

# Exercise 1

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## 1 Fundamentals of Differential Equations

### 1.a Difference between ordinary derivative, partial derivative, and material (total) derivative

- **Ordinary derivative** ( $\frac{d}{dt}$ ): Describes the rate of change of a function with respect to one variable. It is used for functions depending on a single variable, such as  $f(t)$ .
- **Partial derivative** ( $\frac{\partial}{\partial t}$ ): Describes the rate of change of a multivariable function with respect to one of its variables, while holding other variables constant. This is often used in multivariable functions such as  $f(x, t)$ , where we can find  $\frac{\partial f}{\partial t}$  while  $x$  remains fixed.
- **Material (total) derivative** ( $\frac{D}{Dt}$ ): is a measure of the rate of change of a physical quantity (like velocity or temperature) experienced by an observer moving with the fluid. It combines both local and convective rates of change as, for example, in a function  $f(x, t)$ ,  $\frac{D}{Dt} = \frac{\partial}{\partial t} + u \frac{\partial}{\partial x}$  for some velocity field  $u$ .

### 1.b Ordinary and partial differential equations

- **Ordinary Differential Equations (ODEs)**: These involve derivatives with respect to a single variable. For example,  $\frac{dy}{dt} = y$  is an ODE.
- **Partial Differential Equations (PDEs)**: These involve partial derivatives with respect to multiple variables. For instance, the heat equation  $\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2}$  is a PDE.

### 1.c Order of a differential equation

The order of a differential equation is the highest order of derivative present in the equation.

- **First-order ODE**:  $\frac{dy}{dt} = ky$ .
- **Second-order PDE**: The wave equation  $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$ .
- **Third-order ODE**:  $\frac{d^3 y}{dt^3} + \frac{d^2 y}{dt^2} + y = 0$ .

### 1.d Linear and non-linear differential equations

- **Linear Differential Equations**: These have terms that are linear in the unknown function and its derivatives. For example,  $\frac{dy}{dt} + 3y = 0$  is linear.
- **Non-linear Differential Equations**: These have terms that are non-linear in the unknown function or its derivatives. For instance, the Navier-Stokes equation  $\rho \left( \frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = -\nabla p + \mu \nabla^2 \mathbf{u} + \mathbf{f}$  is non-linear. This non-linearity arises due to the convective term  $\mathbf{u} \cdot \nabla \mathbf{u}$ , which represents the interaction of the velocity field with itself. Specifically,  $\mathbf{u} \cdot \nabla \mathbf{u}$  is non-linear because it involves the product of the velocity field  $\mathbf{u}$  with its own gradient  $\nabla \mathbf{u}$ .

### 1.e Initial value problem (IVP) and boundary value problem (BVP)

- **Initial Value Problem (IVP)**: A problem that requires solving a differential equation with specified initial conditions, such as  $y(0) = y_0$ , in time.
- **Boundary Value Problem (BVP)**: A problem where the solution to a differential equation is sought within a specified range, with conditions, usually Dirichlet or Neumann, given at the boundaries of the range, like  $u(0) = 0$  and  $u(1) = 1$ .

### 1.f Parabolic and elliptic PDE examples and their conditions

The difference between parabolic and elliptic PDEs can be defined through the computation of a discriminant  $\Delta = b^2 - 4ac$ , where  $a$ ,  $b$ , and  $c$  are coefficients from the second-order PDE of the form  $a \frac{\partial^2 u}{\partial x^2} + b \frac{\partial^2 u}{\partial x \partial y} + c \frac{\partial^2 u}{\partial y^2} + \dots = 0$ . If  $\Delta = 0$ , the PDE is parabolic, and if  $\Delta < 0$ , the PDE is elliptic.

- **Parabolic PDE**: The heat equation  $\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2}$  is parabolic and typically requires both initial and boundary conditions.
- **Elliptic PDE**: Laplace's equation  $\nabla^2 u = 0$  is elliptic and usually requires boundary conditions but not initial conditions, since it does not depend on time.

## 2 Order Reduction

The governing equation of the damped oscillator is given by:

$$m \frac{d^2 y(t)}{dt^2} + b \frac{dy(t)}{dt} + cy(t) = 0 \quad (1)$$

with initial conditions:

$$y(0) = s_0, \quad \frac{dy(0)}{dt} = v_0.$$

We aim to transform this second-order ODE into a system of first-order differential equations. To reduce a second-order ODE to a system of first-order ODEs we can introduce new variables to represent the derivatives of the function  $y(t)$ . In particular, we define:

$$y_1(t) = y(t)$$

and introduce a new variable  $y_2(t)$  to represent the first derivative of  $y(t)$ :

$$y_2(t) = \frac{dy(t)}{dt}.$$

Since  $\frac{dy_2(t)}{dt} = \frac{d^2 y(t)}{dt^2}$ , we can substitute this into the original equation to obtain:

$$m \frac{dy_2(t)}{dt} + by_2(t) + cy_1(t) = 0. \quad (2)$$

In this way, we can express the problem as two coupled first-order differential equations:

$$\begin{cases} \frac{dy_1(t)}{dt} &= y_2(t), \\ \frac{dy_2(t)}{dt} &= -\frac{b}{m}y_2(t) - \frac{c}{m}y_1(t). \end{cases}$$

However, we also need to rewrite the initial conditions for  $y(t)$  and  $\frac{dy(t)}{dt}$ :

- $y_1(0) = s_0$ ,
- $y_2(0) = v_0$ .

This approach allows us to solve the system using methods suited for first-order differential equations, enabling easier numerical or analytical analysis.

## 3 Blasius Equation

### Part (a): Convert the Blasius Equation to a System of First-Order ODEs

The Blasius equation is given by:

$$f''' + \frac{1}{2}ff'' = 0 \quad (3)$$

with  $f' = \frac{u}{U_\infty}$ . Three boundary conditions are necessary to solve this equation:

- $\eta = 0$ :  $f' = f = 0$  (no-slip condition)
- $\eta \rightarrow \infty$ :  $f' = 1$  (free outer flow)

We aim to transform this third-order ODE into a system of first-order differential equations. To reduce a third-order ODE to a system of first-order ODEs we can introduce new variables to represent the derivatives of the function  $f(\eta)$ . In particular, we define:

$$y_1 = f, \quad y_2 = f' = \frac{df}{d\eta}, \quad y_3 = f'' = \frac{d^2 f}{d\eta^2}$$

Then, the derivatives of these variables with respect to  $\eta$  are:

$$\frac{dy_1}{d\eta} = y_2, \quad \frac{dy_2}{d\eta} = y_3$$

Now, we can substitute this into the original Blasius equation to obtain:

$$\frac{dy_3}{d\eta} = -\frac{1}{2}y_1y_3$$

In this way, we can express the problem as three coupled first-order differential equations:

$$\begin{cases} \frac{dy_1}{d\eta} = y_2 \\ \frac{dy_2}{d\eta} = y_3 \\ \frac{dy_3}{d\eta} = -\frac{1}{2}y_1y_3 \end{cases}$$

with boundary conditions:

- At  $\eta = 0$ :  $y_1 = 0$ ,  $y_2 = 0$
- As  $\eta \rightarrow \infty$ :  $y_2 = 1$

#### Part (b): Providing an Initial Condition for $f''(0)$

To solve this problem as an initial value problem, we need an initial value for  $f''(0)$ . However, the boundary condition  $y_2(\infty) = 1$  is specified at infinity, making it impractical to impose this condition directly at a finite point. To address this, we can use an iterative approach:

1. Guess an initial value for  $f''(0)$ .
2. Integrate the system of equations from  $\eta = 0$  to a sufficiently large value of  $\eta$  where  $y_2(\eta)$  approaches a constant. To solve this system numerically, we use the Runge-Kutta method of fourth order (RK4) that allows to approximate solutions to ordinary differential equations.

For a step size  $h$ , the RK4 method computes the next values  $y_{i+1}$  as follows:

$$\begin{aligned} k_1 &= h \cdot f(\eta_i, y_i) \\ k_2 &= h \cdot f\left(\eta_i + \frac{h}{2}, y_i + \frac{k_1}{2}\right) \\ k_3 &= h \cdot f\left(\eta_i + \frac{h}{2}, y_i + \frac{k_2}{2}\right) \\ k_4 &= h \cdot f(\eta_i + h, y_i + k_3) \end{aligned}$$

The next value of the solution is updated by:

$$y_{i+1} = y_i + \frac{1}{6}(k_1 + 2k_2 + 2k_3 + k_4)$$

3. Check if  $y_2(\eta)$  approaches 1 as  $\eta \rightarrow \infty$ . If  $y_2(\eta)$  is not close to 1, adjust the initial guess for  $y_3(0)$  iterate this process until the condition  $y_2(\infty) = 1$  (or close to it) is satisfied within a desired tolerance.

This iterative approach allows us to find an appropriate initial condition for  $y_3(0) = f''(0)$  that satisfies the boundary condition at infinity.

#### References

- [1] *CFD Repository*,  
Available at: <https://github.com/GiuseppePisante/CFD.git>
- [2] *GitHub Copilot*,  
GitHub. Available at: <https://github.com/features/copilot>