# Theoretical Quantum Optics

# Problem Sheet

Lecturer: Prof. Dr. Igor Lesanovsky

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**Sheet:** 9 **Hand-out:** 7.12.23 **Hand-in:** 14.12.23

### Problem 22. Photon number conservation by beam splitters

The action of a beam splitter on two modes with annihilation operators a and b can be described by the unitary operator:

$$U(\theta) = e^{i\theta(a^{\dagger}b + ab^{\dagger})}. (22.1)$$

Show that the beam splitter conserves the total photon number  $n \equiv a^{\dagger}a + b^{\dagger}b$ .

## Problem 23. Beam splitter and the 2-mode Wigner function

Show that under the transformation

$$\begin{pmatrix} c \\ d \end{pmatrix} = U \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} \cos \theta & e^{i\varphi} \sin \theta \\ -\sin \theta e^{-i\varphi} & \cos \theta \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix}, \tag{23.1}$$

the P, Q and W functions transform as follows:

$$F_{\text{out}}(\alpha, \beta) = F_{\text{in}}(\alpha', \beta'), \quad \text{with} \quad \begin{pmatrix} \alpha' \\ \beta' \end{pmatrix} = U^{-1} \begin{pmatrix} \alpha \\ \beta \end{pmatrix},$$
 (23.2)

where F stands for any of the functions P, Q, and W.

Hint: Show that incoming coherent states  $|\alpha'\beta'\rangle = D_a(\alpha')D_b(\beta')|00\rangle_{ab}$  are transformed into outgoing coherent states  $|\alpha\beta\rangle_{cd}$  with labels linked by equation (23.2).

#### Problem 24. Phase shifter

A common optical element is a phase shifter. It simply means an operation that corresponds classically to a phase shift of the complex amplitude E of the electric field,  $E \to Ee^{-i\varphi}$ . In quantum mechanics a phase shifter is described by the unitary transformation  $U(\varphi) \equiv e^{-ia^{\dagger}a\varphi}$ .

a. Show that  $U(\varphi)$  transforms the annihilation operators as

$$a \to U(\varphi)aU^{\dagger}(\varphi) = ae^{i\varphi}.$$
 (24.1)

**b.** Show that a coherent state is transformed as

$$|\alpha\rangle \to U(\varphi) |\alpha\rangle = |\alpha e^{-i\varphi}\rangle$$
 (24.2)

**c.** Show the action of the phase shifter on a squeezed vacuum state,  $|\xi,0\rangle = S(\xi)|0\rangle$  with

$$S(\xi) = \exp\left(-\frac{\xi}{2}(a^{\dagger})^2 + \frac{\xi^*}{2}a^2\right),$$
 (24.3)

$$|\xi,0\rangle \to U(\varphi)|\xi,0\rangle = |\xi e^{-2i\varphi},0\rangle.$$
 (24.4)

### Problem 25. Photon-subtracted squeezed vacuum state

The photon-subtracted squeezed vacuum state is given by

$$|\psi\rangle = \frac{1}{N_{\psi}} aS(\xi) |0\rangle = S(\xi) |1\rangle. \tag{25.1}$$

- a. Show the second equality in (25.1) holds up to a phase factor. You need not determine the normalization factor  $N_{\psi}$ .
- **b.** Calculate the Wigner function for  $\tilde{\rho} \equiv \rho_1 = |1\rangle \langle 1|$ . Hint: Start with the Wigner function of the vacuum state  $W_0(\alpha, \alpha^*) = \frac{2}{\pi} e^{-2|\alpha|^2}$  and apply the second half of exercise 15.

Using the transformation identities from Sheet 2 Excercise 4

$$S(\xi)aS^{\dagger}(\xi) = a\cosh(r) - a^{\dagger}e^{i\theta}\sinh(r) \equiv a(\xi)$$
  

$$S(\xi)a^{\dagger}S^{\dagger}(\xi) = a^{\dagger}\cosh(r) - ae^{-i\theta}\sinh(r) \equiv a^{\dagger}(\xi),$$
(25.2)

where  $\xi = re^{i\theta}$ , we have

$$\rho(a, a^{\dagger}) \equiv S(\xi)\tilde{\rho}(a, a^{\dagger})S^{\dagger}(\xi) = \tilde{\rho}(a(\xi), a^{\dagger}(\xi))$$
(25.3)

for any function  $\tilde{\rho}(a, a^{\dagger})$  that can be expanded as a power series. The linearity property of the Wigner function then implies

$$W_{\rho}(\alpha, \alpha^*) = W_{\tilde{\rho}}(\alpha(\xi), \alpha^*(\xi)), \tag{25.4}$$

where  $\alpha(\xi) = \alpha \cosh(r) - \alpha^* e^{i\theta} \sinh(r)$  and  $\alpha^*(\xi) = \alpha^* \cosh(r) - \alpha e^{-i\theta} \sinh(r)$ , and  $W_{\rho}(W_{\tilde{\rho}})$  is the Wigner function for the state  $\rho$  ( $\tilde{\rho}$ ), respectively.

- c. Determine the Wigner function for the photon-subtracted squeezed vacuum state (25.1) using your result from **b**. and equation (25.4). For what values of  $|\alpha(\xi)|^2$  does the Wigner function become negative?
- d. Plot the Mandel Q-parameter given in exercise 19 by using the results

$$\langle \psi | (a^{\dagger})^{2} a^{2} | \psi \rangle = \frac{3t^{2} (3 + 2t^{2})}{(1 - t^{2})^{2}}$$

$$\langle \psi | a^{\dagger} a | \psi \rangle = \frac{1 + 2t^{2}}{1 - t^{2}},$$
(25.5)

with  $t = \tanh(r) \in (0,1)$ . For approximately what values of r is  $Q_M$  negative? Compare your results to those from  $\mathbf{c}$ .