Peskin Solutions: Chapter 9

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1 Problem 9.1

(a)

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}^2 + (\partial_\mu\phi^* - ieA_\mu\phi^*)(\partial^\mu\phi + ieA^\mu\phi) - m^2\phi^*\phi = \mathcal{L}_A + \mathcal{L}_\phi + \mathcal{L}_I$$

The $\mathcal{L}_A$ is just free E-M field, so the propogator is the propogator of photon.

The $\mathcal{L}_\phi = \partial_\mu\phi^*\partial^\mu\phi - m^2\phi^*\phi = \partial_\mu(\phi^*\partial^\mu\phi) - \phi^*\partial_\mu(\partial^\mu\phi) - m^2\phi^*\phi$, because the differential term in the Lagrangian density makes no difference, we got $\mathcal{L}_\phi = -\phi^*\partial_\mu(\partial^\mu\phi) - m^2\phi^*\phi = \phi^*(-\partial^2 - m^2)\phi = \phi^*\hat{T}\phi$.

With generating functional method, we have $\mathcal{L}_\phi + \eta^*\phi + \phi^*\eta$ in the $Z[J]$, then do a shift $\phi \to \phi' = \phi + \hat{T}^{-1}\eta$, we got $\mathcal{L}_\phi + \eta^*\phi + \phi^*\eta = \mathcal{L}_{\phi'} - \eta^*\hat{T}^{-1}\eta$. If G is the Green function of $\hat{T}$, then $\mathcal{L}_{\phi'} - \eta^*\hat{T}^{-1}\eta = \mathcal{L}_{\phi'} - \eta^*G * \eta$, after two functional derivatives, we will find the propogator is exactly the G.

So the propogator of ϕ and ϕ^* is $\frac{i}{p^2-m^2+i\epsilon}$. (How to calculate the Green function of $\hat{T}$? - Check Eq.(2.57) in Peskin)

Then comes to the vertex, $\mathcal{H}_I = -\mathcal{L}_I$ (P. 289 in Peskin), theoretically we should check Eq.(4.31) and do the contraction to get Feynman rules, but here we can just look at $exp[i\int\mathcal{L}_I]$, here $\mathcal{L}_I = ieg^{\mu\nu}(\partial_\mu\phi^*A_\nu\phi - A_\mu\phi^*\partial_\nu\phi) + e^2g^{\mu\nu}A_\mu\phi^*A_\nu\phi$, then $i\mathcal{L}_I = i*i*ieg^{\mu\nu}(-i\partial_\mu\phi^*A_\nu\phi + A_\mu\phi^*i\partial_\nu\phi) + ie^2g^{\mu\nu}A_\mu\phi^*A_\nu\phi$.

There are three terms, let's throw those fields away and turn $i\partial\phi$ to $p_\phi\phi$, $-i\partial\phi^*$ to $p_{\phi^*}\phi^*$, here ps are along particle/anti-particle lines, besides, the third term has two A fields, which are commutative, so there should be a factor 2 for the $AA\phi^*\phi$ vertex.

So,

$$For\ \phi^*A\phi : -ie(p+p')^\mu$$

$$For\ AA\phi^*\phi : 2ie^2g^{\mu\nu}$$

Theoretically,

$$<\phi\phi^*|S|\gamma> = <\phi\phi^*|T\int d^4x i\mathcal{L}_I|\gamma> = <\phi\phi^*|T\int d^4x(-ie)g^{\mu\nu}(-i\partial_\mu\phi^*A_\nu\phi + A_\mu\phi^*i\partial_\nu\phi)|\gamma>$$

and

$$< \phi\phi^*|S|\gamma\gamma > = < \phi\phi^*|T\int d^4x i\mathcal{L}_I|\gamma\gamma > = < \phi\phi^*|T\int d^4x(ie^2)g^{\mu\nu}A_\mu\phi^*A_\nu\phi|\gamma\gamma >$$

give the Feynman rules of two kinds of vertex with contraction.

(b)

With Eq.(4.84), m_e is ignored, then,

$$(\frac{d\sigma}{d\Omega})_{c.m.} = \frac{|\vec{p}_\phi|}{32(2\pi)^2E_e^2\cdot 2E_e}\frac{1}{4}\Sigma|\mathcal{M}(ee\to\phi^*\phi)|^2$$

The outlines of ϕ and ϕ^* are 1, the Feynman diagram looks similar to the diagram in P.131.

$$i\mathcal{M} = (-ie)^2\bar{v}(k_2)\gamma^\mu u(k_1)\frac{-ig_{\mu\nu}}{s+i\epsilon}(p_1-p_2)^\nu = ie^2\bar{v}(k_2)(\not{p_1}-\not{p_2})u(k_1)\frac{1}{s+i\epsilon}$$

$$\frac{1}{4}\Sigma_{spin}|\mathcal{M}|^2 = \Sigma_{spin}\frac{e^4}{4s^2}\bar{v}(k_2)(\not{p_1}-\not{p_2})u(k_1)\bar{u}(k_1)(\not{p_1}-\not{p_2})v(k_2)$$

$$= \Sigma_{spin}\frac{e^4}{4s^2}tr(v(k_2)\bar{v}(k_2)(\not{p_1}-\not{p_2})u(k_1)\bar{u}(k_1)(\not{p_1}-\not{p_2}))$$

$$= \frac{e^4}{4s^2}tr(\not{k_2}(\not{p_1}-\not{p_2})\not{k_1}(\not{p_1}-\not{p_2}))$$

$$= \frac{e^4}{4s^2}\left[8\left(k_1\cdot p_1 - k_1\cdot p_2\right)\left(k_2\cdot p_1 - k_2\cdot p_2\right) - 4\left(k_1\cdot k_2\right)\left(p_1-p_2\right)^2\right]$$

Choose a specific frame,

$$k_1 = (E,0,0,E), \quad p_1 = (E, p\sin\theta, 0, p\cos\theta)$$

$$k_2 = (E,0,0,-E), \quad p_2 = (E, -p\sin\theta, 0, -p\cos\theta)$$

$$\frac{1}{4}\Sigma_{spin}|\mathcal{M}|^2 = \frac{e^4p^2}{2E^2}\sin^2\theta$$

$ee\to\mu\mu$ is Eq.(5.11)
So,

$$\left(\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega}\right)_{\mathrm{CM}} = \frac{1}{2(2E)^2}\frac{p}{8(2\pi)^2E}\left(\frac{1}{4}\sum|\mathcal{M}|^2\right) = \frac{\alpha^2}{8s}\left(1-\frac{m^2}{E^2}\right)^{3/2}\sin^2\theta$$

(c)

Two diagrams because there are two kinds of vertex which are listed in (a). Because the sign in the $\mathcal{L}_I$ between two vertex terms is +, the order of fermion fields in the $<|S|>$ expression are same, which is $<|\phi^*\phi|>$, so the two diagrams should be added.

$$i\Pi_1^{\mu\nu} = e^2 \int \frac{d^4k}{(2\pi)^4}(2k+q)^\mu \frac{1}{k^2-m^2+i\epsilon}(2k+q)^\nu \frac{1}{(k+q)^2-m^2+i\epsilon}$$

$$i\Pi_2^{\mu\nu} = -2e^2 g^{\mu\nu}\int \frac{d^4k}{(2\pi)^4}\frac{1}{(k+q)^2-m^2+i\epsilon}$$

add togeter, get

$$i\Pi^{\mu\nu} = -e^2\int \frac{d^4k}{(2\pi)^4}\frac{2g^{\mu\nu}(k^2-m^2)-(2k+q)^\mu(2k+q)^\nu}{(k^2-m^2)((k+q)^2-m^2)}$$

$$\frac{1}{(k^2-m^2)((k+q)^2-m^2)} = \int_0^1 dx\frac{1}{[(k+(1-x)q)^2+xq^2-x^2q^2-m^2]^2}$$

change the variable, $l = k+(1-x)q$, with Eq.(6.45),

$$numerator = g^{\mu\nu}l^2 + 2g^{\mu\nu}(1-x)^2q^2 - 2g^{\mu\nu}m^2 - (2x-1)^2q^\mu q^\nu$$

do the Wick rotation, $l^0 = il_E^0$ and $l^i = l_E^i$, so we have $d^4l = id^4l_E$ and $l^2 = -l_E^2$,

$$i\Pi^{\mu\nu} = -ie^2\int_0^1 dx\int\frac{d^4l_E}{(2\pi)^4}\frac{-g^{\mu\nu}l_E^2+2g^{\mu\nu}(1-x)^2q^2-2g^{\mu\nu}m^2-(2x-1)^2q^\mu q^\nu}{[l_E^2+m^2+x^2q^2-xq^2]^2}$$

$$= -ie^2\int_0^1 dx\int\frac{d^4l_E}{(2\pi)^4}[\frac{-g^{\mu\nu}l_E^2}{(l_E^2+\Delta)^2}+\frac{2g^{\mu\nu}(1-x)^2q^2-2g^{\mu\nu}m^2-(2x-1)^2q^\mu q^\nu}{(l_E^2+\Delta)^2}]$$

use dimensional regularization, with Eq.(7.85) and Eq.(7.86),

$$i\Pi^{\mu\nu} = -ie^2\int_0^1 dx[(2g^{\mu\nu}(1-x)^2q^2-2g^{\mu\nu}m^2-(2x-1)^2q^\mu q^\nu)\frac{1}{(4\pi)^{d/2}}\frac{\Gamma\left(2-\frac{d}{2}\right)}{\Gamma(2)}\left(\frac{1}{\Delta}\right)^{2-\frac{d}{2}}$$

$$-g^{\mu\nu}\frac{1}{(4\pi)^{d/2}}\frac{d}{2}\frac{\Gamma\left(2-\frac{d}{2}-1\right)}{\Gamma(2)}\left(\frac{1}{\Delta}\right)^{2-\frac{d}{2}-1}]$$

$$= -ie^2\int_0^1 dx\frac{1}{(4\pi)^{d/2}}(\frac{1}{\Delta})^{2-d/2}\Gamma(2-d/2)[(2g^{\mu\nu}(1-x)^2q^2-(2x-1)^2q^\mu q^\nu)-g^{\mu\nu}\frac{d}{2-d}(x^2q^2-xq^2)]$$

set $d = 4 - \epsilon$ with $\epsilon \to 0$,

$$i\Pi^{\mu\nu} = \frac{-ie^2}{(4\pi)^2}\int_0^1 dx(\frac{\epsilon}{2} - log\Delta - \gamma + log(4\pi))[(g^{\mu\nu}(2x-2)(2x-1)q^2 - (2x-1)^2 q^\mu q^\nu)]$$

Because $\int_0^1 dx(\frac{2}{\epsilon} - log\Delta - \gamma + log(4\pi))(2x-1) = \int_0^1 dx\frac{2}{\epsilon}(2x-1) = 0$, we have

$$i\Pi^{\mu\nu} = \frac{-ie^2}{(4\pi)^2}\int_0^1 dx(\frac{\epsilon}{2} - log\Delta - \gamma + log(4\pi))(2x-1)^2[(g^{\mu\nu}q^2 - q^\mu q^\nu)]$$

with MS-bar scheme,

$$\Pi(q^2) = \frac{-\alpha}{4\pi}\int_0^1 dx(-log\Delta)(2x-1)^2$$

If we adopt $-q^2 >> m^2$,

$$\Pi(q^2) = \frac{-\alpha}{4\pi}\int_0^1 dx(-log(x-x^2) - log(-q^2))(2x-1)^2 \to \frac{-\alpha}{12\pi}log(-q^2)$$

while looking at Eq.(7.90), $\int_0^1 dxx(1-x) = \frac{1}{6}$, we know for $e+e-$ pair,

$$\Pi(q^2) \to \frac{-\alpha}{3\pi}log(-q^2)$$

which is four times as our results.

2 Problem 9.1

(a)