Self-consistent GW in Practice

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Outline

- Electronic structure problem
- Overview of real material simulations in UGF2
- scGW in UGF2
- Examples: Silicon



Electronic structure simulation of real materials

- Molecules
 - Finite-size system
- Solids
 - Infinite periodic unit cells
 - Atoms in a unit cell + Translational vectors
- Born-Oppenheimer approximation
 - Stationary nuclei
 - Neglect the kinetic energy of nuclei
 - Coulomb repulsion energy between nuclei can be view as a constant
- Electronic Hamiltonian of an N-electron system with M nuclei

$$H = \sum_{i=1}^{N} -\frac{1}{2} \nabla_i^2 + \sum_{i=1}^{N} \sum_{A=1}^{M} -\frac{Z_A}{r_{iA}} + \sum_{i=1}^{N} \sum_{j>i}^{N} \frac{1}{r_{ij}}$$

One-body part

two-body part

Electronic Hamiltonian in second quantization

Certain type of basis could benefit certain electronic structure methods

a) Position basis:
$$H = \int d\mathbf{r} \psi^{\dagger}(\mathbf{r}) [-\frac{1}{2}\nabla^{2} + V(\mathbf{r})] \psi(\mathbf{r}) + \frac{1}{2} \int d\mathbf{r} \int d\mathbf{r}' U(\mathbf{r} - \mathbf{r}') \psi^{\dagger}(\mathbf{r}) \psi^{\dagger}(\mathbf{r}') \psi(\mathbf{r}') \psi(\mathbf{r}')$$

b) Localized atomic basis $\chi_i^{\mathbf{R}}(\mathbf{r})$: \mathbf{R} : unit cell index, i: orbital index

$$H = \sum_{\mathbf{R}\mathbf{R}'} \sum_{ij} (H_0)_{ij}^{\mathbf{R}\mathbf{R}'} c_i^{\mathbf{R}\dagger} c_j^{\mathbf{R}} + \frac{1}{2} \sum_{\mathbf{R}_1 \mathbf{R}_2 \mathbf{R}_3 \mathbf{R}_4} \sum_{ijkl} U_{ijkl}^{\mathbf{R}_1 \mathbf{R}_2 \mathbf{R}_3 \mathbf{R}_4} c_i^{\mathbf{R}_1\dagger} c_k^{\mathbf{R}_3\dagger} c_l^{\mathbf{R}_4} c_j^{\mathbf{R}_2}$$

$$(H_0)_{ij}^{\mathbf{R}\mathbf{R}'} = \int d\mathbf{r} \chi_i^{\mathbf{R}^*}(\mathbf{r}) [-\frac{1}{2} \nabla^2 + V(\mathbf{r})] \chi_j^{\mathbf{R}'}(\mathbf{r})$$

$$U_{ijkl}^{\mathbf{R}_1 \mathbf{R}_2 \mathbf{R}_3 \mathbf{R}_4} = \int d\mathbf{r} \int d\mathbf{r}' \chi_i^{\mathbf{R}_1^*}(\mathbf{r}) \chi_j^{\mathbf{R}_2}(\mathbf{r}) U(\mathbf{r} - \mathbf{r}') \chi_k^{\mathbf{R}_3^*}(\mathbf{r}') \chi_l^{\mathbf{R}_4}(\mathbf{r}')$$

$$U_{ijkl}^{\mathbf{R}_1 \mathbf{R}_2 \mathbf{R}_3 \mathbf{R}_4} = \int d\mathbf{r} \int d\mathbf{r}' \chi_i^{\mathbf{R}_1^*}(\mathbf{r}) \chi_j^{\mathbf{R}_2}(\mathbf{r}) U(\mathbf{r} - \mathbf{r}') \chi_k^{\mathbf{R}_3^*}(\mathbf{r}') \chi_l^{\mathbf{R}_4}(\mathbf{r}')$$

c) Bloch localized atomic basis $\chi_i^{\mathbf{k}}(\mathbf{r})$: $\chi_i^{\mathbf{k}}(\mathbf{r}) = \sum_{\mathbf{R}} \chi_i^{\mathbf{R}}(\mathbf{r}) e^{i\mathbf{k}\mathbf{R}}, \ \chi_i^{\mathbf{k}}(\mathbf{r} + \mathbf{R}) = \chi_i^{\mathbf{k}}(\mathbf{r})$

$$H = \sum_{\mathbf{k}} \sum_{ij} (H_0)_{ij}^{\mathbf{k}} c_i^{\mathbf{k}\dagger} c_j^{\mathbf{k}} + \frac{1}{2N_k} \sum_{\mathbf{k}\mathbf{k}'\mathbf{q}} \sum_{ijkl} U_{i \ j \ k}^{\mathbf{k} \cdot -\mathbf{q} \ \mathbf{k}' \ \mathbf{k}' + \mathbf{q}} c_i^{\mathbf{k}\dagger} c_k^{\mathbf{k}'\dagger} c_l^{\mathbf{k}' + \mathbf{q}} c_j^{\mathbf{k} - \mathbf{q}}$$

$$(H_0)_{ij}^{\mathbf{k}} = \int_{\Omega} d\mathbf{r} \chi_i^{\mathbf{k}^*}(\mathbf{r}) [-\frac{1}{2} \nabla^2 + V(\mathbf{r})] \chi_j^{\mathbf{R}'}(\mathbf{r})$$

$$U_{i \ j \ k}^{\mathbf{k} \cdot -\mathbf{q} \ \mathbf{k}' \ \mathbf{k}' + \mathbf{q}} = \int_{\Omega} d\mathbf{r} \int d\mathbf{r}' \chi_i^{\mathbf{k}^*}(\mathbf{r}) \chi_j^{\mathbf{k} - \mathbf{q}}(\mathbf{r}) U(\mathbf{r} - \mathbf{r}') \chi_k^{\mathbf{k}'^*}(\mathbf{r}') \chi_l^{\mathbf{k}' + \mathbf{q}}(\mathbf{r}')$$

UGF2: MBPT for periodic systems

Github: https://github.com/CQMP/UGF2

c) Bloch localized atomic basis $\chi_i^{\mathbf{k}}(\mathbf{r})$: $\chi_i^{\mathbf{k}}(\mathbf{r}) = \sum_{\mathbf{R}} \chi_i^{\mathbf{R}}(\mathbf{r}) e^{i\mathbf{k}\mathbf{R}}, \ \chi_i^{\mathbf{k}}(\mathbf{r} + \mathbf{R}) = \chi_i^{\mathbf{k}}(\mathbf{r})$ $H = \sum_{\mathbf{k}} \sum_{ij} (H_0)_{ij}^{\mathbf{k}} c_i^{\mathbf{k}\dagger} c_j^{\mathbf{k}} + \frac{1}{2N_k} \sum_{\mathbf{k}\mathbf{k}'\mathbf{q}} \sum_{ijkl} U_{i\ j\ k}^{\mathbf{k}\ \mathbf{k}-\mathbf{q}\ \mathbf{k}'} \frac{\mathbf{k}' + \mathbf{q}}{l} c_i^{\mathbf{k}\dagger} c_k^{\mathbf{k}'\dagger} c_l^{\mathbf{k}'+\mathbf{q}} c_j^{\mathbf{k}-\mathbf{q}}$

$$H = \sum_{\mathbf{k}} \sum_{ij} (H_0)_{ij}^{\mathbf{k}} c_i^{\mathbf{k}\dagger} c_j^{\mathbf{k}} + \frac{1}{2N_k} \sum_{\mathbf{k}\mathbf{k}'\mathbf{q}} \sum_{ijkl} U_i^{\mathbf{k}} \int_{j}^{\mathbf{k}-\mathbf{q}} \int_{k}^{\mathbf{k}'} \mathbf{k}' + \mathbf{q} c_i^{\mathbf{k}\dagger} c_k^{\mathbf{k}'\dagger} c_l^{\mathbf{k}'+\mathbf{q}} c_j^{\mathbf{k}-\mathbf{q}}$$

- Once the Hamiltonians have been computed, the UGF2 program
 - Performs many-body perturbation theory (MBPT) using different MBPT solvers $\mathscr{F}^{\mathrm{MBPT}}$

$$\text{MBPT self-consistent loop:} \qquad \qquad G_{ij}^{\mathbf{k}}(i\omega_n) = [(i\omega_n + \mu)S^{\mathbf{k}} - H_0^{\mathbf{k}} - \Sigma^{\mathbf{k}}(i\omega_n)]_{ij}^{-1} \qquad \qquad \\ \Sigma_{ij}^{\mathbf{k}} = \mathscr{F}^{\mathrm{MBPT}}[G] \qquad \qquad \qquad \\ \Sigma_{ij}^{\mathbf{k}} = \mathscr{F}^{\mathrm{MBPT}}[G]$$

- $\mathscr{F}^{\mathrm{MBPT}}$: HF, GF2, GW etc
- Assumes Bloch Gaussian-type orbitals as the one-particle basis
- Assumes decomposed three-index Coulomb integrals $V_{i,j}^{\mathbf{k}\mathbf{k}-\mathbf{q}}(Q)$

$$U_{i \ j \ k}^{\mathbf{k} \ \mathbf{k} - \mathbf{q} \ \mathbf{k}' \ \mathbf{k}' + \mathbf{q}} = \sum_{Q} V_{i \ j}^{\mathbf{k} \mathbf{k} - \mathbf{q}}(Q) V_{k \ l}^{\mathbf{k}' \mathbf{k}' + \mathbf{q}}(Q)$$

- Uses intermediate numerical representations (IR) and sparse-sampling technique to represent dynamic objects on the imaginary axes

- Run many-body perturbation theory (MBPT) using UGF2
 - C++ and CUDA
 - Input: Matrix elements stored in the HDF5 output files from UGF2/script/init_data_df.py
 - Output: Green's function and Self-energy

• Define the problem

- Atoms in the primitive unit cell
- Translational vectors
- Gaussian basis set
- **-** *k*-mesh

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Compute matrix elements of the Hamiltonian

- Python scripts using PySCF
- init_data_df.py script in UGF2/scripts/
- Density fitting for the two-electron Coulomb interaction
- Store the Hamiltonian matrix elements in HDF5 files

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- Output: Green's function and Self-energy
- Post-processing: Band structure, Fermi surface etc
 - Python and C++
 - Band interpolation along the high-symmetry k-path
 - Analytical continuation

Building UGF2

Github: https://github.com/CQMP/UGF2

Dependencies

- CMake
- ALPSCore (https://github.com/ALPSCore/ALPSCore)
- Eigen3
- HDF5
- CUDA (optional)

Example of building UGF2 on pauli:

- Load required libraries: module load BuildEnv/intel-2021.3.0 module load alpscore/2.3.1a-mpi
- 2. git clone https://github.com/CQMP/UGF2.git and go to the UGF2 folder
- 3. mkdir build && cd build
- 4. CC=mpicc CXX=mpicxx cmake ..
 - -DCMAKE_CXX_FLAGS='-DEIGEN_USE_BLAS' Use openblas as backends for some -DCMAKE_EXE_LINKER_FLAGS='-lopenblas' linear algebra functions (optional)
 - -DWITH_CUDA=ON] Enable cuda HF/GW solver (optional)
- 5. make -j4

Running UGF2

Basic parameters

- scf_type MBPT solver, e.g. HF, GF2, GW etc
- nel_cell number of electrons per unit cell
- nao number of atomic orbitals per unit cell
- nk number of k-points
- ink number of reduced k-points in the presence of inversion symmetry
- ns number of spins (1: restricted, 2: unrestricted)
- NQ number of auxiliary basis per unit cell
- beta inverse temperature (Ha⁻¹)
- ni number of imaginary-time points in the IR gird
- TNL IR-grid files for fermonic functions
- TNL_B IR-grid files for bosonic functions
- intermax maximum number of MBPT iterations
- E_thr Energy convergence threshold
- input_file input files with one-electron Hamiltonians
- dfintegral_file density-fitted Coulomb integrals for GW, GF2
- HF_dfintegral_file density-fitted Coulomb integrals for HF
- Results HDF5 files to store results
- rst read "Results" and restart the calculation

Try ./UGF2 --help to check the complete parameter list!

Outputs of UGF2/script/ init_data_df.py

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Outputs of UGF2/script/ init_data_df.py

IR representation for dynamic quantities

- ni number of imaginary-time points in the IR gird
- TNL IR-grid files for fermonic functions
- TNL_B IR-grid files for bosonic functions

For all dynamic quantities, a set of IR non-uniform grids and the corresponding Fourier transformation matrices are pre-computed and stored at

/UGF2/MB_analysis/data/ir_grid/lambda_ni.h5

Ex: 1e3_72.h5, 1e4_104.h5, 1e5_136.h5, 1e6_168.h5

Phys. Rev. B 96, 035147 (2017) Phys. Rev. B 101, 035144 (2020)

- lambda: controlled parameter in the IR representation.
 Phys. Rev.
 In principle, lambda ≥ (beta) * (energy window of the system)
- ni: number of imaginary-time points
- lambda_ni.h5 stores both the fermonic and bosonic grids.
- Lower temperature or larger energy window -> larger imaginary grids

Theoretical details can be found in https://green.physics.lsa.umich.edu/mw19/index.php?title=IR Basis Set https://green.physics.lsa.umich.edu/mw19/index.php?title=Sparse Sampling in time and frequency

Hartree-Fock (scf_type=HF, HF_X2C1e)

Computational bottleneck:

HF exchange diagram:
$$K_{ij}^{\mathbf{k}} = \frac{-1}{N_k} \sum_{\mathbf{k}'} \sum_{ab} G_{ab}^{\mathbf{k}'} (\tau = \beta^-) U_{i\ a\ b\ j}^{\mathbf{k}\mathbf{k}'\mathbf{k}'\mathbf{k}}$$

- Non-relativistic or spin-free relativistic HF: scf_type=HF
- Two-component relativistic HF: scf_type=HF_X2C1e X2C=true
- MPI parallelization over the (ink) independent k-points

Self-consistent GW (scf_type=GW, GW_X2C1e)

$$\tilde{W} = \sum_{l=1}^{\tilde{P}_{0,QQ'}^{\mathbf{q}}(i\Omega_n)} \left\{ \underbrace{W + \left(W \right) W}_{l}^{\mathbf{q}} \right\}^2 + \dots \right\} W$$

$$\tilde{P}_{QQ'}^{\mathbf{q}}(i\Omega_n)$$

$$\tilde{P}_{0,QQ'}^{\mathbf{q}}(\tau) = \frac{-1}{N_k} \sum_{l=1}^{\infty} \sum_{l=1}^{\infty} V_{d\ a}^{\mathbf{k},\mathbf{k}+\mathbf{q}}(Q) G_{c\sigma,d\sigma}^{\mathbf{k}}(-\tau) G_{a\sigma\ b\sigma}^{\mathbf{k}+\mathbf{q}}(\tau) V_{b\ c}^{\mathbf{k}+\mathbf{q},\mathbf{k}}(Q')$$

- Non-relativistic or spin-free relativistic GW: scf_type=GW
- Two-component relativistic GW: scf_type=GW_X2C1e X2C=true

Two-level MPI parallelization

- controlled by **ntauspinprocs** number of processes for the second layer of MPI parallelization on τ and spin axes
- (ink) independent q-points
- (ni) independent τ -points

Examples:

Given N MPI processes and ntauspinprocs = x ->

- N/x of processes on the q-axis
- x of processes on the τ -axis

Cuda GW (scf_type=cuGW, cuGW_X2C1e)

$$\tilde{W} = \sum_{\mathbf{W}} \left\{ \underbrace{\mathbf{W} \mathbf{W}}_{0,QQ'}(i\Omega_n) + \left(\mathbf{W} \mathbf{W} \mathbf{W} \right)^2 + \dots \right\} \mathbf{W}$$

$$\tilde{P}_{QQ'}^{\mathbf{q}}(i\Omega_n)$$

$$\tilde{P}_{0,QQ'}^{\mathbf{q}}(\tau) = \frac{-1}{N_k} \sum_{\mathbf{k}} \sum_{\sigma} \sum_{abcd} V_{d\ a}^{\mathbf{k},\mathbf{k}+\mathbf{q}}(Q) G_{c\sigma,d\sigma}^{\mathbf{k}}(-\tau) G_{a\sigma\ b\sigma}^{\mathbf{k}+\mathbf{q}}(\tau) V_{b\ c}^{\mathbf{k}+\mathbf{q},\mathbf{k}}(Q')$$

- Non-relativistic or spin-free relativistic GW: scf_type=cuGW
- Two-component relativistic GW: scf_type=cuGW_X2C1e X2C=true

- Parallelization scheme
 - MPI parallelization over the **q**-axis
 - Asynchronous streams for the summation over **k**-points
 - Batched ZGEMM over the τ-axis controlled by
 nt batch size of τ batch in cuda GW solver

Iterative solvers for self-consistency loop

$$G_{ij}^{\mathbf{k}}(i\omega_n) = [(i\omega_n + \mu)S^{\mathbf{k}} - H_0^{\mathbf{k}} - \Sigma^{\mathbf{k}}(i\omega_n)]_{ij}^{-1}$$

$$\Sigma_{ij}^{\mathbf{k}} = \mathcal{F}^{\mathrm{MBPT}}[G]$$

Parameters for iterative solvers

- damp damping parameter for self-energy and Fock matrix
 (1: no damp, 0: full damp)
- DIIS_size DIIS space size
- DIIS_start iteration where the DIIS solver begins
- DIIS_interval number of iteration between two DIIS extrapolations

More elaborate iterative solvers can be found in https://green.physics.lsa.umich.edu/mw19/index.php?title=DIIS_and_Convergence_Acceleration

Example: Silicon

• Basis set: gth-dzvp-molopt-sr, Pseudopotential: gth-pbe

After running UGF2/script/init_data_df.py...

- input.h5: input file with all the one-electron quantities
- df_int: density-fitted Coulomb integrals for GW/GF2 with Madelung constant correction
- df_hf_int: density-fitted Coulomb integrals for HF

run.sh:

```
#!/bin/bash
#SBATCH -p super,batch
#SBATCH -t 48:00:00
#SBATCH -N 2
#SBATCH -n 128
#SBATCH -c 2
#SBATCH -o Si6_GW.o%j
#SBATCH -J Si6
#SBATCH --exclusive
#OpenMP settings:
export HDF5_USE_FILE_LOCKING=FALSE
export INT_DIR=/pauli-storage/cnyeh/Si/nk6/integrals/
export BinDir=/home/cnyeh/Project/UGF2_CQMP/build/
export IR_DIR=/home/cnyeh/Project/UGF2_CQMP/MB_analysis/data/ir_grid/
date
srun -n 128 -c 2 --cpu_bind=cores $BinDir/UGF2 gw_param --dfintegral_file=$INT_DIR/df_int
--HF_dfintegral_file=$INT_DIR/df_hf_int --input_file=input.h5 --TNL=$IR_DIR/1e6_168.h5
-TNL_B=$IR_DIR/1e6_168.h5 --Results=sim.h5 --ntauspinprocs=2
date
```

gw_param:

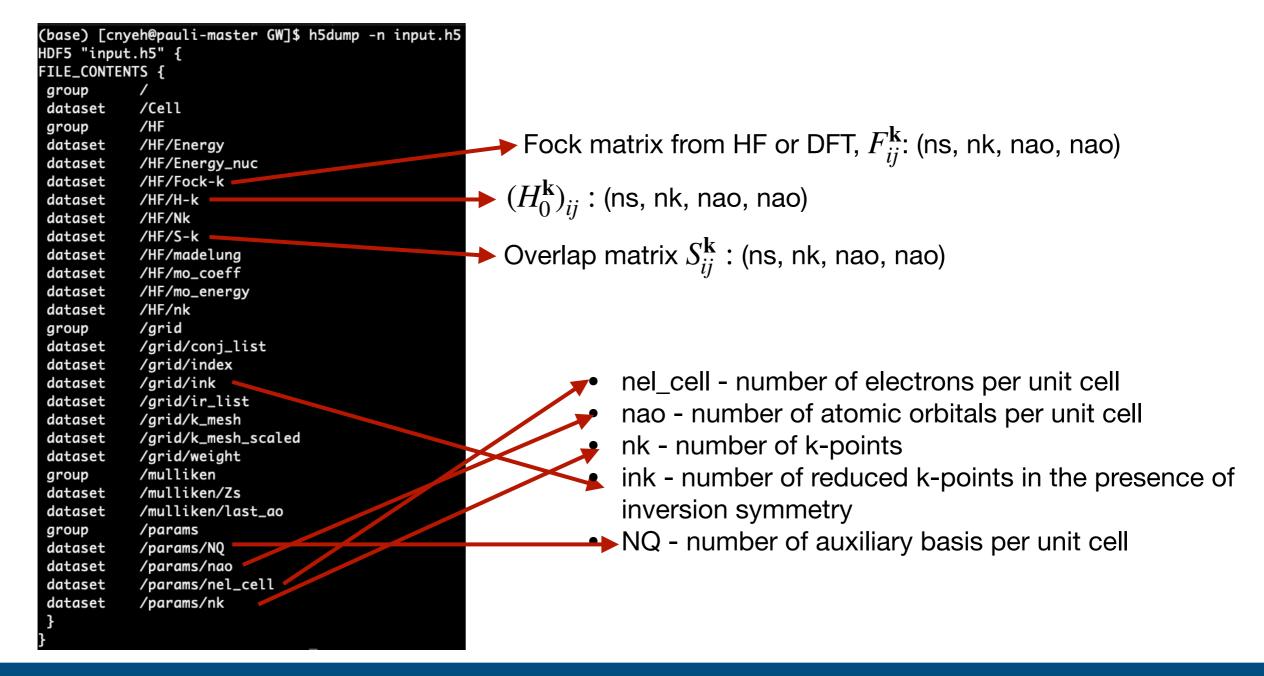
```
nel_cell=8
nao=26
nk=216
ink=112
ns=1
NQ = 124
ni=168
beta=1000
itermax=15
rst=false
scf_type=GW
IR=true
CONST_DENSITY=true
E_thr=1e-5
damp=0.7
DIIS_size=4
DIIS_start=4
DIIS_interval=1
```

Example: Silicon

• Basis set: gth-dzvp-molopt-sr, Pseudopotential: gth-pbe

input.h5:

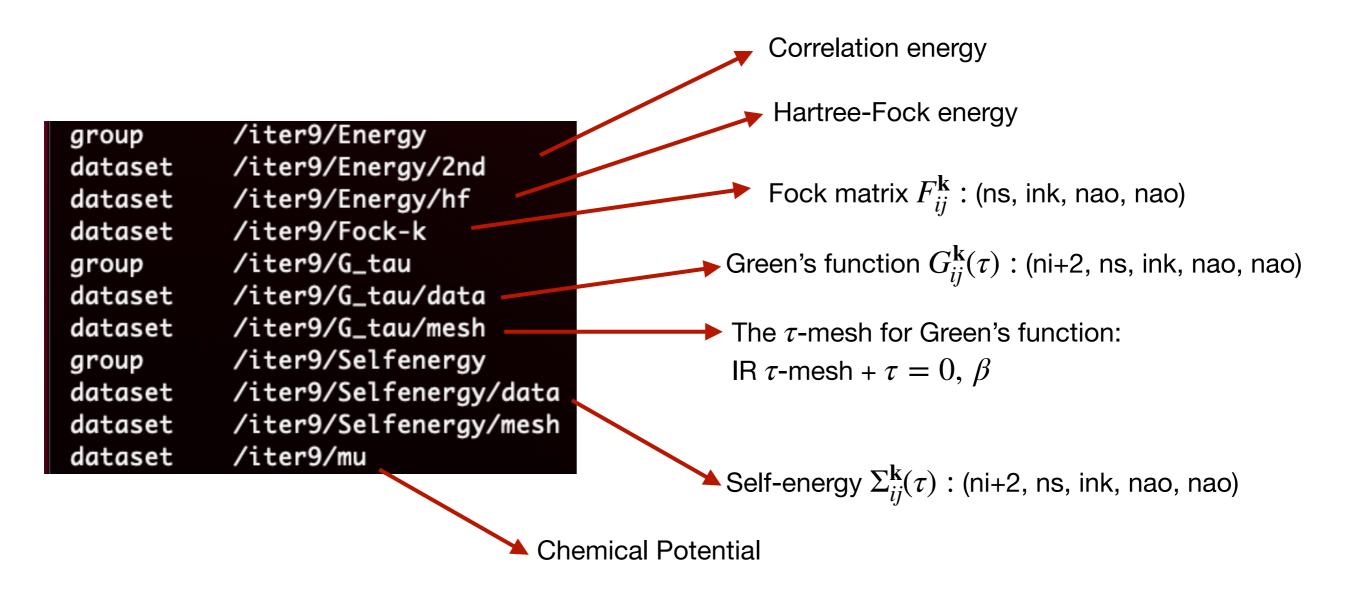
The output file after computing the matrix elements using UGF2/script/init_data_df.py. Most of the GW parameters can be found here.



Example: Silicon

• Basis set: gth-dzvp-molopt-sr, Pseudopotential: gth-pbe

sim.h5: UGF2 output file



Comments

- Define the problem
 - Gaussian basis set → basis set convergence
 - Size of *k*-mesh: → finite-size effect
 Typically, 6x6x6 is good enough for insulator. The larger the unit cell, the less *k*-points needed.
- Compute matrix elements of the Hamiltonian
 - Python scripts using PySCF
 - Density fitting for the two-electron Coulomb interaction

 J. Chem. Phys. 147, 164119 (2017)

 J. Chem. Phys. 154, 131104 (2021)

 → This step could be very time-consuming!
- Run many-body perturbation theory (MBPT) using UGF2
 - GW $\sim O(N_{\tau}N_k^2N_{ao}^4)$
 - 10~20 iterations to reach the self-consistency
 - beta is typically smaller than 1000 (i.e. ~ 315 K)