

Solutions to  
Quantum Field Theory for the Gifted Amateur  
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# Chapter 1

## From Particles to Fields

### 1. WRONG

$$\begin{aligned} G(x_2; x_1) &= \int d^D \mathbf{y} \int d^D \mathbf{p} \int d^D \mathbf{p}' \frac{\theta(x_2^0 - y^0) \theta(y^0 - x_1^0)}{(2\pi)^{2D}} F(|\mathbf{p}|, x_2^0 - y^0) F(|\mathbf{p}|, y^0 - x_1^0) e^{i\mathbf{p} \cdot (\mathbf{x}_2 - \mathbf{y} + \mathbf{y} - \mathbf{x}_1)} \\ &= \int d^D \mathbf{y} \int d^D \mathbf{p} \frac{\theta(x_2^0 - x_1^0)}{(2\pi)^D} F(|\mathbf{p}|, x_2^0 - x_1^0) e^{i\mathbf{p} \cdot (\mathbf{x}_2 - \mathbf{x}_1)} \end{aligned}$$

### 2. Substituting into Eq. 1.7:

$$\int_{-\infty}^{\infty} d^D \mathbf{x} N(t) \exp \left( \frac{im|\mathbf{x}|^2}{2t} - i\mathbf{p} \cdot \mathbf{x} \right) = \theta(t) e^{-if(|\mathbf{p}|)t}$$

This is a product of  $D$  Gaussian integrals, with the evaluation of each  $n$ th component in Cartesian coordinates:

$$\int dx_n \exp \left( \frac{imx_n^2}{2t} - ip_n x_n \right) = \sqrt{\frac{2\pi it}{m}} \exp \left( -\frac{ip_n^2 t}{2m} \right)$$

Which is the propagator in momentum space, and gives the final result:

$$N(t) = \theta(t) \left( \frac{m}{2\pi it} \right)^{D/2} e^{i\psi t}, \quad \psi = \left( \frac{|\mathbf{p}|^2}{2m} - f(|\mathbf{p}|) \right) t$$

### 3. Unsatisfactorily, substitute the solution into $i\partial_t \hat{A} = [\hat{A}, \hat{H}]$ .

$$\begin{aligned} i\partial_t \left[ e^{i\hat{H}t} \hat{A}(0) e^{-i\hat{H}t} \right] &= e^{i\hat{H}t} \hat{A}(0) e^{-i\hat{H}t} \hat{H} - \hat{H} e^{i\hat{H}t} \hat{A}(0) e^{-i\hat{H}t} \\ -\hat{H} e^{i\hat{H}t} \hat{A}(0) e^{-i\hat{H}t} + e^{i\hat{H}t} \hat{A}(0) e^{-i\hat{H}t} \hat{H} &= e^{i\hat{H}t} \hat{A}(0) e^{-i\hat{H}t} \hat{H} - \hat{H} e^{i\hat{H}t} \hat{A}(0) e^{-i\hat{H}t} \end{aligned}$$

This works because the operator commutes with its exponential, easily verified by the Taylor series representation.

4. The ground state wavefunction of the quantum harmonic oscillator is:

$$\phi_0(\mathbf{x}) = \left(\frac{m\omega}{\pi}\right)^{1/4} \exp\left(-\frac{m\omega|\mathbf{x}|^2}{2}\right)$$

The action for the harmonic oscillator in Euclidean time with the boundary conditions  $(\mathbf{0}, 0)$ ,  $(\mathbf{x}, -it)$  is

$$S = \frac{m\omega|\mathbf{x}|^2 \cos -i\omega t}{2 \sin -i\omega t} = -\frac{m\omega|\mathbf{x}|^2 \cosh \omega t}{2 \sinh \omega t}$$

$$\phi_0(\mathbf{x}) \propto \int \mathcal{D}q \lim_{t_E \rightarrow \infty} \exp \left[ -\frac{m\omega|\mathbf{x}|^2 \cosh \omega t}{2 \sinh \omega t} \right]$$

The limit is evaluated via L'Hôpital's rule, giving  $\phi_0(\mathbf{x}) \propto \exp(m\omega|\mathbf{x}|^2/2)$ .

5. The propagator for the harmonic oscillator of frequency  $\omega_0$  coupled to an external source is:

$$G^J(x_2; x_1) = A[\mathbf{x}_cl] \exp(iS), \quad S = \int dt \frac{m}{2} [\dot{x}^2 - \omega^2 x^2 + J(t)x]$$

The equation of motion is:

$$\ddot{x} + \omega^2 x = \frac{J(t)}{m}$$

Solving this with boundary conditions  $(x_i, t_i)$ ,  $(x_f, t_f)$ :

$$x(t) = Ae^{i\omega t} + Be^{-i\omega t} + \int_{t_i}^{t_f} dt' G(t, t') \frac{J(t')}{m}, \quad G(t, t') = - \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \frac{e^{-i\omega(t-t')}}{\omega^2 - \omega_0^2}$$

where  $G(t, t')$  is the Green function, evaluated as follows:

$$G(t, t') =$$