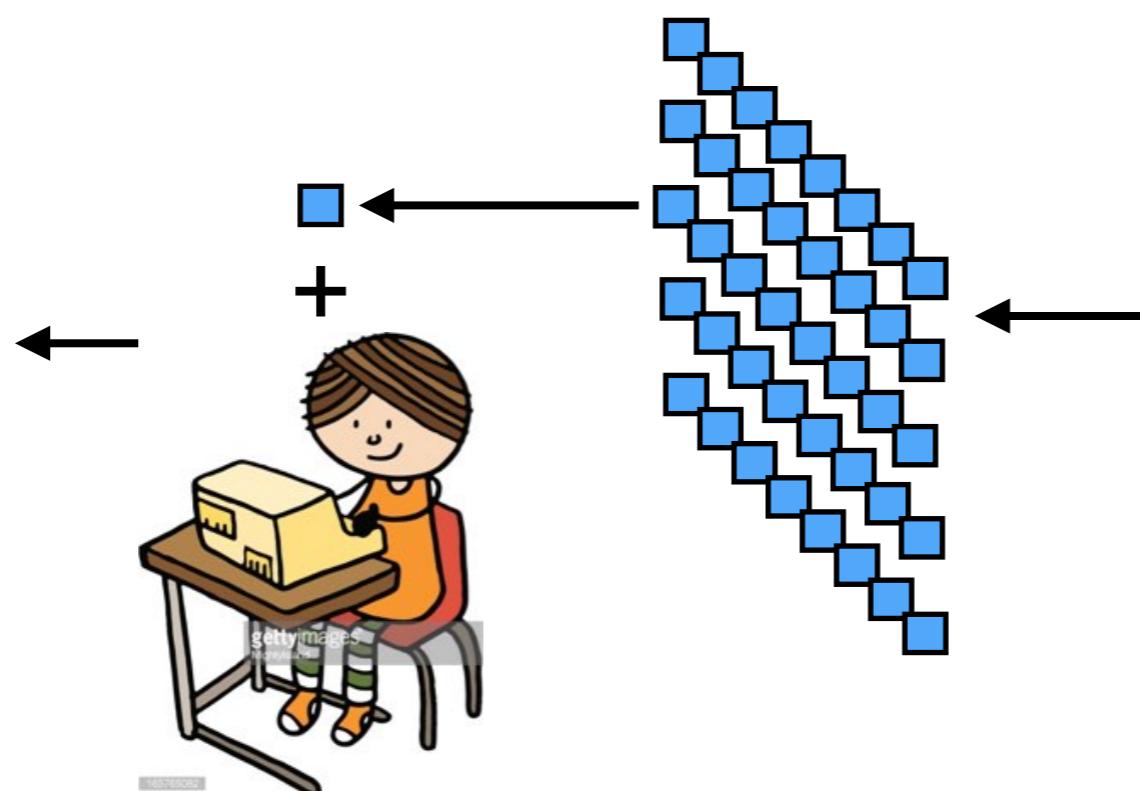
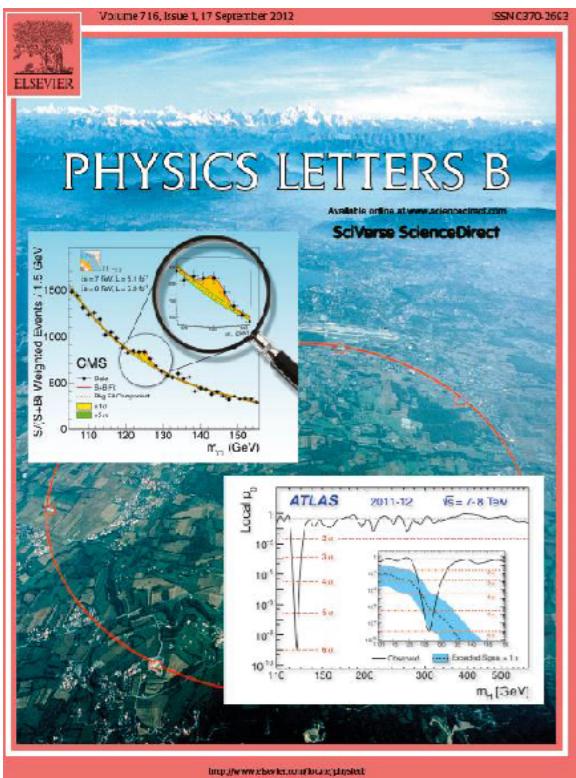
- Data reduction usually aims for factors of 100 or more (more than shown here !)

# Ingredients to the ATLAS physics program



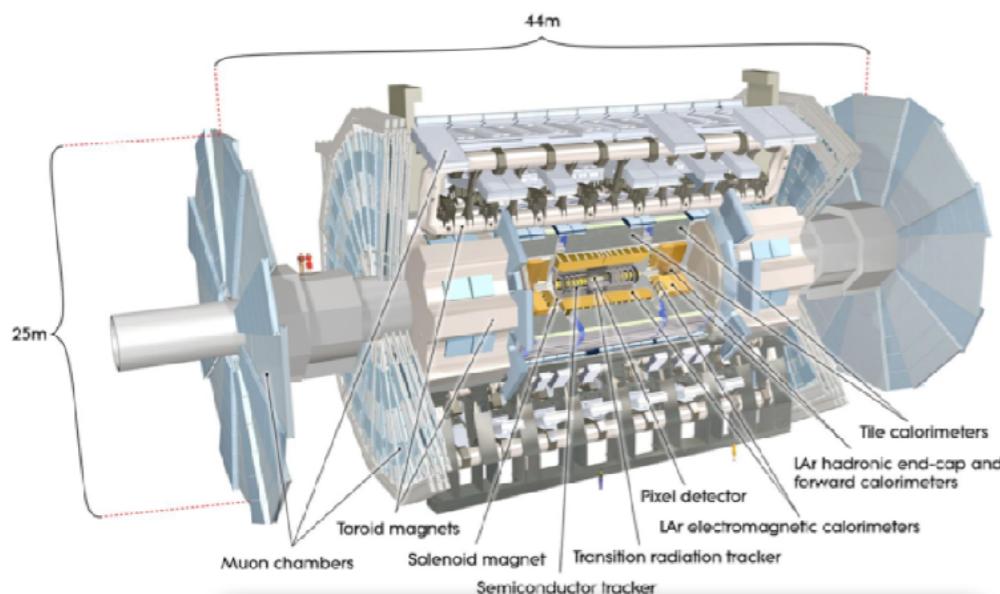
- We make lots of reduced samples of both data and simulation, which all need to be replicated around the world - a computing challenge !

# The best computing model

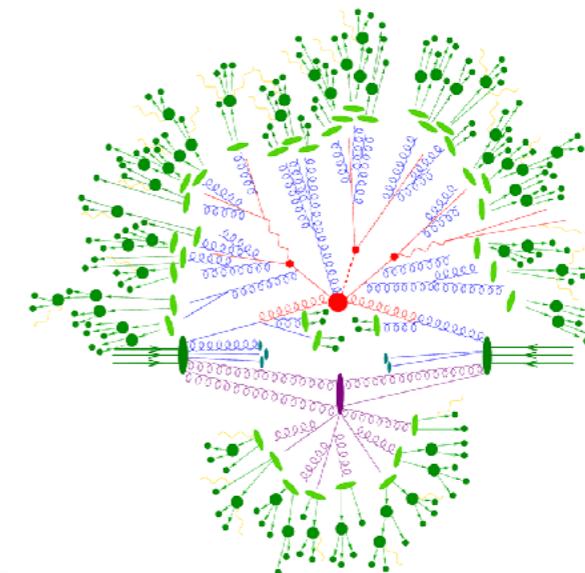


- How to most efficiently do this across the whole physics program making the best use of computing resources and the best use of people's time is an important question

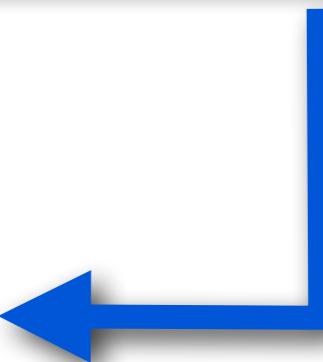
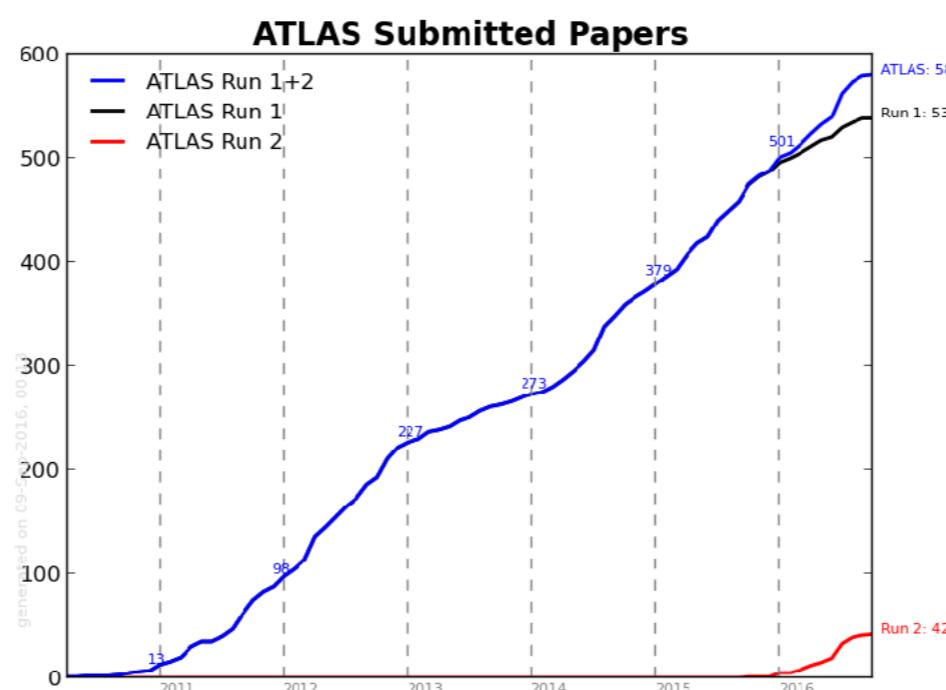
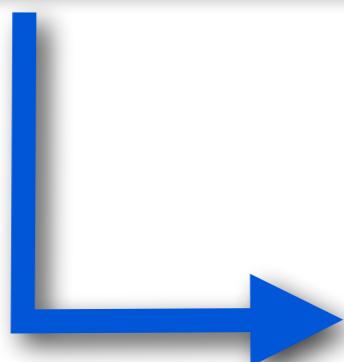
# Now you know how to do exabyte-scale physics analysis!



Exabytes of Data

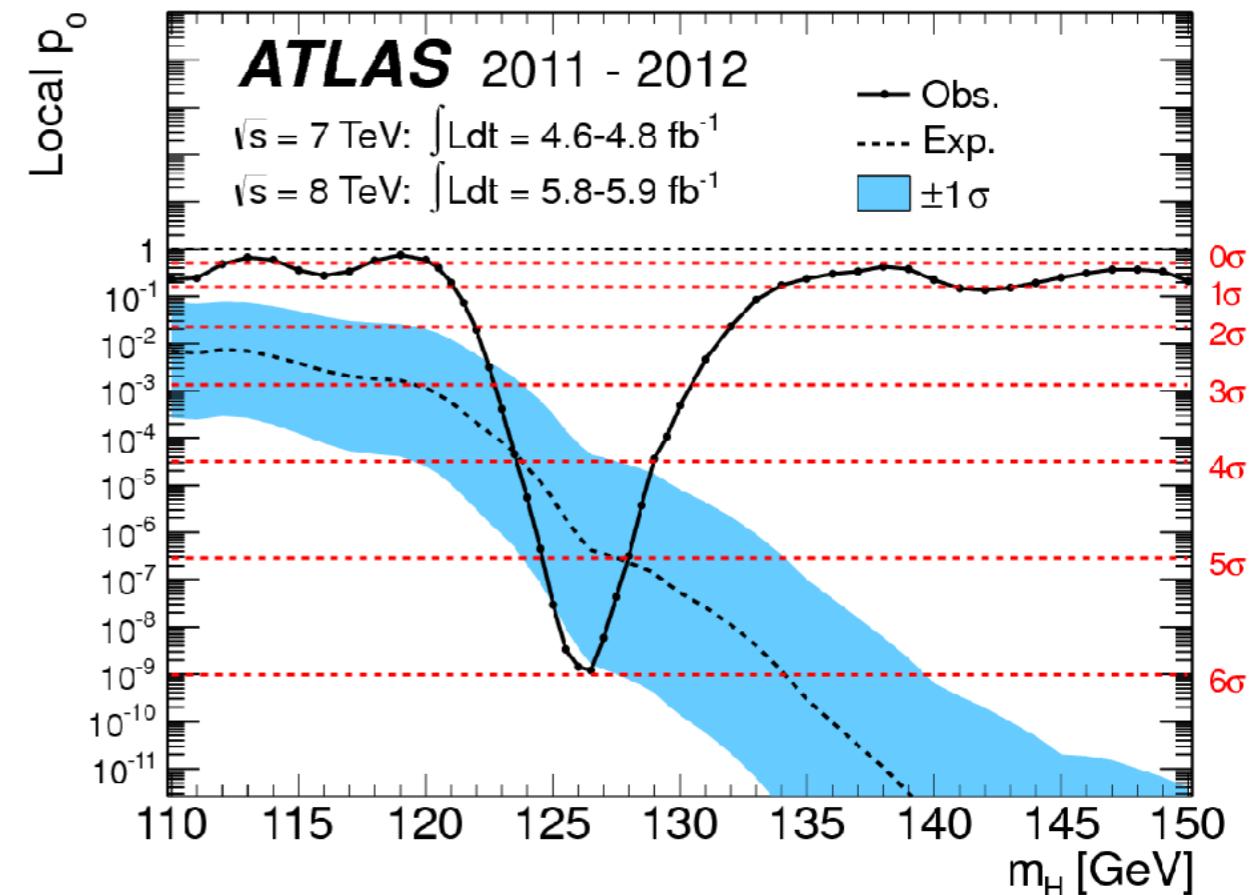
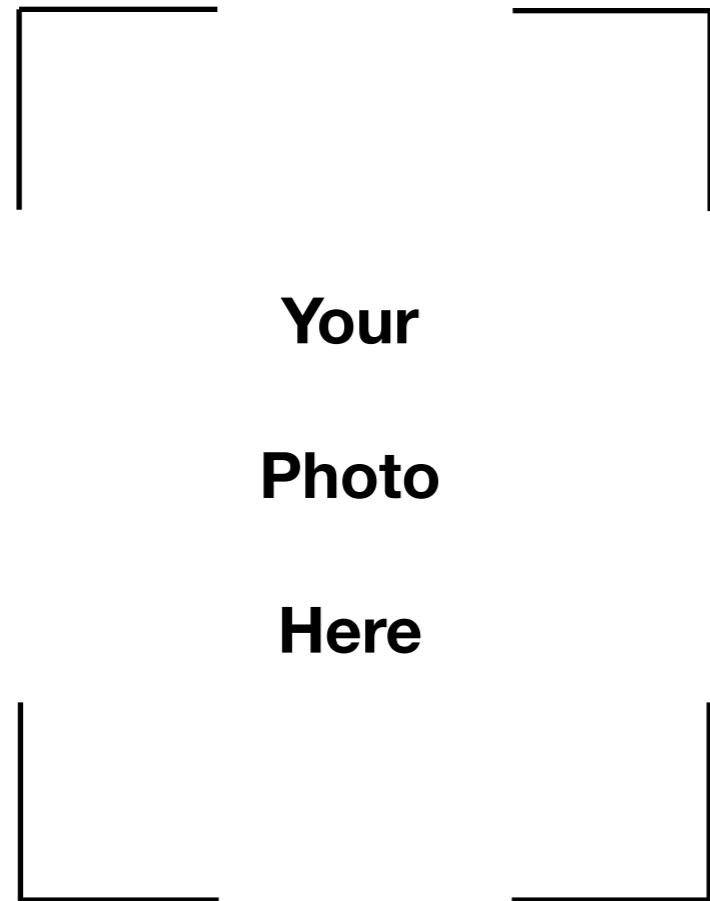


Exabytes of Simulation



Publish!

# Now it's over to you !



- Our future computing needs outstrip our computing resources
  - and computing gets more heterogeneous and complicated
  - and we want to be as environmentally-responsible as possible
- So you have work to do - good luck and have fun!

# Contact details

- I am usually based at Geneva Observatory in Versoix, but will be here at CERN Wednesday 28th through Friday 30th June.
  - I will be available for Q&A every afternoon from 3pm-4pm in restaurant 1, feel free to send questions to my email
- email: [paul.laycock@unige.ch](mailto:paul.laycock@unige.ch)