Task2_PeriodicMultilayer_Bandgap

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1 Task 2: 1D Periodic Multilayer Dielectric Slab, Eigenmode formulation

1.1 Maxwell Equations (time domain)

$$\nabla \times E = -\mu \frac{\partial H}{\partial t} \tag{1}$$

$$\nabla \times H = \epsilon \frac{\partial E}{\partial t} \tag{2}$$

1.2 Time Domain → Frequency Domain

Using the Fourier transform we get:

$$E(r,t) \to E(r,\omega)$$
 (3)

$$\frac{\partial E}{\partial t} \to i\omega E \tag{4}$$

Thus our wave equation transforms into:

$$\nabla E = -i\omega\mu H \tag{5}$$

$$\nabla E = -i\omega \epsilon E \tag{6}$$

1.3 1D Coordinates, Periodic

$$H_{y_i}, E_{z_{i+\frac{1}{2}}}, i \in \mathcal{N} \tag{7}$$

$$H_{y_0} = H_{y_n} e^{ik_x a} (8)$$

$$E_{z_{0+\frac{1}{2}}} = E_{z_{n+\frac{1}{2}}} e^{ik_x a} \tag{9}$$

$$\frac{\partial E_z}{\partial x} = -i\omega \mu H_y \tag{10}$$

$$\frac{\partial H_y}{\partial x} = i\omega \epsilon E_z \tag{11}$$

1.4 Discretization

$$\frac{\partial E_{zi}}{\partial x} \approx \frac{E_{zi+1} - E_{zi-1}}{\Delta x} \tag{12}$$

$$\frac{\partial E_{z0}}{\partial x} \approx \frac{E_{z1} - E_{zn}e^{ik_x a}}{\Delta x} \tag{13}$$

$$\frac{\partial E_{zn}}{\partial x} \approx \frac{E_{z0}e^{ik_x a} - E_{zn-1}}{\Delta x} \tag{14}$$

Forward difference

$$-i\omega\mu_i E_{z_{i+\frac{1}{2}}} \approx \frac{H_{y_{i+1}} - H_{y_i}}{\Delta x} \tag{15}$$

Backward difference

$$-i\omega\epsilon_i H_{y_i} \approx \frac{E_{z_{i+\frac{1}{2}}} - E_{z_{i-\frac{1}{2}}}}{\Delta x} \tag{16}$$

1.5 Wave equation

$$H_{y_i} = \frac{i}{\omega \mu_i} \frac{\partial E_{z_i}}{\partial x} \tag{17}$$

$$E_{zi} = \frac{-i}{\omega \epsilon_i} \frac{\partial H_{y_i}}{\partial x} \tag{18}$$

$$\frac{\partial}{\partial x} \frac{i}{\omega \mu_i} \frac{\partial}{\partial x} E_{z_i} = i\omega \epsilon_i E_{z_i} \tag{19}$$

$$\frac{1}{\epsilon_i} \frac{\partial}{\partial x} \frac{1}{\mu_i} \frac{\partial}{\partial x} E_{z_i} = \omega^2 E_{z_i} \tag{20}$$

or

$$\frac{1}{\mu_i} \frac{\partial}{\partial x} \frac{1}{\epsilon_i} \frac{\partial}{\partial x} H_{y_i} = \omega^2 H_{y_i} \tag{21}$$

1.6 Matrix form

$$H_{y_i} = \frac{i}{\omega \mu_i} \frac{\partial E_z}{\partial x} \approx \frac{i}{\omega \mu_i} \frac{E_{z_{i+\frac{1}{2}}} - E_{z_{i-\frac{1}{2}}}}{\Delta x}$$
 (22)

$$\begin{bmatrix} H_{y_0} \\ H_{y_1} \\ \vdots \\ H_{y_{n-1}} \\ H_{y_n} \end{bmatrix} = \frac{i}{\omega \Delta x} \begin{bmatrix} \frac{1}{\mu_0} & 0 & 0 & \cdots & 0 \\ 0 & \frac{1}{\mu_1} & 0 & \cdots & 0 \\ 0 & 0 & \frac{1}{\mu_2} & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & 0 \\ 0 & 0 & \cdots & 0 & \frac{1}{\mu_n} \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 & \cdots & -e^{ik_x a} \\ -1 & 1 & 0 & \cdots & 0 \\ 0 & -1 & 1 & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & 0 \\ 0 & 0 & \cdots & -1 & 1 \end{bmatrix} \begin{bmatrix} E_{z_0 + \frac{1}{2}} \\ E_{z_{1+\frac{1}{2}}} \\ \vdots \\ E_{z_{n-\frac{1}{2}}} \\ E_{z_{n+\frac{1}{2}}} \end{bmatrix}$$
(23)

$$E_{z_{i+\frac{1}{2}}} = \frac{-i}{\omega \epsilon_i} \frac{\partial H_{y_i}}{\partial x} \approx \frac{-i}{\omega \epsilon_i} \frac{H_{y_{i+1}} - H_{y_i}}{\Delta x}$$
 (24)

$$\begin{bmatrix} E_{z_{0+\frac{1}{2}}} \\ E_{z_{1+\frac{1}{2}}} \\ \vdots \\ E_{z_{n-\frac{1}{2}}} \\ E_{z_{n+\frac{1}{2}}} \end{bmatrix} = \frac{i}{\omega \Delta x} \begin{bmatrix} \frac{1}{\epsilon_{0}} & 0 & 0 & \dots & 0 \\ 0 & \frac{1}{\epsilon_{1}} & 0 & \dots & 0 \\ 0 & 0 & \frac{1}{\epsilon_{2}} & \dots & \vdots \\ \vdots & \vdots & \vdots & \ddots & 0 \\ 0 & 0 & \dots & 0 & \frac{1}{\epsilon_{n}} \end{bmatrix} \begin{bmatrix} -1 & 1 & 0 & \dots & 0 \\ 0 & -1 & 1 & \dots & 0 \\ 0 & 0 & -1 & \dots & \vdots \\ \vdots & \vdots & \vdots & \ddots & 1 \\ e^{ik_{x}a} & 0 & \dots & 0 & -1 \end{bmatrix} \begin{bmatrix} H_{y_{0}} \\ H_{y_{1}} \\ \vdots \\ H_{y_{n-1}} \\ H_{y_{n}} \end{bmatrix}$$
(25)

$$\omega^2 E_z = \frac{1}{\epsilon} \frac{\partial}{\partial x} \frac{1}{\mu} \frac{\partial}{\partial x} E_z \tag{26}$$

$$\omega^{2} \begin{bmatrix} E_{z_{0+\frac{1}{2}}} \\ E_{z_{1+\frac{1}{2}}} \\ \vdots \\ E_{z_{n-\frac{1}{2}}} \\ E_{z_{n+\frac{1}{2}}} \end{bmatrix} = \frac{i}{\Delta x} \begin{bmatrix} \frac{1}{\epsilon_{0}} & 0 & 0 & \dots & 0 \\ 0 & \frac{1}{\epsilon_{1}} & 0 & \dots & 0 \\ 0 & 0 & \frac{1}{\epsilon_{2}} & \dots & \vdots \\ \vdots & \vdots & \vdots & \ddots & 0 \\ 0 & 0 & \dots & 0 & \frac{1}{\epsilon_{n}} \end{bmatrix} \begin{bmatrix} -1 & 1 & 0 & \dots & 0 \\ 0 & -1 & 1 & \dots & 0 \\ 0 & 0 & -1 & \dots & \vdots \\ \vdots & \vdots & \vdots & \ddots & 1 \\ e^{ik_{x}a} & 0 & \dots & 0 & -1 \end{bmatrix} \underbrace{\frac{i}{\Delta x}} \begin{bmatrix} \frac{1}{\mu_{0}} & 0 & 0 & \dots & 0 \\ 0 & \frac{1}{\mu_{1}} & 0 & \dots & 0 \\ 0 & 0 & \frac{1}{\mu_{2}} & \dots & \vdots \\ \vdots & \vdots & \vdots & \ddots & 0 \\ 0 & 0 & \dots & 0 & \frac{1}{\mu_{n}} \end{bmatrix} \begin{bmatrix} 1 & 0 \\ -1 & 1 \\ 0 & -1 \\ \vdots & \vdots \\ 0 & 0 \end{bmatrix}$$

or, simplifing, and replacing $i + \frac{1}{2}$ with i:

$$\omega^{2} \begin{bmatrix} E_{z0} \\ E_{z1} \\ \vdots \\ E_{zn-1} \\ E_{zn} \end{bmatrix} = -\frac{1}{\Delta x^{2}} \begin{bmatrix} \frac{1}{\epsilon_{0}} & 0 & 0 & \cdots & 0 \\ 0 & \frac{1}{\epsilon_{1}} & 0 & \cdots & 0 \\ 0 & 0 & \frac{1}{\epsilon_{2}} & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & 0 \\ 0 & 0 & \cdots & 0 & \frac{1}{\epsilon_{n}} \end{bmatrix} \begin{bmatrix} -1 & 1 & 0 & \cdots & 0 \\ 0 & -1 & 1 & \cdots & 0 \\ 0 & 0 & -1 & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & 1 \\ e^{ik_{x}a} & 0 & \cdots & 0 & -1 \end{bmatrix} \begin{bmatrix} \frac{1}{\mu_{0}} & 0 & 0 & \cdots & 0 \\ 0 & \frac{1}{\mu_{1}} & 0 & \cdots & 0 \\ 0 & 0 & \frac{1}{\mu_{1}} & 0 & \cdots & 0 \\ 0 & 0 & \frac{1}{\mu_{2}} & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & 0 \\ 0 & 0 & \cdots & 0 & \frac{1}{\mu_{n}} \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & -1 & 1 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & 0 \\ 0 & 0 & \cdots & 0 & \frac{1}{\mu_{n}} \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & -1 & 1 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & 0 \\ 0 & 0 & \cdots & 0 & \frac{1}{\mu_{n}} \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & -1 & 1 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & 0 \\ 0 & 0 & \cdots & 0 & \frac{1}{\mu_{n}} \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{1}{\mu_{1}} & 0 & \cdots & 0 \\ 0 & 0 & \frac{1}{\mu_{n}} & \cdots & 0 \\ 0 & 0 & \cdots & 0 & \frac{1}{\mu_{n}} \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{1}{\mu_{1}} & 0 & \cdots & 0 \\ 0 & 0 & 0 & \frac{1}{\mu_{n}} & \cdots & 0 \\ 0 & 0 & \cdots & 0 & \frac{1}{\mu_{n}} & \cdots & 0 \\ 0 & 0 & \cdots & 0 & \frac{1}{\mu_{n}} & \cdots & 0 \\ 0 & 0 & \cdots & 0 & \frac{1}{\mu_{n}} & \cdots & 0 \\ 0 & 0 & \cdots & 0 & \frac{1}{\mu_{n}} & \cdots & 0 \\ 0 & 0 & \cdots & 0 & \frac{1}{\mu_{n}} & \cdots & 0 \\ 0 & 0 & \cdots & 0 & \frac{1}{\mu_{n}} & \cdots & 0 \\ 0 & 0 & \cdots & 0 &$$

1.7 Boundary Conditions

In this case the boundary conditions are periodic, satisfied by the e^{ik_xa} elements in the resulting matrix

2 Implementation

2.1 Setup grid

```
In [102]: n=200  # number of grid nodes

dx=1/n  # discretization step, domain size = 1

eps1 = 13  # Layer 1
 eps2 = 1  # Layer 2

num_eigs = 6  # Solver for first # eigenmodes
```

2.2 Build matrix

2.3 Solve for eigenmodes

```
In [117]: kk0 = 2*sp.pi*np.linspace(-0.5,0.5,300) # k-vector in medium with eps1-ep
k = np.zeros((num_eigs, kk0.size), dtype=complex)
for ik in range(kk0.size):
    k0=kk0[ik]

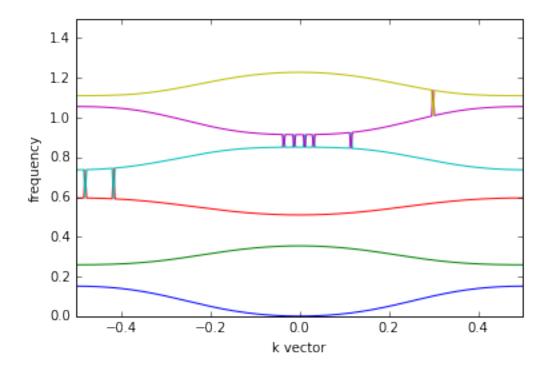
M_1[n-1, 0] = np.cos(k0*1)/dx # to impose periodic boundary condition
```

```
M_2[0, n-1] = -np.cos(k0*1)/dx
M = -eps*M_2*M_1
kt = k0/np.sqrt((eps1+eps2)/2); # target k=w/c
k2, V = linalg.eigs(M, k=num_eigs, M=None, sigma=kt**2)
k[:,ik] = np.sqrt(k2) # k=w/c
```

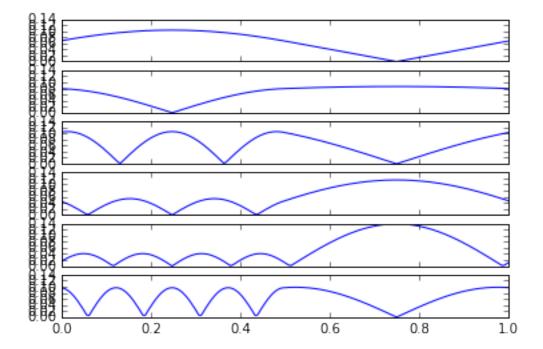
2.4 Plot eigenmodes

2.4.1 Bands

Out[126]: (0.0, 1.5)



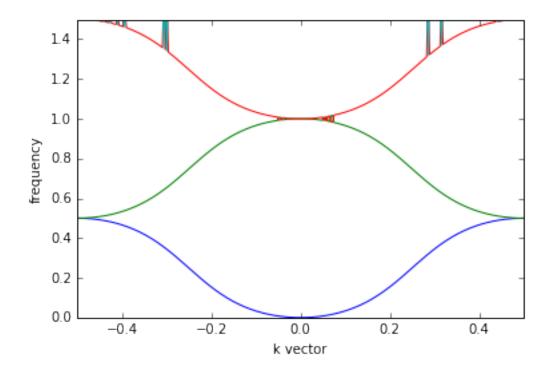
2.4.2 Field distribution

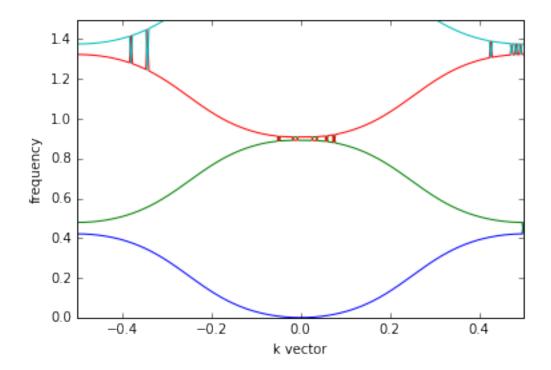


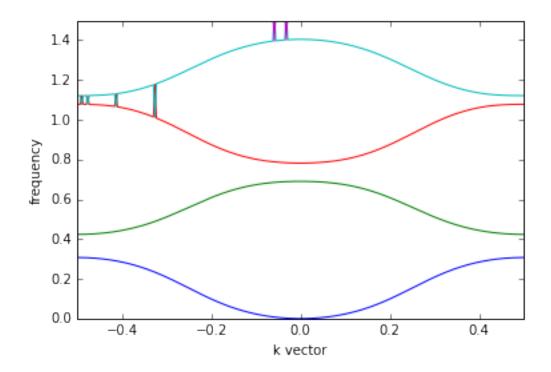
2.5 Permittivity Contrast Bandgap dependence

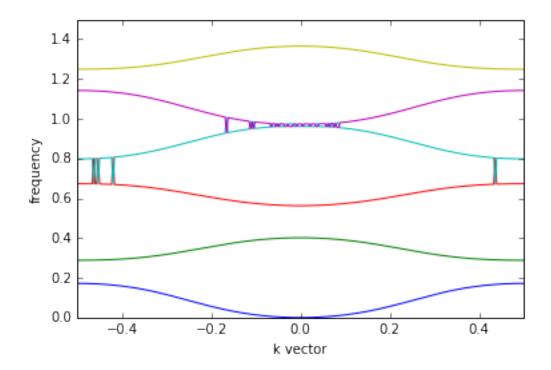
```
# Backward
              diag = np.ones(n) * 1/dx
              up\_diag = np.ones(n) * -1/dx
              M_2 = sp.sparse.dia_matrix(([up_diag, diag], [-1, 0]), [n,n])
              M_1 = sp.sparse.lil_matrix(M_1)
              M_2 = sp.sparse.lil_matrix(M_2)
              kk0 = 2*sp.pi*np.linspace(-0.5, 0.5, 300) # k-vector in medium with eps
              k = np.zeros((num_eigs, kk0.size), dtype=complex)
              for ik in range(kk0.size):
                  k0=kk0[ik]
                  M_1[n-1, 0] = np.cos(k0*1)/dx # to impose periodic boundary cond.
                  M_2[0, n-1] = -np.cos(k0*1)/dx
                  M = -eps * M 2 * M 1
                  kt = k0/np.sqrt((eps1+eps2)/2); # target k=w/c
                  k2, V = linalg.eigs(M, k=num\_eigs, M=None, sigma=kt**2)
                  k[:,ik] = np.sqrt(k2) \# k=w/c
              return kk0/(2*sp.pi), np.real(k/(2*sp.pi))
2.5.1 \epsilon_1 = 1, \epsilon_2 = 1
In [130]: k0, k = build_bandgap_diagram(1, 1, num_eigs)
          for i in range(num_eigs):
              plt.hold(True)
              plt.plot(k0, k[i,:],'-')
          plt.xlabel("k vector")
          plt.ylabel("frequency")
          plt.xlim([-0.5, 0.5])
          plt.ylim([0.0, 1.5])
Out[130]: (0.0, 1.5)
```

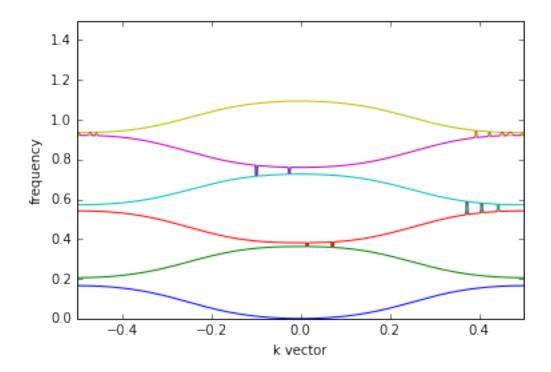
 $M_1 = \text{sp.sparse.dia_matrix}(([up_diag, diag], [1, 0]), [n,n])$

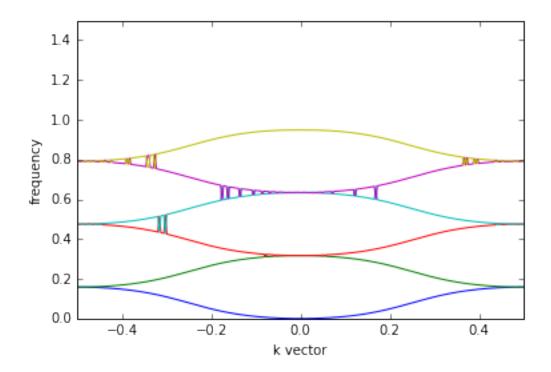












In []: