1) For a physical dipole with -q at \$\vec{r}\$ and +q at \$\vec{r}\$+\$\vec{d}\$

$$U = 9V(\vec{r}+\vec{d}) - V(\vec{r})$$

$$= 9\left[-\int_{\vec{r}}^{\vec{r}} + \vec{d}\right]$$

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$$= -\vec{p} \cdot \vec{e}$$

$$U = -\vec{P_1} \cdot \vec{E_2}$$

BW
$$\vec{E}_2 = \frac{1}{4\pi\epsilon_0} \frac{1}{23} \left[3(\vec{P}_2 \cdot \hat{p}) \hat{p} - \vec{P}_2 \right]$$

$$: U = \frac{1}{4\pi\epsilon_0} \frac{1}{23} \left[(\vec{P}_1 \cdot \vec{P}_2) - 3 (\vec{P}_1 \cdot \hat{r}) (\vec{P}_2 \cdot \hat{r}) \right]$$

$$P_V = \begin{cases} P_0 & 0 < r < a \\ 0 & r > a \end{cases}$$

For
$$2 < \alpha$$
. $\mathcal{E} = (4\pi 2^2) = \frac{4}{3}\pi 2^3$

$$\Rightarrow \mathcal{E} = \frac{9 \cdot 2}{3 \cdot 6}$$

$$V = -\int \vec{E} \cdot d\vec{i}$$

$$= -\int \frac{\rho_0 \hbar}{3\epsilon} dz$$

$$= -\frac{\rho_0}{3\epsilon} \left[\frac{\hbar^2}{2} \right]^{\frac{1}{2}} + C_1$$

$$= -\frac{\rho_0 \hbar^2}{6\epsilon} + C_1$$

For
$$2 > a$$
, $\epsilon \in (4\pi 2^{2}) = e^{6} \frac{4}{3}\pi a^{3}$

$$\Rightarrow E = \frac{e^{6}a^{3}}{3\epsilon_{0}2^{2}}$$

$$f \cdot V = -\int \vec{E} \cdot d\vec{i}$$

$$= -\int \frac{e^{0} a^{3}}{3 \epsilon_{0} k^{2}} dr$$

$$= -\frac{e^{0} a^{3}}{3 \epsilon_{0}} \left[-\frac{1}{k} \right] + C_{2}$$

$$= \frac{e^{0} a^{3}}{3 \epsilon_{0} k} + C_{2}$$

Boundary condition

As
$$k \to \infty$$
 $V \longrightarrow 0$... $C_2 \stackrel{\text{\tiny deg}}{=} 0$

Also
$$\forall$$
 at $\mathcal{P} = a$ $\vee(a^+) = \vee(a^-)$

$$-\frac{\rho_0 a^2}{6\epsilon} + C_1 = \frac{\rho_0 a^3}{3\epsilon_0 a}$$

$$\Rightarrow -\frac{\rho_0 a^2}{6\epsilon} + C_1 = \frac{\rho_0 a^2}{3\epsilon_0}$$

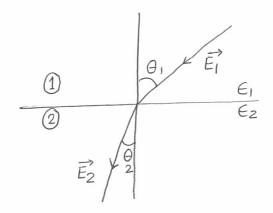
$$\Rightarrow C_1 = \frac{\rho_0 a^2}{6\epsilon} \left[1 + \frac{2\epsilon}{\epsilon_0}\right]$$

$$= \frac{\rho_0 a^2}{6\epsilon_0 \epsilon_2} \left(1 + 2\epsilon_2\right)$$

$$V(0) = C_1 = \frac{R_0 a^n}{6 \epsilon_0 \epsilon_R} (1 + 2 \epsilon_R)$$

$$V(a) = \frac{\rho_0 a^{\nu}}{3\epsilon_0}$$

(3)



$$E_{11}$$
 is continuous $\Rightarrow E_{1x} = E_{2x}$

DI is continuous (with Tp=0)

$$\therefore D_{1y} = D_{2y}$$

$$\Rightarrow \mathcal{E}_{1} \mathcal{E}_{1y} = \mathcal{E}_{2} \mathcal{E}_{2y}$$

$$\frac{\tan \theta_2}{\tan \theta_1} = \frac{\frac{E_1 x}{E_1 x}}{\frac{E_2 x}{E_2 y}} = \frac{\frac{E_1 x}{E_1 x}}{\frac{E_2 x}{E_2 y}}$$

$$= \frac{\frac{E_1 y}{E_2 y}}{\frac{E_2 y}{E_2 y}}$$

$$= \frac{\frac{E_2 y}{E_1 x}}{\frac{E_2 y}{E_2 y}}$$

Applying Ampere's law -

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 \text{ Iencl}$$

$$\Rightarrow B.2\PiP = \mu_0 \text{ I} \frac{\PiP^2}{\Pi a^2}$$

$$\Rightarrow B.2\PiP = \mu_0 \text{ I} \frac{P^2}{a^2}$$

$$\Rightarrow B = \mu_0 \text{ IP}$$

$$= \frac{\mu_0 \text{ IP}}{2\Pi a^2}$$

$$\vec{B} = \frac{\mu_0 \text{ IP}}{2\Pi a^2}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J}$$

$$\Rightarrow \vec{J} = \frac{1}{\mu_0} (\vec{\nabla} \times \vec{B})$$

$$= \frac{1}{\mu_0} \frac{1}{\ell} \left[\frac{2}{2\ell} (\ell + \frac{\mu_0 I \ell}{2 \pi a^{\nu}}) \right] \hat{2}$$

$$= \frac{1}{\mu_0} \frac{1}{\ell} \frac{\mu_0 I}{2 \pi a^{\nu}} 2\ell \hat{2}$$

$$= \frac{I}{\pi a^{\nu}} \hat{2}$$

$$\vec{B} = \vec{\nabla} \times \vec{A}$$

$$\vec{B} = \vec{\nabla} \times \vec{A}$$

$$= \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} & -4xyz \end{vmatrix}$$

$$= \hat{i} \left(-4\alpha z - 0 \right) - \hat{j} \left(-4yz - 0 \right) + \hat{k} \left(y^2 - \alpha^2 \right)$$

$$= -4\alpha z \hat{i} + 4yz \hat{j} + (y^2 - \alpha^2) \hat{k}$$

$$\vec{B}(-1,2,5) = -4(-5)\hat{i} + 4 \times 10\hat{j} + (4-1)\hat{k}$$

$$= 20\hat{i} + 40\hat{j} + 3\hat{k} \quad \frac{Wb}{m^{2}}$$

$$\oint \vec{B} \cdot d\vec{s}$$

$$= \oint \vec{A} \cdot d\vec{i} = \phi_1 + \phi_2 + \phi_3 + \phi_4$$

$$\ell$$

Now
$$\phi_i = \left[\int_{y=-1}^{4} (x^3y^2 + y^3x^2 - 4xyz^2) \cdot dy^2 \right]$$

$$= \begin{cases} 4 & y^{2} x & dy \\ y = -1 & x = 1 \\ 2 = 1 & 2 = 1 \end{cases}$$

$$= \frac{4}{3} + \frac{1}{3} = \frac{65}{3}$$

$$\frac{2}{3}$$

$$\begin{aligned}
& \Phi_{2} = \left[\int_{\alpha_{2}}^{\alpha_{2}} (x^{2}y^{2} + y^{2}x^{2}) - 4xy^{2}k \right] \cdot dx^{2} \right] \\
& = -\int_{\alpha_{2}}^{\alpha_{2}} x^{2}y \, dx \\
& = -4 \left[\frac{\alpha_{3}^{3}}{3} \right]_{1}^{0} \\
& = +4 \times \frac{1}{3} \\
& = +\frac{4}{3} \end{aligned}$$

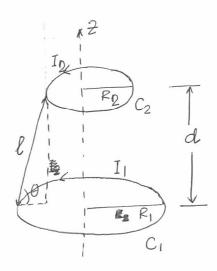
$$\begin{aligned}
& \Phi_{3} = \left[\int_{\alpha_{2}}^{-1} (x^{2}y^{2} + y^{2}x^{2}) - 4xy^{2}k \right] \cdot dy^{2} \right] \\
& = 0
\end{aligned}$$

$$\begin{aligned}
& \Phi_{4} = \left[\int_{\alpha_{2}}^{-1} (x^{2}y^{2} + y^{2}x^{2}) - 4xy^{2}k \right] \cdot (dx^{2}) \\
& = \int_{\alpha_{2}}^{-1} (x^{2}y^{2} + y^{2}x^{2}) - 4xy^{2}k \right] \cdot (dx^{2}) \\
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& = \int_{\alpha_{2}}^{-1} (x^{2}y^{2} + y^{2}) + 4xy^{2}k \cdot (dx^{2}) \\
& = \int_{\alpha_$$

: Flux through the surface =
$$\phi = \frac{65}{3} + \frac{4}{3} + 0 - \frac{1}{3}$$

= $\frac{68}{3}$ Wb.





We consider two coils C_1 and C_2 of radii R_1 and R_2 respectively, carrying currents I_1 and I_2 respectively. Let d be the separation between the centers of the coils.

Force per unit length on either coil is

$$F_0 = \frac{\mu_0 I_1 I_2}{2\pi \ell}$$

By symmetry, the component of Fo perpendicular to the z-arris will cancel.

The total force along z-arris
$$F = \frac{\mu_0 I_1 I_2}{2\pi \ell} S_{1n} \Theta \times 2\pi R_2 \qquad R_1 \simeq R_2$$

$$= 2\pi R_2 \times \frac{\mu_0 I_1 I_2}{2\pi \ell} \times \frac{d}{\ell}$$

$$F = \frac{\mu_0 I_1 I_2}{d^2 + (R_1 - R_2)^2}$$

Force is marinum when,

$$\frac{dF}{dd} = 0$$

$$\Rightarrow \frac{d^{r} + (R_{1} - R_{2})^{2} - d \cdot 2d}{(d^{r} + (R_{1} - R_{2})^{r})^{2}} = 0$$

$$\Rightarrow (R_{1} - R_{2})^{r} = d^{r}$$

$$\Rightarrow d = R_{1} - R_{2}$$