$$Q_{chsubs0} = Q_0 \exp(\frac{V_{gs} - V_{th}}{nv_t})$$
(3.1.1a)

where Q_0 is

$$Q_0 = \sqrt{\frac{q\varepsilon_{si}N_{ch}}{2\phi_s}}v_t \exp(-\frac{V_{off}}{nv_t})$$
(3.1.1b)

$$Q_{chs0} = C_{ox}(V_{gs} - V_{th})$$
(3.1.2)

In both Eqs. (3.1.1a) and (3.1.2), the parameters $Q_{chsubs0}$ and Q_{chs0} are the channel charge densities at the source for very small Vds. To form a unified expression, an effective $(V_{gs}-V_{th})$ function named V_{gsteff} is introduced to describe the channel charge characteristics from subthreshold to strong inversion

$$V_{gsteff} = \frac{2 n v_t \ln \left[1 + \exp(\frac{V_{gs} - V_{th}}{2 n v_t}) \right]}{1 + 2 n Cox \sqrt{\frac{2\Phi_s}{q\varepsilon_{si}N_{ch}}} \exp(-\frac{V_{gs} - V_{th} - 2V_{off}}{2 n v_t})}$$
(3.1.3)

The unified channel charge density at the source end for both subthreshold and inversion region can therefore be written as

$$Q_{chs0} = C_{ox}V_{gsteff}$$
 (3.1.4)

Figures 3-1 and 3-2 show the smoothness of Eq. (3.1.4) from subthreshold to strong inversion regions. The V_{gsteff} expression will be used again in subsequent sections of this chapter to model the drain current.