

# Flavour Physics: A Taster

CERN Summer Student Lecture Programme 2023

Lecture 3 of 3: Flavour in the LHC era and beyond

**17-19 July 2023**

**Mark Williams**  
**University of Edinburgh**



**THE UNIVERSITY  
of EDINBURGH**

# Introduction

Yesterday we covered several key flavour physics ideas:

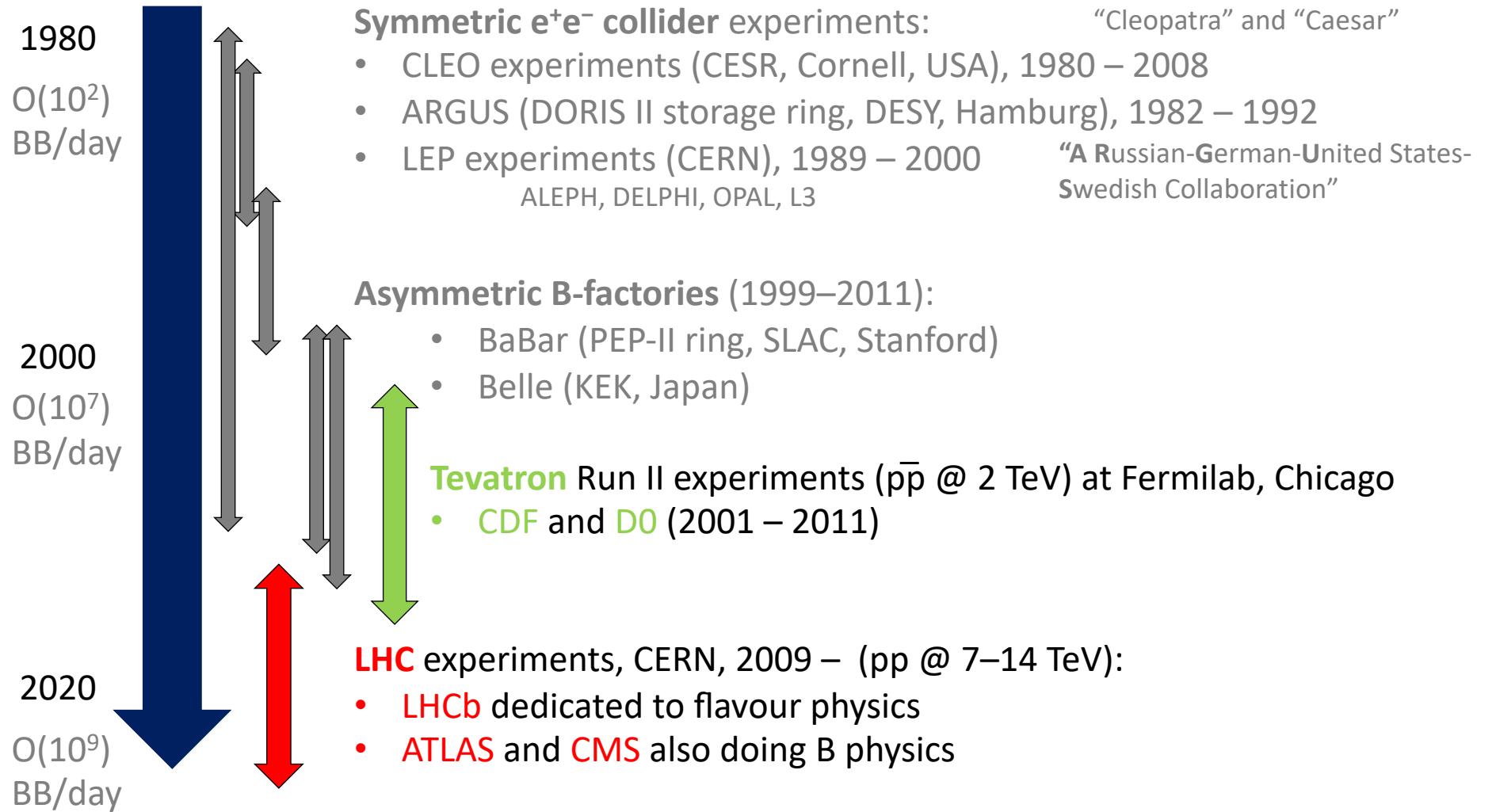
- CP violation in the SM (quark sector)
- Unitarity triangle(s)
- Measuring CKM phases
- B-factory measurements of  $\beta$  and  $\alpha$

Today we will discuss b (and c) physics in the LHC era (and beyond):

- Hadron colliders vs B-factories
- Mixing and CP violation in  $B_s^0$  and  $D^0$  mesons
- CKM angle gamma
- Rare decays and lepton universality
- The future

# Part I: Flavour physics at hadron colliders

# Overview of b experiments



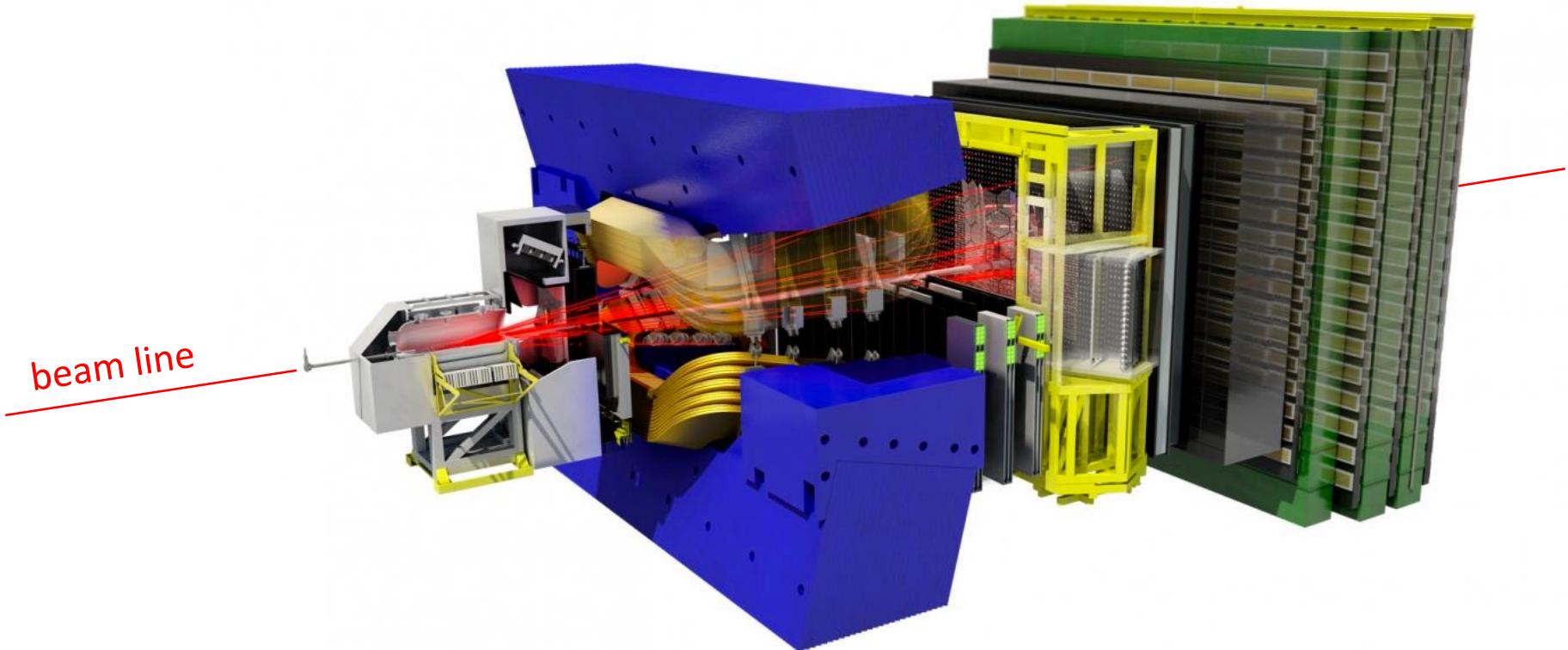
# Flavour physics at hadron colliders

	B Factories	Hadron colliders
	<i>Belle (1999-2010)</i> <i>BaBar (1999-2008)</i>	<i>Tevatron (&lt;2 TeV, 1983–2011)</i> <i>LHC (&lt;14 TeV, 2008–)</i>
Collision environment	Asymmetric $e^+e^- \rightarrow Y(4S)$	$p p$ or $p\bar{p}$ (also ions...)
Flavour tagging (initial $B^0$ or $\bar{B}^0$ )	Clean! Pure $B\bar{B}$ event ✓	Busy! Proton remnants give background particles
	Excellent ✓ (30% ‘tagging power’)	Challenging (~5%)

# Flavour physics at hadron colliders

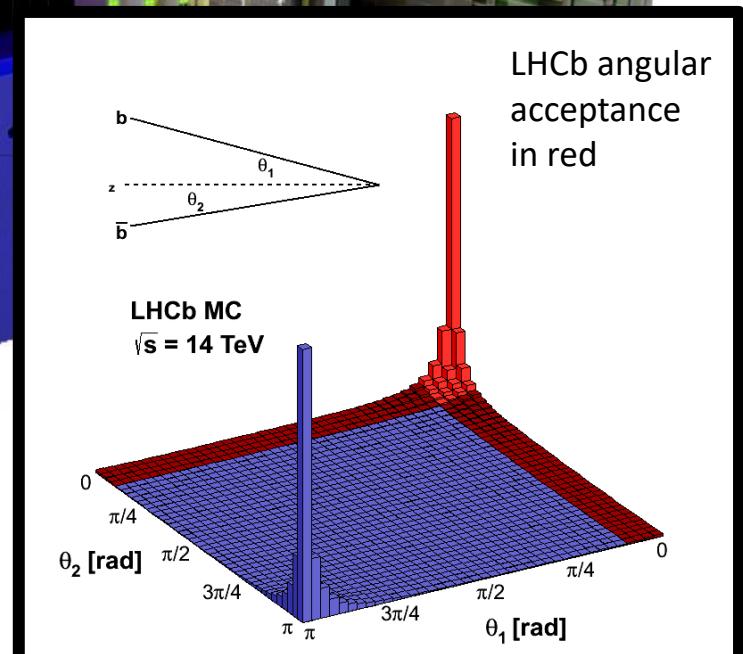
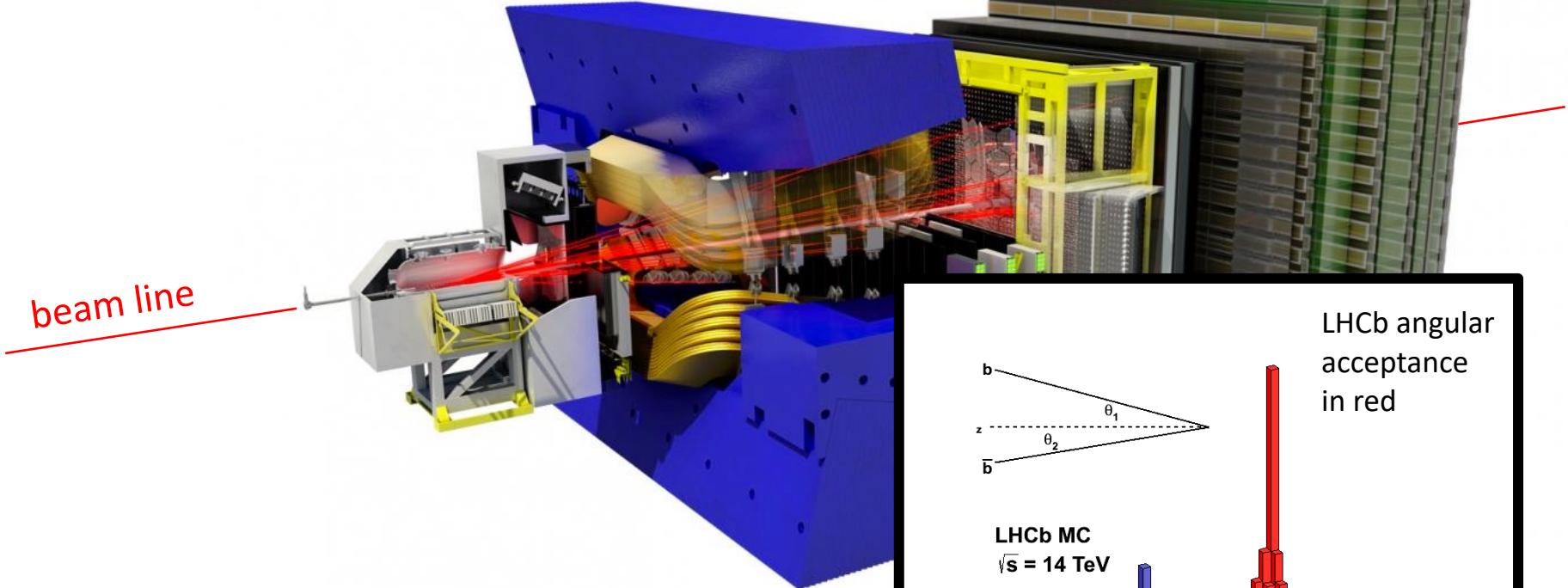
	B Factories	Hadron colliders
	<i>Belle (1999-2010)</i> <i>BaBar (1999-2008)</i>	<i>Tevatron (&lt;2 TeV, 1983–2011)</i> <i>LHC (&lt;14 TeV, 2008–)</i>
Collision environment	Asymmetric $e^+e^- \rightarrow Y(4S)$  Clean! Pure $B\bar{B}$ event ✓	pp or $p\bar{p}$ (also ions...)  Busy! Proton remnants give background particles
Flavour tagging (initial $B^0$ or $\bar{B}^0$ )	Excellent ✓ (30% ‘tagging power’)	Challenging (~5%)
Production $\sigma(B)$	1 nb	~100–500 $\mu b$ ✓
B hadron boost	Small ( $\beta\gamma \approx 0.5$ )	Large ( $\beta\gamma \approx 100$ ) ✓
B hadrons created	$B^+B^-$ (50%), $B^0\bar{B}^0$ (50%)	$B^\pm$ (40%), $B^0$ (40%), $B_s^0$ (10%) b baryons (10%) ✓

# LHCb experiment (v2010-2018)



# LHCb experiment

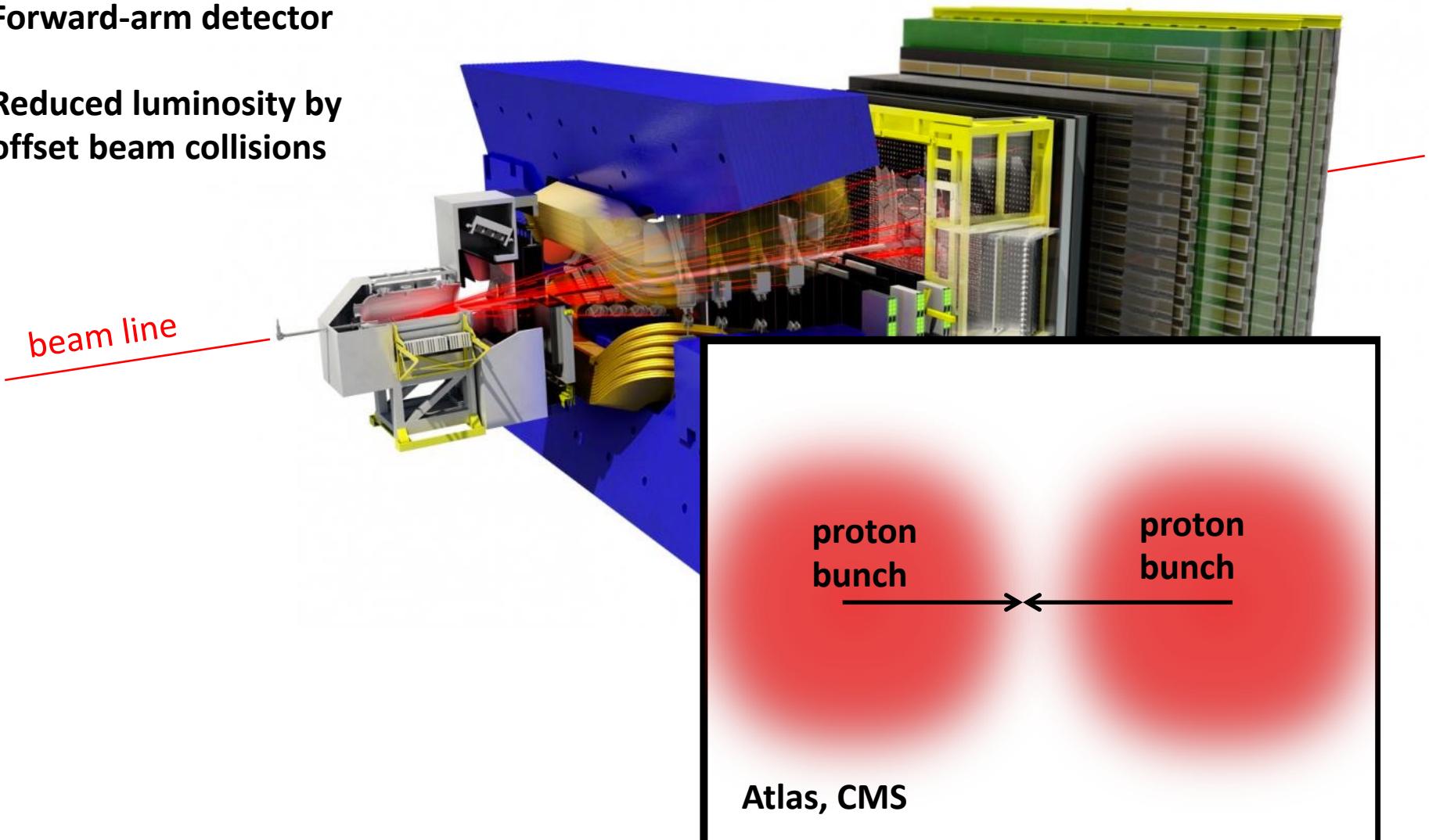
## Forward-arm detector



# LHCb experiment

Forward-arm detector

Reduced luminosity by  
offset beam collisions

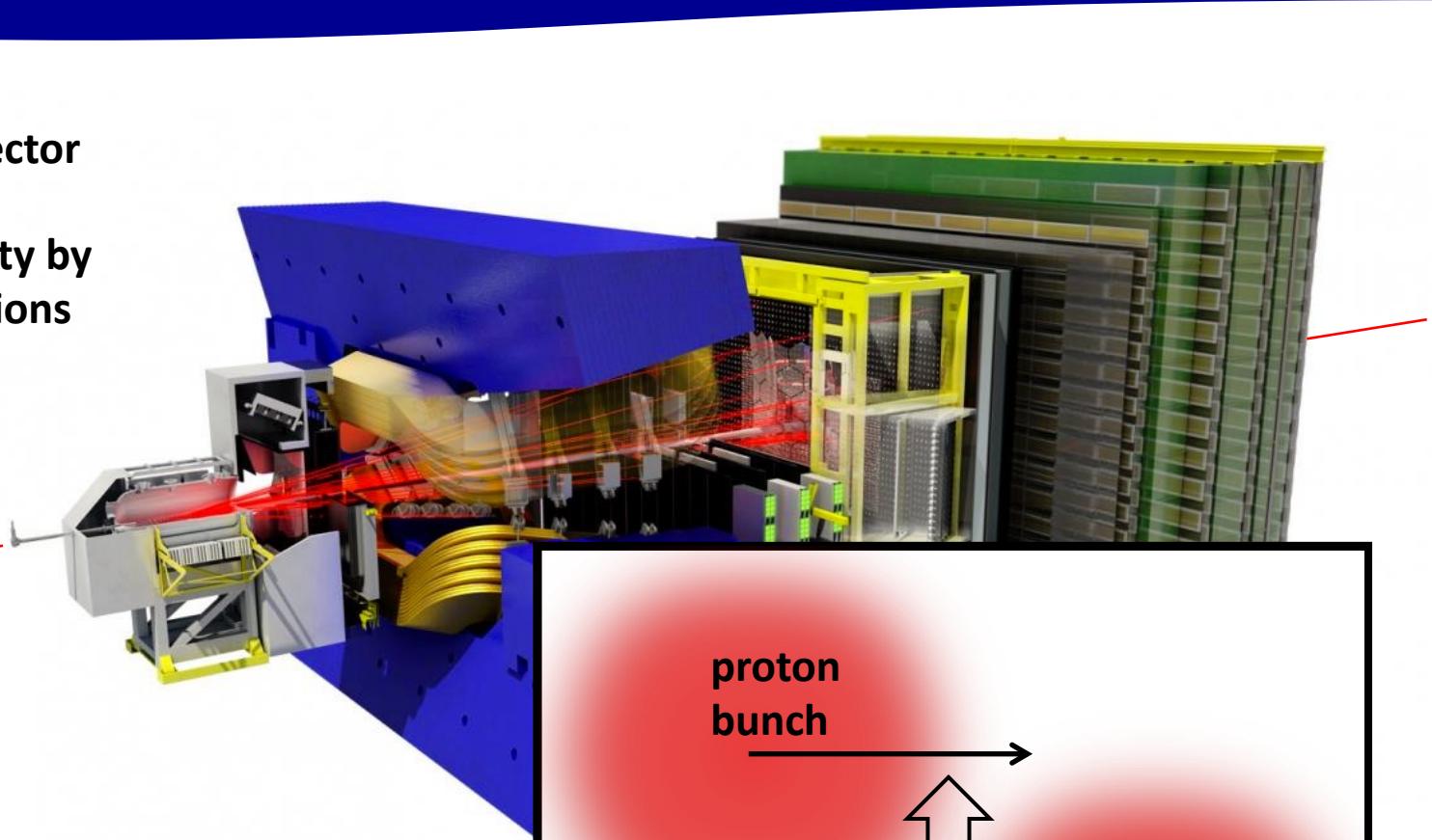


# LHCb experiment

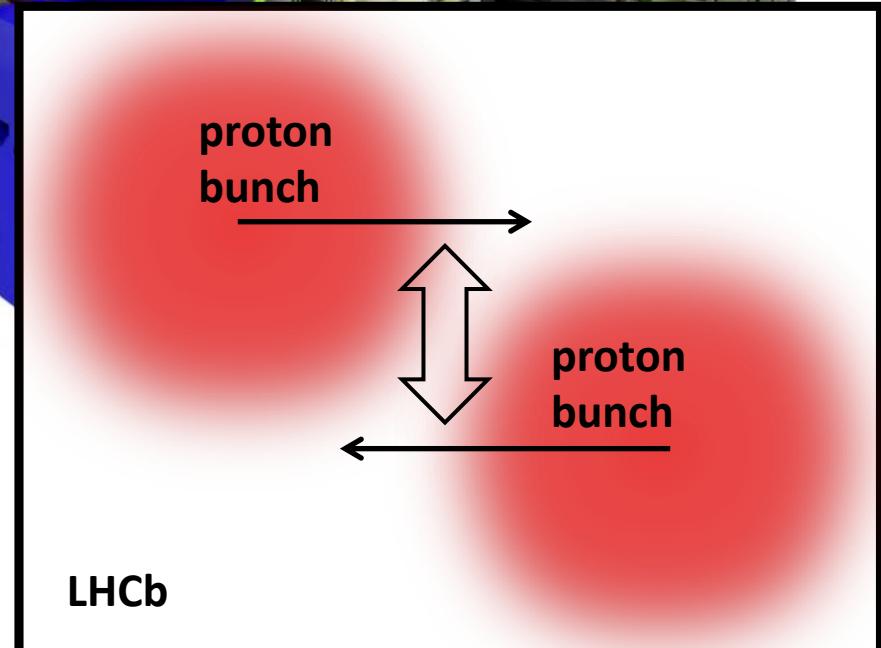
Forward-arm detector

Reduced luminosity by  
offset beam collisions

beam line



Beams move closer as  
intensity reduces over time  
⇒ luminosity levelling

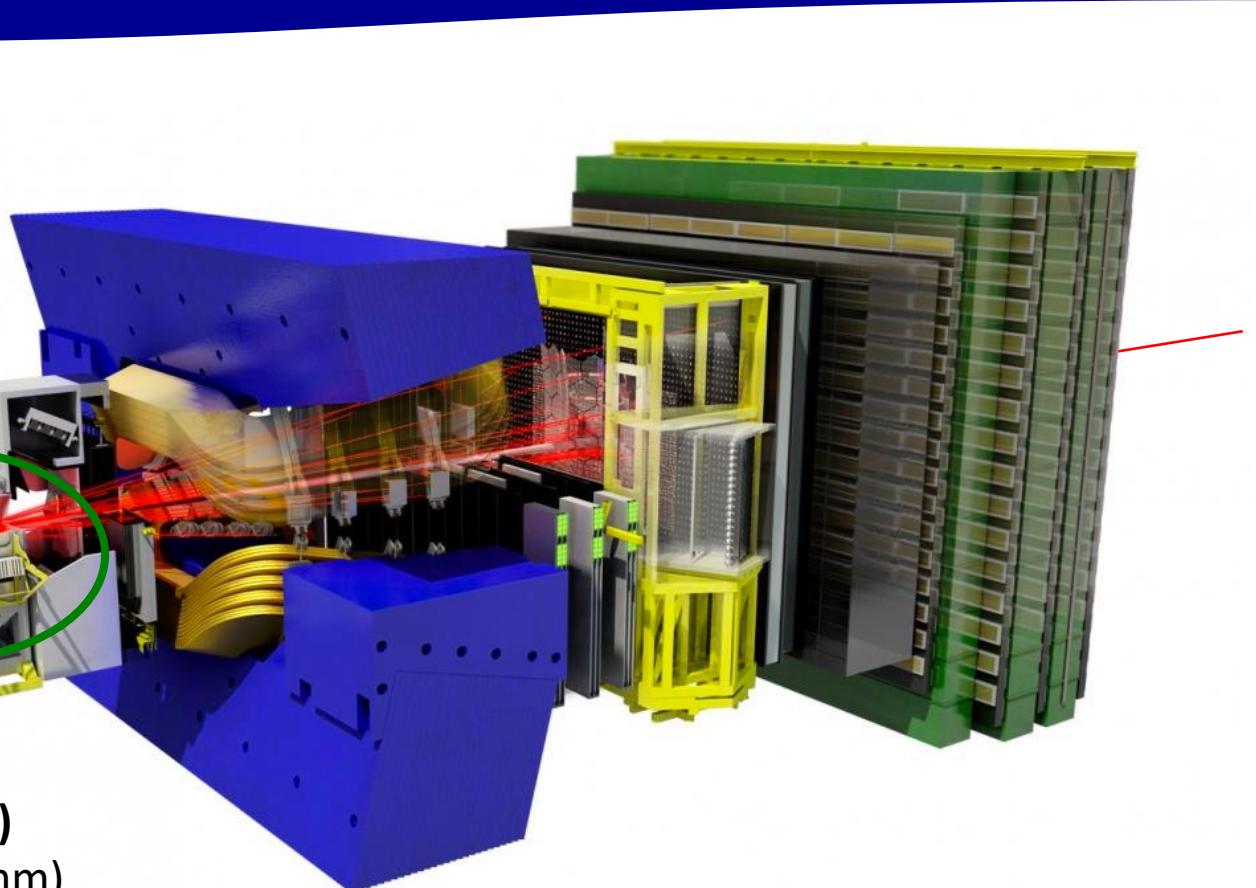
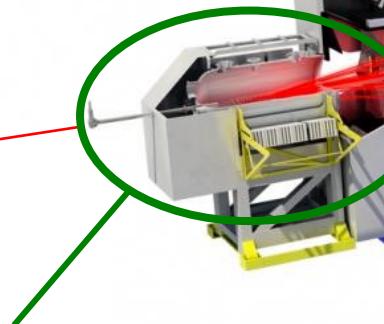


# LHCb experiment

Forward-arm detector

Reduced luminosity by  
offset beam collisions

beam line



Precise Vertex Location (VELO)

Very close to the collisions (8mm)

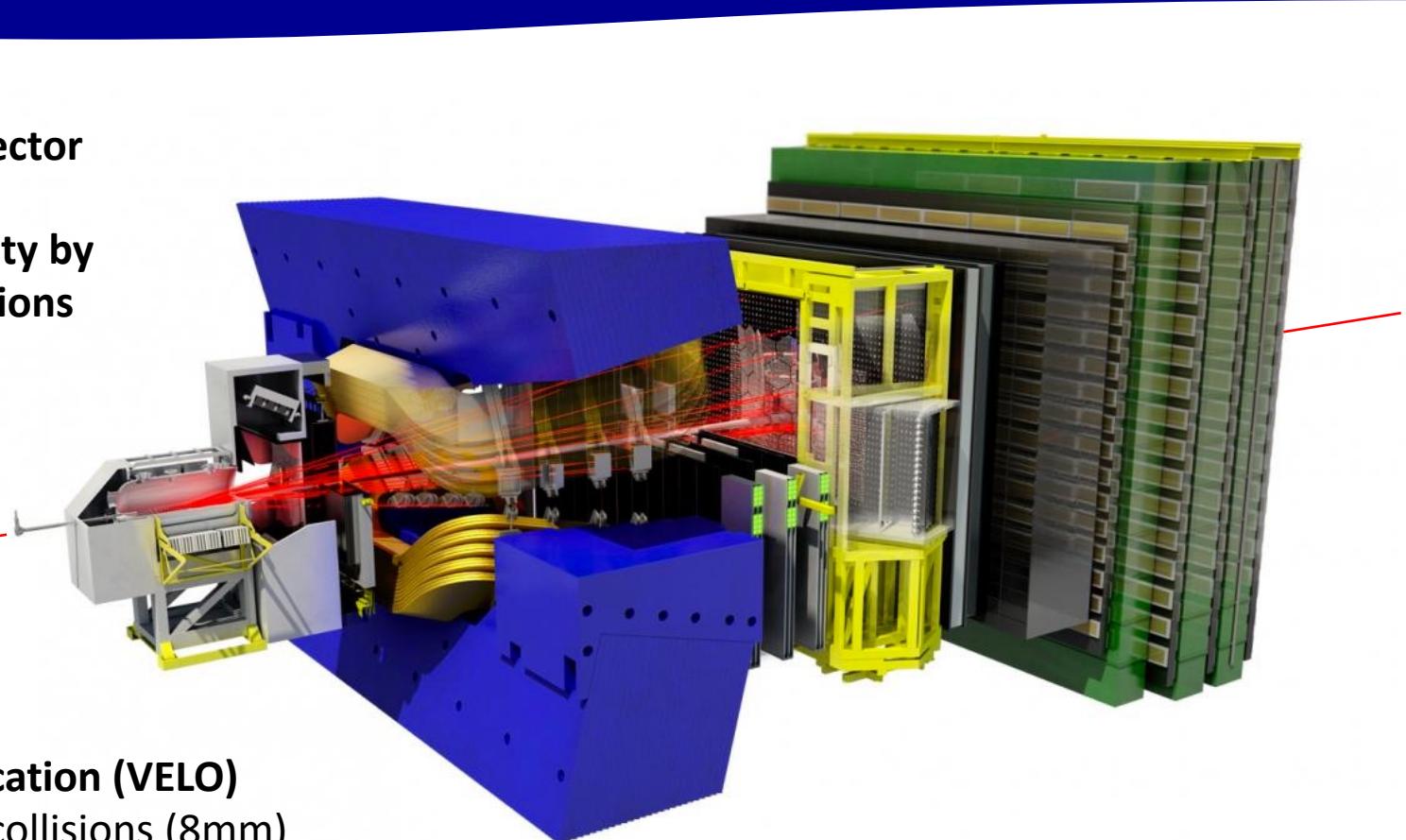
→ must be moved away for safety  
every time beam is injected (!)

# LHCb experiment

Forward-arm detector

Reduced luminosity by  
offset beam collisions

beam line



Precise Vertex Location (VELO)

Very close to the collisions (8mm)  
→ must be moved away for safety  
every time beam is injected (!)

**Excellent particle identification** using Cherenkov  
radiation to measure particle speed

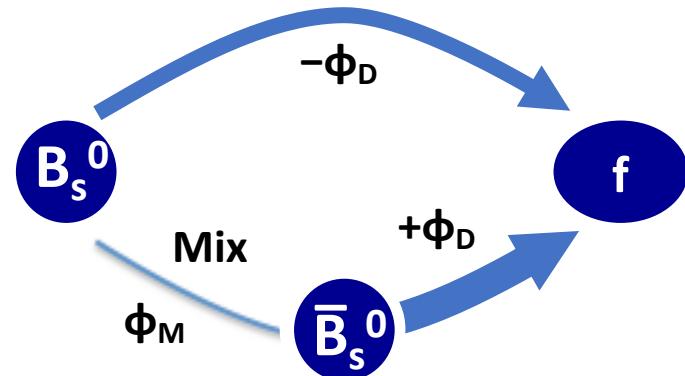
**Powerful software-based trigger** – make  
decisions using full event reconstruction

# Part IIa: CP violation in $B_s^0$ mesons

# CP violation in $B_s^0$ mixing

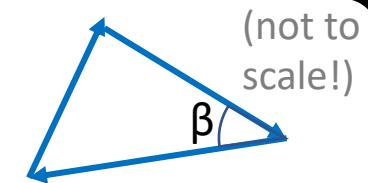
$B_s^0$  equivalent of  $\sin(2\beta)$  measurement

For  $B_s^0$  system, we are dealing with a different unitarity triangle from  $B^0$



$B^0$  case:  $V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$   $\Rightarrow$  angles  $\alpha, \beta, \gamma$

$O(\lambda^3)$        $O(\lambda^3)$        $O(\lambda^3)$



$B_s^0$  case:  $V_{us}V_{ub}^* + V_{cs}V_{cb}^* + V_{ts}V_{tb}^* = 0$   $\Rightarrow$  angles  $\alpha_s, \beta_s, \gamma_s$

$O(\lambda^4)$        $O(\lambda^2)$        $O(\lambda^2)$



In fact we measure  $\phi_s = \phi_M - 2\phi_D$       BUT      for tree-level decays  $b \rightarrow c\bar{c}s$ :

$$\boxed{\phi_s = -2\beta_s}$$

$\Rightarrow$  Just as with  $\sin(2\beta)$  we need to pick a 'golden mode'

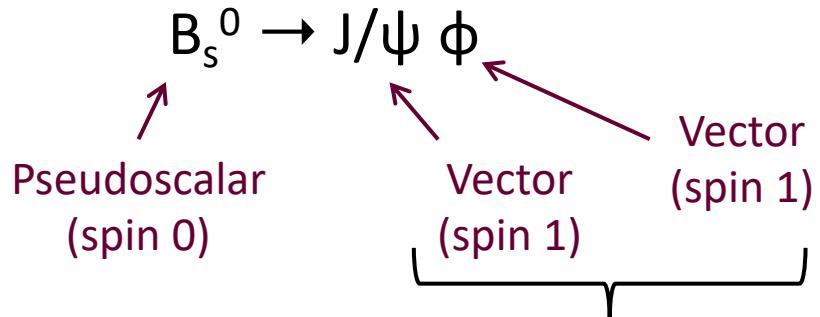
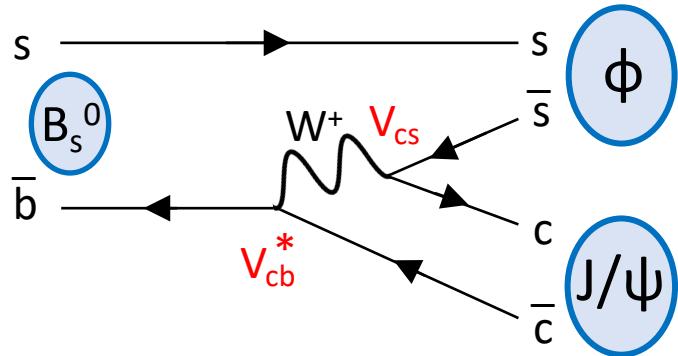
# CP violation in $B_s^0 \rightarrow J/\psi \phi$

Decay similar to  $B^0 \rightarrow J/\psi K_s^0$   
with spectator quark exchange  $d \rightarrow s$

Extra challenge – final state is not CP eigenstate  
⇒ need to analyse angular distributions to  
disentangle three CP states

In SM,  $\phi_s$  is not independent variable  
⇒ highly constrained by CKM mechanism  
(only 4 free parameters)

$$\begin{aligned} -2\beta_s &= -2 \arg(-(V_{ts} V_{tb}^*)/(V_{cs} V_{cb}^*)) \\ &= -0.0369 {}^{+0.0007}_{-0.0010} \text{ rad (SM)} \end{aligned}$$



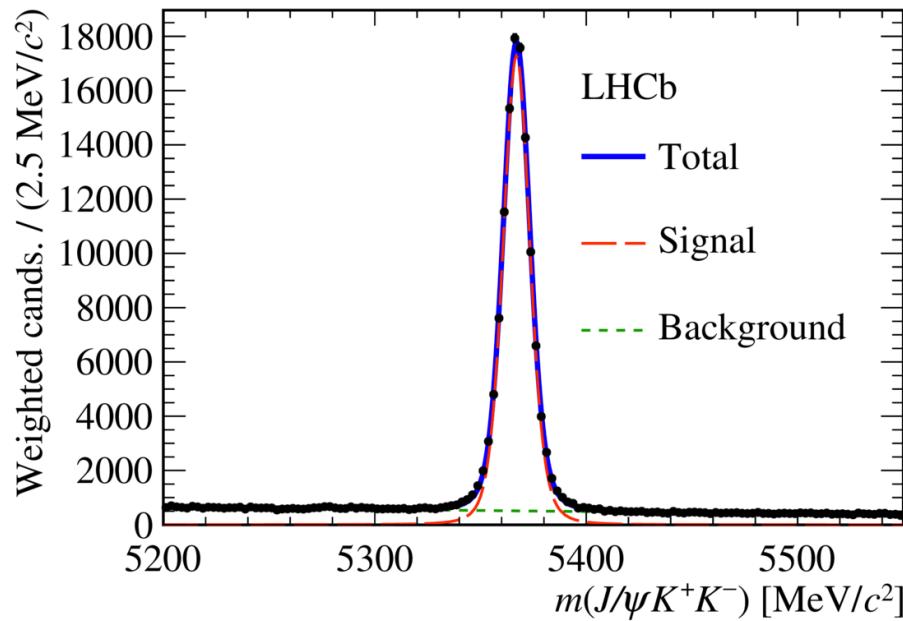
- ↑↑: L=2, CP even, amplitude  $A_{||}(t)$
- ↑↓: L=0, CP even,  $A_0(t)$
- ↑→: L=1, CP odd,  $A_\perp(t)$

# CP violation in $B_s^0 \rightarrow J/\psi\phi$

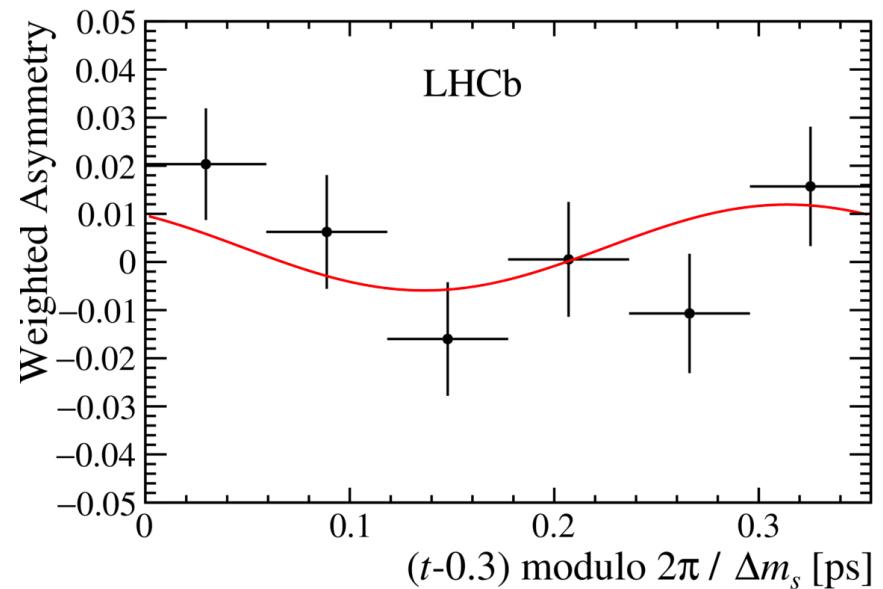
Clean decay mode:  $J/\psi \rightarrow \mu^+\mu^-$ ,  $\phi \rightarrow K^+K^-$

⇒ muons provide good trigger signature  
⇒ ATLAS & CMS join the party!

<https://doi.org/10.1140/epjc/s10052-019-7159-8>



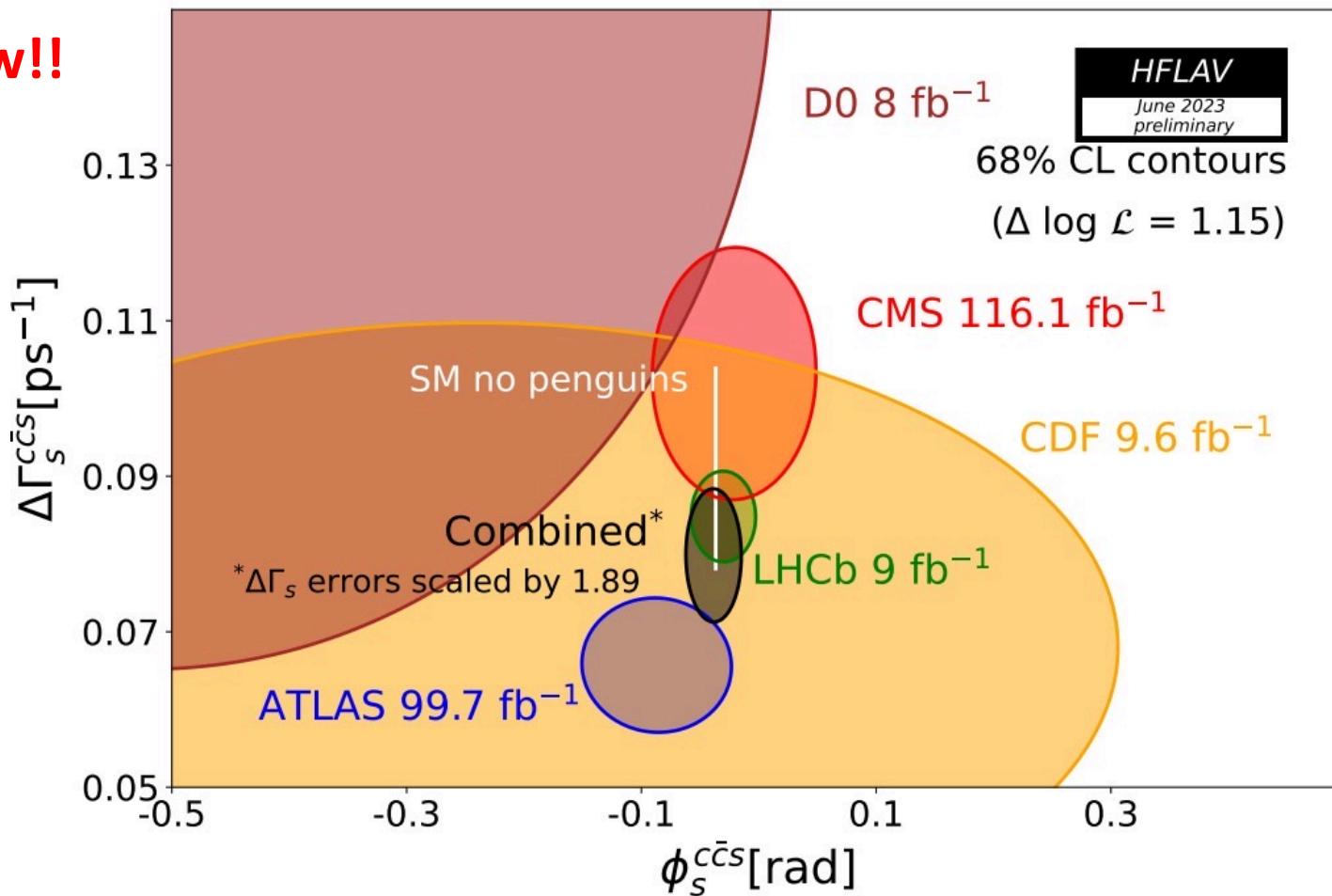
High-purity signal



Asymmetry vs decay time [c.f.  $\sin(2\beta)$ ]  
No significant asymmetry ⇒  $\phi_s \approx 0$

# CP violation in $B_s^0 \rightarrow J/\psi\phi$

New!!



**World averages:**  $\Phi_s = -0.039 \pm 0.016$  rad  
(provisional)

$\Delta\Gamma_s = 0.082 \pm 0.005$  ps<sup>-1</sup>

<https://indico.cern.ch/event/1281612/>  
CERN Seminar 13 June 2023

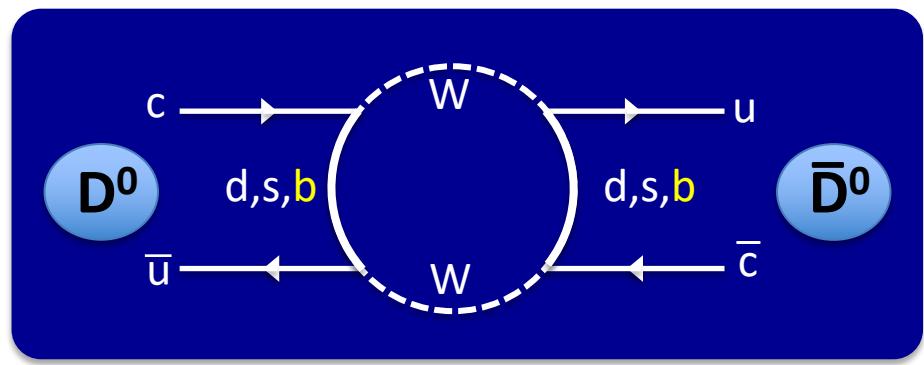
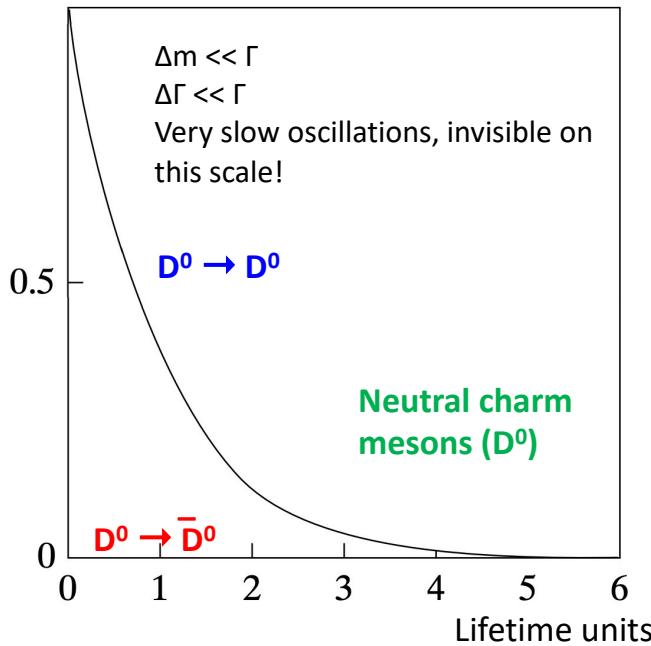
# Part IIb: Charm mixing and CPV

# Charm mixing and CP violation

We've covered strange ( $K^0$ ) and beauty ( $B^0, B_s^0$ ) physics: what about the charm quark?

Neutral charm meson ( $D^0$ ) can oscillate

but... both  $\Delta m$  and  $\Delta\Gamma$  are tiny  
⇒ very hard to observe oscillations



Why is charm mixing so suppressed?

Combination of:

- **CKM suppression** (contribution of b-quark loop suppressed by  $V_{cb}V_{ub} \sim \lambda^2 \lambda^3$ )
- **GIM suppression** (d and s quarks have similar masses, so amplitudes nearly cancel)

The final frontier in meson mixing

# Charm mixing

Remember the ‘master’ equation derived for the  $B^0$  case (but general):

$$D^0 \text{ at } t=0: \quad \Gamma(D(t) \rightarrow f) \propto e^{-\Gamma t} \\ \times [ \cosh(\Delta\Gamma t/2) + A_{CP}^{\text{dir}} \cos(\Delta mt) + A_{\Delta\Gamma} \sinh(\Delta\Gamma t/2) + A_{CP}^{\text{mix}} \sin(\Delta mt) ]$$

$$\bar{D}^0 \text{ at } t=0: \quad \Gamma(\bar{D}(t) \rightarrow f) \propto e^{-\Gamma t} \\ \times [ \cosh(\Delta\Gamma t/2) - A_{CP}^{\text{dir}} \cos(\Delta mt) + A_{\Delta\Gamma} \sinh(\Delta\Gamma t/2) - A_{CP}^{\text{mix}} \sin(\Delta mt) ]$$

For charm, both  $\Delta m$  and  $\Delta\Gamma$  are small:

$$x = \Delta m / \Gamma < 1\%$$
$$y = \Delta\Gamma / 2\Gamma < 1\%$$

# Charm mixing

Remember the ‘master’ equation derived for the  $B^0$  case (but general):

$$D^0 \text{ at } t=0: \quad \Gamma(D(t) \rightarrow f) \propto e^{-\Gamma t} \times [ \cancel{\cosh(\Delta\Gamma t/2)} + A_{CP}^{\text{dir}} \cancel{\cos(\Delta m t)} + A_{\Delta\Gamma} \cancel{\sinh(\Delta\Gamma t/2)} + A_{CP}^{\text{mix}} \cancel{\sin(\Delta m t)} ] \\ \cancel{1 + \frac{1}{2}(2y\Gamma t)^2} \quad \cancel{1 - \frac{1}{2}(x\Gamma t)^2} \quad \cancel{y\Gamma t} \quad \cancel{x\Gamma t}$$

$$\bar{D}^0 \text{ at } t=0: \quad \Gamma(\bar{D}(t) \rightarrow f) \propto e^{-\Gamma t} \times [ \cancel{\cosh(\Delta\Gamma t/2)} - A_{CP}^{\text{dir}} \cancel{\cos(\Delta m t)} + A_{\Delta\Gamma} \cancel{\sinh(\Delta\Gamma t/2)} - A_{CP}^{\text{mix}} \cancel{\sin(\Delta m t)} ] \\ \cancel{1 + \frac{1}{2}(2y\Gamma t)^2} \quad \cancel{1 - \frac{1}{2}(x\Gamma t)^2} \quad \cancel{y\Gamma t} \quad \cancel{x\Gamma t}$$

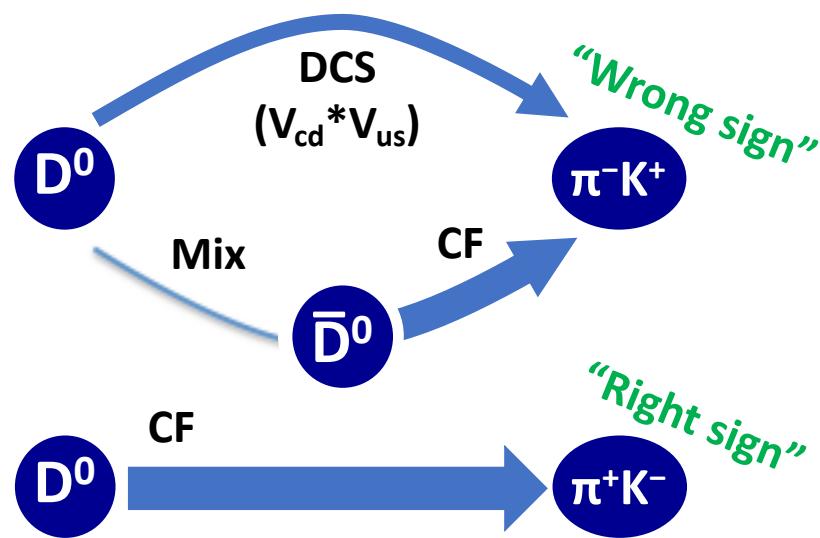
For charm, both  $\Delta m$  and  $\Delta\Gamma$  are small:

$$x = \Delta m / \Gamma \quad < 1\% \\ y = \Delta\Gamma / 2\Gamma \quad < 1\%$$

⇒ Quadratic time dependence is very good approximation

# Charm mixing: Wrong-sign K $\pi$

## Charm mixing and CPV



Pick decay with two amplitudes that can interfere (one with oscillation)

Plot ratio to non-oscillated decay  $D^0 \rightarrow K^-\pi^+$

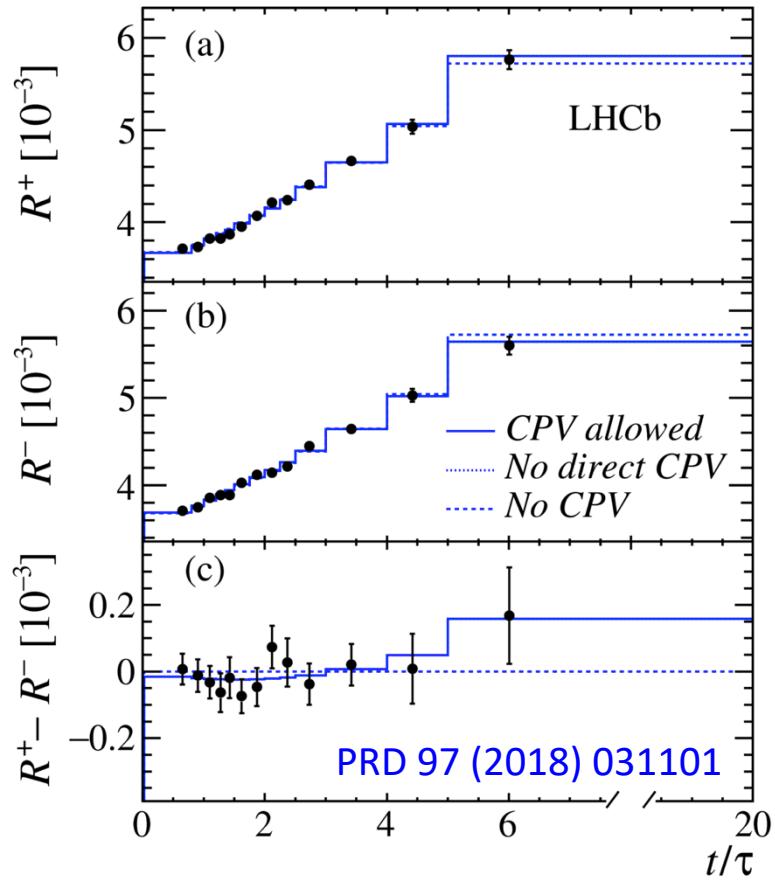
- Mixing  $\Rightarrow$  quadratic time dependence
- CP violation  $\Rightarrow$  different for  $D^0$  and  $\bar{D}^0$

$$R(t) = R_D + \sqrt{R_D} y' \left( \frac{t}{\tau} \right) + \frac{x'^{-2} + y'^{-2}}{4} \left( \frac{t}{\tau} \right)^2$$

↑                      ↑                      ↑  
DCS                  Interference term      Mix+CF

# Charm mixing: Wrong-sign $K\pi$

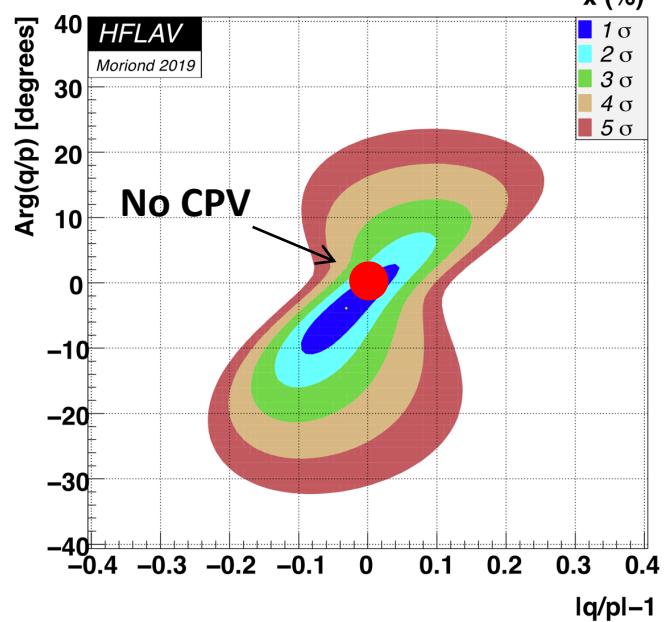
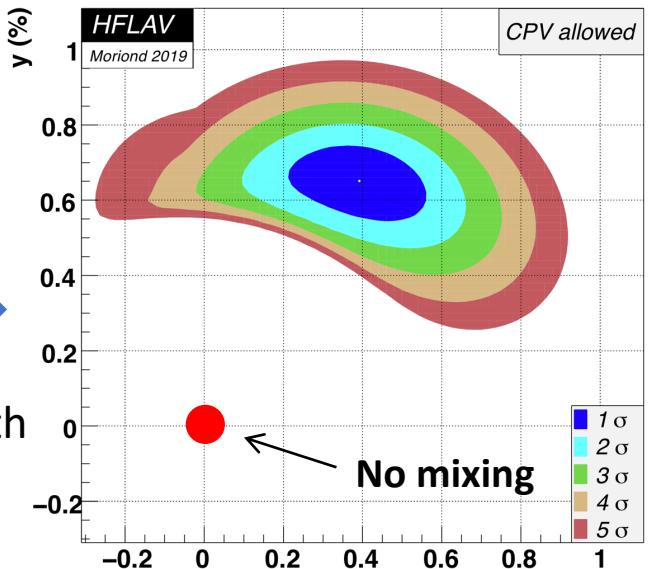
## Charm mixing and CPV



Combine  
with other  
measurements

⇒  $x \leq 0$  excluded with  
3.1 $\sigma$  significance

No sign of CPV  
in charm mixing  
or interference  
⇒ need more  
precision!



# Charm mixing: state-of-the-art

Recently LHCb published analysis of  
'Golden mode'  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$   
(see back-up slides)

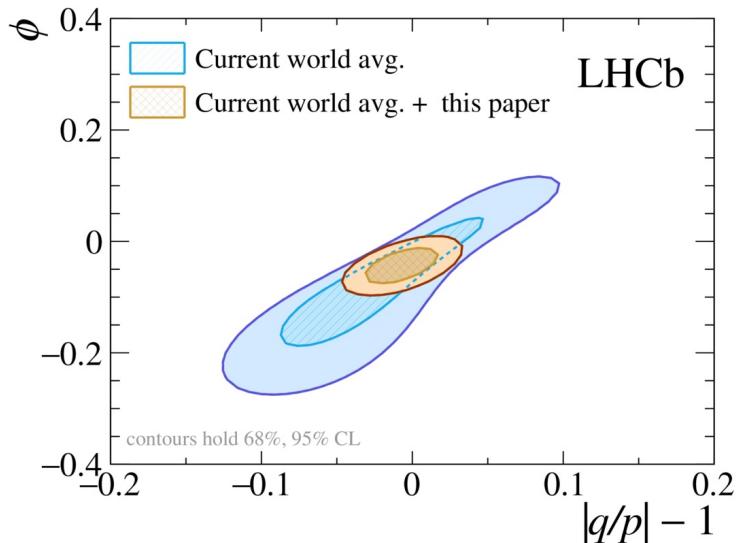
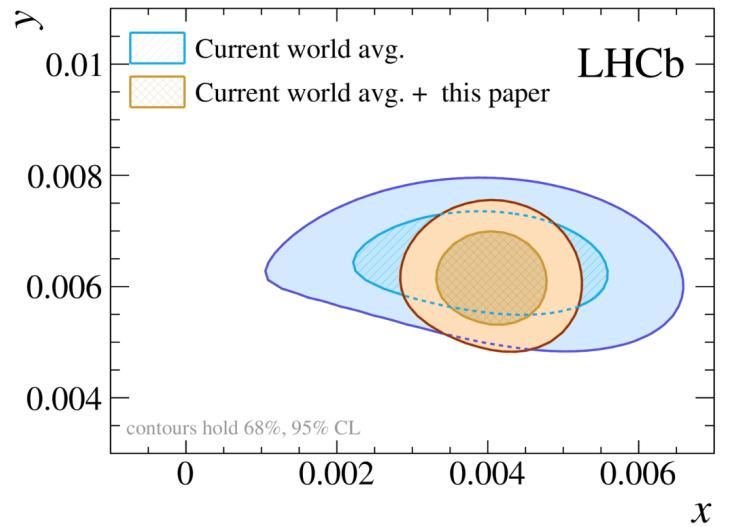
Large impact on world-averages for mixing  
and CP violation parameters

First measurement of non-zero  $x$  (and  $\Delta m$ )  
( $>7\sigma$  significance)

Oscillation period  $\sim 630$ ps  
( $D^0$  lifetime 0.4ps)

[arXiv:2106.03744](https://arxiv.org/abs/2106.03744)

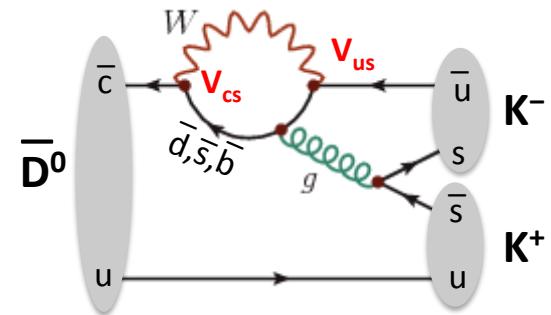
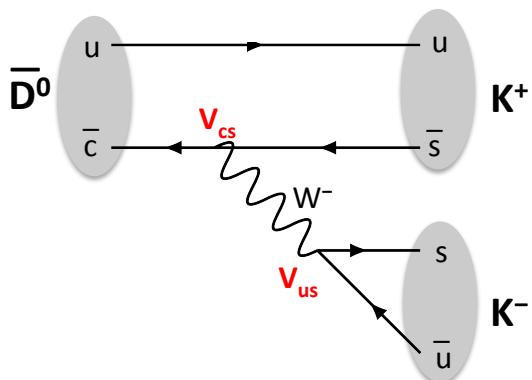
Phys. Rev. Lett. 127, 111801



# CP violation in charm decays

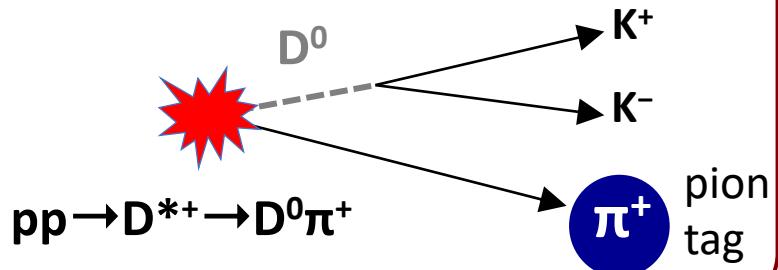
Ingredients:

(1) Two amplitudes with same final state  $\Rightarrow$  interference

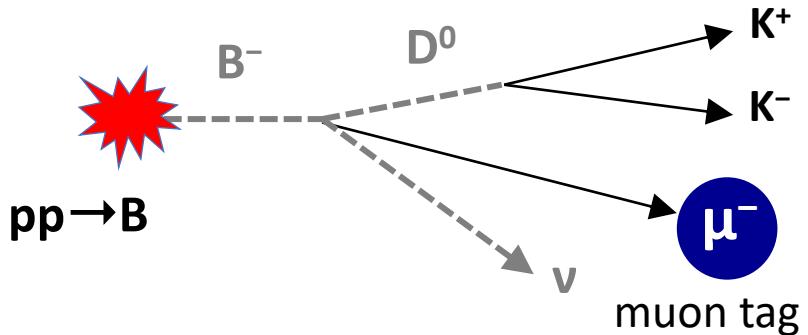


(2) Knowledge of flavour ( $D^0$  or  $\bar{D}^0$ ) at production

**$\pi$ -tagged (“prompt charm”)**



**$\mu$ -tagged (“charm from B”)**



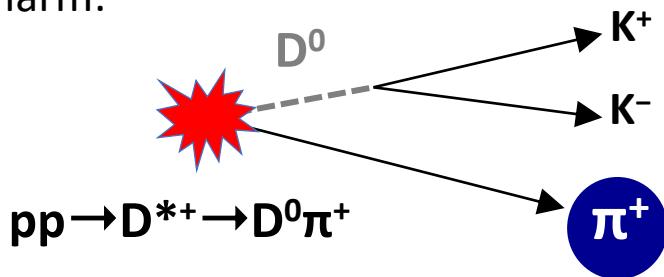
# CP violation in charm decays

Ingredients:

$$A_{CP}(D^0 \rightarrow f) = \frac{\Gamma(D^0 \rightarrow f) - \Gamma(\bar{D}^0 \rightarrow f)}{\Gamma(D^0 \rightarrow f) + \Gamma(\bar{D}^0 \rightarrow f)}$$

(3) Detailed knowledge of production and detector asymmetries

e.g. for prompt charm:



$$A_{Raw}(KK) = \mathbf{A_{CP}(KK)} + A_{Prod}(p\bar{p} \rightarrow D^*) + A_{Det}(\pi_{tag})$$

$$A_{Raw}(\pi\pi) = \mathbf{A_{CP}(\pi\pi)} + A_{Prod}(p\bar{p} \rightarrow D^*) + A_{Det}(\pi_{tag})$$

# CP violation in charm decays

Ingredients:

$$A_{CP}(D^0 \rightarrow f) = \frac{\Gamma(D^0 \rightarrow f) - \Gamma(\bar{D}^0 \rightarrow f)}{\Gamma(D^0 \rightarrow f) + \Gamma(\bar{D}^0 \rightarrow f)}$$

(3) Detailed knowledge of production and detector asymmetries

$$A_{Raw}(KK) = A_{CP}(KK) + A_{Prod}(pp \rightarrow D^*) + A_{Det}(\pi_{tag})$$

$$A_{Raw}(\pi\pi) = A_{CP}(\pi\pi) + A_{Prod}(pp \rightarrow D^*) + A_{Det}(\pi_{tag})$$

*OR* Clever method to eliminate them...

$$\Rightarrow A_{Raw}(KK) - A_{Raw}(\pi\pi) = A_{CP}(KK) - A_{CP}(\pi\pi) \equiv \Delta A_{CP}$$

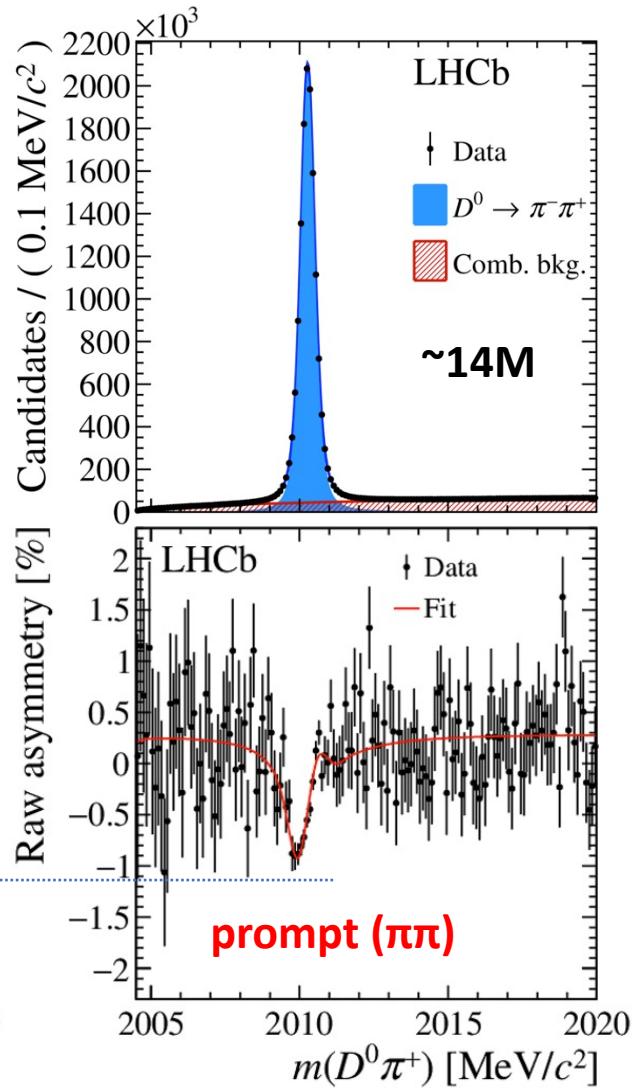
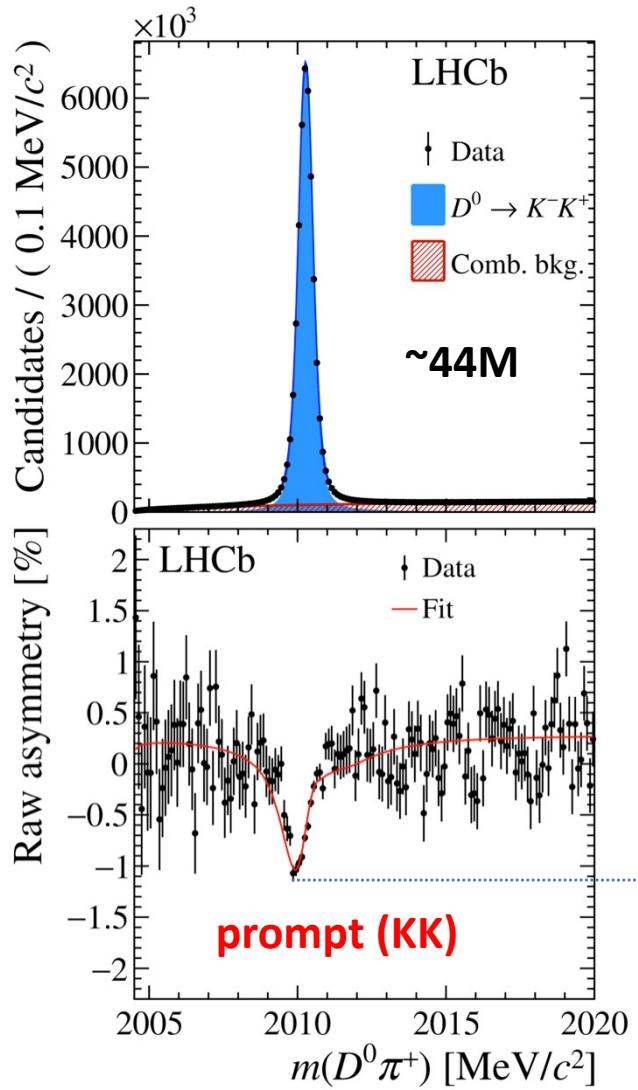
# CP violation in charm decays

Huge (>10M) signal samples,  
high purity

$$\Delta A_{CP} = (-0.154 \pm 0.029)\%$$

Inconsistent with CP  
symmetry at  $5.3\sigma$   
significance  $\Rightarrow$  discovery!

Q: Is this from SM?  
A: Not yet clear!



PRL 122 (2019) 211803

Part IIc: CKM angle  $\gamma$

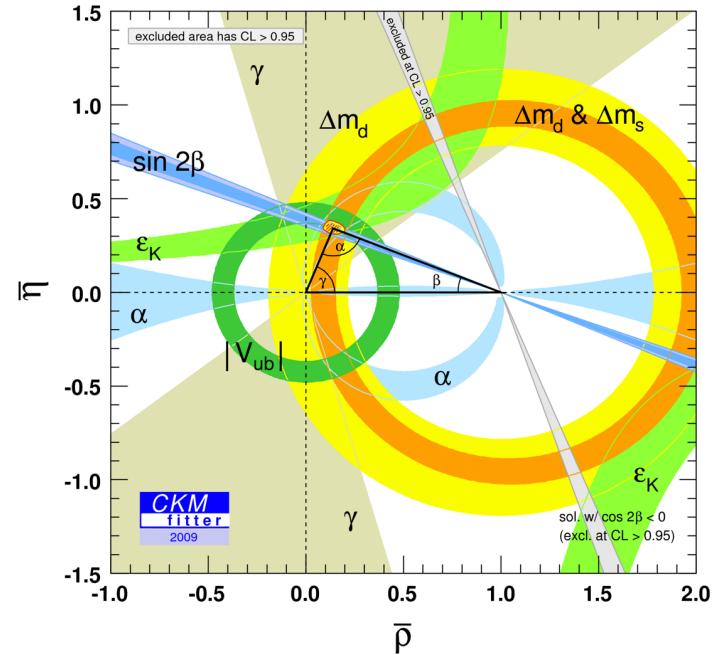
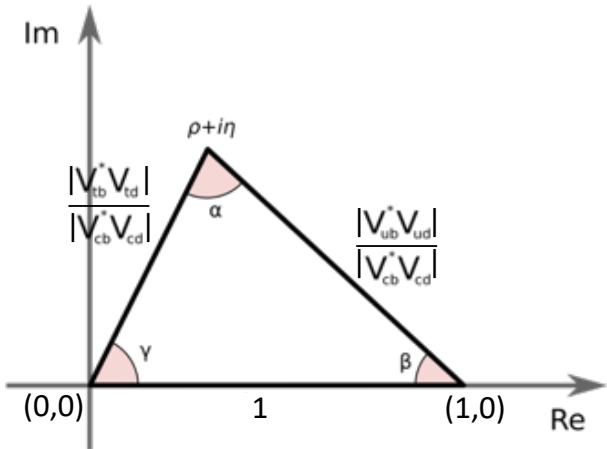
# CKM angle $\gamma$

At the start of the LHC era:

- The least experimentally constrained angle
- The most precisely predicted angle (negligible theory errors)

$$\sigma_{\text{theory}}(\gamma) \approx 10^{-7} \text{ rad} \quad \text{https://arxiv.org/abs/1308.5663}$$

**Reminder:**



$$\begin{aligned}\beta &= \phi_1 = \arg \left( -\frac{V_{cd} V_{cb}^*}{V_{td} V_{tb}^*} \right) \\ \alpha &= \phi_2 = \arg \left( -\frac{V_{td} V_{tb}^*}{V_{ud} V_{ub}^*} \right) \\ \gamma &= \phi_3 = \arg \left( -\frac{V_{ud} V_{ub}^*}{V_{cd} V_{cb}^*} \right)\end{aligned}$$

# CKM angle $\gamma$

$$\begin{aligned}\beta &= \phi_1 = \arg \left( -\frac{V_{cd} V_{cb}^*}{V_{td} V_{tb}^*} \right) \\ \alpha &= \phi_2 = \arg \left( -\frac{V_{td} V_{tb}^*}{V_{ud} V_{ub}^*} \right) \\ \gamma &= \phi_3 = \arg \left( -\frac{V_{ud} V_{ub}^*}{V_{cd} V_{cb}^*} \right)\end{aligned}$$

b $\rightarrow$ cW transitions, with B $^0$  mixing  
(e.g. B $^0 \rightarrow J/\psi K_S^0$ )

b $\rightarrow$ uW transitions, with B $^0$  mixing  
(e.g. B $^0 \rightarrow \pi^+ \pi^-$ )

No top loop needed! – can extract in tree-level decays (b $\rightarrow$ cW vs b $\rightarrow$ uW)  
⇒ Very clean SM test

Measure  $\gamma$  in tree-level processes ⇒ precise SM benchmark

Measure  $\gamma$  in loop processes ⇒ sensitive to NP

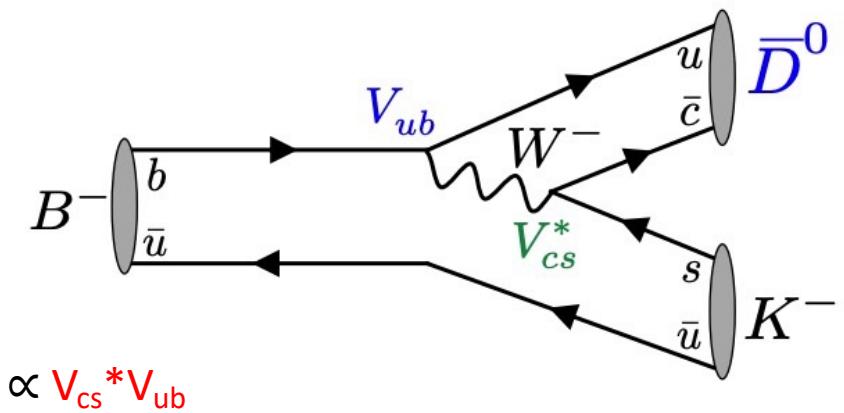
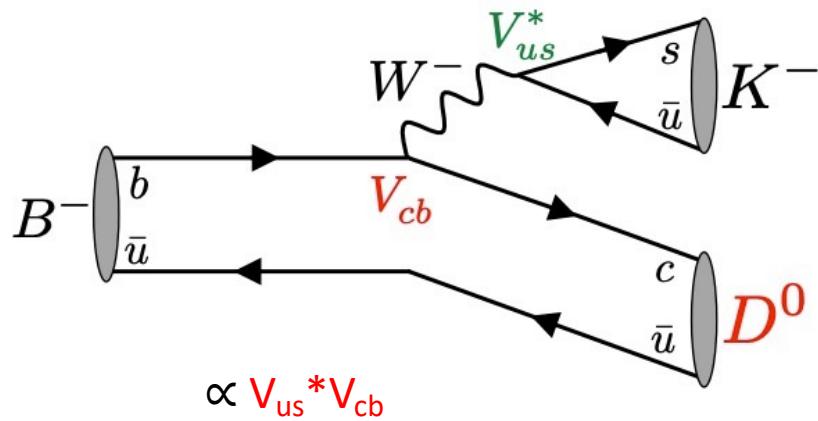
Precise  $\gamma$  studies were a major motivation for building the LHCb experiment

# Measuring $\gamma$

Require interference between  $b \rightarrow cW$  and  $b \rightarrow uW$

$$\gamma = \phi_3 = \arg \left( -\frac{V_{ud} V_{ub}^*}{V_{cd} V_{cb}^*} \right)$$

Textbook case is  $B^\pm \rightarrow \bar{D}^0 K^\pm$



Transitions have different final states ( $D^0$  vs  $\bar{D}^0$ )

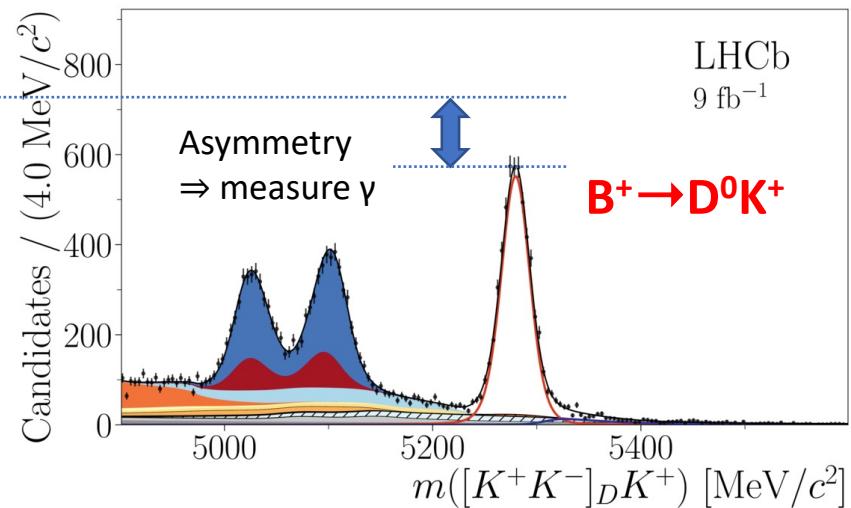
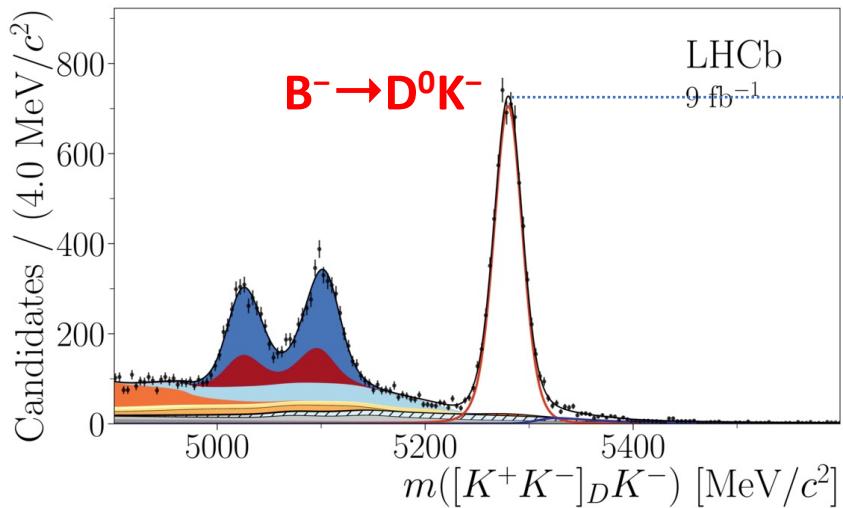
Interference if  $D^0$  and  $\bar{D}^0$  decay to **same final state  $f$**

Many different methods and decay channels – best results from combination

# Measuring $\gamma$

Sensitivity to  $\gamma$  without time-dependent analysis – see asymmetries in yields!

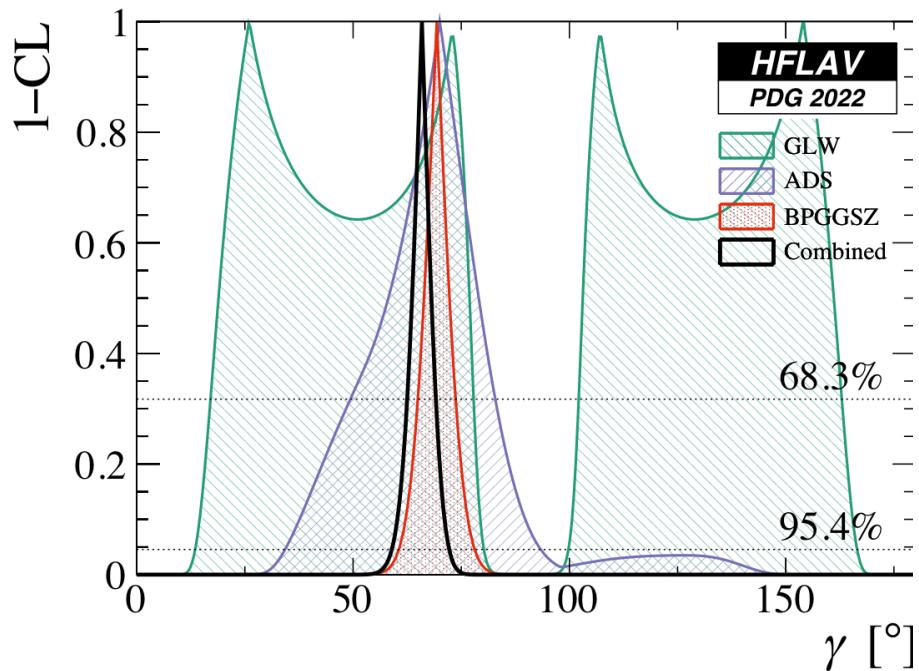
⇒ Can convert measured yields into precise  $\gamma$  measurement



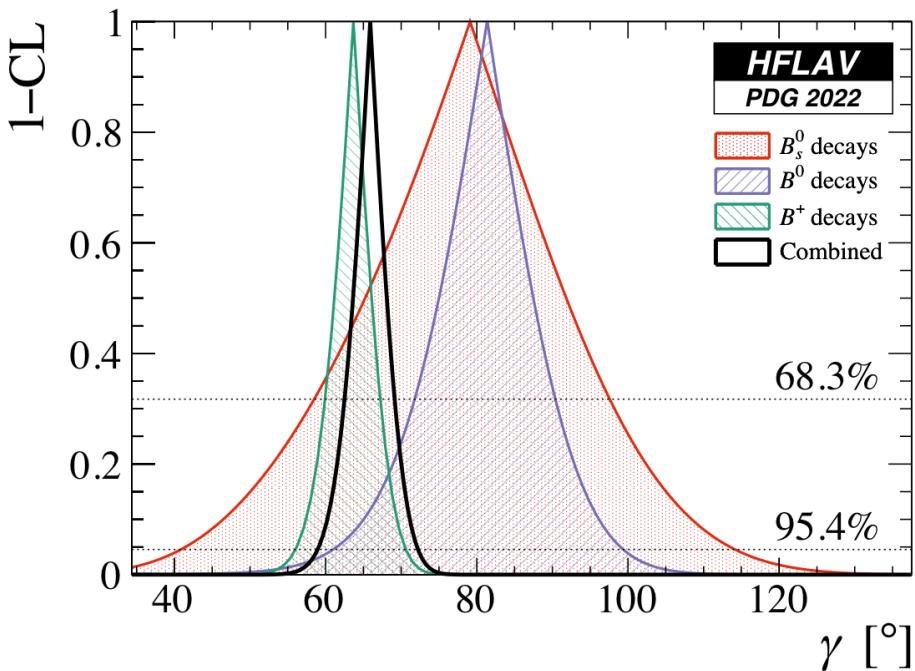
[https://doi.org/10.1007/JHEP04\(2021\)081](https://doi.org/10.1007/JHEP04(2021)081) (2021)

# Combining all $\gamma$ measurements

Split by method:



Split by B species:



Combining all measurements:

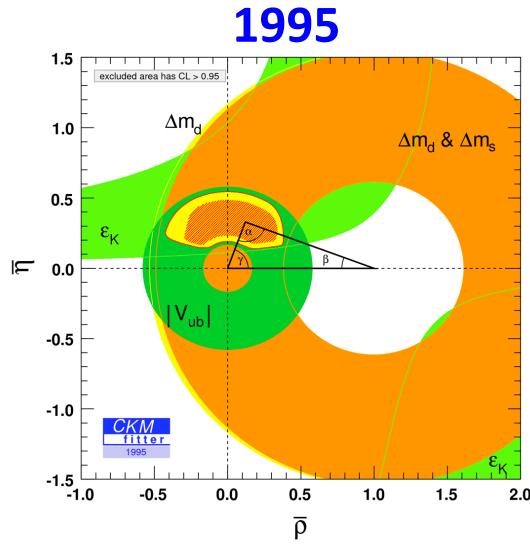
$$\gamma = (65.9^{+3.3}_{-3.5})^\circ$$

Indirect constraints:

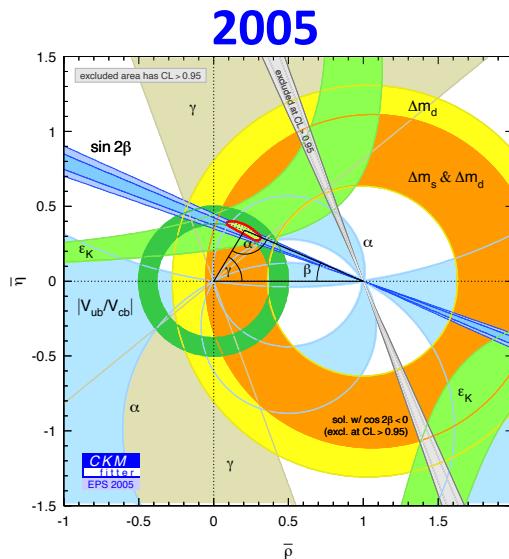
$$\gamma = (65.5^{+1.1}_{-2.7})^\circ$$

i.e. all other unitarity triangle measurements

# Unitarity triangle fits through time

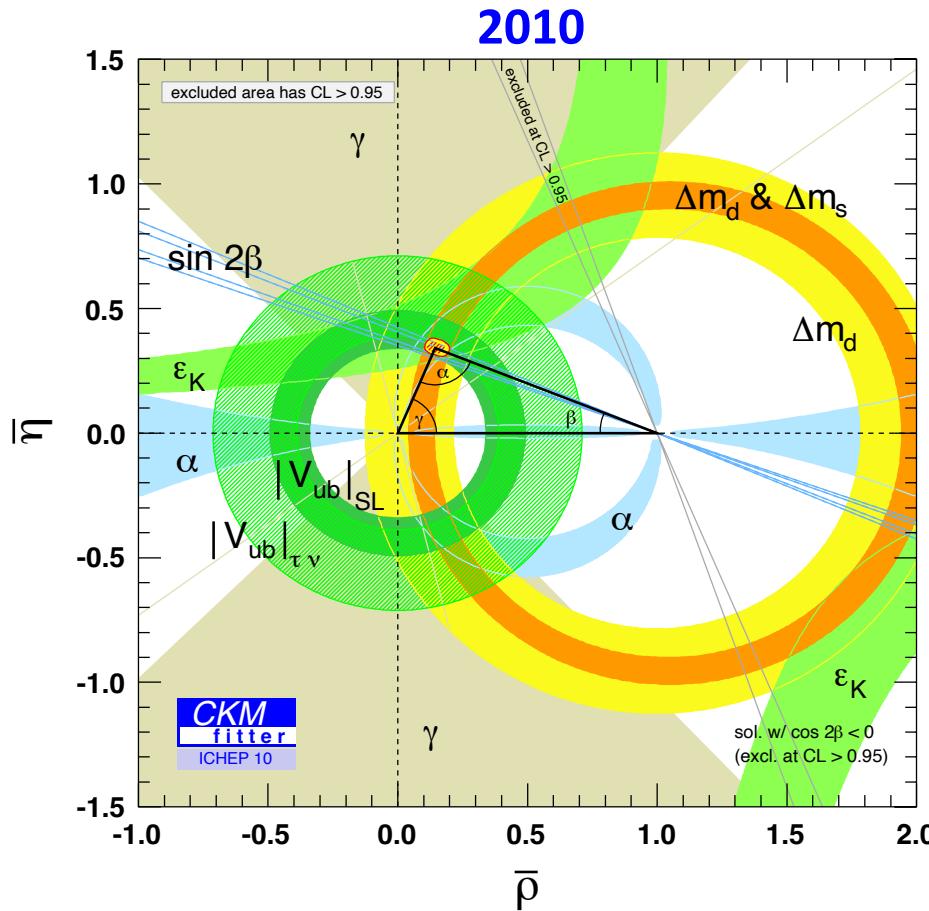
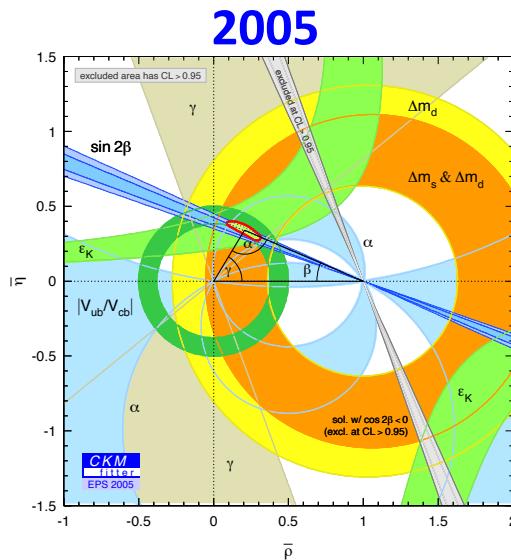
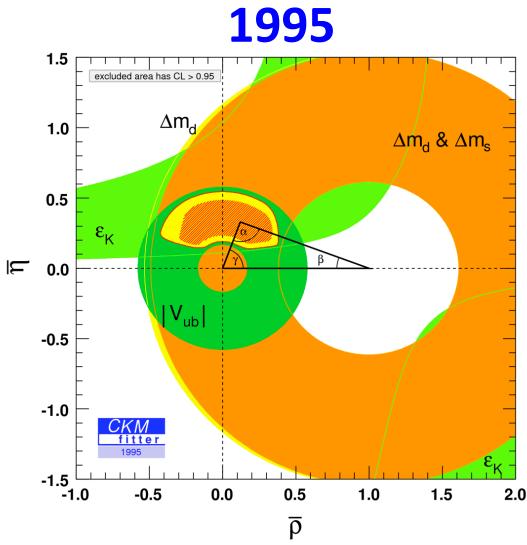


Before B-factories: minimal constraints



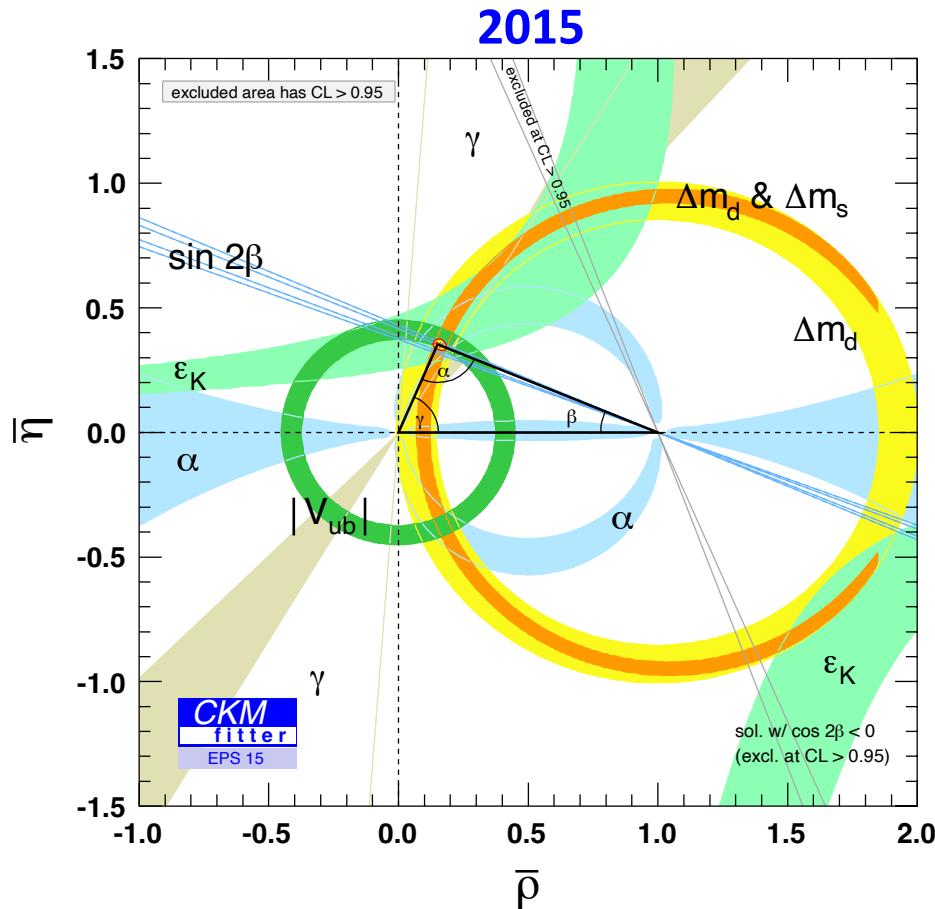
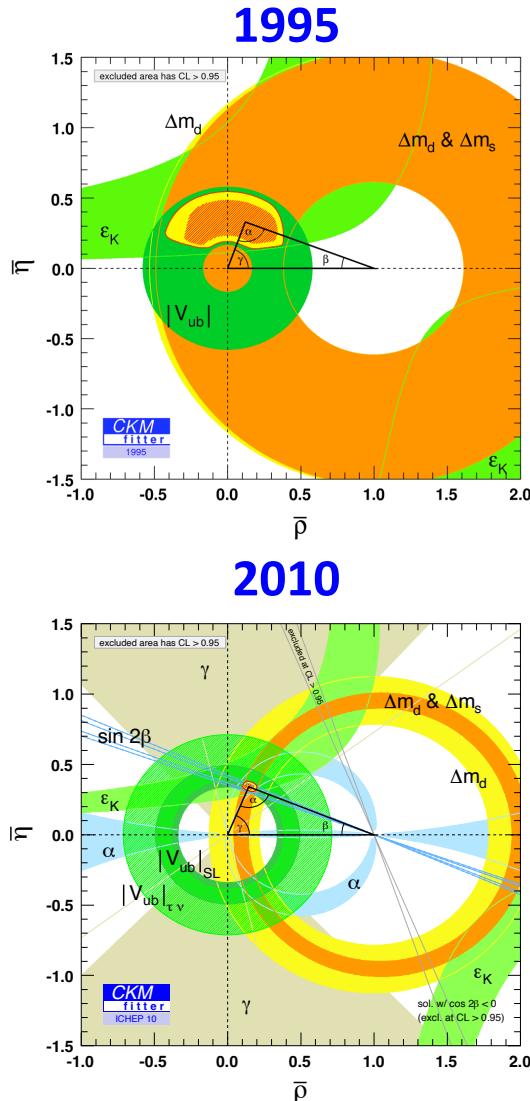
+ B-factory measurements of CKM angles

# Unitarity triangle fits through time



+  $B_s^0$  mixing from Tevatron  
+ more B-factory inputs on angles

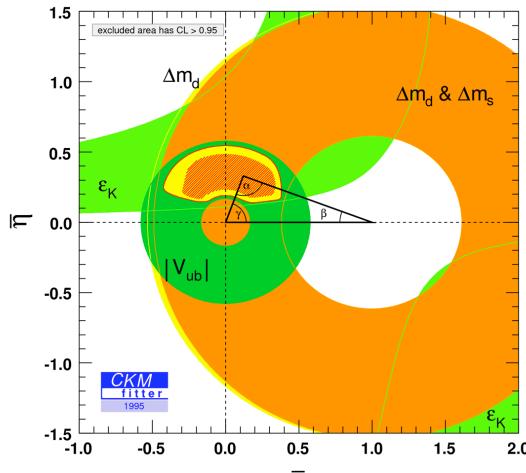
# Unitarity triangle fits through time



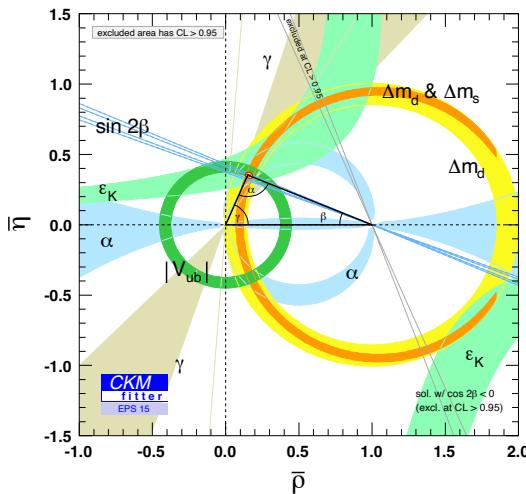
+ LHCb starts to deliver

# Unitarity triangle fits through time

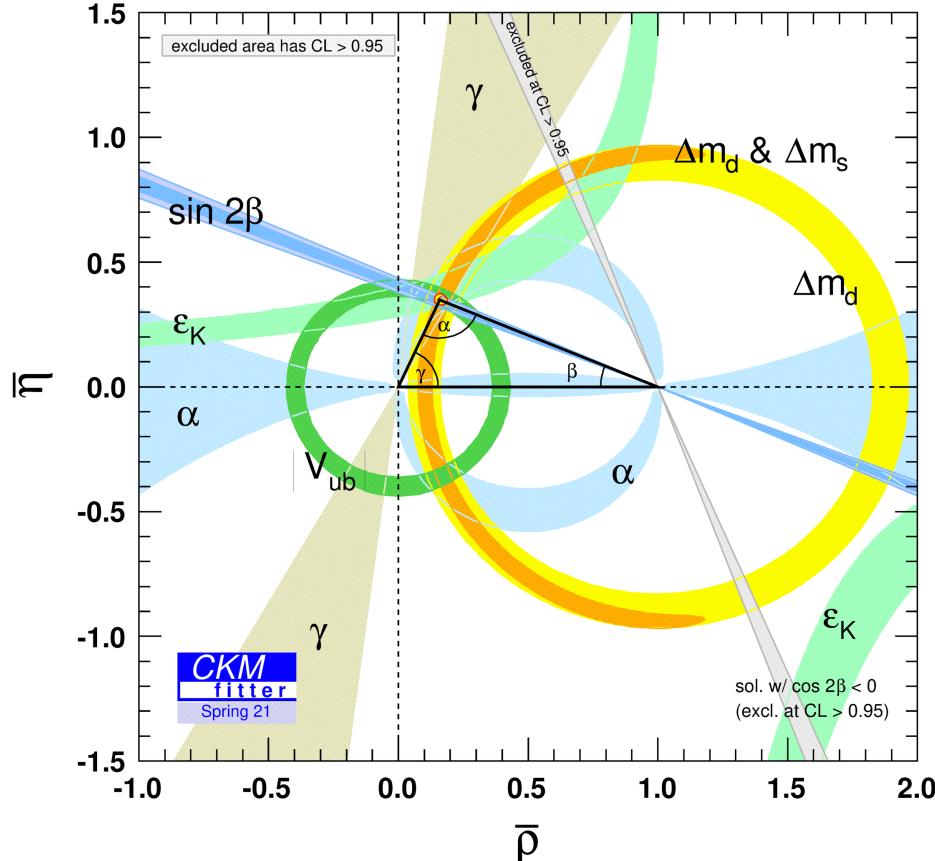
1995



2015

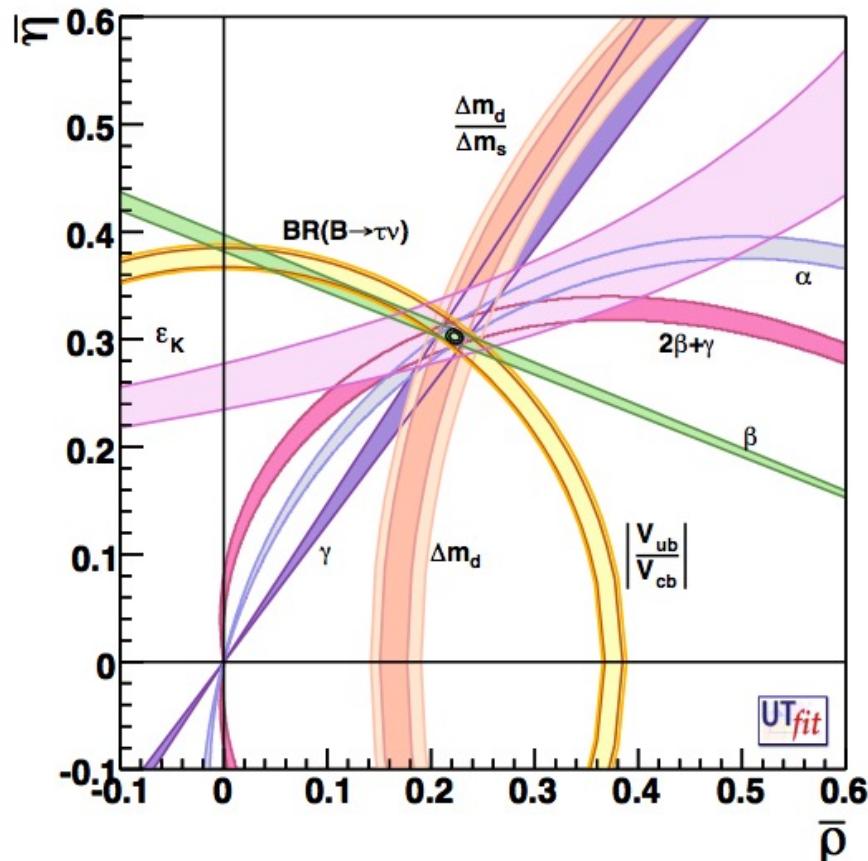


2021



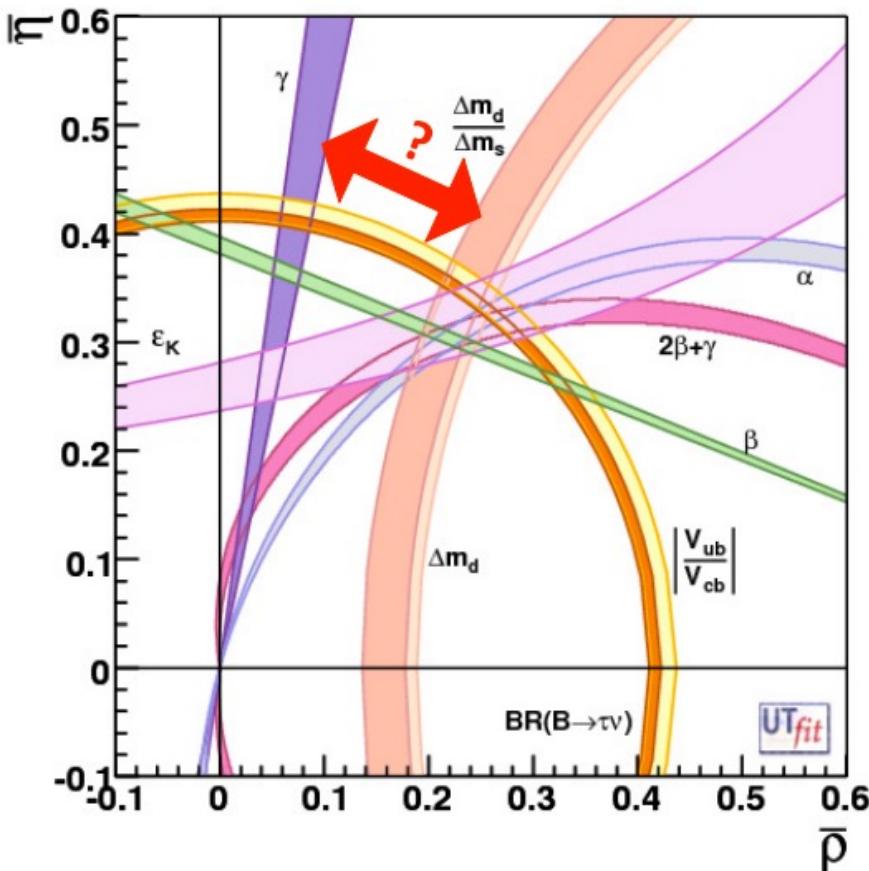
# Future of CKM?

The “nightmare”



<https://arxiv.org/abs/0710.3799>

The “dream”



# Part III: Rare decays and lepton universality

# Rare b decays

Rare decays helped to shape the SM. Can they show us the way beyond it?

Studies of rare b decays are a key part of LHC physics programme

- Both overall rates, and properties (e.g. angular distributions) of rare processes can be influenced by New Physics

✓ LHC is a b-factory! Huge numbers of b quarks produced.  
⇒ Great place to look for rare decays

⚠ LHC is a busy environment  
⇒ Essential to understand (and reject / account for) backgrounds

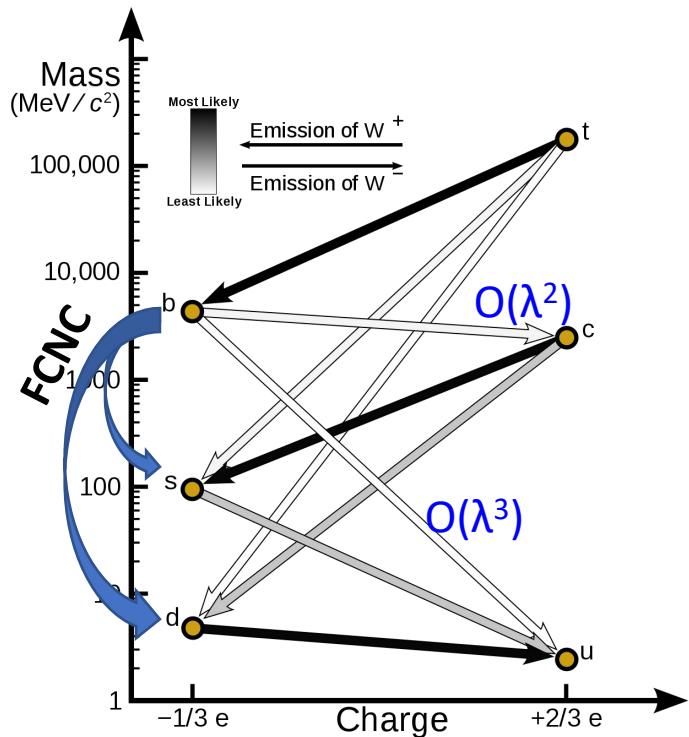
**Muon modes are clean and easy to trigger: low(ish)-hanging fruit**

# Rare b decays

In CKM picture, all b decays are ‘rare’ since they change generation

Rare in this context really means not  $b \rightarrow c$

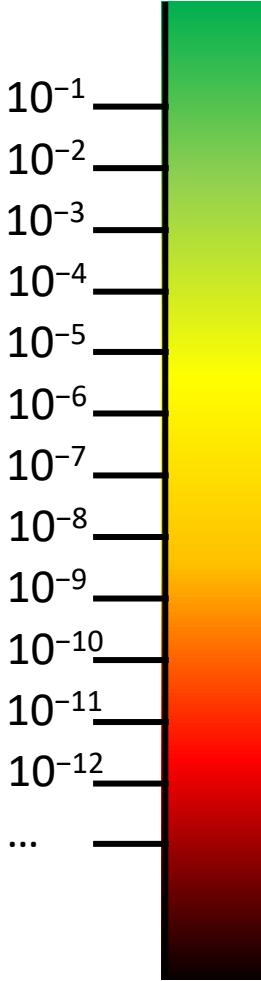
**Flavour changing neutral currents** (FCNC) are even more suppressed, since direct transitions are forbidden ( $b \rightarrow s$ ,  $b \rightarrow d$ )  
⇒ no tree-level SM diagrams



(From Lecture 2)

# Rare b decays

Branching  
fraction



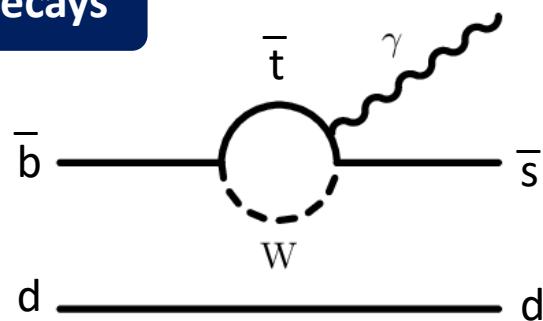
$B^0 \rightarrow D^- \mu^+ \nu$  ( $2.3 \times 10^{-2}$ )

$B^0 \rightarrow J/\psi K^0$  ( $8.9 \times 10^{-4}$ )

$B^0 \rightarrow K^{*0} \gamma$  ( $4.2 \times 10^{-5}$ )

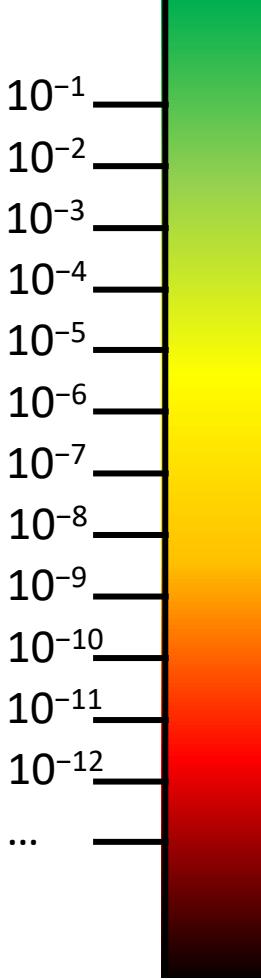
$B^0 \rightarrow \omega^* \gamma$  ( $4.4 \times 10^{-7}$ )

**Radiative decays**



# Rare b decays

Branching  
fraction



$B^0 \rightarrow D^- \mu^+ \nu$  ( $2.3 \times 10^{-2}$ )

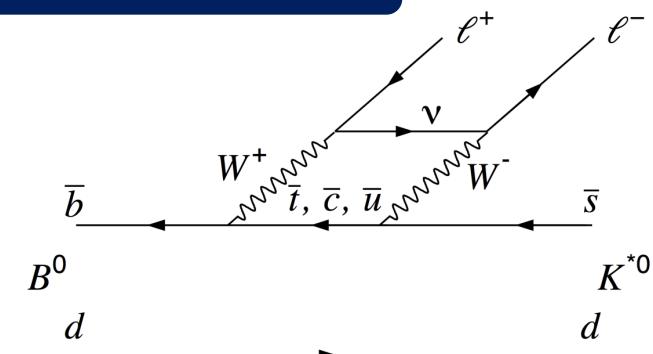
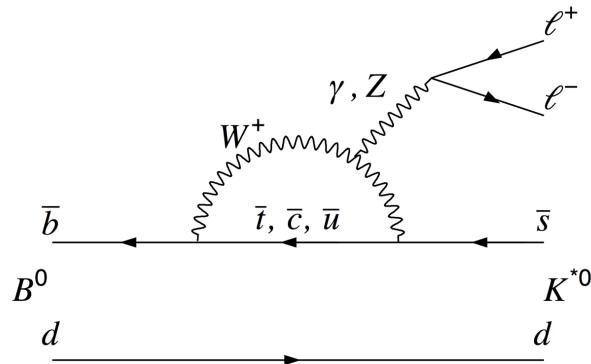
$B^0 \rightarrow J/\psi K^0$  ( $8.9 \times 10^{-4}$ )

$B^0 \rightarrow K^{*0} \gamma$  ( $4.2 \times 10^{-5}$ )

$B^0 \rightarrow \omega^* \gamma$  ( $4.4 \times 10^{-7}$ )

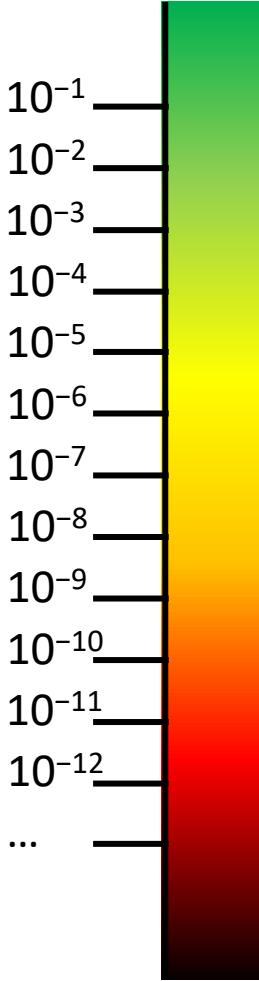
$B^0 \rightarrow K^{*0} \mu^+ \mu^-$  ( $9.9 \times 10^{-7}$ )

Electroweak penguins



# Rare b decays

Branching  
fraction



$B^0 \rightarrow D^- \mu^+ \nu$  ( $2.3 \times 10^{-2}$ )

$B^0 \rightarrow J/\psi K^0$  ( $8.9 \times 10^{-4}$ )

$B^0 \rightarrow K^{0*} \gamma$  ( $4.2 \times 10^{-5}$ )

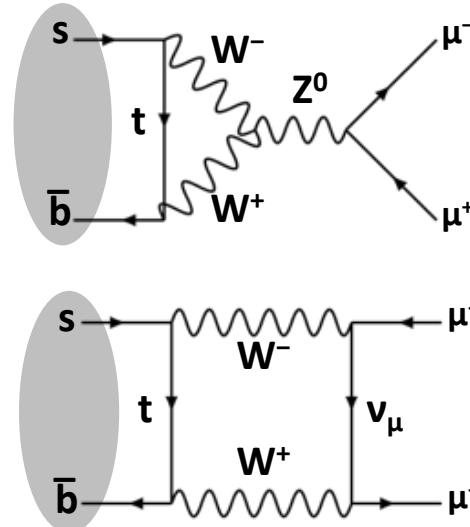
$B^0 \rightarrow \omega^* \gamma$  ( $4.4 \times 10^{-7}$ )

$B^0 \rightarrow K^{0*} \mu^+ \mu^-$  ( $9.9 \times 10^{-7}$ )

$B_s^0 \rightarrow \mu^+ \mu^-$  ( $3.5 \times 10^{-9}$ )

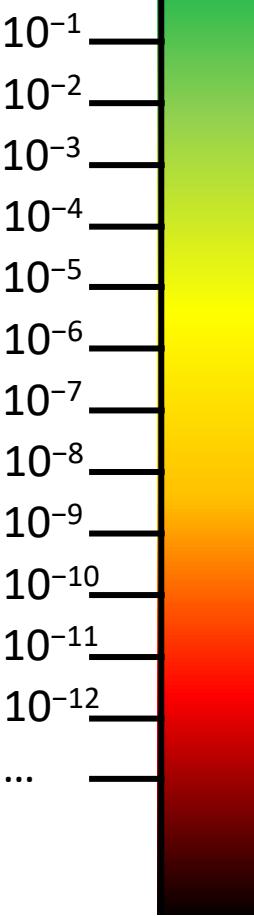
$B^0 \rightarrow \mu^+ \mu^-$  ( $< 2 \times 10^{-10}$ )

Box diagrams



# Rare b decays

Branching  
fraction



$B^0 \rightarrow D^- \mu^+ \nu$  ( $2.3 \times 10^{-2}$ )

$B^0 \rightarrow J/\psi K^0$  ( $8.9 \times 10^{-4}$ )

$B^0 \rightarrow K^{0*} \gamma$  ( $4.2 \times 10^{-5}$ )

$B^0 \rightarrow \omega^* \gamma$  ( $4.4 \times 10^{-7}$ )

$B^0 \rightarrow K^{0*} \mu^+ \mu^-$  ( $9.9 \times 10^{-7}$ )

$B_s^0 \rightarrow \mu^+ \mu^-$  ( $3.5 \times 10^{-9}$ )

$B^0 \rightarrow \mu^+ \mu^-$  ( $< 2 \times 10^{-10}$ )

Lepton number violating

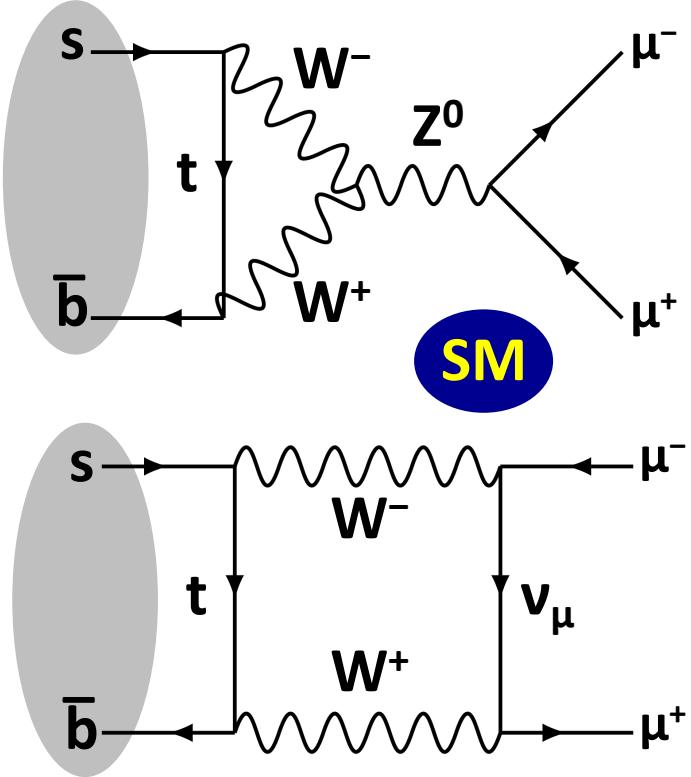
Lepton family violating

Baryon number violating ...

# $B_s^0 \rightarrow \mu^+ \mu^-$

Similar to  $\mu \rightarrow e\gamma$  from lecture 1: highly suppressed in SM with precise prediction.

$$\text{Br}(B_s^0 \rightarrow \mu^+ \mu^-) = 3.3 \times 10^{-9}$$



# $B_s^0 \rightarrow \mu^+ \mu^-$

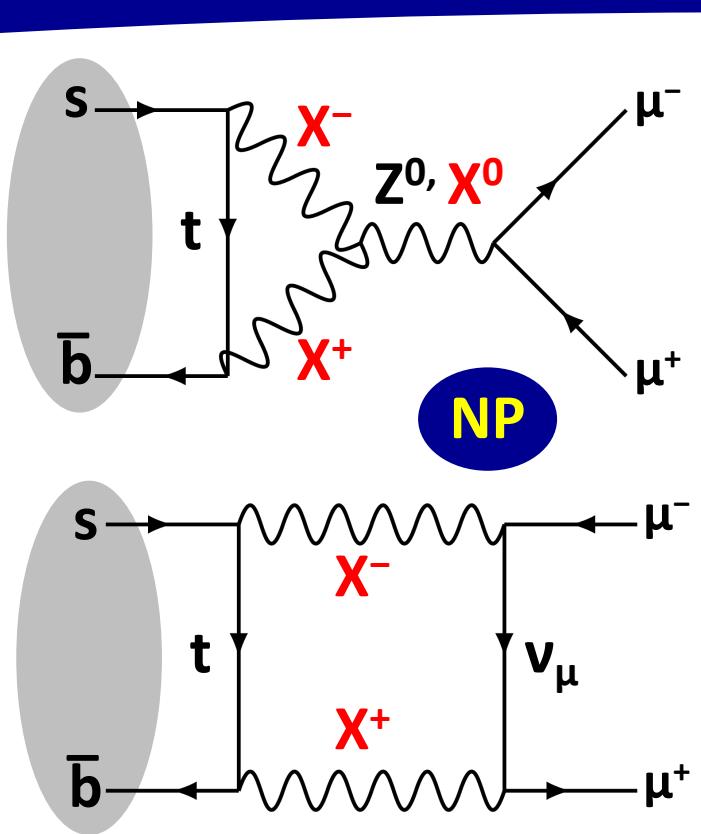
Similar to  $\mu \rightarrow e\gamma$  from lecture 1: highly suppressed in SM with precise prediction.

$$\text{Br}(B_s^0 \rightarrow \mu^+ \mu^-) = 3.3 \times 10^{-9}$$

Almost all NP theories predict enhancement (or suppression)

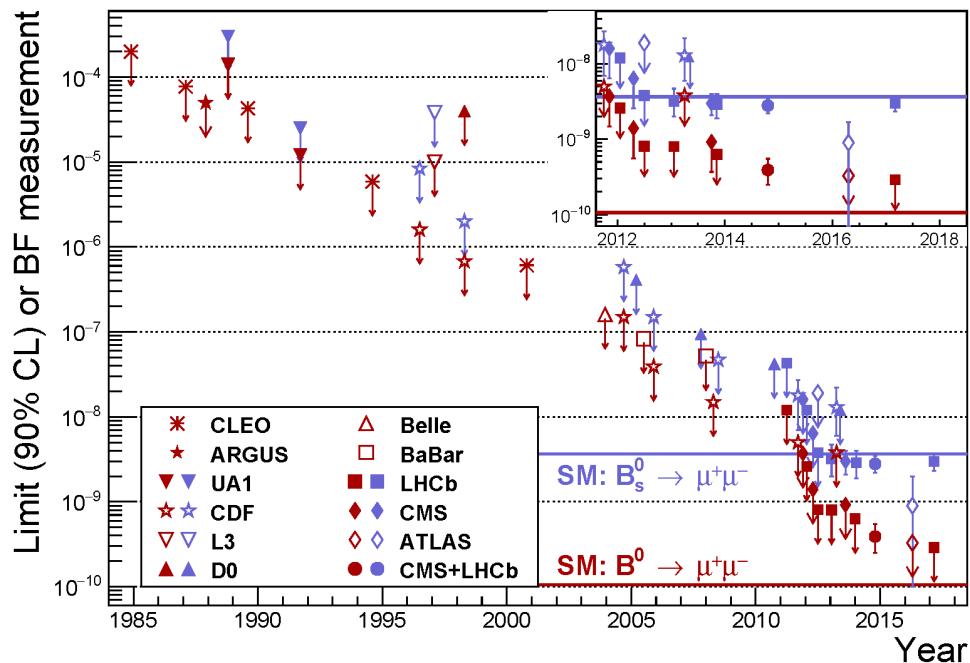
Especially strong dependence on SUSY parameter, e.g.

$$\text{Br}(B_s^0 \rightarrow \mu^+ \mu^-) \propto \tan^6 \beta$$



# $B_s^0 \rightarrow \mu^+\mu^-$

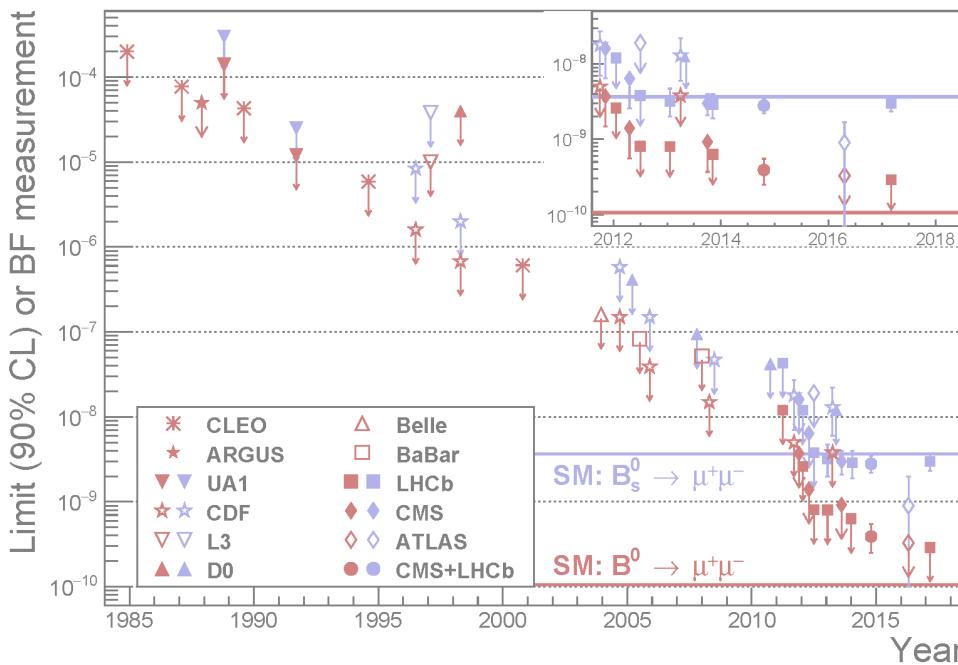
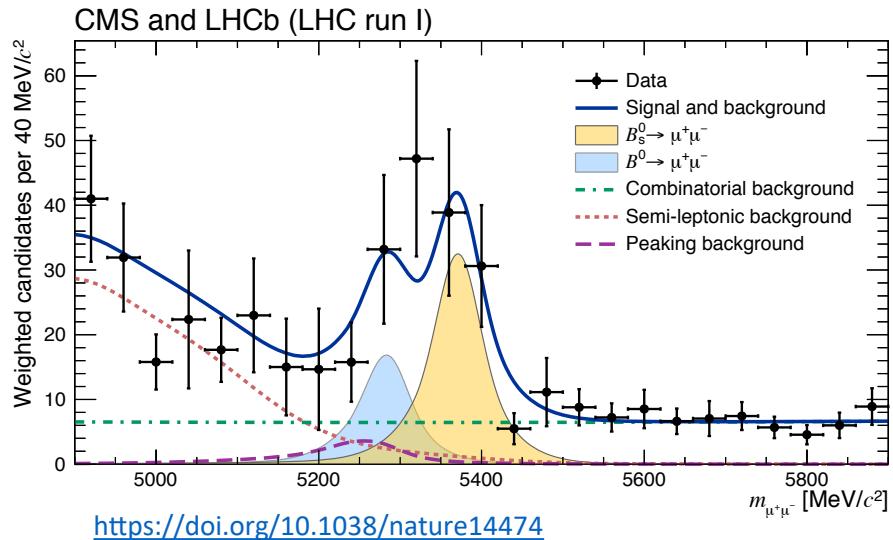
For a long time – improving limits  
Then in 2013 something changed...



# $B_s^0 \rightarrow \mu^+ \mu^-$

For a long time – improving limits  
Then in 2013 something changed...

By combining results from CMS and LHCb, reached “ $5\sigma$ ” standard for claiming observation



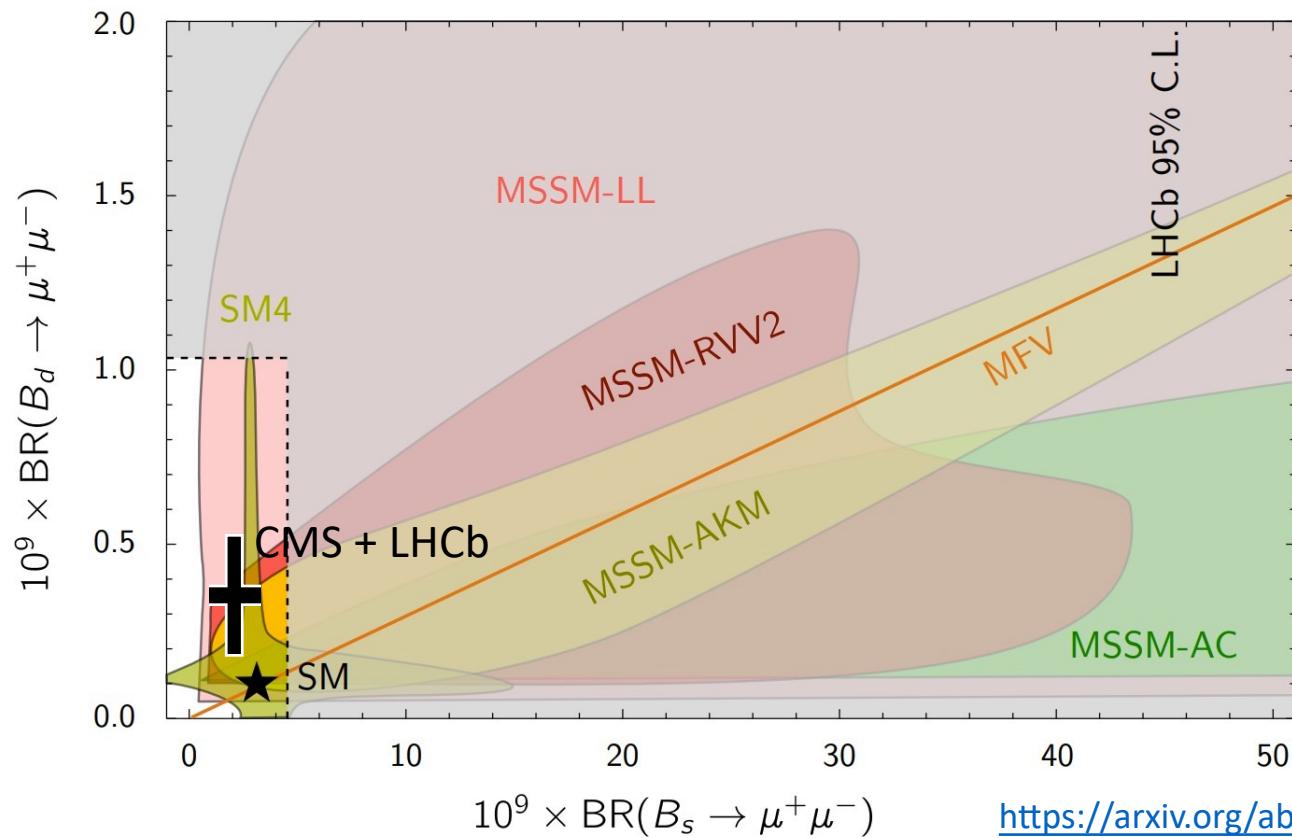
# $B_s^0 \rightarrow \mu^+ \mu^-$

Measurements in agreement with SM  
(some tension for  $B^0$  rate)

Killed a lot of SUSY parameter space

$$\text{Br}(B_s^0 \rightarrow \mu^+ \mu^-) = (2.8 \pm 0.7) \times 10^{-9}$$

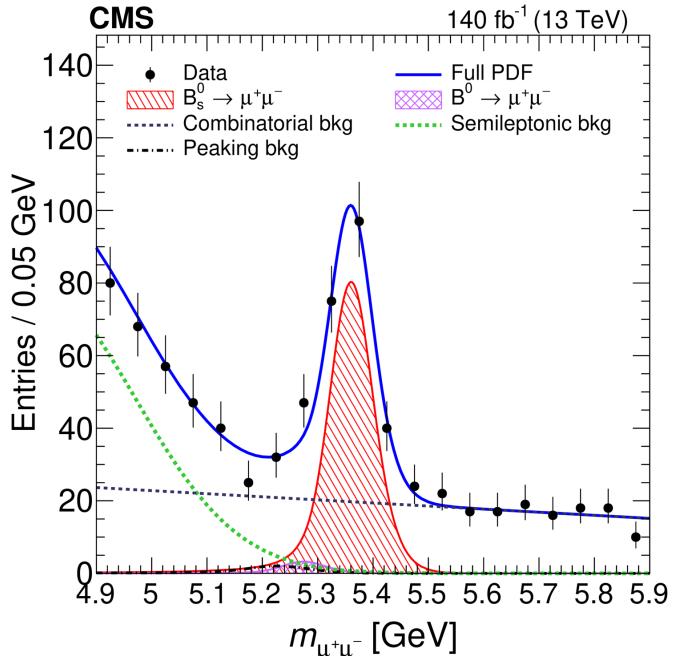
$$\text{Br}(B_d^0 \rightarrow \mu^+ \mu^-) = (3.9 \pm 1.5) \times 10^{-10}$$



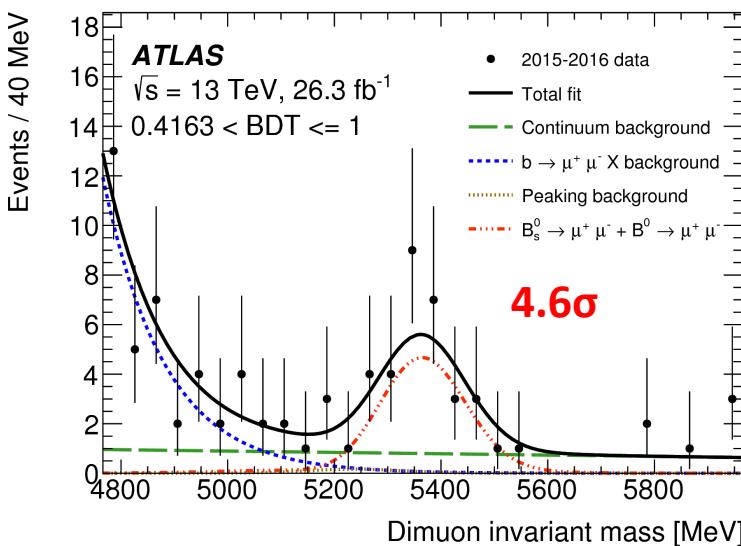
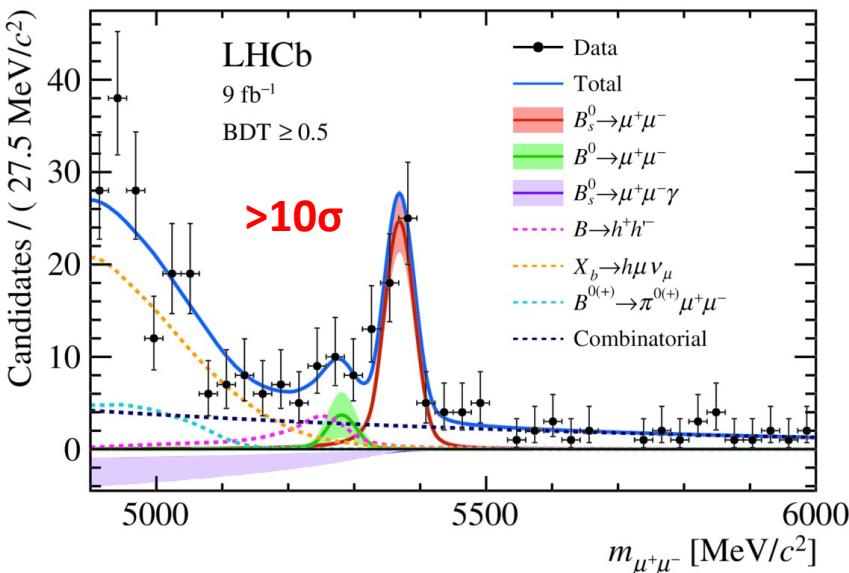
<https://arxiv.org/abs/1205.6094>

# $B_s^0 \rightarrow \mu^+\mu^-$ latest

High-significance observations from LHCb, CMS (full Run 1+2), & Atlas (partial Run 1+2)



<https://doi.org/10.1103/PhysRevLett.128.041801>  
<https://doi.org/10.1016/j.physletb.2023.137955>  
[https://doi.org/10.1007/JHEP04\(2019\)098](https://doi.org/10.1007/JHEP04(2019)098)

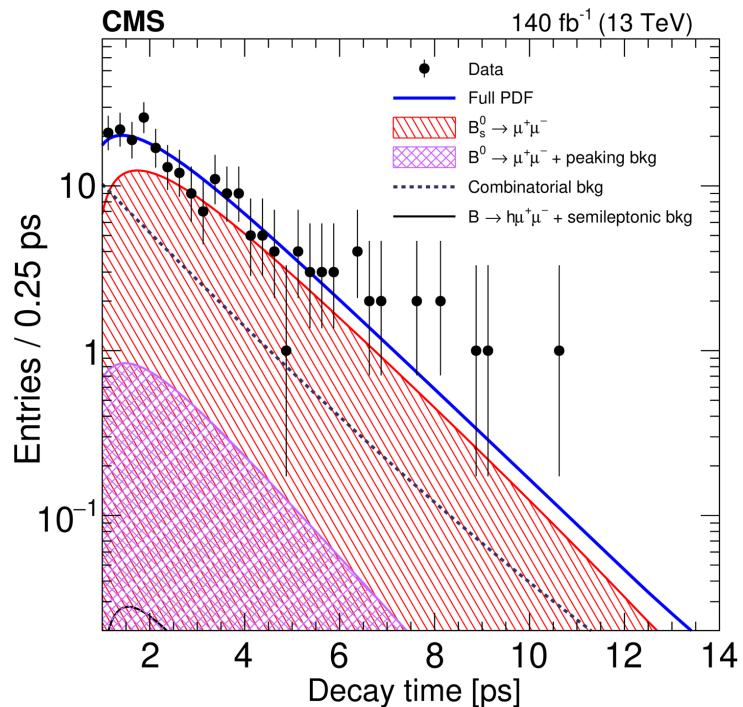
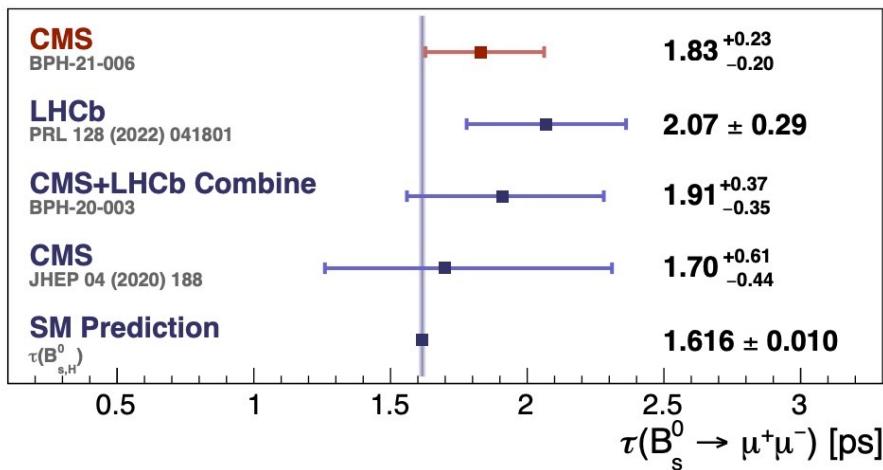


# $B_s^0 \rightarrow \mu^+\mu^-$ latest

With larger signal yields, can also measure  $B_s^0 \rightarrow \mu\mu$  lifetime

In SM, only heavy ( $H$ ) mass eigenstate can decay to  $\mu\mu \Rightarrow$  another probe of new physics

$$\begin{aligned} (\text{SM}) \quad \tau_H &= 1.620 \pm 0.007 \text{ ps} \\ &\tau_L = 1.423 \pm 0.005 \text{ ps} \end{aligned}$$



At current precision, cannot rule-out either  $\tau_H$  or  $\tau_L$  case:

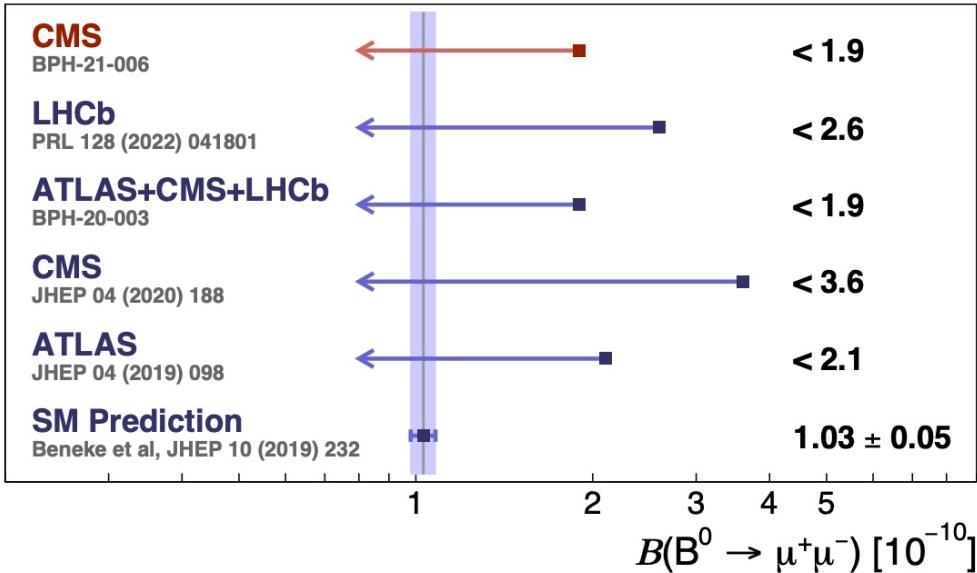
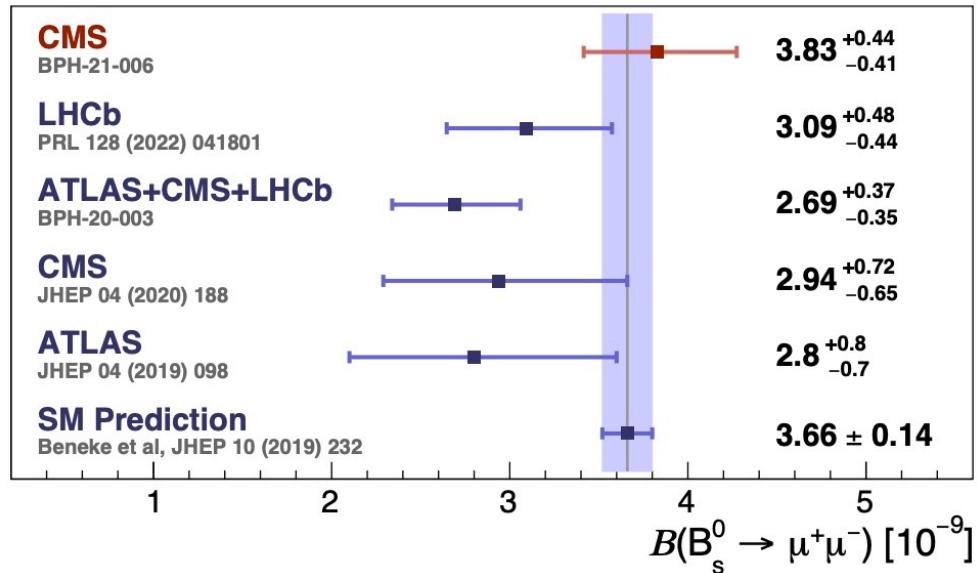
$\Rightarrow$  Need Run 3 data to fully test this

# $B_s^0 \rightarrow \mu^+\mu^-$ latest

Discovering  $B_s^0 \rightarrow \mu\mu$  is one of the most significant achievements of the LHC experiments

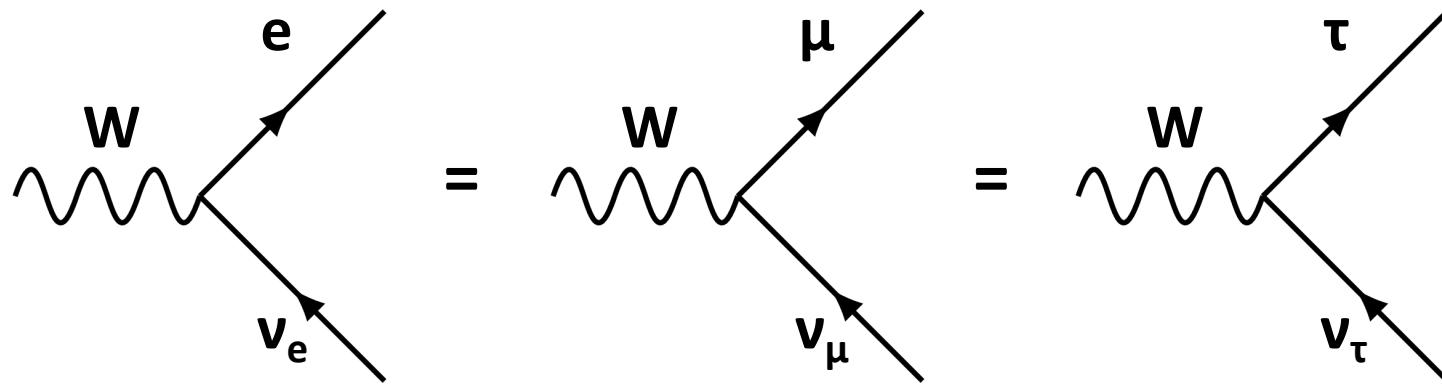
Provides very tight constraints on possible new physics models

Next target is  $B^0 \rightarrow \mu\mu$   
⇒ stay tuned in Run 3!



# Lepton universality

Lepton universality: *weak interaction acts equally regardless of lepton flavour*



Pillar of standard model – any deviation can **only** be caused by new physics

Theoretically clean...

...Experimentally challenging...

# Lepton universality: “ $R_{K^{(*)}}$ ”

Rare decays may be sensitive to new physics which doesn't respect lepton universality.

Branching ratios for FCNC  $b \rightarrow s\mu^+\mu^-$  are consistently lower than expected in SM

- Could be due to hard-to-calculate QCD effects
- If so, should also see low BRs for electron modes

Measure ratio of decay rates of muons  
and electron channels in  $B^+ \rightarrow K^+ l^+ l^-$

$$R_K = \frac{\int_{q_{\min}^2}^{q_{\max}^2} \frac{d\Gamma[B^+ \rightarrow K^+ \mu^+ \mu^-]}{dq^2} dq^2}{\int_{q_{\min}^2}^{q_{\max}^2} \frac{d\Gamma[B^+ \rightarrow K^+ e^+ e^-]}{dq^2} dq^2}$$

$\approx 1$  in SM  
(very reliable prediction)

# Lepton universality: “ $R_{K^{(*)}}$ ”

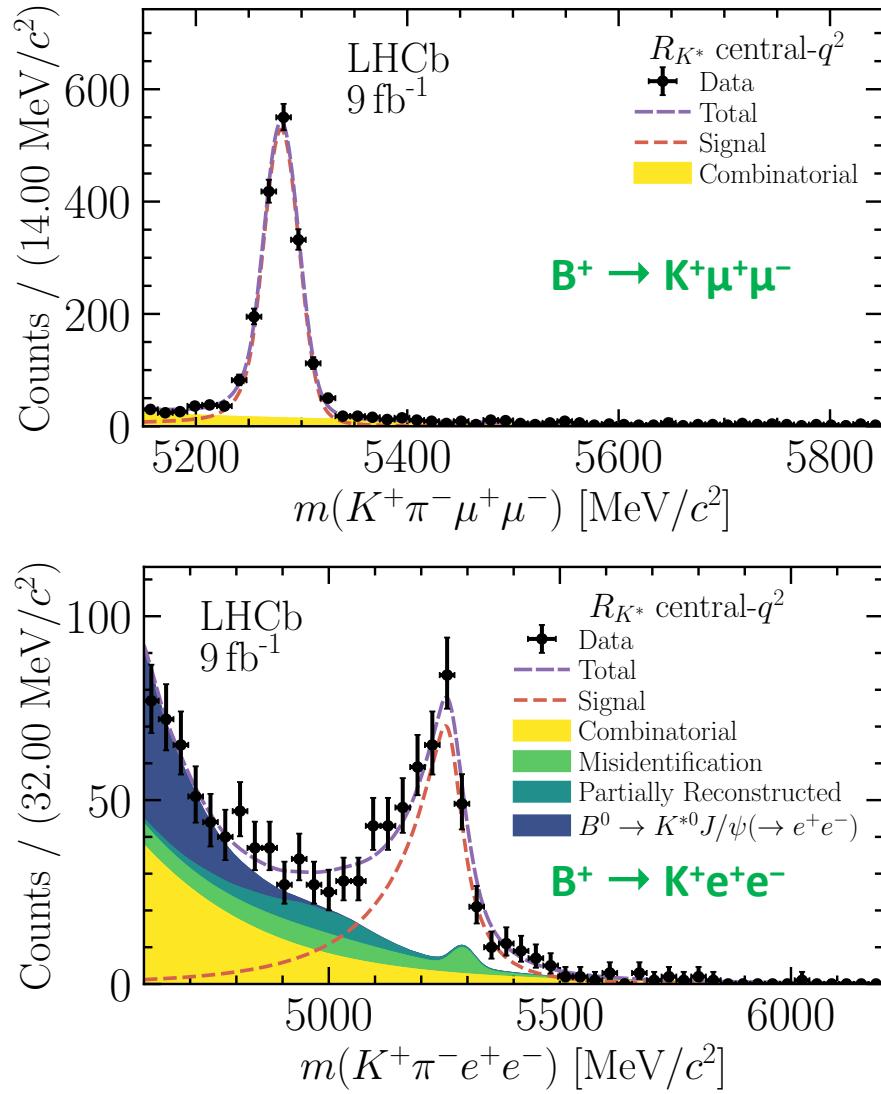
Challenge: electrons and muons interact differently with matter  
⇒ totally different signatures in LHCb

**Muon:** track + muon system hits

**Electron:** track + EM calorimeter shower

Electrons also suffer from *Bremsstrahlung* – lose energy in detector.

<https://arxiv.org/abs/2212.09153> (sub. to PRD)



# Latest $R_K$ (and $R_{K^*}$ ) results

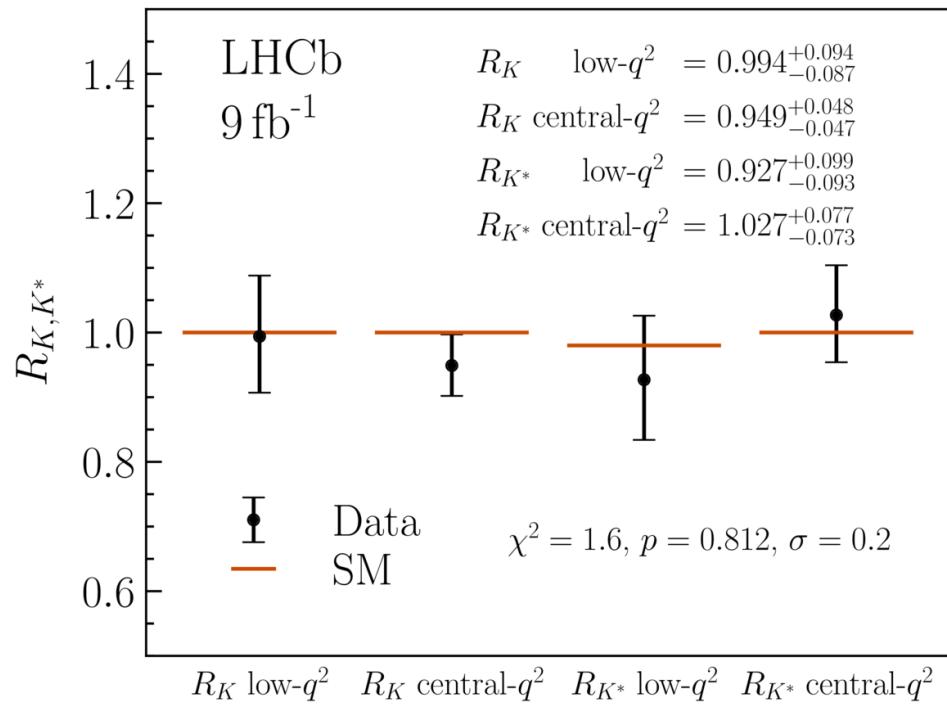
Initial LHCb measurements hinted at non-universality,  $R_K$  below SM by  $\sim 3\sigma$

Latest result is a simultaneous measurement of  $R_K$  and  $R_{K^*}$  with all Run 1-2 data

Now compatible with the SM after applying tighter cuts to reject tricky backgrounds from mis-ID particles

⇒ In these channels, the hint of disagreement with the SM has gone away

**Discovering BSM physics requires a deep knowledge of detector effects and backgrounds!  
(including estimates of ‘unknown unknowns’)**



<https://arxiv.org/abs/2212.09153> (sub. to PRD)

# “The anomalies”

<https://www.nikhef.nl/~pkoppenb/anomalies.html>

Is a picture emerging of a broken Standard Model?

Could it be due to issues with calculations  
(e.g. hadronic effects?)

Or... is this just cherry-picking statistical fluctuations?

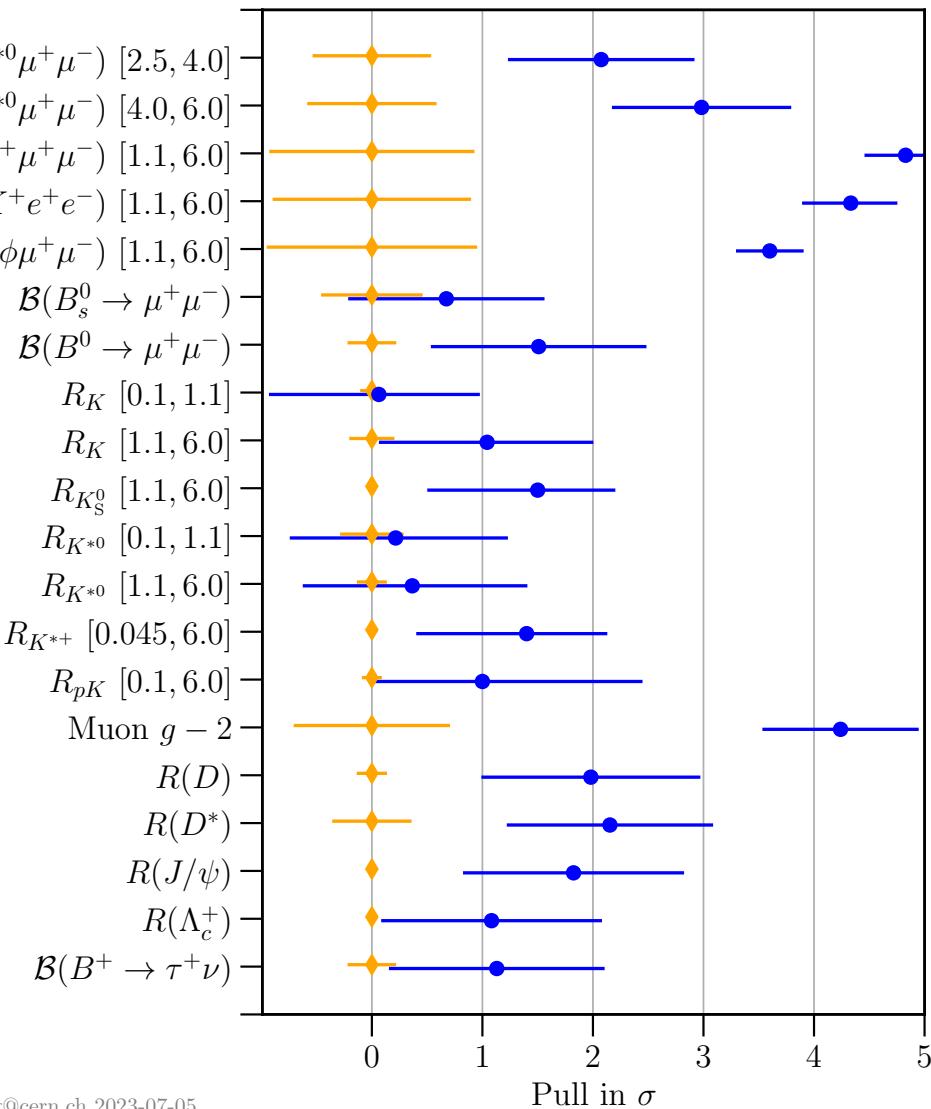
(T.S. Eliot)

Standard Model

This is the way the world ends  
This is the way the world ends  
This is the way the world ends  
Not with a bang but a whimper.

Direct observation

Combination of high-precision measurements



## Part IV: The future of flavour

# The future of flavour

The need for more precision...

*"A special search at Dubna was carried out by Okonov and his group. They did not find a single  $K_L^0 \rightarrow \pi^+ \pi^-$  event among 600 decays into charged particles (Anikira et al., JETP 1962). At that stage the search was terminated by the administration of the lab. The group was unlucky."*

- L. Okun

Remember,  $B(K_L^0 \rightarrow \pi^+ \pi^-) = 0.2\%$

Most measurements limited by statistical precision  $\Rightarrow$  need **more data**

Also need better control over systematics  $\Rightarrow$  **better detectors**

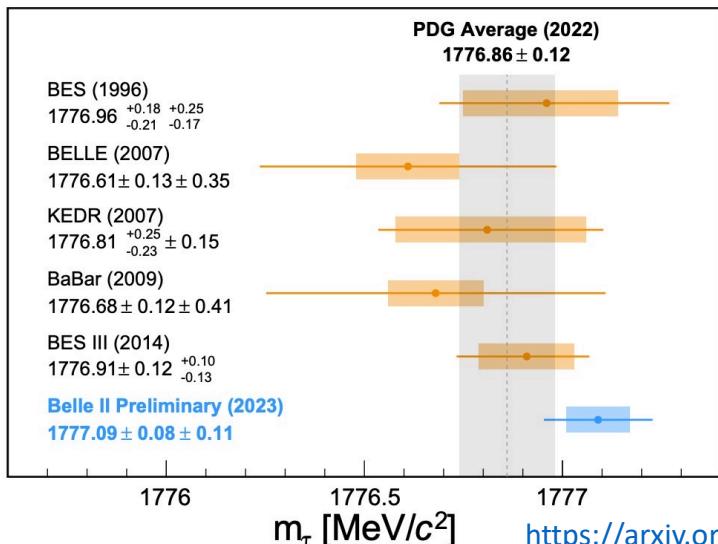
Many new experiments, or upgrades, planned in coming years...

# Belle II (2018 -)

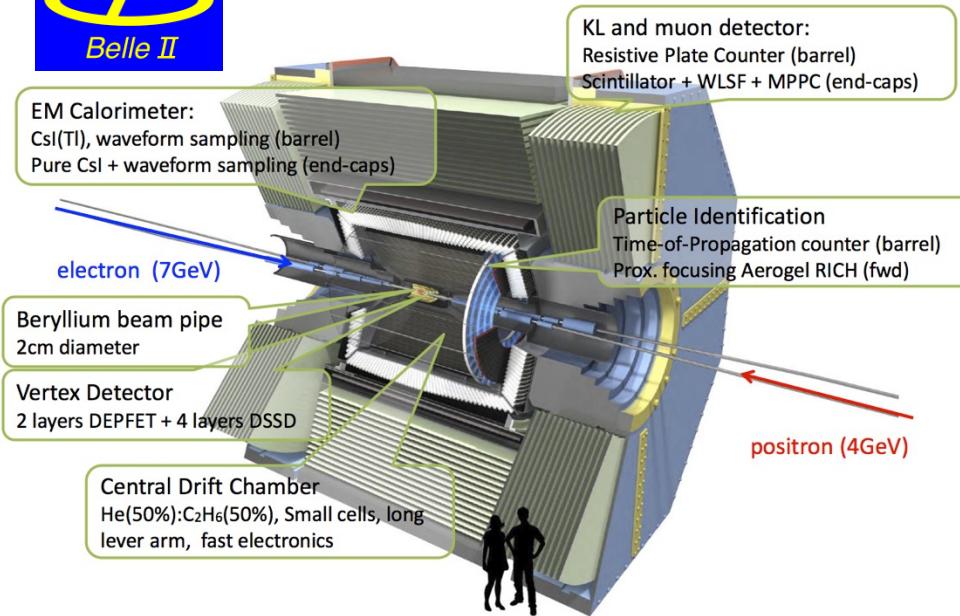
Will collect 40x more data than Belle  
(already a world record luminosity!)

Major accelerator and detector  
upgrades to reach  $50 \text{ ab}^{-1}$

First physics run with complete detector  
started in March 2019



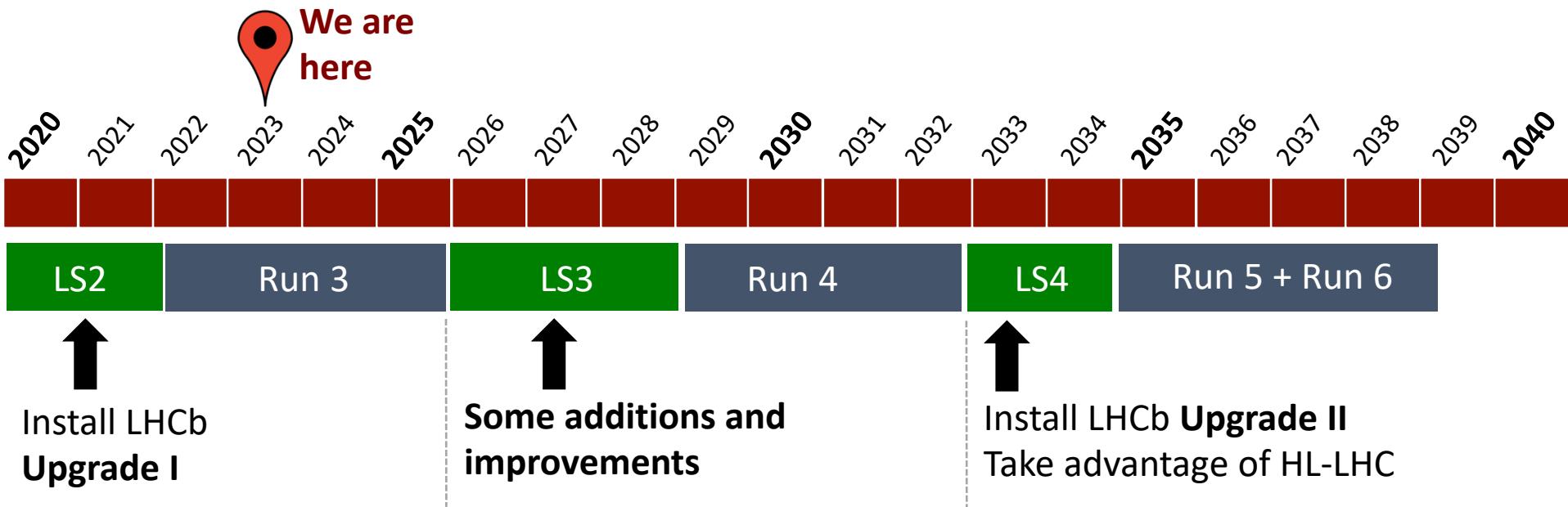
## Belle II Detector



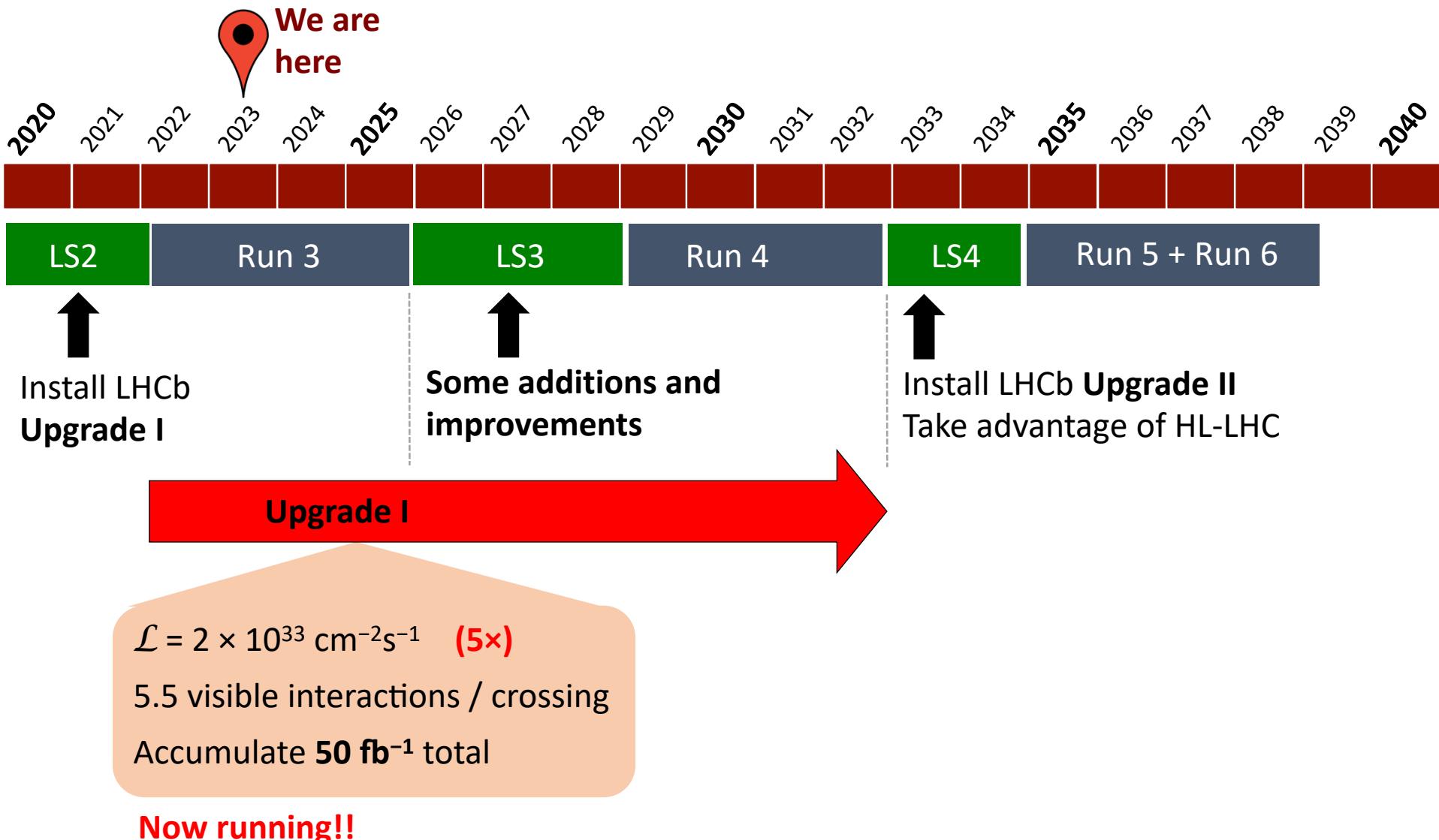
Already surpassing original Belle precision  
in several areas (with fraction of data)

Complementary to LHCb programme

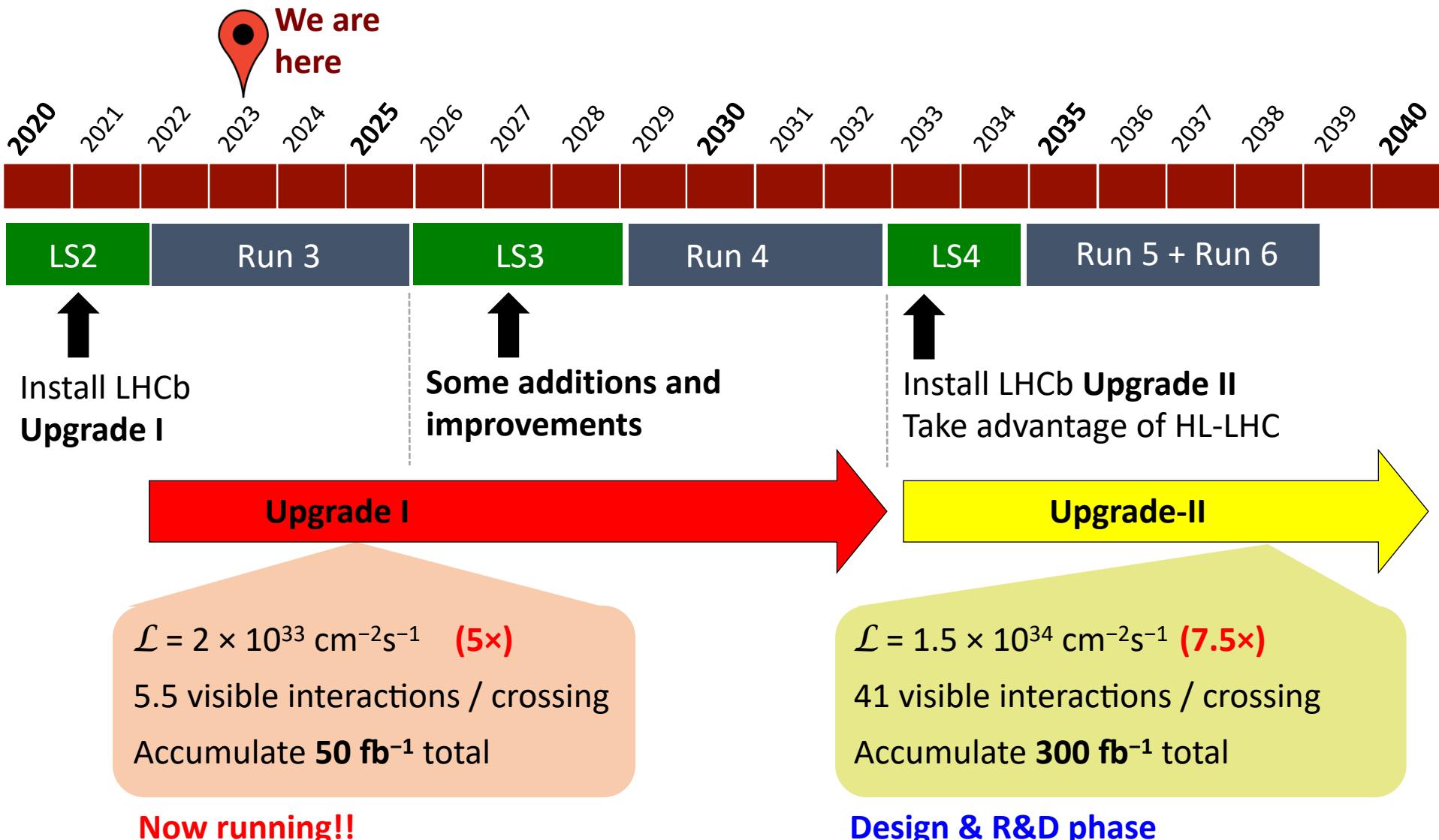
# LHCb Upgrades (2022 – 2040)



# LHCb Upgrades (2022 – 2040)



# LHCb Upgrades (2022 – 2040)



# Kaon physics

Last frontier in kaon physics – observe  $K \rightarrow \pi \nu \bar{\nu}$

Highest CKM suppression of  $s \rightarrow d$  coupling

⇒ measurement sensitive to  $|V_{td}|$  - compare results with B mixing

$$\text{SM} \left[ \begin{array}{l} \mathcal{B}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = (9.1 \pm 0.7) \times 10^{-11} \\ \mathcal{B}(K_L^0 \rightarrow \pi^0 \nu \bar{\nu}) = (3.0 \pm 0.3) \times 10^{-11} \end{array} \right]$$

<https://arxiv.org/abs/1503.02693>

# Kaon physics

Last frontier in kaon physics – observe  $K \rightarrow \pi \nu \bar{\nu}$

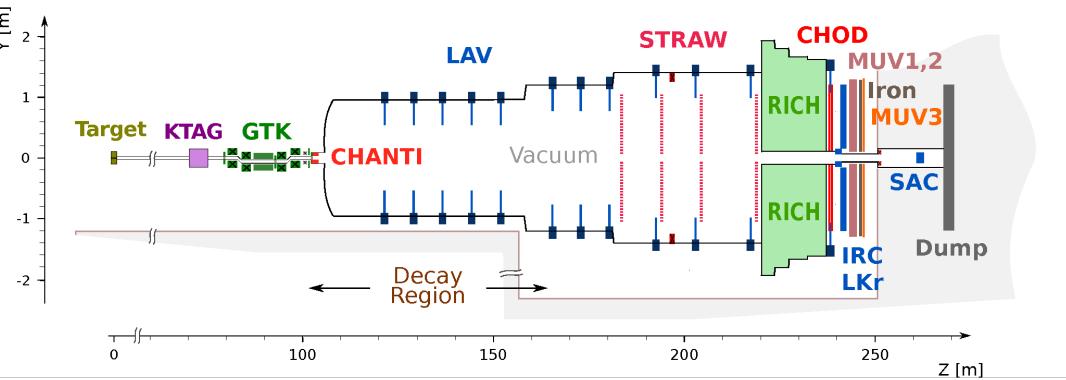
Highest CKM suppression of  $s \rightarrow d$  coupling  
⇒ measurement sensitive to  $|V_{td}|$  - compare results with B mixing

SM  $\begin{cases} B(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = (9.1 \pm 0.7) \times 10^{-11} \\ B(K_L^0 \rightarrow \pi^0 \nu \bar{\nu}) = (3.0 \pm 0.3) \times 10^{-11} \end{cases}$

<https://arxiv.org/abs/1503.02693>



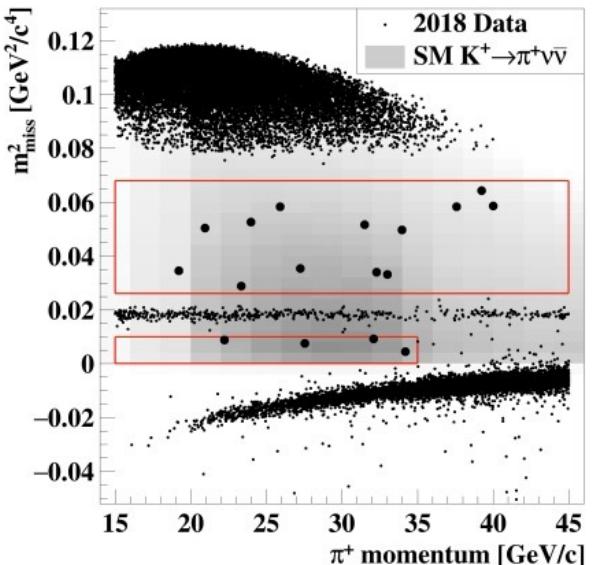
NA62 @ CERN  
(2015 –)



With 2016–18 data,  $3.4\sigma$  evidence:

$$BR(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = (10.6^{+4.0}_{-3.4} |_{\text{stat}} \pm 0.9 |_{\text{syst}}) \times 10^{-11} \text{ at } 68\% \text{ CL}$$

Will reconstruct  $\sim 100 K^+ \rightarrow \pi^+ \nu \bar{\nu}$  at SM rate ⇒ 10% precision on  $|V_{td}|$



# Summary

## Flavour physics is a powerful and versatile tool

- Challenge SM predictions ⇒ see NP indirectly

History tells us that this is often the gateway to major discoveries

## A huge field, covering many areas, with many experiments past, present, future

- Many important ones not covered (including Atlas/CMS, Tevatron, BES-III, CLEO)

## Future looks bright

- Many new and upgraded experiments coming
- Are we on the verge of a breakthrough?

## Experiments

<http://lhcb-public.web.cern.ch/>  
<https://www.belle2.org/>  
<https://mu2e.fnal.gov/>  
<http://muon-g-2.fnal.gov/>  
<https://www.psi.ch/mu3e/>

## Resources

<https://hflav.web.cern.ch/>  
<http://ckmfitter.in2p3.fr/>  
<http://www.utfit.org/>  
<http://pdglive.lbl.gov/>



@LHCbExperiment  
@LHCbPhysics  
@BelleIICollab  
@QuarkWilliams

# Enjoy your stay!



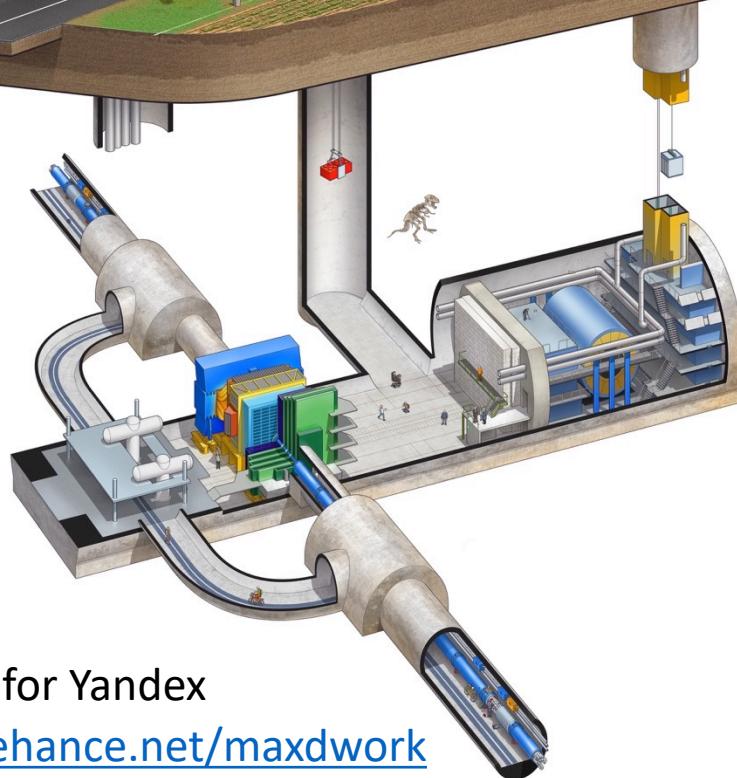
Max Degtyarev

<https://www.behance.net/maxdwork>

# Extra Slides



- Current/future muon experiments
- CPV in  $B_s^0$  mixing
- The ‘golden mode’ for charm mixing:  
 $D^0 \rightarrow K_S^0 \pi^+ \pi^-$
- Example methods for measuring  $\gamma$
- CPV in decay: “ $K\pi$  problem”
- Operator product expansion and radiative B decays
- $R_K$  measurement



Max Degtyarev for Yandex  
<https://www.behance.net/maxdwork>

# CP violation in charm decays (2022)

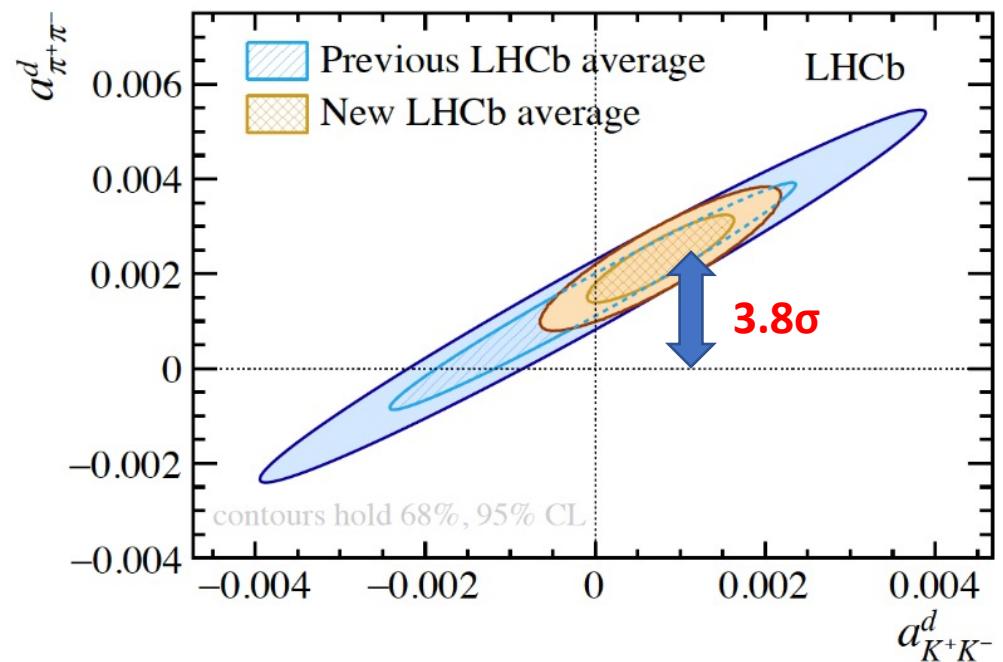
Separate measurement of CP asymmetry in  $D^0 \rightarrow K^+K^-$

⇒ allows CP asymmetries in both channels to be measured, with constraint from  $\Delta A_{CP}$

$$A_{CP}(D^0 \rightarrow \pi^+\pi^-) = (23.2 \pm 6.1) \times 10^{-4}$$

3.8 $\sigma$  from zero

⇒ First evidence of CP violation in a specific charm quark decay channel



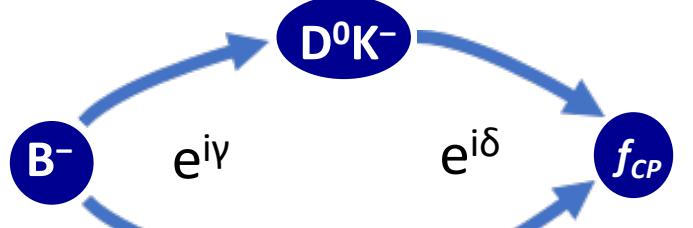
<https://cerncourier.com/a/lhcbs-digs-deeper-in-cp-violating-charm-decays/>

<https://agenda.infn.it/event/28874/contributions/169315/>

# Measuring $\gamma$ in tree decays

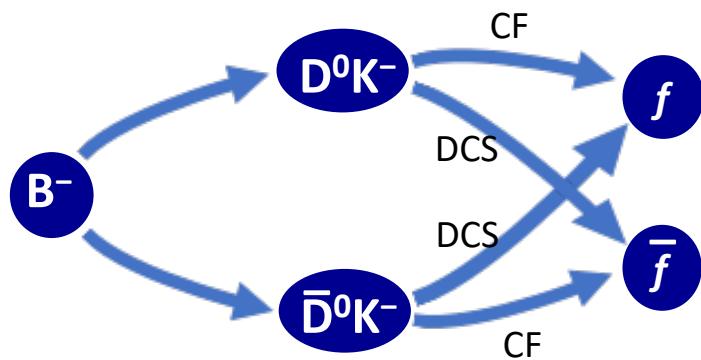
## (1) 'GLW'

- CP eigenstates e.g.  $f = \pi^+\pi^-, K^+K^-$
- [[https://doi.org/10.1016/0370-2693\(91\)90034-N](https://doi.org/10.1016/0370-2693(91)90034-N)]
- [[https://doi.org/10.1016/0370-2693\(91\)91756-L](https://doi.org/10.1016/0370-2693(91)91756-L)]



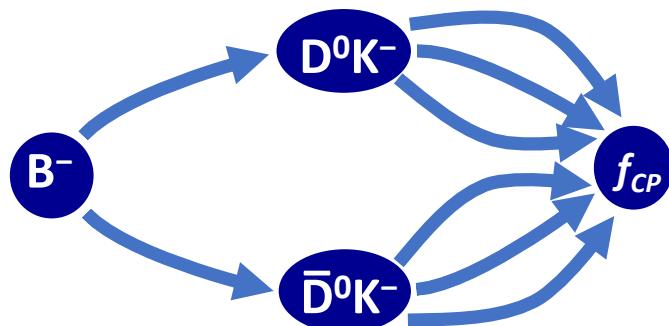
## (2) 'ADS'

- Cabibbo-Favoured or Doubly-Cabibbo-Suppressed decays e.g.  $f = K^-\pi^+$
- [<https://doi.org/10.1103/PhysRevD.63.036005>]
- [<https://doi.org/10.1103/PhysRevLett.78.3257>]

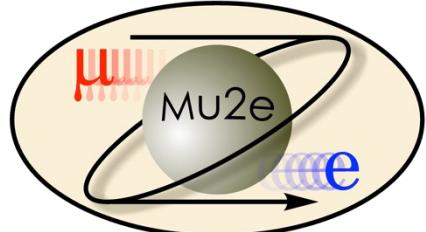


## (3) 'BPGGSZ'

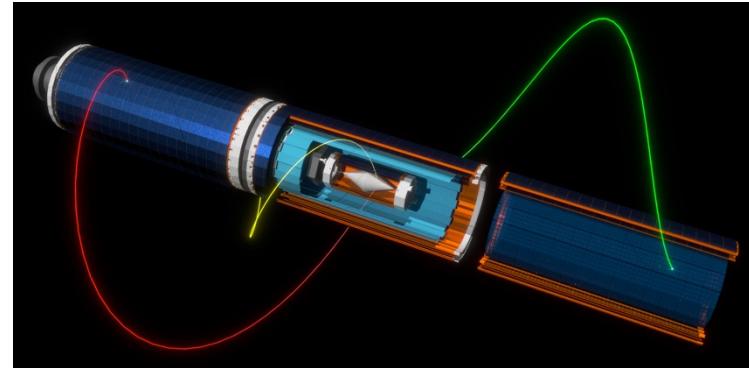
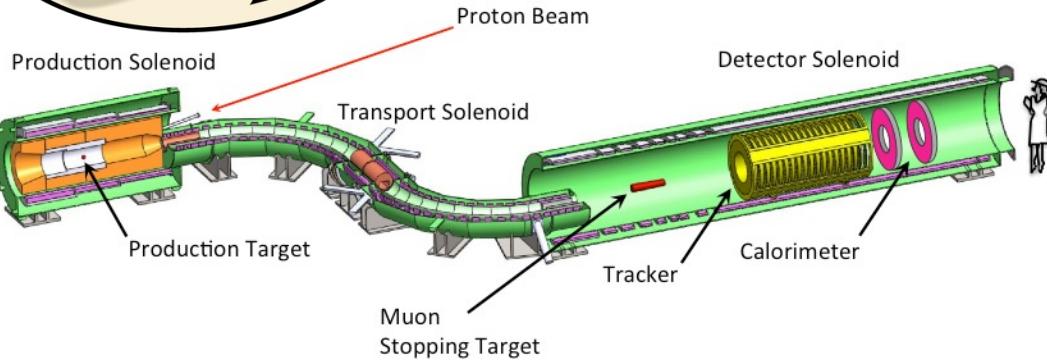
- 3-body final states e.g.  $f = K_S^0\pi^+\pi^-$
- Reached via intermediate resonances
- [<https://doi.org/10.1103/PhysRevD.68.054018>]



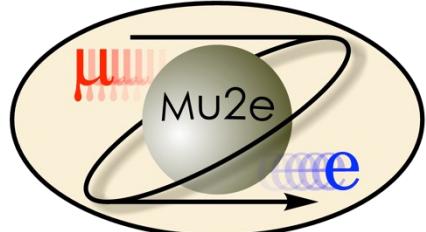
# Muon physics



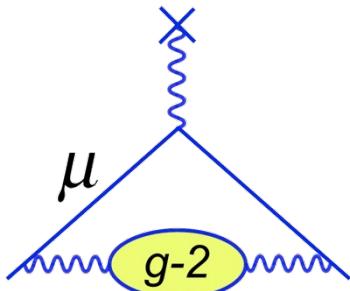
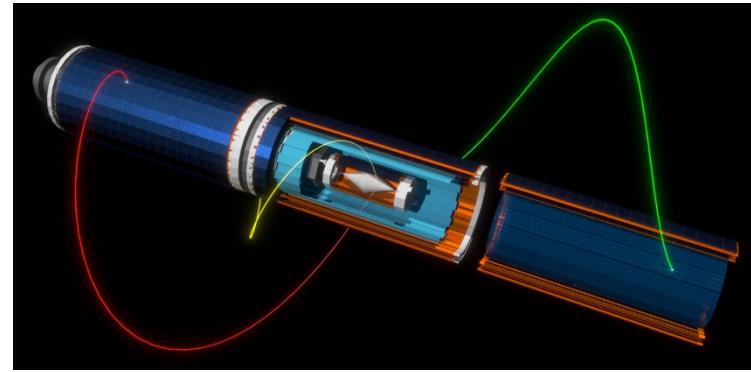
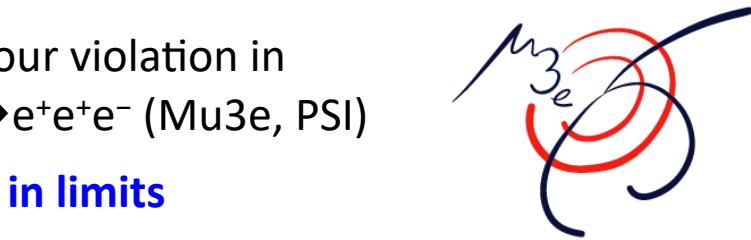
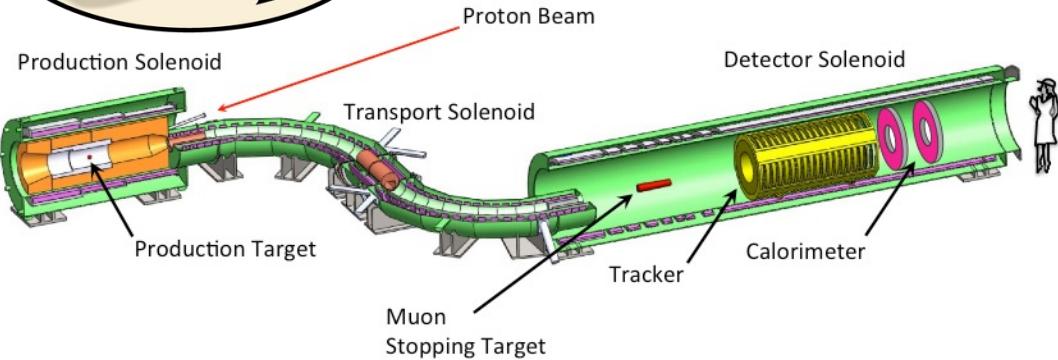
Search for charged lepton flavour violation in  
 $\mu^- \rightarrow e^-$  (**Mu2e**, FNAL) and  $\mu^+ \rightarrow e^+ e^+ e^-$  (**Mu3e**, PSI)  
⇒  $\times 10^4$  improvement in limits



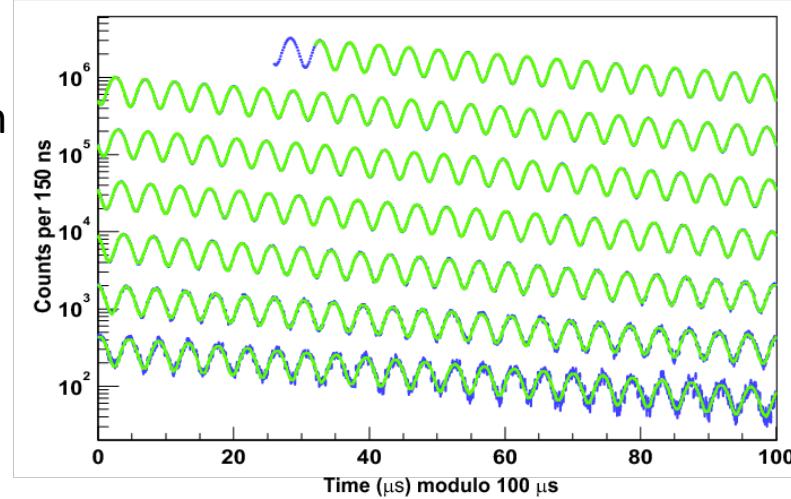
# Muon physics



Search for charged lepton flavour violation in  
 $\mu^- \rightarrow e^-$  (Mu2e, FNAL) and  $\mu^+ \rightarrow e^+ e^+ e^-$  (Mu3e, PSI)  
⇒  **$\times 10^4$  improvement in limits**



**g-2** experiment (FNAL)  
⇒ Precise measurement of muon magnetic moment  
Previous exp. result disagrees with prediction at  $\sim 3.5\sigma$   
⇒  **$4\times$  improvement in precision**  
**Now a  $4\sigma$  discrepancy (4/2021)**



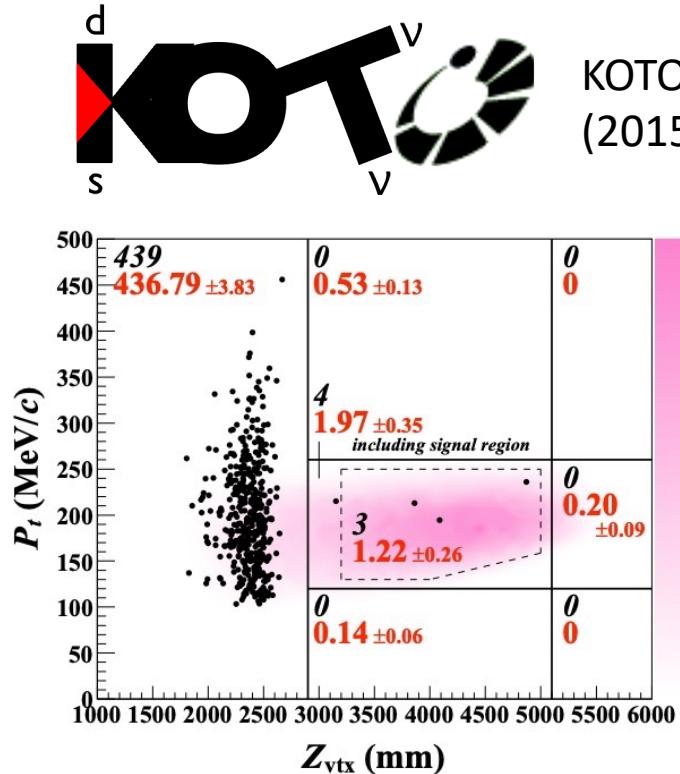
# Kaon physics

Last frontier in kaon physics – observe  $K \rightarrow \pi vv$

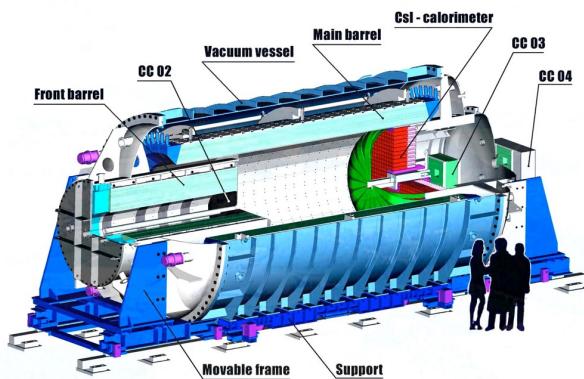
Highest CKM suppression of  $s \rightarrow d$  coupling  
⇒ measurement sensitive to  $|V_{td}|$  - compare results with B mixing

$$\text{SM} \begin{cases} B(K^+ \rightarrow \pi^+ vv) = (9.1 \pm 0.7) \times 10^{-11} \\ B(K_L^0 \rightarrow \pi^0 vv) = (3.0 \pm 0.3) \times 10^{-11} \end{cases}$$

<https://arxiv.org/abs/1503.02693>



Measure  $B(K_L^0 \rightarrow \pi^0 vv)$  to 10% precision  
(@ SM rate)



- All data from 2016–18 analysed
- 3 signal-like events observed, consistent with BG
  - $B(K_L^0 \rightarrow \pi^0 vv) < 490 \times 10^{-11}$  @90% CL

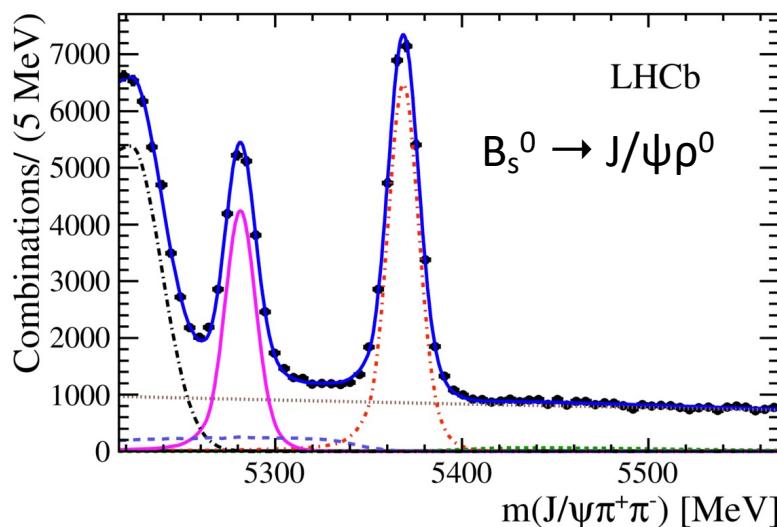
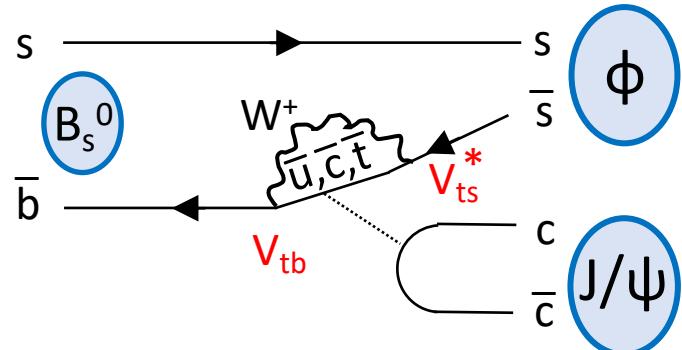
10× better than previous limit

# Penguin pollution in $B_s^0 \rightarrow J/\psi \phi$

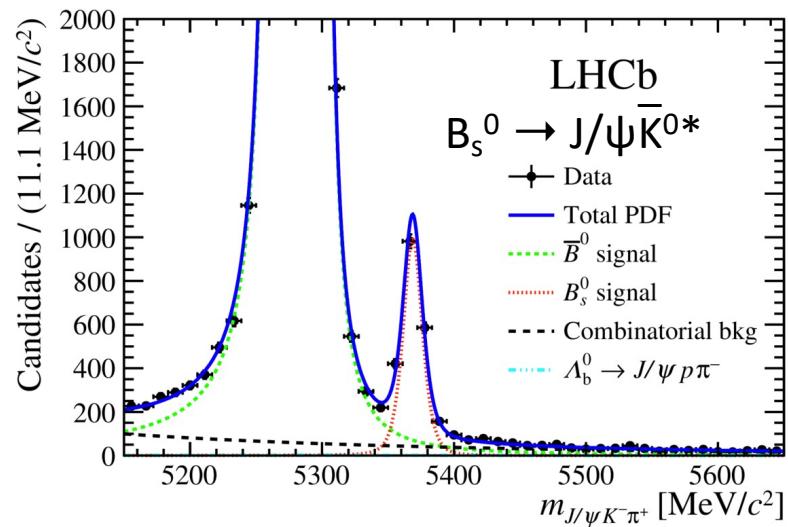
Penguin pollution breaks equality  $\phi_s = -2\beta_s$   
 ⇒ can mimic effects of new physics

Strategy is to study in dedicated channels to set limits on the size:

$\delta\phi_s = [-0.018, 0.021]$  rad at 95% confidence



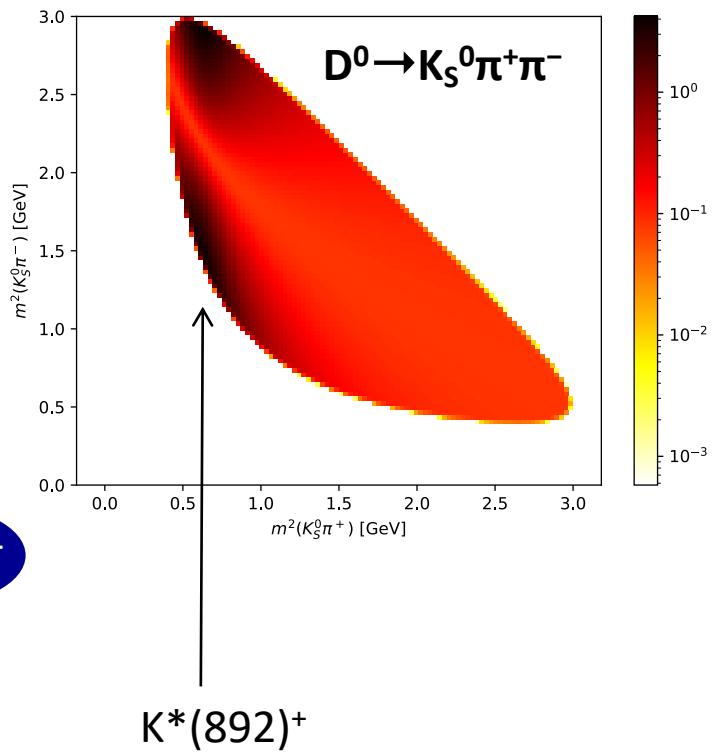
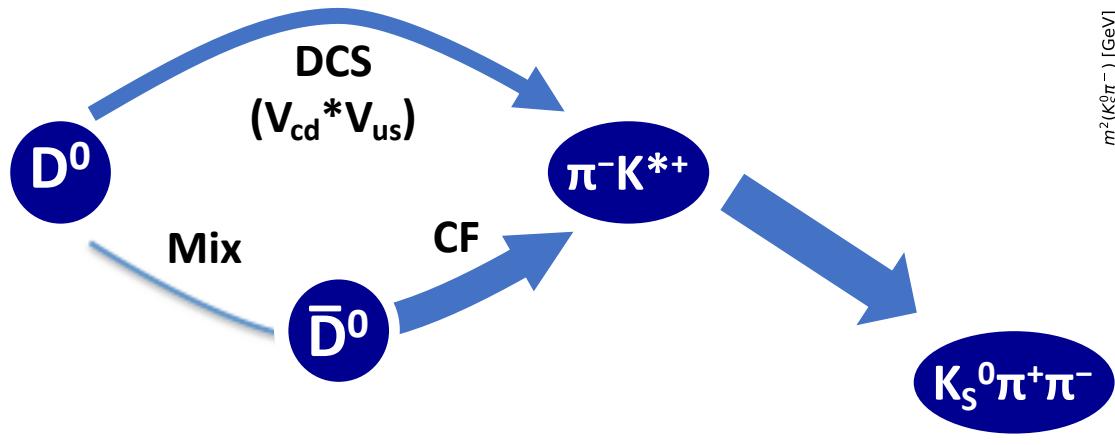
<https://doi.org/10.1016/j.physletb.2015.01.008>



[https://doi.org/10.1007/JHEP11\(2015\)082](https://doi.org/10.1007/JHEP11(2015)082)

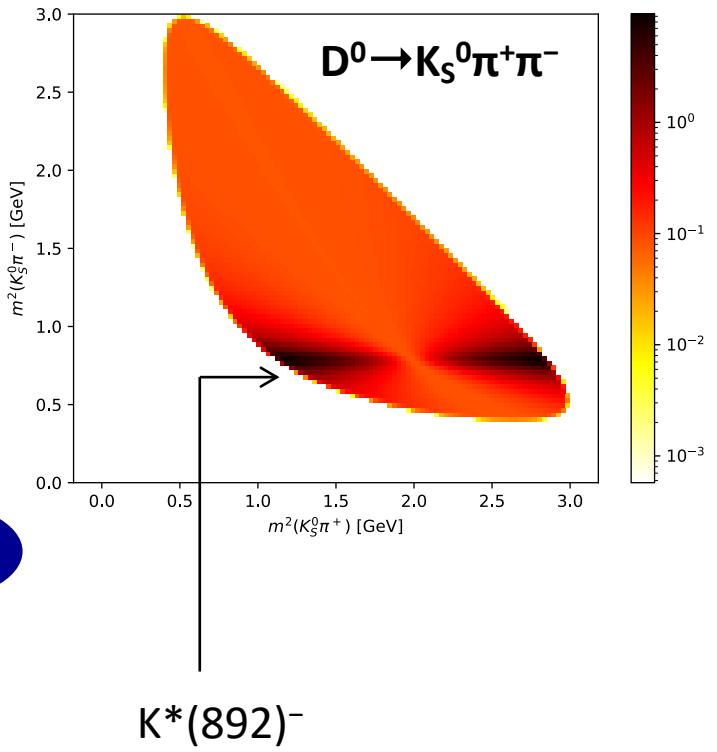
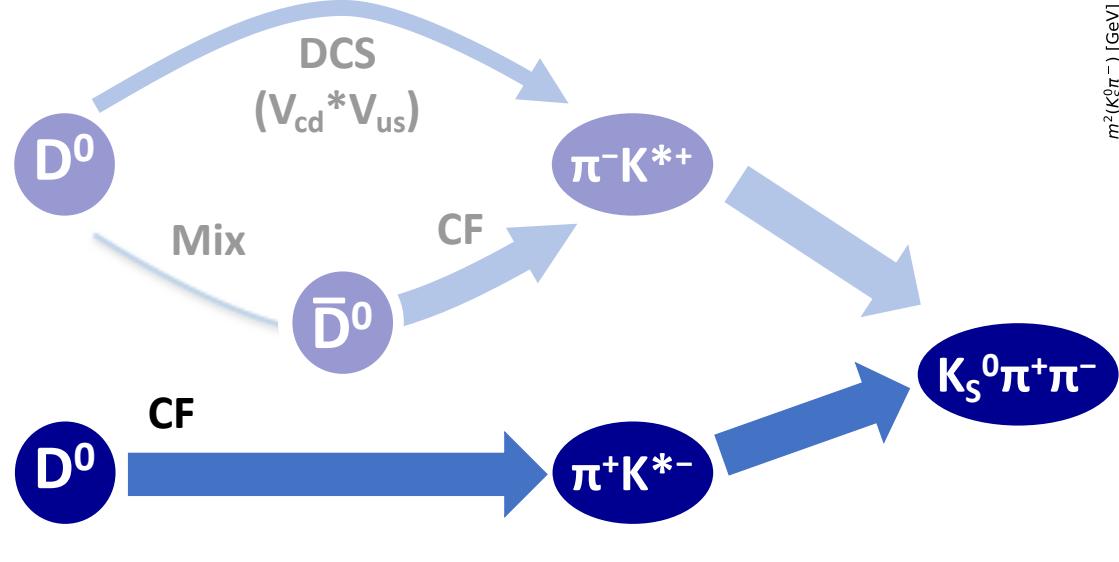
# Charm mixing: ‘Golden mode’ $D^0 \rightarrow K_S^0 \pi^+ \pi^-$

‘Right-sign’ (CF) and ‘wrong-sign’ (DCS or oscillated) decays to *same final state*



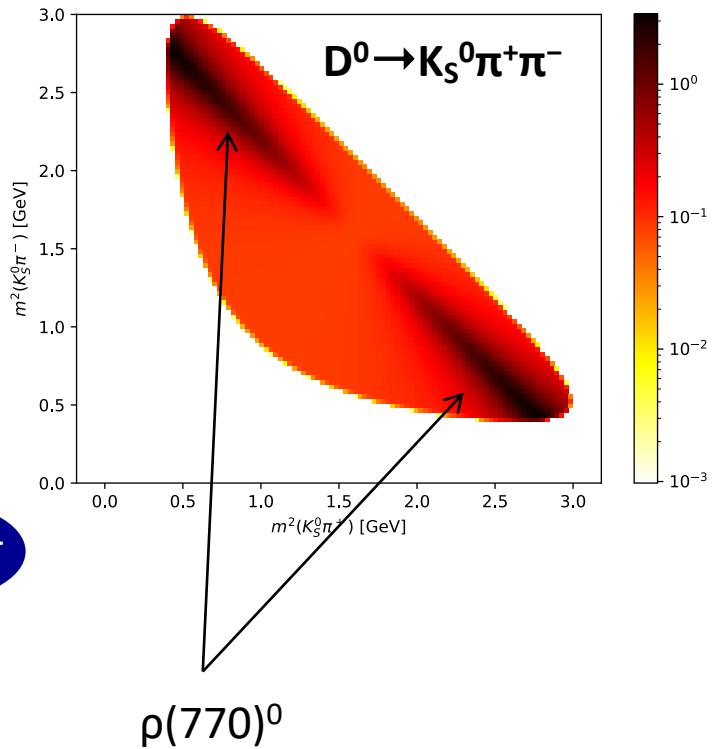
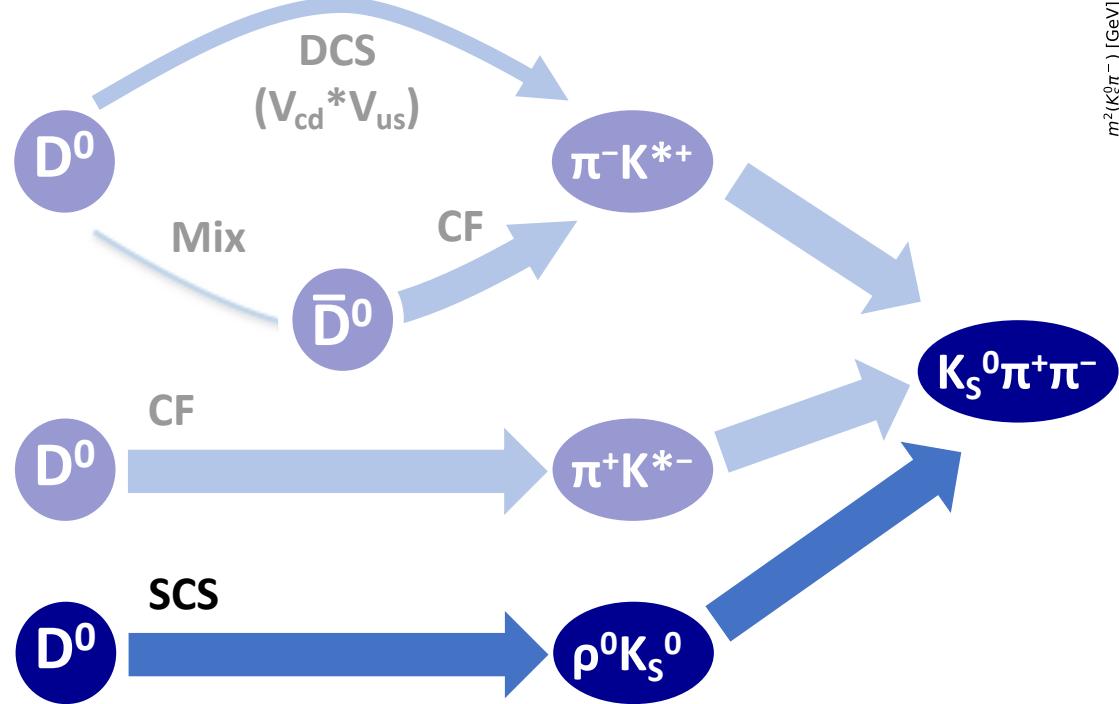
# Charm mixing: ‘Golden mode’ $D^0 \rightarrow K_S^0 \pi^+ \pi^-$

‘Right-sign’ (CF) and ‘wrong-sign’ (DCS or oscillated) decays to *same final state*



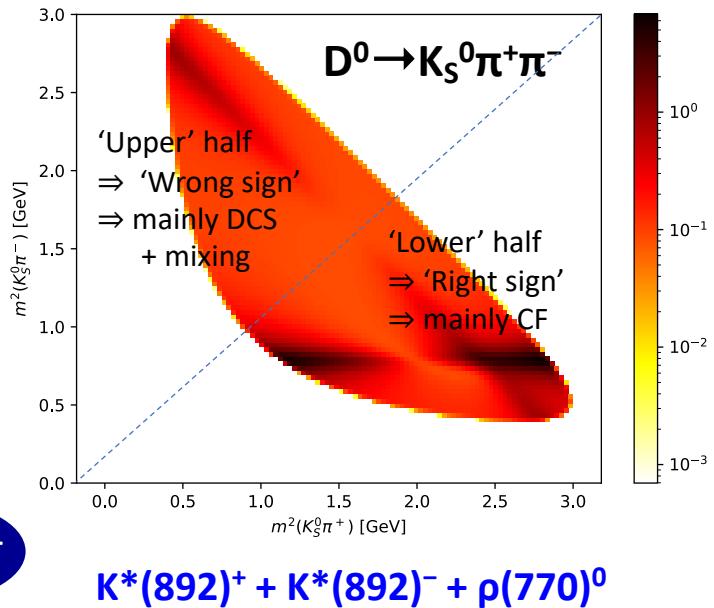
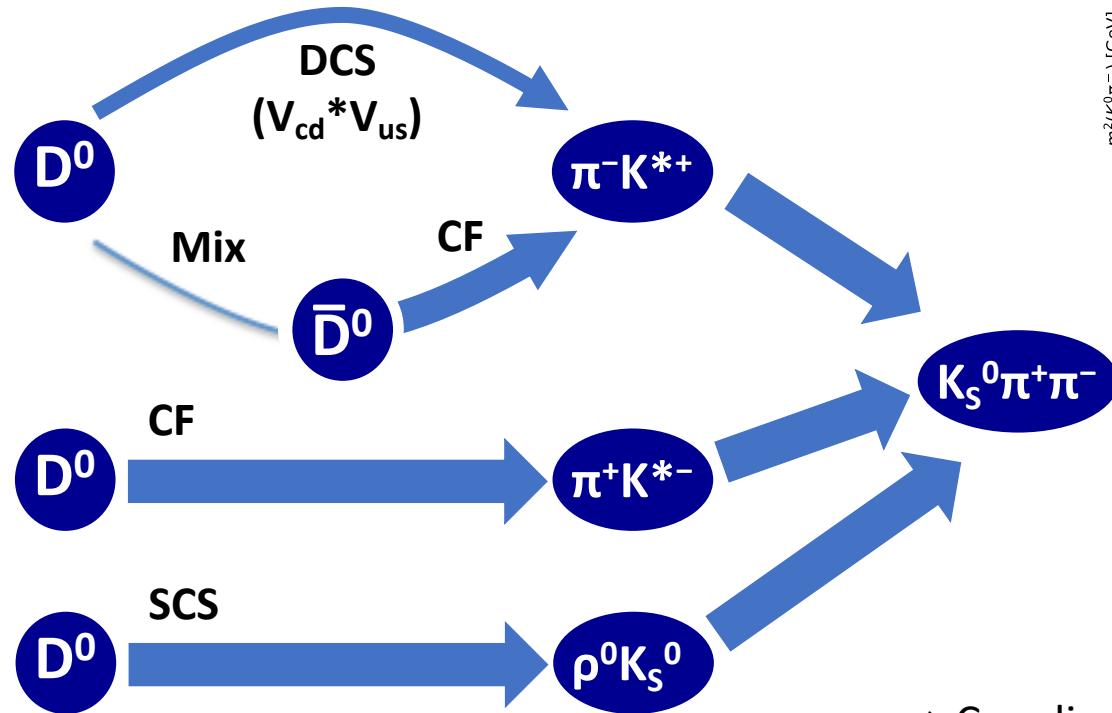
# Charm mixing: ‘Golden mode’ $D^0 \rightarrow K_S^0 \pi^+ \pi^-$

‘Right-sign’ (CF) and ‘wrong-sign’ (DCS or oscillated) decays to *same final state*



# Charm mixing: ‘Golden mode’ $D^0 \rightarrow K_S^0 \pi^+ \pi^-$

‘Right-sign’ (CF) and ‘wrong-sign’ (DCS or oscillated) decays to *same final state*



$\Rightarrow$  Can directly measure all four mixing and CPV parameters  $x, y, |q/p|, \arg(q/p)$

Requires **time and phase-space** dependent analysis

# Charm mixing: ‘Golden mode’ $D^0 \rightarrow K_S^0 \pi^+ \pi^-$

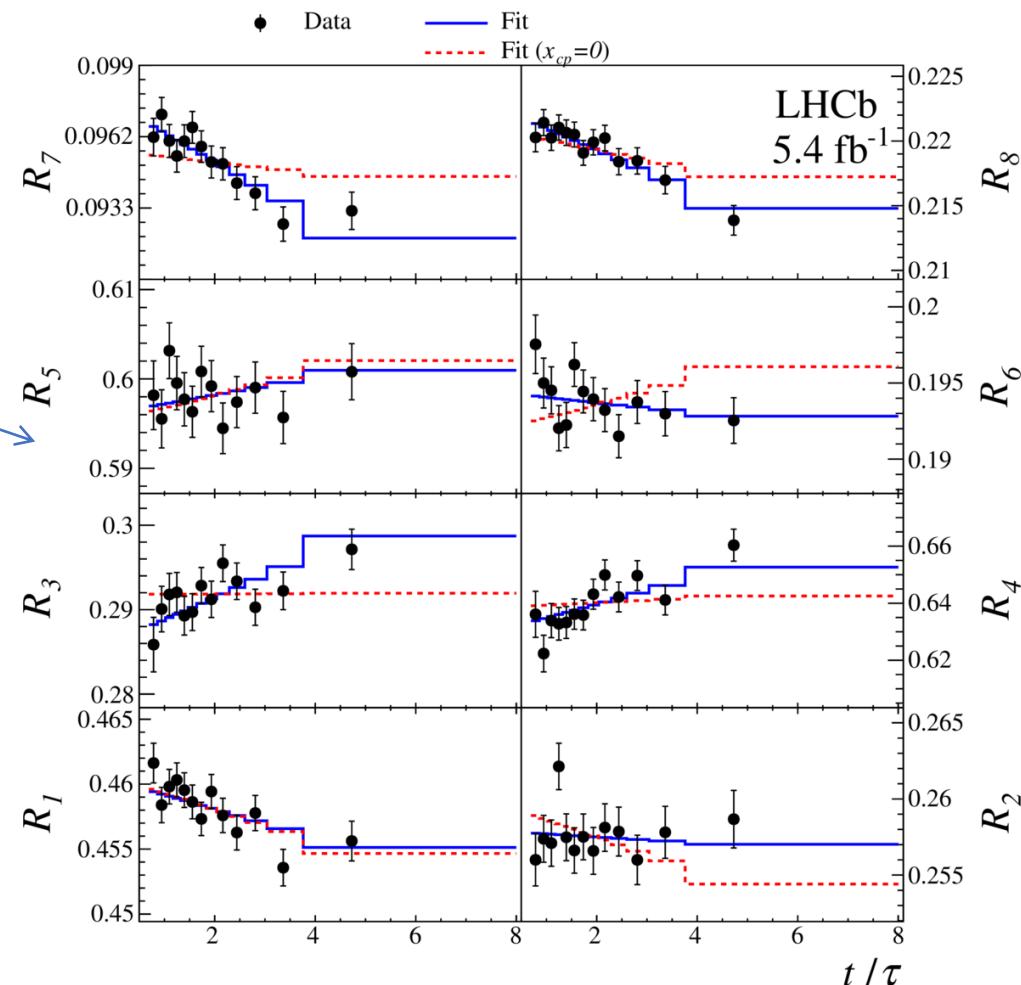
Plot ratio of candidates in upper/lower half of Dalitz plane  
 ⇒ 8 ratios in Dalitz bins  $R_1 - R_8$   
 ⇒ Fit to extract parameters

**Clear time dependence from mixing**

$$x \equiv \Delta m / \Gamma = [0.397 \pm 0.046 \pm 0.029]\% \\ y \equiv \Delta \Gamma / 2\Gamma = [0.459 \pm 0.120 \pm 0.085]\%$$

First measurement of non-zero  $x$   
 (>7σ significance)

Oscillation period ~630ps  
 ( $D^0$  lifetime 0.4ps)

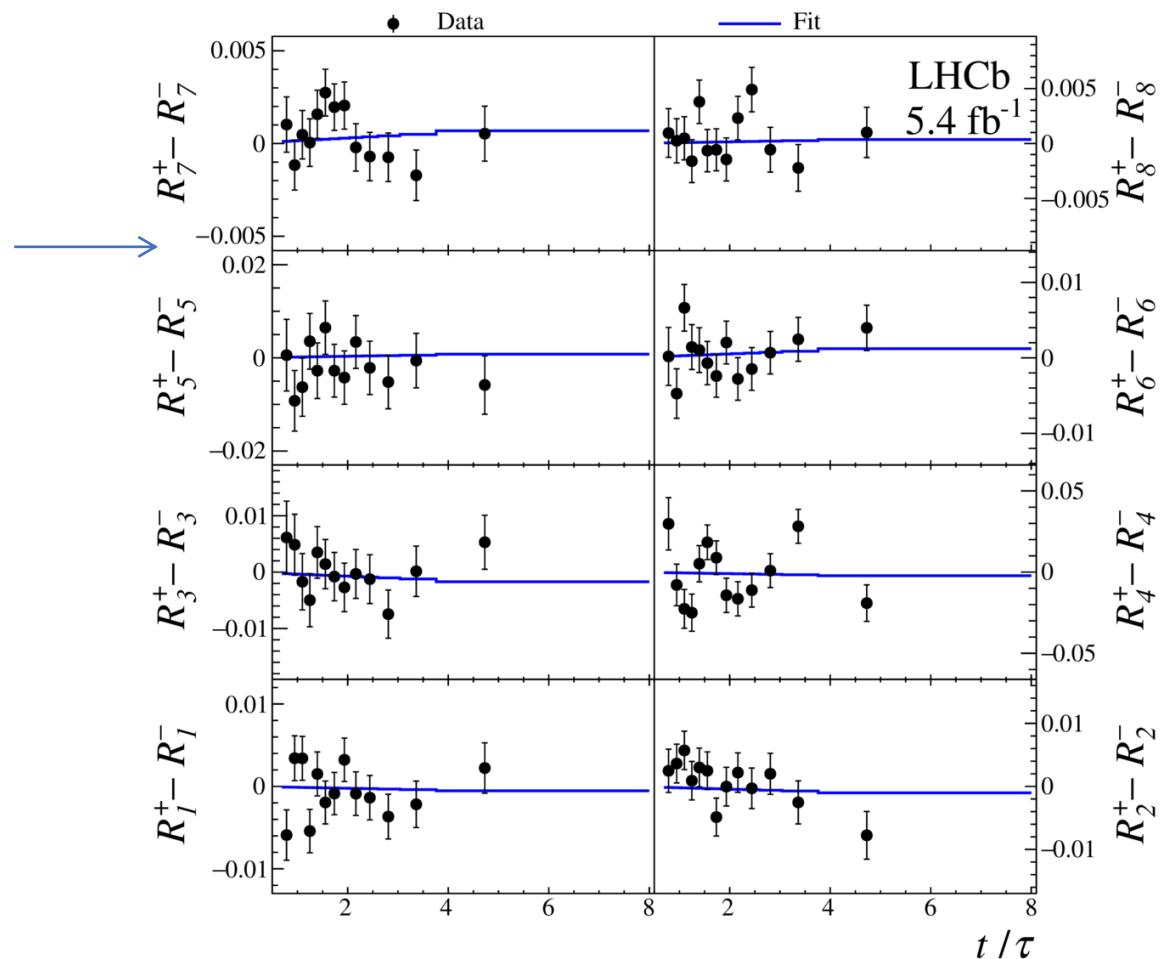


arXiv:2106.03744

Phys. Rev. Lett. 127, 111801

# Charm mixing: ‘Golden mode’ $D^0 \rightarrow K_S^0 \pi^+ \pi^-$

No significant differences  $D^0$  vs  $\bar{D}^0$   
⇒ no evidence for CP violation



arXiv:2106.03744

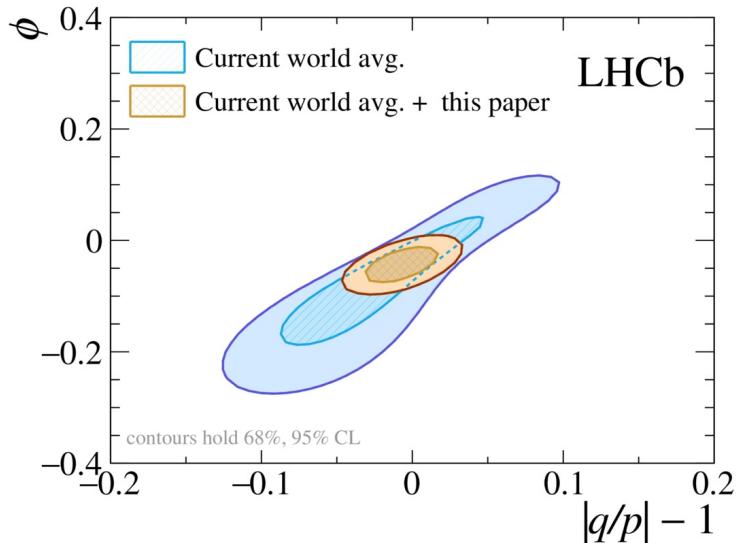
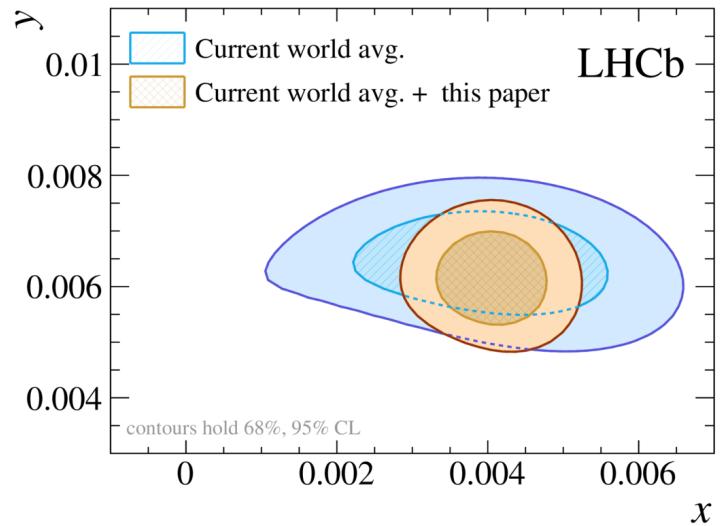
Phys. Rev. Lett. 127, 111801

# Charm mixing: ‘Golden mode’ $D^0 \rightarrow K_S^0 \pi^+ \pi^-$

Combine with all previous measurements

⇒ **Significant improvements in WA** for both mixing and CPV parameters

⇒ Hence the ‘Golden Mode’ for charm



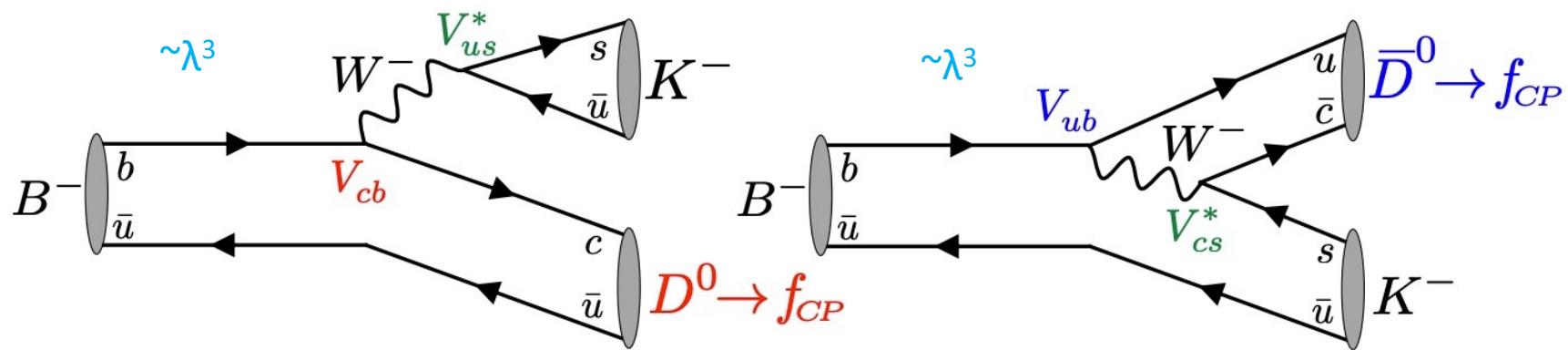
# Measuring $\gamma$ : GLW method

Colour favoured:  $b \rightarrow c W$

- Magnitude  $|F|$
- Weak phase  $\phi_F$ , strong phase  $\delta_F$

Colour suppressed:  $b \rightarrow u W$

- Magnitude  $|S|$
- Weak phase  $\phi_S$ , strong phase  $\delta_S$



Amplitudes to final state  $f_{CP}$ :

$$B^-: A_f = |F| e^{i(\delta_F + \phi_F)} + |S| e^{i(\delta_S + \phi_S)}$$

$$B^+: A_f = |F| e^{i(\delta_F - \phi_F)} + |S| e^{i(\delta_S - \phi_S)}$$

Under CP operation:

- Weak phases change sign
- Strong phases unchanged

# Measuring $\gamma$ : GLW method

'Trick': Weak phase difference  $\Phi_F - \Phi_S = \gamma$  always

while  $\delta_B \equiv \delta_F - \delta_S$  and  $r_B = |S|/|F|$  depend on the decay

( $F^+$ : fraction of  $D^0$  decay to  $CP=+1$  eigenstate)

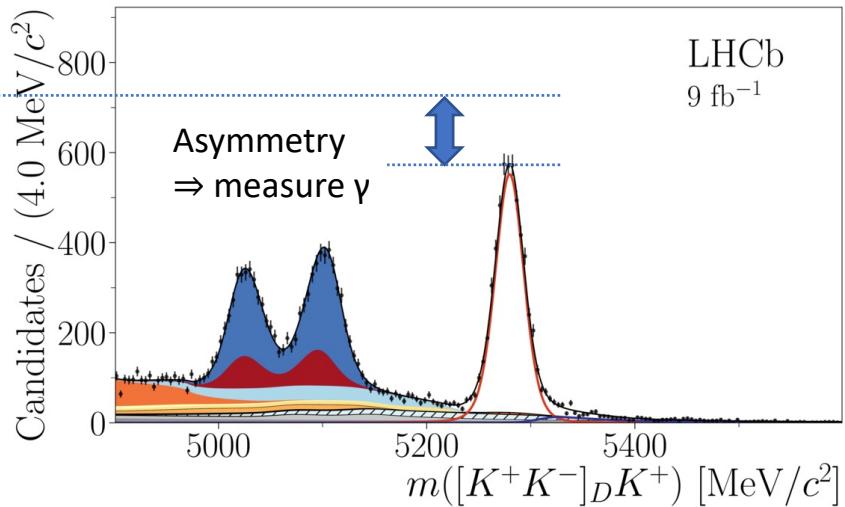
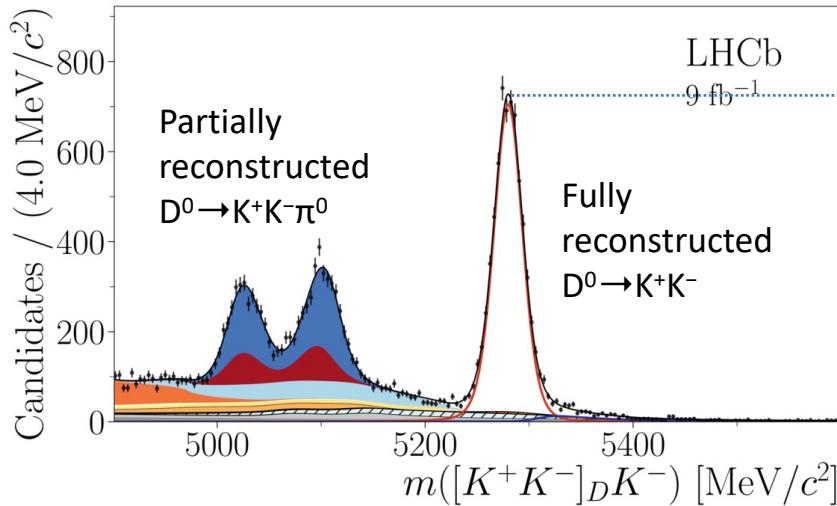
## GLW observables

Asymmetry ( $B^+$  vs  $B^-$ ):

$$A_{CP} = \frac{\pm 2r_B(2F^+ + 1) \sin(\delta_B) \sin(\gamma)}{1 + r_B^2 \pm 2r_B(2F^+ + 1) \cos(\delta_B) \cos(\gamma)}$$

Total rate:

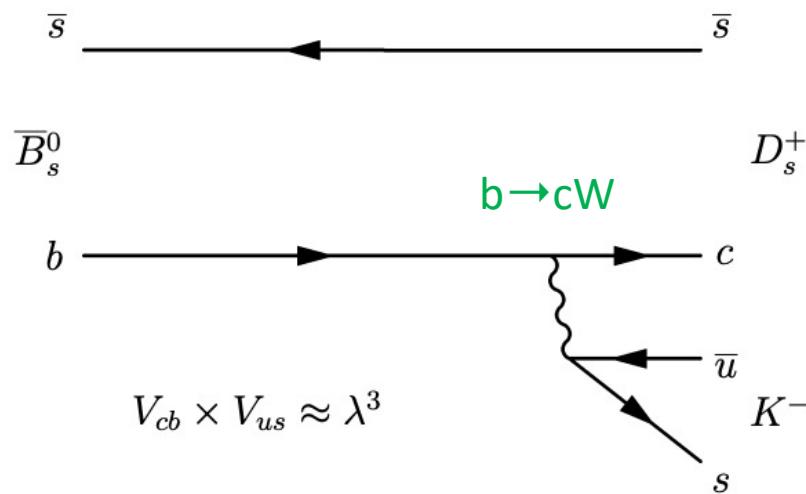
$$R_{CP} = 1 + r_B^2 \pm 2r_B(2F^+ + 1) \cos(\delta_B) \cos(\gamma)$$



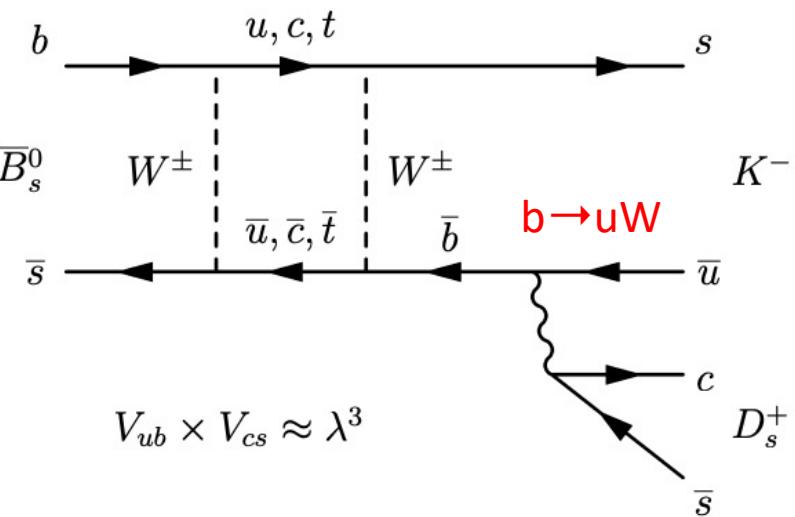
# Measuring $\gamma$ with $B_s^0$ mixing

Can also measure  $\gamma$  with  $B_s^0$  meson mixing providing interference [compare  $\sin(2\beta)$ ]

$\bar{B}_s^0 \rightarrow D_s^+ K^-$  (tree)



$\bar{B}_s^0 \rightarrow D_s^+ K^-$  (loop, with  $B_s^0$  mixing)



- Weak phase difference is  $(\gamma - 2\beta_s)$   $\Rightarrow$  need input from  $B_s^0$  measurements
- Need time-dependent analysis to observe oscillations

# Measuring $\gamma$ with $B_s^0$ mixing

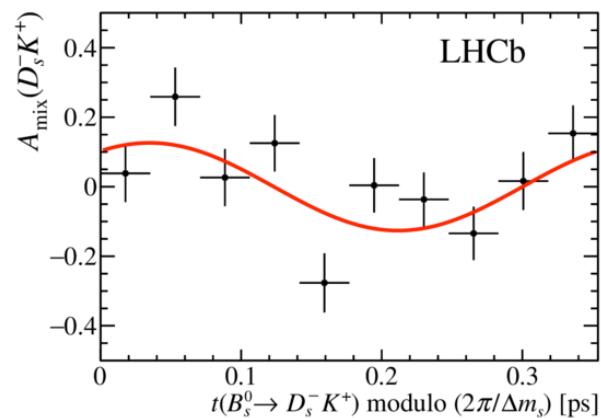
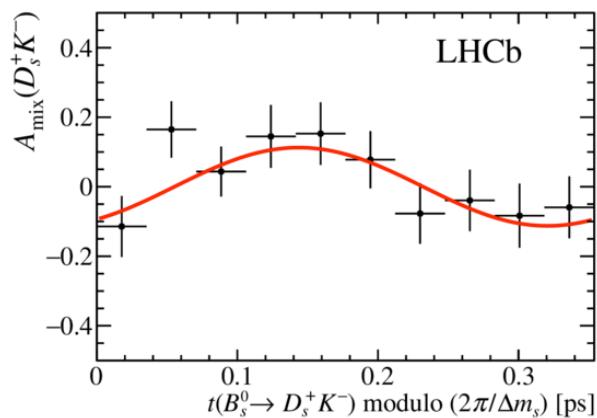
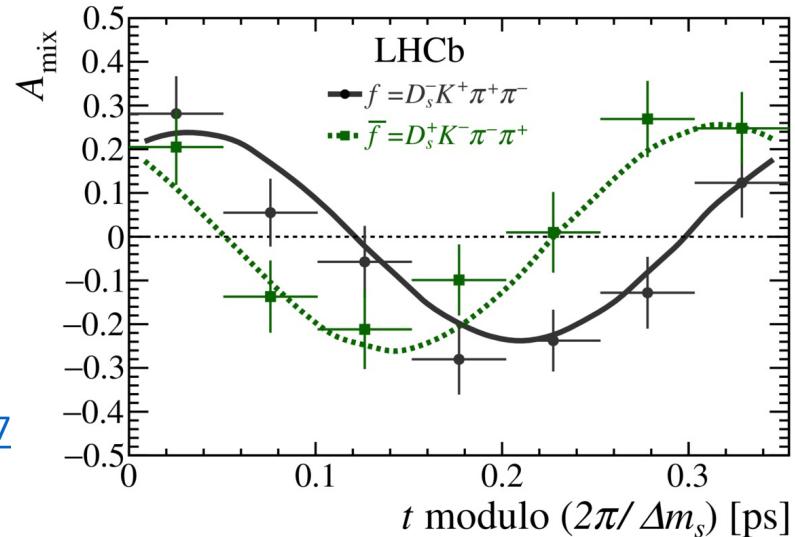
Fit for time-dependent asymmetry



$$\gamma = (44 \pm 12)^\circ \text{ modulo } 180^\circ$$

Full LHCb Run 1+2 (9fb<sup>-1</sup>)

[https://doi.org/10.1007/JHEP03\(2021\)137](https://doi.org/10.1007/JHEP03(2021)137)



$$B_s^0 \rightarrow D_s^- K^+ \quad \text{modulo } 180^\circ$$

$$\gamma = (128^{+17}_{-22})^\circ$$

LHCb Run 1 (3fb<sup>-1</sup>)

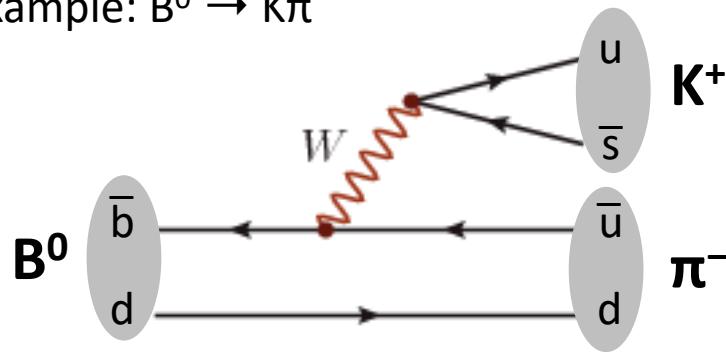
[https://doi.org/10.1007/JHEP03\(2018\)059](https://doi.org/10.1007/JHEP03(2018)059)

# CPV in B decay: “K $\pi$ problem”

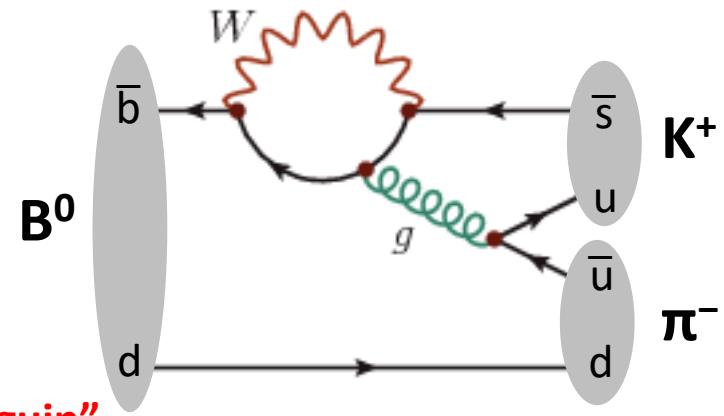
CPV in decay

Remember – need  $\geq 2$  interfering processes  
(preferably with similar magnitudes)

Example:  $B^0 \rightarrow K\pi$

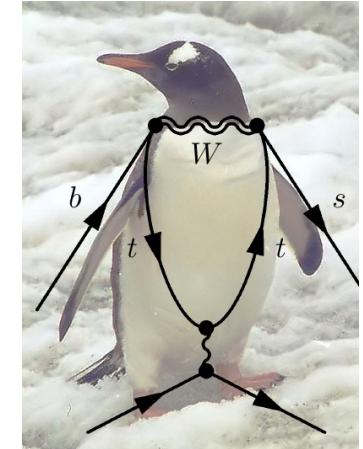


Tree-level



“Penguin”

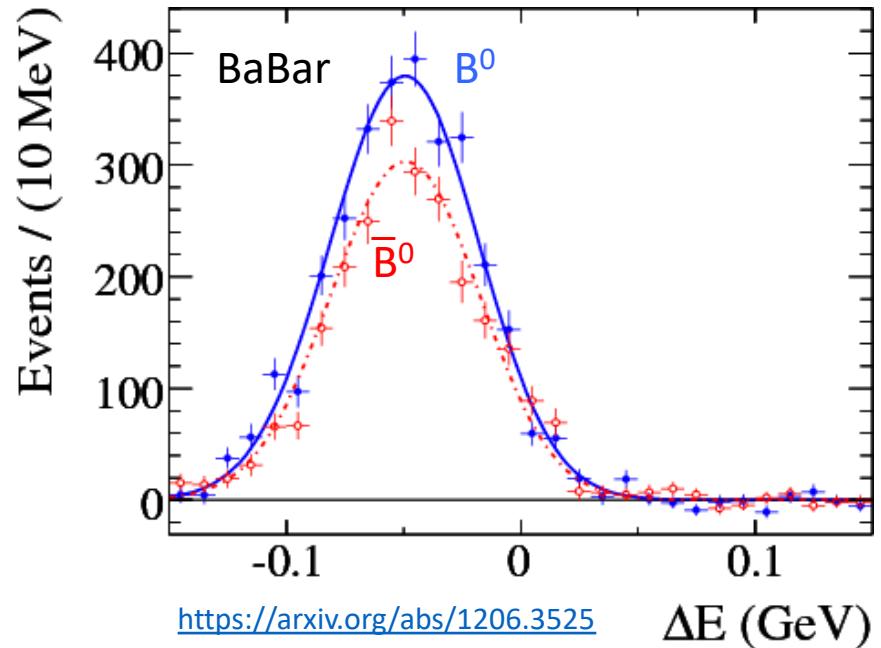
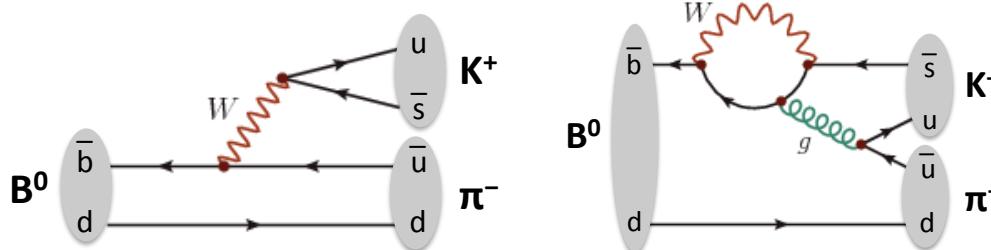
Weak phase between amplitudes =  $\gamma$



# CPV in B decay: “Kπ problem”

## CPV in decay

Example:  $B^0 \rightarrow K\pi$



Problem: CPV should be same for corresponding  $B^+$  decay  
(only difference is spectator quark), but...

$$A_{CP}(B^0 \rightarrow K^+ \pi^-) = (-8.4 \pm 0.4)\%$$

$$A_{CP}(B^+ \rightarrow K^+ \pi^0) = (+4.0 \pm 2.1)\%$$

Sign of new physics?  
Could be from subtle QCD effects...

Averages  
from HF-LAV

# CPV in B decay: “K $\pi$ problem”

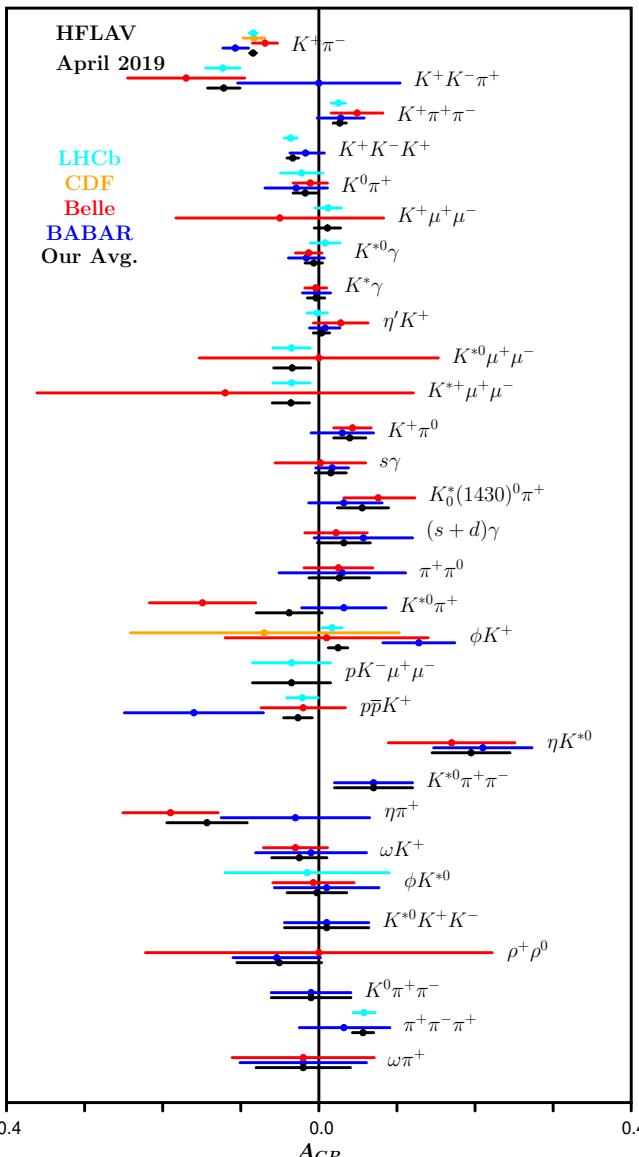
## CPV in decay

Only small sample of results  
(for  $B^0$  and  $B^+$ )

A lot of measurements of a lot of modes!  
Most consistent with CP symmetry

⇒ Remember: only 1 CP violating phase in SM!

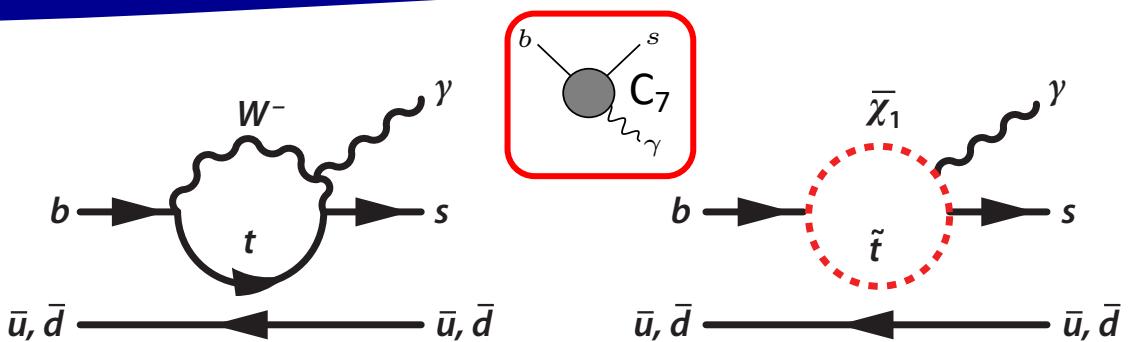
$A_{CP}$  of Most Precisely Measured Modes



<https://hflav-eos.web.cern.ch/hflav-eos/rare/April2019/ACP/index.html>

# Radiative penguins: $b \rightarrow s\gamma$

Radiative penguins give access to new physics, via Wilson Coefficient  $C_7$



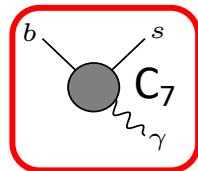
Two approaches:

- *Exclusive*, to specific final state (e.g.  $B^0 \rightarrow K^{0*}\gamma$ )  
⇒ experimentally easier, theoretically messier
- *Inclusive*, including any strange hadrons in final state ( $B^0 \rightarrow X_s\gamma$ )  
⇒ experimentally challenging (one for the B factories!), theoretically clean

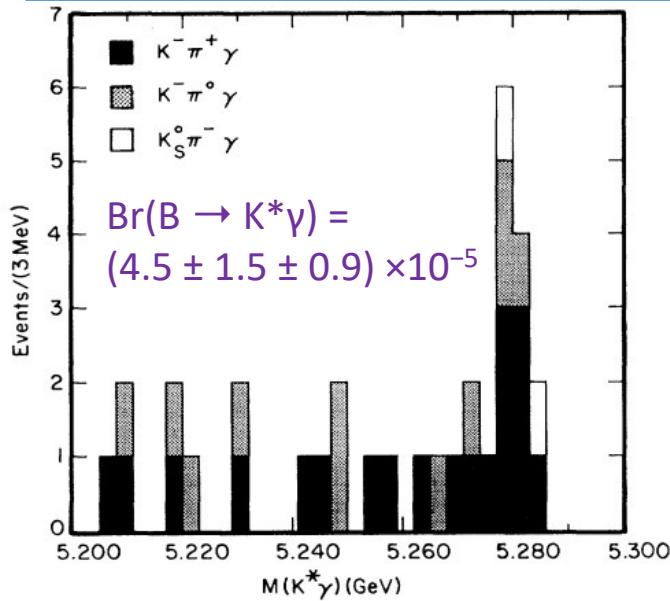
Can also study  $b \rightarrow d\gamma$  decays (e.g.  $B \rightarrow \rho\gamma$ )  
⇒ further suppressed by  $|V_{td}/V_{ts}|^2 \approx (\lambda^3/\lambda^2)^2 \approx 0.05$

# Radiative penguins: Exclusive

$B^{(0,\pm)} \rightarrow K^{(0,\pm)*}\gamma$  discovered by CLEO-II in 1993



<https://doi.org/10.1103/physrevlett.71.674>

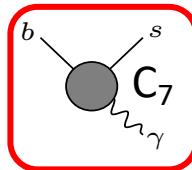
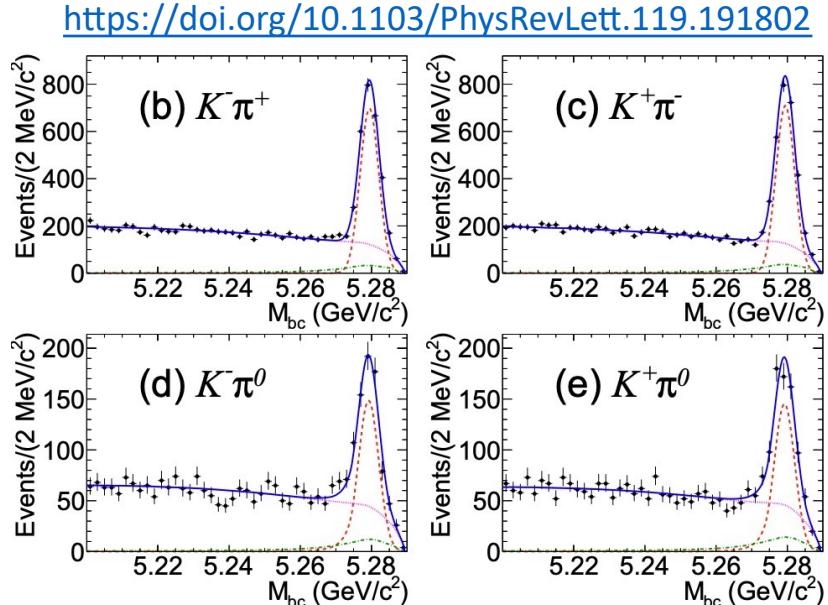


# Radiative penguins: Exclusive

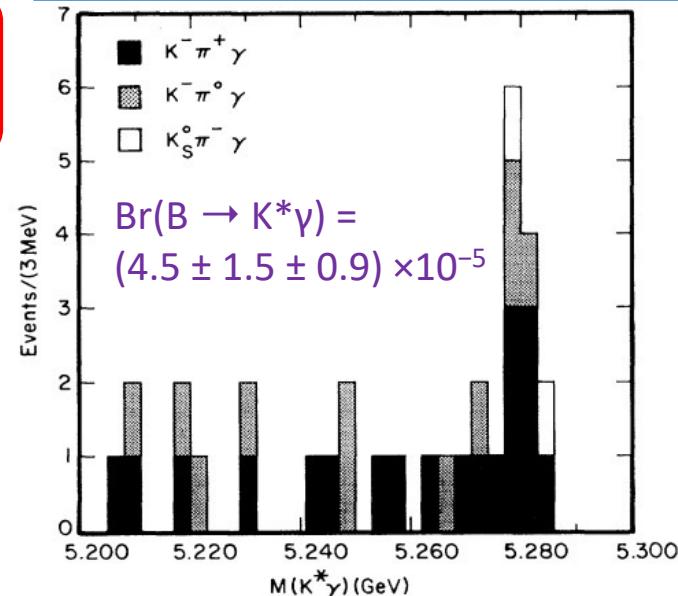
$B^{(0,\pm)} \rightarrow K^{(0,\pm)*}\gamma$  discovered by CLEO-II in 1993

The B-factories later made high-precision measurements

e.g. Belle (2017):



<https://doi.org/10.1103/physrevlett.71.674>



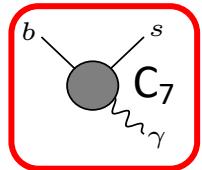
Current world-averages:

$$Br(B^0 \rightarrow K^0\gamma) = (4.18 \pm 0.25) \times 10^{-5}$$

$$Br(B^\pm \rightarrow K^\pm\gamma) = (3.92 \pm 0.22) \times 10^{-5}$$

Theory predictions in range  $3.5 - 7.0 \times 10^{-5}$   
Easier to calculate CP or isospin asymmetries

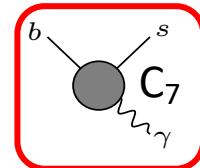
# Radiative penguins: Inclusive



Much trickier experimentally! Final state is **photon** ( $\gamma$ ) + **anything strange** ( $X_s$ )

No perfect method – several options (fully-inclusive, semi-inclusive, summed exclusive, ...) each with pros and cons.

# Radiative penguins: Inclusive

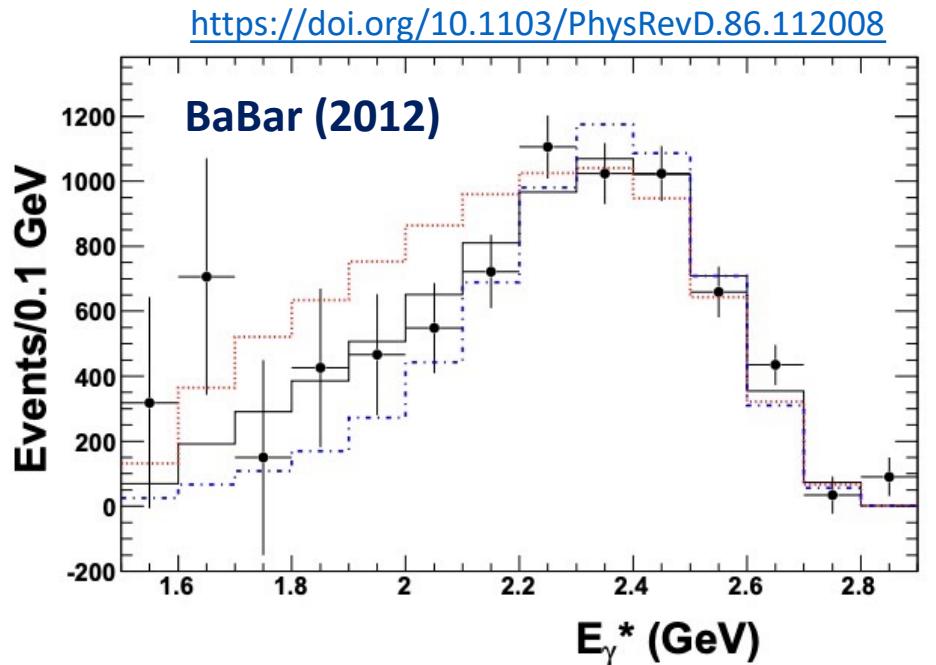


Much trickier experimentally! Final state is **photon** ( $\gamma$ ) + **anything strange** ( $X_s$ )

No perfect method – several options (fully-inclusive, semi-inclusive, summed exclusive, ...) each with pros and cons.

Need clean environment of B-factory:

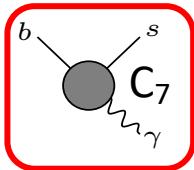
- Look for (relatively) high energy photon
- Identify  $X_s$  topology using multivariate tools trained on simulation
- Plot **photon energy spectrum** after subtracting backgrounds
- Input from theory to convert to BR measurement



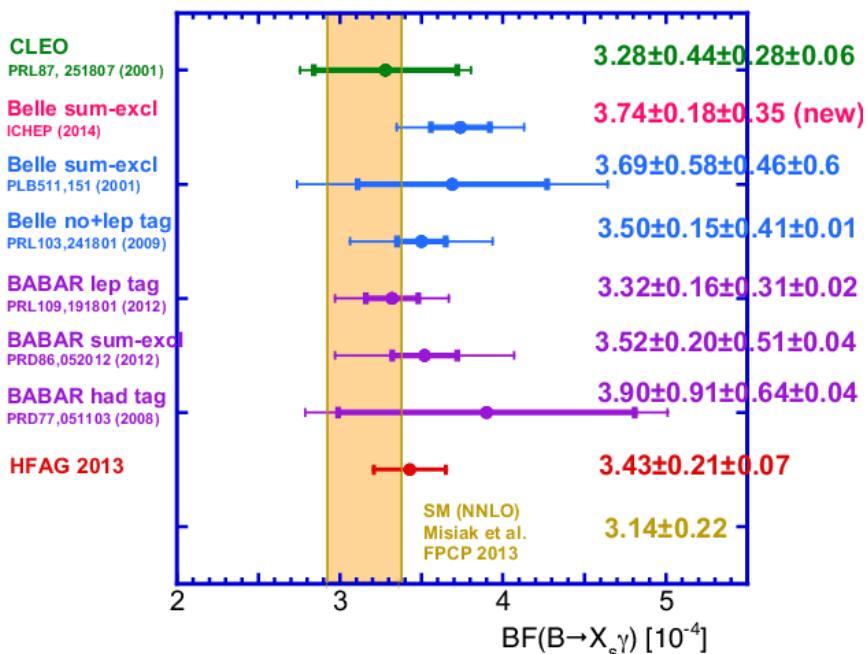
Data (markers) versus three theory models (histograms)

# Radiative penguins: Inclusive

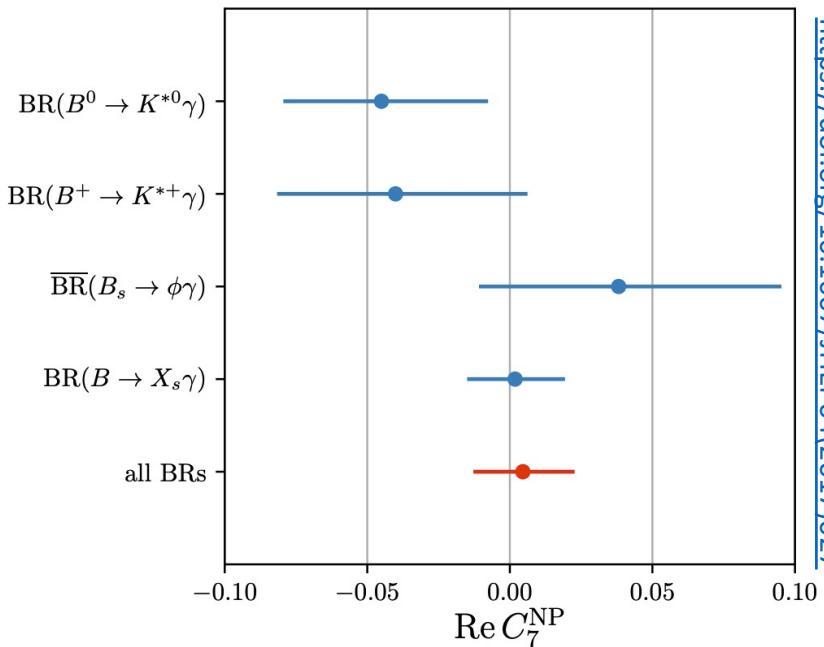
Can combine results from all experiments and use to constrain  $C_7$



*Measurements in good agreement*



*Constraints on new physics in  $C_7$*



$$\text{Latest exp.: } \text{Br}(B \rightarrow X_s \gamma) = (3.32 \pm 0.15) \times 10^{-4}$$

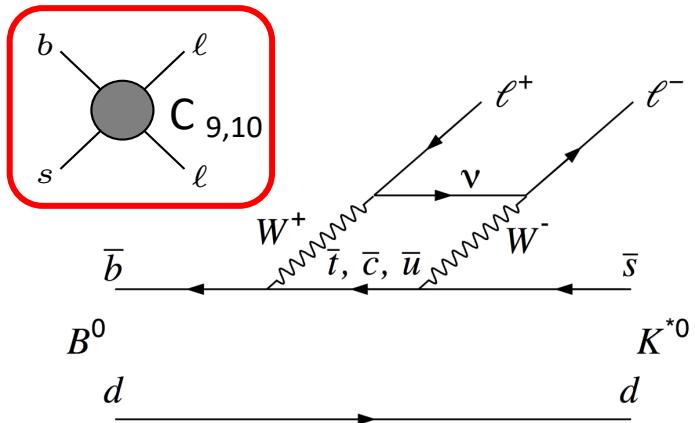
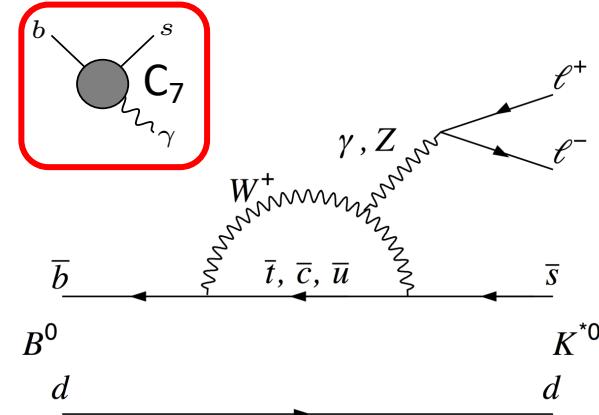
$$\text{Latest SM: } \text{Br}(B \rightarrow X_s \gamma) = (3.36 \pm 0.23) \times 10^{-4}$$

# Electroweak penguins: $b \rightarrow sl^+l^-$

Experimentally easier: two charged leptons + hadrons

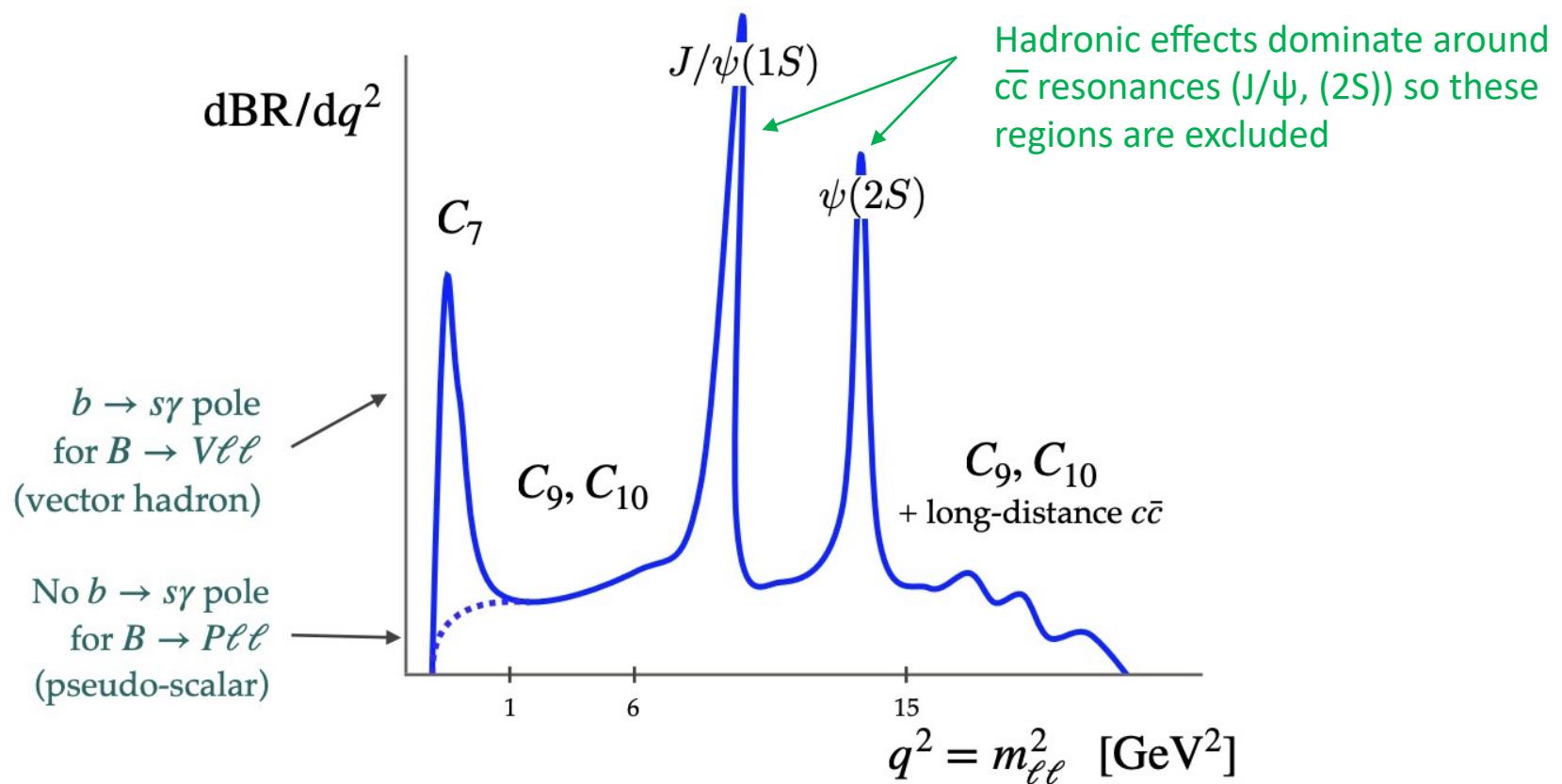
'Golden mode' here is  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$

Much richer decay structure – many observables  
which can constrain new physics models...



# Electroweak penguins: $b \rightarrow sl^+l^-$

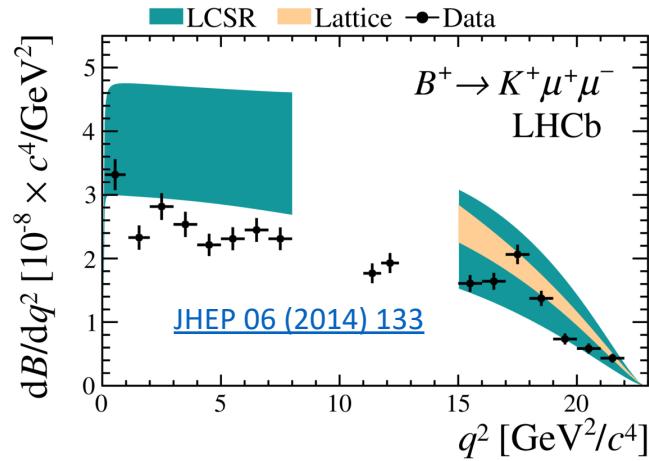
Decay dynamics depend strongly on dimuon mass (= momentum transfer,  $q^2$ )



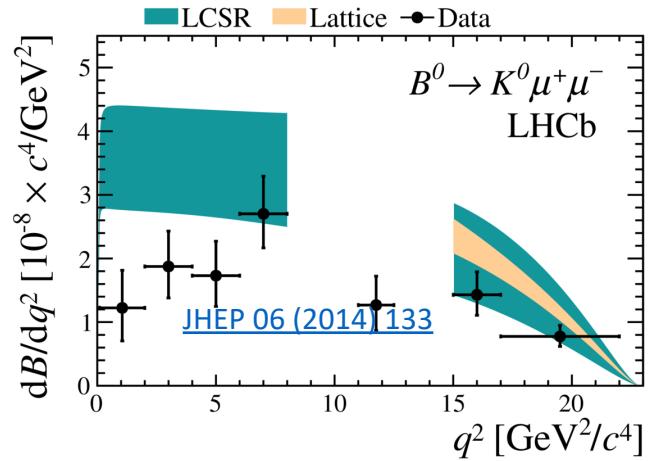
# Electroweak penguins: $b \rightarrow sl^+l^-$

Branching ratios for different  $b \rightarrow s\mu^+\mu^-$  channels tend to undershoot SM prediction

$B^+ \rightarrow K^+\mu^+\mu^-$

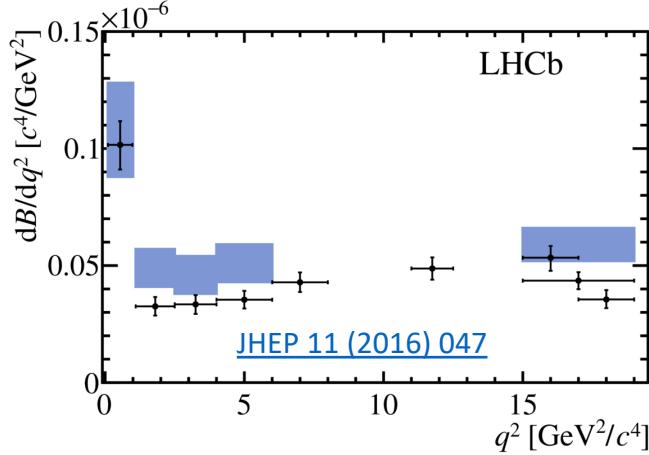


$B^0 \rightarrow K^0\mu^+\mu^-$

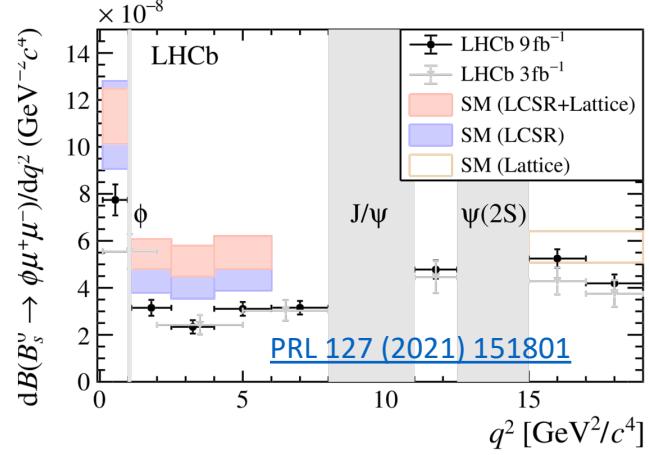


$B^0 \rightarrow K^0\mu^+\mu^-$

$B^0 \rightarrow K^0*\mu^+\mu^-$



$B_s^0 \rightarrow \phi\mu^+\mu^-$

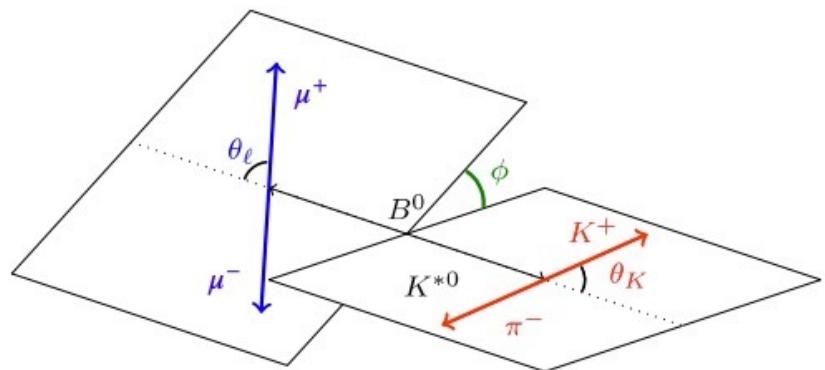


# Electroweak penguins: $b \rightarrow sl^+l^-$

For  $B \rightarrow K^*\mu\mu$ , can do more:

$\Rightarrow K^*$  is a vector meson

$\Rightarrow$  Decay rate depends on three angles



$$\frac{1}{d(\Gamma + \bar{\Gamma})/dq^2} \left. \frac{d^4(\Gamma + \bar{\Gamma})}{dq^2 d\vec{\Omega}} \right|_P = \frac{9}{32\pi} \left[ \frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K \right.$$

$$+ \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell$$

$S_1 \equiv F_L$  (longitudinal polarization fraction of  $K^*$ )

$$- F_L \cos^2 \theta_K \cos 2\theta_\ell + S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi$$

$S_6 \propto A_{FB}$  (forward-backward asymmetry of dimuon system)

$$+ S_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi + S_5 \sin 2\theta_K \sin \theta_\ell \cos \phi$$

$$+ \frac{4}{3} A_{FB} \sin^2 \theta_K \cos \theta_\ell + S_7 \sin 2\theta_K \sin \theta_\ell \sin \phi$$

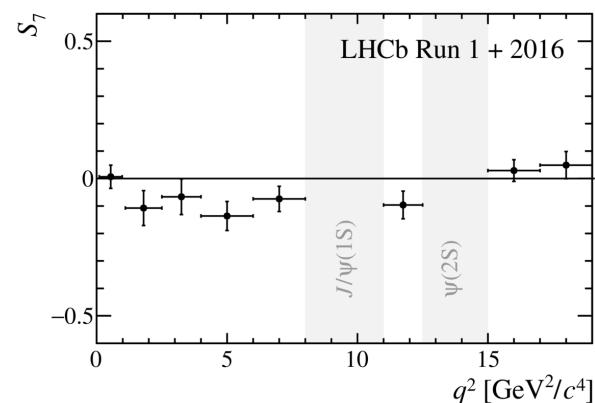
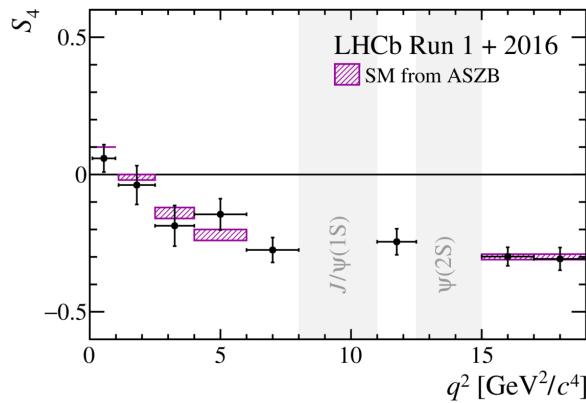
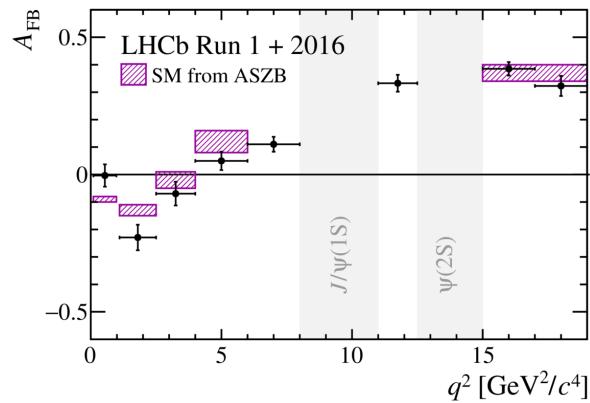
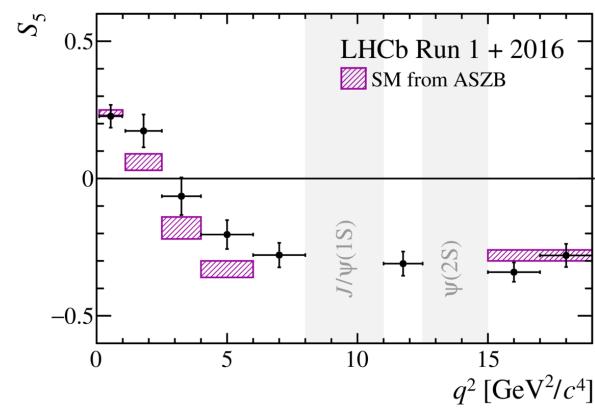
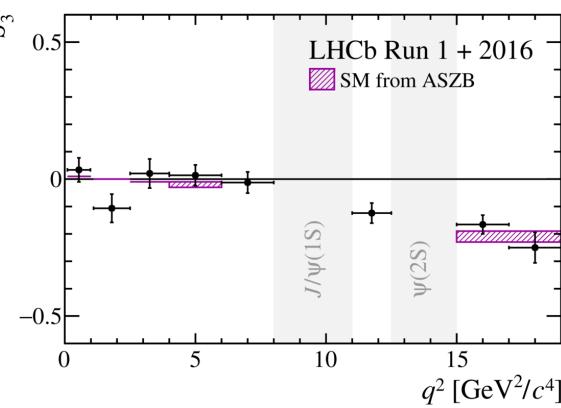
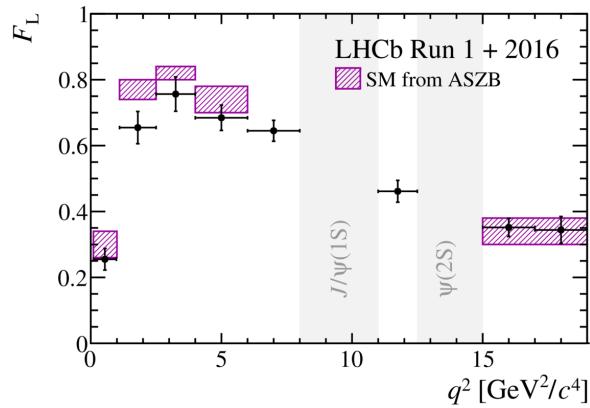
$$+ S_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi + S_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi \Big]$$

Measure the observables:  
 $F_L$ ,  $A_{FB}$ ,  $S_i$

+ also examine CP asymmetries (where  $S_i \rightarrow A_i$ )

# Electroweak penguins: $b \rightarrow sl^+l^-$

At first glance, most observables in agreement with SM ...  
... but hampered by large hadronic uncertainties in SM predictions



[PRL 125 \(2020\) 011802](#)

# Electroweak penguins: $b \rightarrow sl^+l^-$

Build new observables designed to ensure cancellation of major SM uncertainties

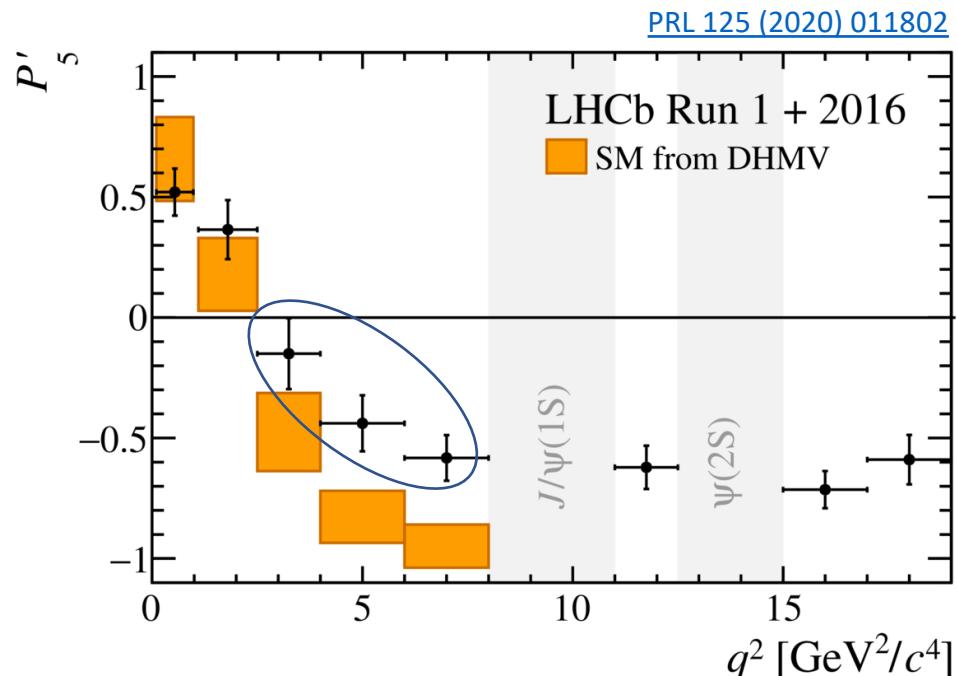
$$P_1 = \frac{2S_3}{(1 - F_L)} = A_T^{(2)},$$

$$P_2 = \frac{2}{3} \frac{A_{FB}}{(1 - F_L)},$$

$$P_3 = \frac{-S_9}{(1 - F_L)},$$

$$P'_{4,5,8} = \frac{S_{4,5,8}}{\sqrt{F_L(1 - F_L)}},$$

$$P'_6 = \frac{S_7}{\sqrt{F_L(1 - F_L)}}.$$



Largest discrepancy in  $P'_5$  variable, at low  $q^2$

Prompted a lot of interest from theoretical community

Crucial to check with other experiments...

# Electroweak penguins: $b \rightarrow sl^+l^-$

Seems to be a consistent disagreement between experiment and theory, but need more precision

Full Run 1+2 data yet to be analysed

