

Modeling and Control of a Single Joint

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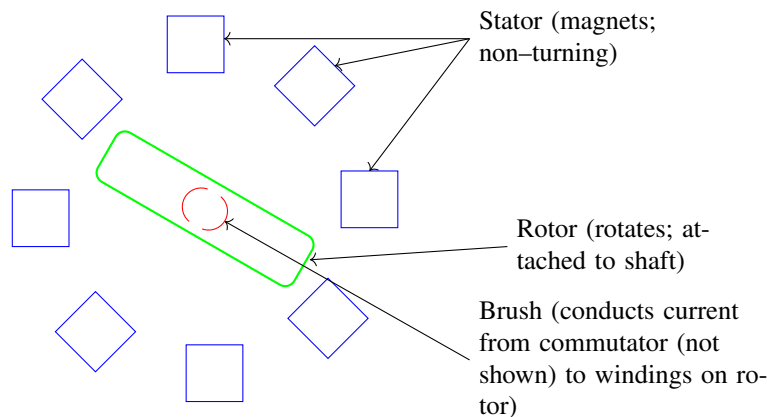
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1 Intro

We're going to talk about how to model (and, maybe, control) the dynamics of a single joint at the actuator level. This will include the dynamics of an electric motor, any gear-based transmission, joint flexibility, and friction. We're going to skip the modeling of linear actuators (electric or otherwise) because in practice rotational motion via electric motors covers 99% of what we see in practice. Note that some robotic arms are hydraulic, specifically arms designed to work in the subsea; the modeling and control of hydraulic actuators is its own thing that we could easily spend an entire class (or at least a couple of lectures) talking about. For more info on hydraulic actuation and control, see the excellent textbook *Hydraulic Control Systems* by Noah Manning and Roger Fales (John Wiley & Sons 2019). Pneumatic actuator control, which is also sometimes used, is even more complex than hydraulic control due to the compressibility of air (or whatever other gas you might consider using). As a consequence, accurate high bandwidth control of pneumatics is very much an open research problem.

2 Electric Motor

Here's a cross section of a simple brushed DC motor:



Typically, the torque produced by a motor like this varies sinusoidally as the angle of the rotor moves with respect to the fixed stator magnets; this is called “torque ripple”. For a decent motor the torque ripple will be perhaps 5% of the average torque. The physics that allow a motor to generate a torque are described by:

$$F = qV \times B$$

where a charge q moving through a magnetic field B with velocity V experiences a force F . The charges are the electrons moving through the windings in the rotor, and the magnetic field is generated by the (usually permanent) magnets in the stator. Usually, the torque of a motor is related to the current flowing through the windings via the “motor torque constant”:

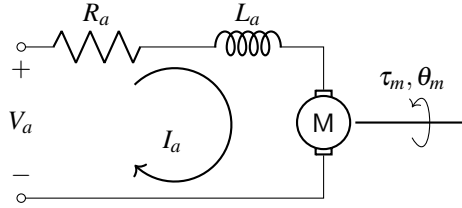
$$\tau_m = K_m I_a \quad (1)$$

As we might recall from a few lectures ago, the torque constant isn’t actually a constant – so this is very definitely an approximation.

When a motor is moving, it acts as a generator and it develops a voltage across its armature. A second motor constant called the “back emf constant” describes the generated voltage as a function of rotational velocity:

$$V = K_e \dot{\theta}_M \quad (2)$$

Here is the electrical circuit of a motor:



where L_a is the armature inductance, R_a is the armature resistance, and V_a is the voltage. The dynamics of this circuit are:

$$L_a \dot{I}_a + R_a I_a = V_a - K_e \dot{\theta}_M$$

If we assume that the inductance of the circuit is negligible, this ceases to be a differential equation and we can derive a “torque–speed” curve as follows:

$$V_a = R_a I_a + K_e \dot{\theta}_m$$

and we can substitute Equation 1 for i_a and then solve for τ_m to get

$$\tau_M = \frac{K_m}{R_a} (V_a - K_e \dot{\theta}_M)$$

Plotting this, we get:

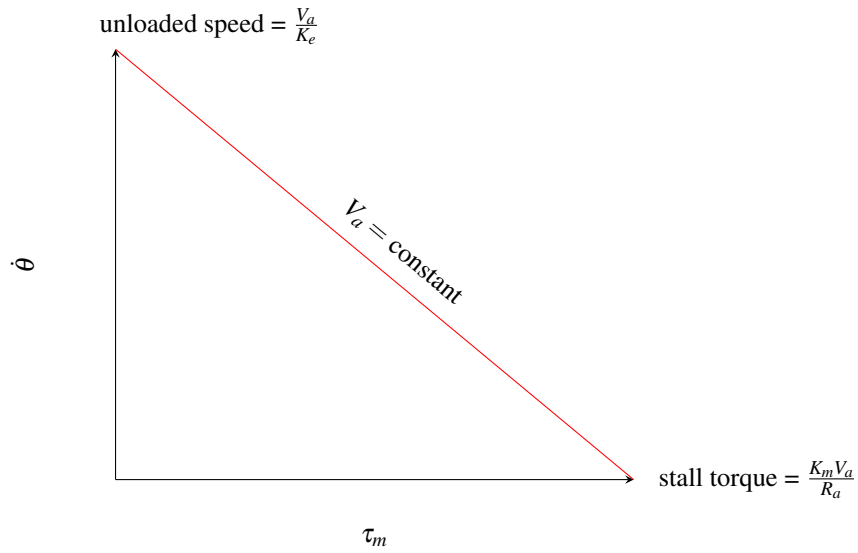
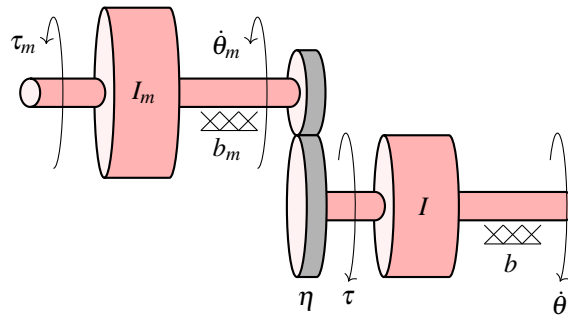


Figure 1: A torque–speed curve

3 Transmission

Motors are (almost) always high speed, low torque devices (i.e. the unloaded speed of a small electric motor will generally be in the range of 500-2000 RPMs). Robots “want” high torque and low speed. Traditionally, we use a transmission to trade speed for torque. Here’s a simplified drawing of a geartrain transmission:



where the gear ratio, η , is greater than 1 (often, considerably greater; gear ratios of 500 aren’t uncommon, and the FRENDA arms used on the RSGS and OSAM-1 missions have gear ratios of 3000-5000 depending on the joint in question). Gear trains amplify

torque and reduce speed:

$$\begin{aligned}\tau &= \eta \tau_m \\ \dot{\theta} &= \frac{1}{\eta} \dot{\theta}_m\end{aligned}\tag{3}$$

Let's figure out the torque balance at the rotor:

$$\begin{aligned}\sum_{\text{rotor}} \tau &= I_m \ddot{\theta}_m \\ I_m \ddot{\theta}_m &= \tau_m - b_m \dot{\theta}_m - \underbrace{\left(\frac{1}{\eta} \right) (I \ddot{\theta} + b \dot{\theta})}_{\substack{\text{inertial viscosity at link} \\ \text{using } \dot{\theta} = \frac{1}{\eta} \dot{\theta}_m}} \\ I_m \ddot{\theta}_m &= \tau_m - b_m \dot{\theta}_m - \frac{1}{\eta} \left(I \frac{1}{\eta} \ddot{\theta}_m + b \frac{1}{\eta} \dot{\theta}_m \right) \\ \tau_m &= \left(I_m + \frac{I}{\eta^2} \right) \ddot{\theta}_m + \left(b_m + \frac{b}{\eta^2} \right) \dot{\theta}_m\end{aligned}$$

We can use equations 3 to write this in terms of “motor variables”:

$$\tau_m = \underbrace{\left(I_m + \frac{I}{\eta^2} \right)}_{\text{“Effective inertia”}} \ddot{\theta}_m + \underbrace{\left(b_m + \frac{b}{\eta^2} \right)}_{\text{“Effective damping”}} \dot{\theta}_m$$

or in terms of “load variables”:

$$\tau = \underbrace{\left(I + \eta^2 I_m \right)}_{\text{“Effective inertia”}} \ddot{\theta} + \underbrace{\left(b + \eta^2 b_m \right)}_{\text{“Effective damping”}} \dot{\theta}$$

Notice that η^2 term — this implies that when the gear ratio is even moderately large, the dynamics of the geartrain can dominate the dynamics of the links.

3.1 Example: Space Systems Lab Ranger Arm

Ranger was a neutral buoyancy prototype for a free-flying robotic satellite servicer built by the SSL in the 1990s. It became a flight program that, tragically, NASA ultimately decided not to fly. The Ranger servicing arms were among the first flight-traceable arms intended for satellite servicing. For the Ranger shoulder actuator, we had

$$\begin{aligned}I_m &= 0.00048 \text{ Kg}\cdot\text{m}^2 \\ I &= 20 \text{ Kg}\cdot\text{m}^2 \\ \eta &= 201\end{aligned}$$

so

$$I_{\text{eff}} = \underbrace{(201)^2(0.00048)}_{19.39} + 20 = 39.39 \text{ Kg}\cdot\text{m}^2$$

4 Flexibility (Unmodeled)

So far the joint model has assumed rigid:

- Gearing
- Shafts
- Bearings
- Links

Of course in reality everything is flexible. In this case gearing tends to have flexible modes of 5–25 Hz, which dominates; everything else typically has modes above 25 Hz. If we added flexibility to the model, it would need to increase the order of the equation:

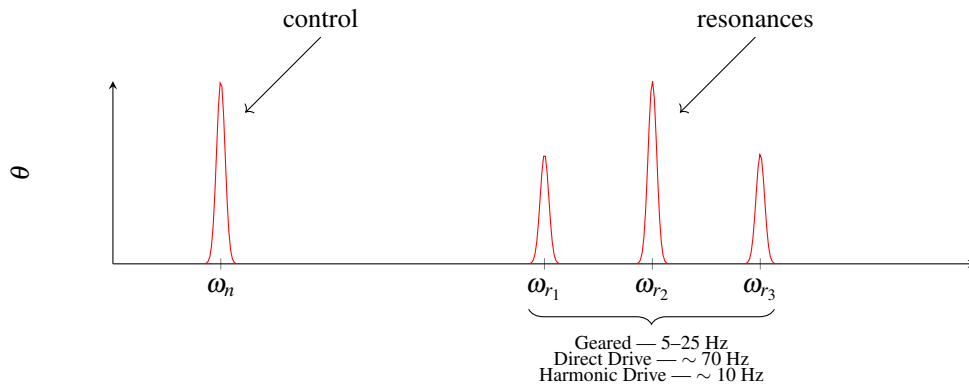
$$\tau_m = I_{\text{eff}}\ddot{\theta} + b_{\text{eff}}\dot{\theta}$$

which is second order.

For high stiffnesses, the natural frequencies of the unmodeled resonances are quite high. In order to not excite these resonances, you should design your control law to follow the closed-loop natural frequency condition:

$$\omega_n \leq \frac{1}{2} \omega_{\text{res}}$$

where ω_{res} is the *lowest* structural resonance frequency.

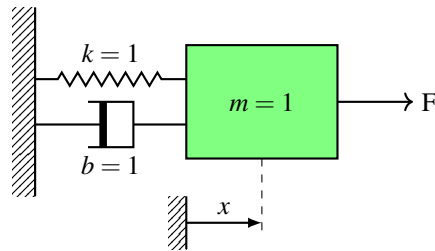


For the Ranger arm,

$$\begin{aligned}
 J_m &= 0.00048 \text{ Kg-m}^2 \\
 J_{\text{link}} &= 0.5 \text{ Kg-m}^2 \\
 \eta &= 201 \\
 \omega_r &= \sqrt{\frac{K_t}{J_{\text{eff}}}} \\
 J_{\text{eff}} &= J_{\text{link}} + \eta^2 J_m = 19.39 + 0.5 \\
 \omega_r &= 52.8 \text{ rad/sec} \\
 \rightarrow f_r &= 8.3 \text{ Hz}
 \end{aligned}$$

4.1 Example 9.7 from the textbook

Here's a spring-mass-damper system:



and we are asked to find a decoupled joint controller such that $\zeta = 1$, $K_{\text{tot}} = \text{MAX}$, and the controller does not excite any unmodelled dynamics (e.g. flexible modes). So the dynamics are

$$\ddot{x} + \dot{x} + x = F$$

The controller will have the form

$$F = \overset{0}{\cancel{\ddot{x}_d}} + (\ddot{x} + x) - K_v \dot{x} - K_p x$$

and the closed loop dynamics are

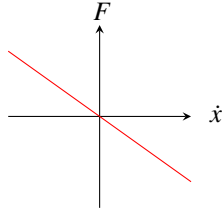
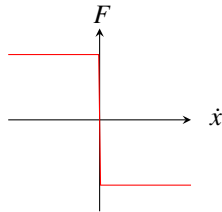
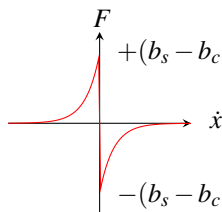
$$\ddot{x} + K_v \dot{x} + K_p x = 0$$

We need to find K_p and K_v :

$$\begin{aligned}
 K_p &= \omega_n^2 = \left(\frac{1}{2} \omega_{\text{res}}\right)^2 = \left(\frac{1}{2} 8\right)^2 = 16 \\
 K_v &= 2\zeta \omega_n = 2(1)(4) = 8
 \end{aligned}$$

5 Friction

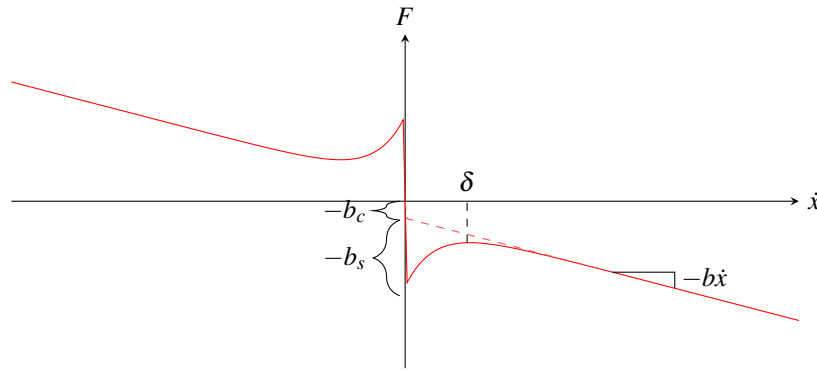
Friction is a significant disturbance present in every robotic joint (and indeed, in every mechanism.) Friction is also a *very* complex phenomenon that is very difficult to accurately model. We generally divide friction into at least three different types:

Type	Equation	Graph
Viscous, or wet	$F = -b_v \dot{x}$	
Coloumb, or dry	$F = -b_c \text{sign}(\dot{x})$	
Stiction	$F = \begin{cases} -F_c & \dot{x} = 0 \text{ and } F_c < F_s \\ F_z & \text{otherwise} \end{cases}$ $F_z = -(b_s - b_c)e^{-(\dot{x} /\delta)} \text{sign}(\dot{x})$	

where the parameters in the stiction model are b_s , the static coefficient of friction; F_c is the nominal control input to the joint, and δ is the “Stribeck velocity”, below which the joint is in the stick–slip regime. Note that this model is an attempt to model stiction deterministically, but actual stick–slip is nearly impossible to model except probabilistically, in the sense when operating the joint near the Stribeck velocity the actual force or torque needed to “break” away from a stick phenomenon and transition to a slip phenomenon is randomly distributed around F_c , as is the distance over which the joint will “slip” before entering a “stick” regime again. In addition, stick–slip can change quite dramatically with joint load. Stick–slip can be particularly challenging when using compliance control because in most cases, tasks requiring compliance require low

Cartesian velocities along most, if not all, axes, so you are almost always operating below the Stribeck velocity. This leads to what appears to be “hunting” behavior.

There are many more complex models of friction, and particularly stiction. The model we have used here is from Márton and Lantos, “Control of mechanical systems with Stribeck friction and backlash,” *Systems & Control Letters* 58.2 (2009), pp. 141-147. For more information on friction models a good reference is Waiboer, Aarts, and Jonker, “Velocity dependence of joint friction in robotic manipulators with gear transmissions,” *ECCOMAS Thematic Conference Multibody Dynamics* (2005). Adding all of these together, we get a decent model for the total friction known as a Stribeck curve:



where the mathematical model is

$$F = -b_c \text{sign}(\dot{x}) - (b_s - b_c) e^{-(|\dot{x}|/\delta)} \text{sign}(\dot{x}) - b\dot{x}$$

6 Backlash

Backlash refers to the “slop” in a gear train where one of the two gears moves before engaging the other gear. Backlash on the inside of a control loop (e.g. where you are measuring the output side of the geartrain) leads to limit cycles. Backlash on the outside of a control loop, where you are measuring the input side of the geartrain, leads to position errors. There is a strict tradeoff between backlash, friction, and compliance. Harmonic gears, for instance, tend to have low backlash but high friction and compliance. Planetary gears tend to have higher backlash but lower friction and compliance.

