

1. Übungsblatt

Aufgabe 1: Single-mode case

The components E_x and B_y of the quantized electric field and magnetic induction of a single standing mode in a resonator in vacuum are given by

$$E_x(z, t) = E_0(ae^{-i\omega t} + a^\dagger e^{i\omega t}) \sin(kz) \quad (1)$$

$$B_y(z, t) = B_0 \frac{1}{i} (ae^{-i\omega t} - a^\dagger e^{i\omega t}) \cos(kz), \quad (2)$$

where $E_0 = \sqrt{\frac{\hbar\omega}{\epsilon_0 V}}$, and $B_0 = E_0/c$, are the field amplitudes for a single photon, $k = \pi n/L$ for $n \in \mathbb{N}$ is the wave number of the mode, L the length of the resonator, V its volume, and c the speed of light. In vacuum, $\mathbf{B} = \mu_0 \mathbf{H}$, where \mathbf{H} is the magnetic field. The creation and annihilation operators a and a^\dagger satisfy the usual commutation relations known from the quantized harmonic oscillator: $[a, a^\dagger] = 1$.

- Calculate the commutator $[E_x(z, t), B_y(z', t')]$.
- Deduce the Heisenberg uncertainty relation between $E_x(z, t)$, and $B_y(z', t')$.
- What are the combinations of times or positions at which E_x and B_y can both have sharp values?
- Calculate the fields at those times or positions and find an explanation for the results in (c).

Aufgabe 2: Multi-mode case for equal times

In the multi-mode case, the full electric and magnetic fields are decomposed into plane waves,

$$\mathbf{E}(\mathbf{r}, t) = \sum_{\mathbf{k}, \lambda} \epsilon_{\mathbf{k}}^{(\lambda)} \mathcal{E}_{\mathbf{k}} a_{\mathbf{k}, \lambda} e^{-i\omega_{\mathbf{k}} t + i\mathbf{k} \cdot \mathbf{r}} + \text{H.c.} \quad (3)$$

and

$$\mathbf{H}(\mathbf{r}, t) = \frac{1}{\mu_0} \sum_{\mathbf{k}, \lambda} \frac{\mathbf{k} \times \epsilon_{\mathbf{k}}^{(\lambda)}}{\omega_{\mathbf{k}}} \mathcal{E}_{\mathbf{k}} a_{\mathbf{k}, \lambda} e^{-i\omega_{\mathbf{k}} t + i\mathbf{k} \cdot \mathbf{r}} + \text{H.c.}, \quad (4)$$

where $\mathcal{E}_{\mathbf{k}} = \sqrt{\frac{\hbar\omega_{\mathbf{k}}}{2\epsilon_0 V}}$ has the dimension of an electric field and $\epsilon_{\mathbf{k}}^{(\lambda)}$ is a unit polarization vector. Here, $\lambda \in \{1, 2\}$ labels the two polarization directions perpendicular to the wavevector \mathbf{k} , whose components satisfy $k_i = 2\pi n_i/L$ with $n_i \in \mathbb{Z}$, and $\omega_{\mathbf{k}} = c|\mathbf{k}|$. The commutation relations between the operators $a_{\mathbf{k}, \lambda}$ and $a_{\mathbf{k}', \lambda'}^\dagger$ are

$$[a_{\mathbf{k}, \lambda}, a_{\mathbf{k}', \lambda'}^\dagger] = [a_{\mathbf{k}, \lambda}^\dagger, a_{\mathbf{k}', \lambda'}] = 0, \quad (5)$$

$$[a_{\mathbf{k},\lambda}, a_{\mathbf{k}',\lambda'}^\dagger] = \delta_{\mathbf{k},\mathbf{k}'} \delta_{\lambda,\lambda'}, \quad (6)$$

indicating that each mode, labeled by \mathbf{k} and λ , can be seen as an independent harmonic oscillator.

a) Verify that the polarization vectors satisfy

$$\epsilon_{\mathbf{k}i}^{(1)} \epsilon_{\mathbf{k}j}^{(1)} + \epsilon_{\mathbf{k}i}^{(2)} \epsilon_{\mathbf{k}j}^{(2)} = \delta_{ij} - \frac{k_i k_j}{k^2}, \quad (7)$$

where i, j represent the Cartesian components and $k^2 = |\mathbf{k}|^2$.

Hint: You may want to write the Cartesian components of the wavevector in spherical coordinates,

$$\mathbf{k} = k(\sin(\theta) \cos(\phi), \sin(\theta) \sin(\phi), \cos(\theta)), \quad (8)$$

and do so as well for the polarization vectors.

b) Show that

$$[E_x(\mathbf{r}, t), H_y(\mathbf{r}', t)] = \frac{\hbar c^2}{2V} \sum_{\mathbf{k}} k_z \left[e^{i\mathbf{k} \cdot (\mathbf{r} - \mathbf{r}')} - e^{-i\mathbf{k} \cdot (\mathbf{r} - \mathbf{r}')} \right] \quad (9)$$

Hint: Write down the field components in the commutator explicitly and use (7).

c) Replace

$$\sum_{\mathbf{k}} \rightarrow \frac{V}{(2\pi)^3} \int d^3 k \quad (10)$$

in equation (9) and consider the limit $V \rightarrow \infty$, to show that in the continuous case

$$[E_x(\mathbf{r}, t), H_y(\mathbf{r}', t)] = -i\hbar c^2 \frac{\partial}{\partial z} \delta^{(3)}(\mathbf{r} - \mathbf{r}'), \quad (11)$$

where

$$\delta^{(3)}(\mathbf{x}) = \frac{1}{(2\pi)^3} \int_{\mathbb{R}^3} e^{i\mathbf{k} \cdot \mathbf{x}} d^3 k \quad (12)$$

is the three-dimensional Dirac-delta function.

d) Generalise (11) to arbitrary $j, k \in \{x, y, z\}$:

$$[E_j(\mathbf{r}, t), H_k(\mathbf{r}', t)] = ? \quad (13)$$

e) Compare your results to the one mode case.