Anisotropic Math

Anton Kodochygov

September 1, 2015

► Eikonal equation (isotropic):

$$|\nabla T| = |p| = U$$

Eikonal equation (isotropic):

$$|\nabla T| = |p| = U$$

Eikonal equation (anisotropic):

$$|\nabla T| = |p| = U(\xi)$$

Eikonal equation (isotropic):

$$|\nabla T| = |p| = U$$

Eikonal equation (anisotropic):

$$|\nabla T| = |p| = U(\xi)$$

Snell's Law (isotropic):

$$sin(\theta_1)/v_1 = sin(\theta_2)/v_2$$

Eikonal equation (isotropic):

$$|\nabla T| = |p| = U$$

Eikonal equation (anisotropic):

$$|\nabla T| = |p| = U(\xi)$$

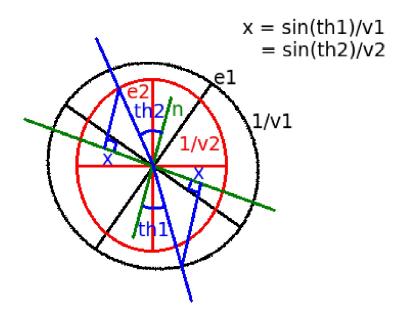
Snell's Law (isotropic):

$$sin(\theta_1)/v_1 = sin(\theta_2)/v_2$$

Snell's Law (anisotropic):

$$\sin(\theta_1)/v_1(\xi_1) = \sin(\theta_2)/v_2(\xi_2)$$

Snell's Law Intuition



Newton's Second Law:

$$\rho \frac{\partial^2 u_i}{\partial t^2} \ = \ c_{ijkl} \frac{\partial^2 u_i}{\partial x_j \partial x_l}$$

Newton's Second Law:

$$\rho \frac{\partial^2 u_i}{\partial t^2} = c_{ijkl} \frac{\partial^2 u_i}{\partial x_j \partial x_l}$$

Christoffel Matrix:

$$G_{ik} = c_{ijkl} n_l n_j$$

Newton's Second Law:

$$\rho \frac{\partial^2 u_i}{\partial t^2} = c_{ijkl} \frac{\partial^2 u_i}{\partial x_j \partial x_l}$$

Christoffel Matrix:

$$G_{ik} = c_{ijkl}n_ln_j$$

Phase Velocities (eigenvectors of G):

$$(G_{ik} - \rho v^2 \delta_{ik}) \bar{U}_k = 0$$



Stress Tensor Simplification

Stress Tensor Simplification

► The 81 element tensor can be represented by a 6x6 matrix

```
\begin{array}{ccccc} c_{ijkl} & \rightarrow & C_{IJ} \\ 11 & \rightarrow & 1 \\ 22 & \rightarrow & 2 \\ 33 & \rightarrow & 3 \\ 23 & \rightarrow & 4 \\ 13 & \rightarrow & 5 \\ 12 & \rightarrow & 6 \end{array}
```

Stress Tensor Simplification

The 81 element tensor can be represented by a 6x6 matrix

$$\begin{array}{cccc} c_{ijkl} & \rightarrow & C_{IJ} \\ 11 & \rightarrow & 1 \\ 22 & \rightarrow & 2 \\ 33 & \rightarrow & 3 \\ 23 & \rightarrow & 4 \\ 13 & \rightarrow & 5 \\ 12 & \rightarrow & 6 \end{array}$$

Furthermore, due to symmetry, we can further simplify this matrix

$$C = \left(\begin{array}{cccccc} C_{11} & C_{12} & C_{13} & 0 & 0 & 0 \\ C_{12} & C_{11} & C_{13} & 0 & 0 & 0 \\ C_{13} & C_{13} & C_{33} & 0 & 0 & 0 \\ 0 & 0 & 0 & C_{44} & 0 & 0 \\ 0 & 0 & 0 & 0 & C_{44} & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{C_{11} - C_{12}}{2} \end{array} \right)$$

Exact Phase Velocities (transverse anisotropy)

Exact Phase Velocities (transverse anisotropy)

$$V_{qP}(\xi) = \sqrt{\frac{C_{11}sin^{2}(\xi) + C_{33}cos^{2}(\xi) + C_{44} + \sqrt{M(\xi)}}{2\rho}}$$

$$V_{qS}(\xi) = \sqrt{\frac{C_{11}sin^{2}(\xi) + C_{33}cos^{2}(\xi) + C_{44} - \sqrt{M(\xi)}}{2\rho}}$$

$$V_{S}(\xi) = \sqrt{\frac{C_{66}sin^{2}(\xi) + C_{44}cos^{2}(\xi)}{\rho}}$$

$$M(\xi) = \left[(C_{11} - C_{44})sin^{2}(\xi) - (C_{33} - C_{44})cos^{2}(\xi) \right]^{2} + (C_{13} + C_{44})^{2}sin^{2}(2\xi)$$

Thomsen's Approximation (Weak Anisotropy $\delta, \gamma, \epsilon << 1$)

Thomsen's Approximation (Weak Anisotropy $\delta, \gamma, \epsilon << 1$)

$$\begin{array}{lcl} V_{qP}(\xi) & \approx & V_{P0}(1+\delta sin^2(\xi)cos^2(\xi)+\epsilon sin^4(\xi)) \\ \\ V_{qS}(\xi) & \approx & V_{S0} \left[1+\left(\frac{V_{P0}}{V_{S0}}\right)^2(\epsilon-\delta)sin^2(\xi)cos^2(\xi)\right] \\ \\ V_{S}(\xi) & \approx & V_{S0}(1+\gamma sin^2(\xi)) \end{array}$$

Thomsen's Approximation (Weak Anisotropy $\delta, \gamma, \epsilon << 1$)

$$\begin{array}{lcl} V_{qP}(\xi) & \approx & V_{P0}(1+\delta sin^2(\xi)cos^2(\xi)+\epsilon sin^4(\xi)) \\ \\ V_{qS}(\xi) & \approx & V_{S0}\left[1+\left(\frac{V_{P0}}{V_{S0}}\right)^2(\epsilon-\delta)sin^2(\xi)cos^2(\xi)\right] \\ \\ V_{S}(\xi) & \approx & V_{S0}(1+\gamma sin^2(\xi)) \end{array}$$

$$\epsilon = \frac{C_{11} - C_{33}}{2C_{33}}$$

$$\delta = \frac{(C_{13} + C_{44})^2 - (C_{33} - C_{44})^2}{2C_{33}(C_{33} - C_{44})}$$

$$\gamma = \frac{C_{66} - C_{44}}{2C_{44}}$$

$$V_{P0} = \sqrt{\frac{C_{33}}{\rho}}$$

$$V_{S0} = \sqrt{\frac{C_{44}}{\rho}}$$

Hamiltonian

$$H = \left(\frac{|\nabla T|^2}{U^2} - 1\right)$$

Hamiltonian

$$H = \left(\frac{|\nabla T|^2}{U^2} - 1\right)$$

Þ

$$\frac{d\vec{x}}{ds} = \frac{\partial H}{\partial \vec{p}} = \frac{\vec{p}}{U^2}$$

Hamiltonian

$$H = \left(\frac{|\nabla T|^2}{U^2} - 1\right)$$

>

$$\frac{d\vec{x}}{ds} = \frac{\partial H}{\partial \vec{p}} = \frac{\vec{p}}{U^2}$$

$$\frac{d\vec{p}}{ds} = \frac{\partial H}{\partial \vec{x}} = \frac{|\vec{p}|^2 \nabla U}{U^3}$$

Hamiltonian

$$H = \left(\frac{|\nabla T|^2}{U^2} - 1\right)$$

Þ

$$\frac{d\vec{x}}{ds} = \frac{\partial H}{\partial \vec{p}} = \frac{\vec{p}}{U^2}$$

$$\frac{d\vec{p}}{ds} = \frac{\partial H}{\partial \vec{x}} = \frac{|\vec{p}|^2 \nabla U}{U^3}$$

$$\frac{dT}{ds} = |\vec{p}\frac{\partial H}{\partial \vec{p}}|$$

Hamiltonian

$$H = \left(\frac{|\nabla T|^2}{U(\xi)^2} - 1\right)$$

$$= \left(\frac{|\nabla T|^2}{U_0^2} (1 + \delta \cos^2(\xi) \sin^2(\xi) + \epsilon \sin^4(\xi))^2 - 1\right)$$

$$= \left(\frac{|\nabla T|^2}{U_0^2} f(\xi) - 1\right)$$

$$\frac{d\vec{x}}{ds} = \frac{\partial H}{\partial \vec{p}}$$

$$= \frac{\vec{p}}{U^2} f(\xi) + \frac{1}{2} \frac{|\nabla T|^2}{U_0^2} \frac{\partial f(\xi)}{\partial \vec{p}}$$

$$= \frac{\vec{p}}{U^2} f(\xi)$$

$$\frac{d\vec{p}}{ds} = \frac{\partial H}{\partial \vec{x}}
= \frac{|\vec{p}|^2 \nabla U}{U^3} f(\xi) + \frac{1}{2} \frac{|\nabla T|^2}{U_0^2} \frac{\partial f(\xi)}{\partial \vec{x}}$$

$$\frac{d\vec{p}}{ds} = \frac{\partial H}{\partial \vec{x}}
= \frac{|\vec{p}|^2 \nabla U}{U^3} f(\xi) + \frac{1}{2} \frac{|\nabla T|^2}{U_0^2} \frac{\partial f(\xi)}{\partial \vec{x}}$$

$$\begin{split} \frac{\partial f(\xi)}{\partial \vec{x}} &= \frac{\partial}{\partial \vec{x}} (1 + \delta \sin^2(\xi) \cos^2(\xi) + \epsilon \sin^4(\xi))^2 \\ &= 2(1 + \delta \sin^2(\xi) \cos^2(\xi) + \epsilon \sin^4(\xi)) \\ &[(\delta(\sin^2(\xi) - \cos^2(\xi)) - 2\epsilon \sin^2(\xi)) \frac{\partial \cos^2(\xi)}{\partial \vec{x}} \\ &+ \sin^2(\xi) \cos^2(\xi) \frac{\partial \delta}{\partial \vec{x}} \\ &+ \sin^4(\xi) \frac{\partial \epsilon}{\partial \vec{x}} \\ \end{bmatrix} \end{split}$$

▶

$$\xi = \mathit{angle}(\vec{\eta}, \vec{p})$$

▶

$$\xi = angle(\vec{\eta}, \vec{p})$$

$$ec{\eta} \cdot ec{p} = |ec{\eta}| |ec{p}| cos(\xi)$$

▶

$$\xi = angle(\vec{\eta}, \vec{p})$$

$$ec{\eta} \cdot ec{p} = |ec{\eta}| |ec{p}| cos(\xi)$$

$$cos^2(\xi) = \left(rac{ec{\eta}\cdotec{p}}{|ec{\eta}||ec{p}|}
ight)^2$$

$$\xi = \mathit{angle}(\vec{\eta}, \vec{p})$$

$$ec{\eta} \cdot ec{p} = |ec{\eta}| |ec{p}| cos(\xi)$$

$$cos^2(\xi) = \left(rac{ec{\eta}\cdotec{p}}{|ec{\eta}||ec{p}|}
ight)^2$$

$$\frac{\partial \cos^2(\xi)}{\partial \vec{x}} = 2 \frac{|\vec{\eta}|^2 |\vec{p}|^2 (\vec{\eta} \cdot \vec{p}) (\frac{\partial \vec{\eta}}{\partial \vec{x}} \cdot \vec{p}) - (\vec{\eta} \cdot \vec{p})^2 |\vec{p}|^2 (\frac{\partial \vec{\eta}}{\partial \vec{x}} \cdot \vec{\eta})}{|\vec{\eta}|^4 |\vec{p}|^4}$$

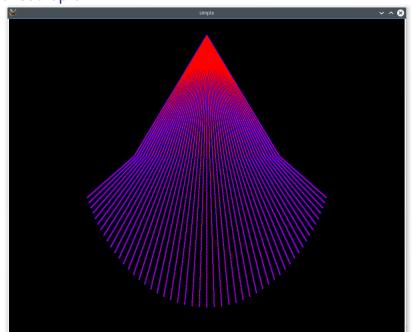
Equations of Motion (anisotropic, last but not least)

Equations of Motion (anisotropic, last but not least)

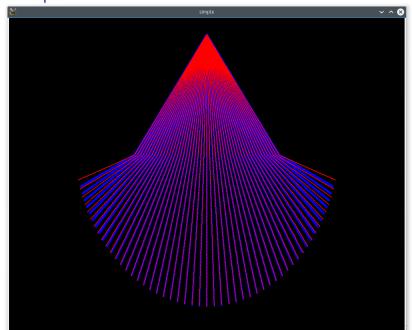
▶

$$\frac{dT}{ds} = |\vec{p}\frac{\partial H}{\partial \vec{p}}|$$

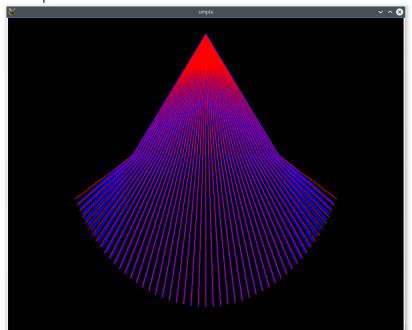
Flat Isotropic



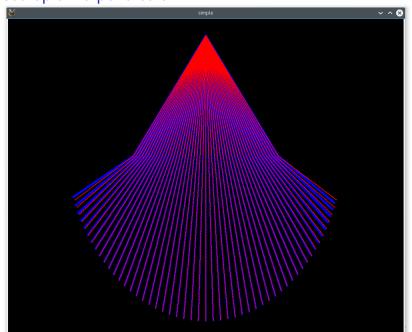
Anisotropic Bottom



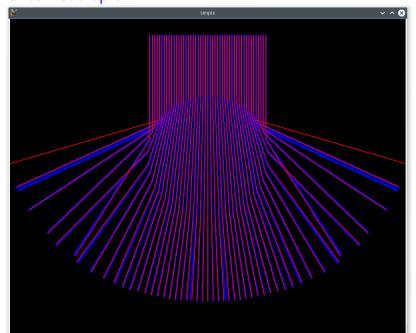
Anisotropic Parallel



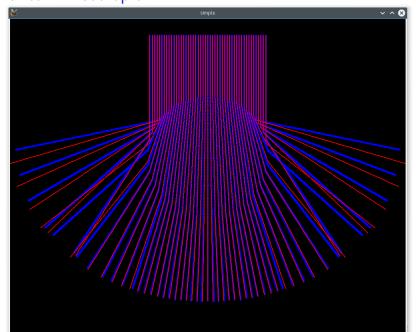
Anisotropic Perpendicular



Spherical Isotropic



Spherical Anisotropic



▶ Isotropic tests seem to closely match ground truth output.

- ▶ Isotropic tests seem to closely match ground truth output.
- Anisotropic tests seem to match ground truth output for small incident angles, but have problems with more severe incident angles.

- ▶ Isotropic tests seem to closely match ground truth output.
- Anisotropic tests seem to match ground truth output for small incident angles, but have problems with more severe incident angles.
 - ▶ This may be due to anisotropic parameter tapering.

- ▶ Isotropic tests seem to closely match ground truth output.
- Anisotropic tests seem to match ground truth output for small incident angles, but have problems with more severe incident angles.
 - This may be due to anisotropic parameter tapering.
 - Alternatively, this may be due to a mistake in my math.

- ▶ Isotropic tests seem to closely match ground truth output.
- Anisotropic tests seem to match ground truth output for small incident angles, but have problems with more severe incident angles.
 - This may be due to anisotropic parameter tapering.
 - Alternatively, this may be due to a mistake in my math.
 - ▶ I have a different anisotropic functor, one based exclusively on Mr. Yingst notes, I'll try it to see if the output improves.

- Isotropic tests seem to closely match ground truth output.
- Anisotropic tests seem to match ground truth output for small incident angles, but have problems with more severe incident angles.
 - This may be due to anisotropic parameter tapering.
 - ▶ Alternatively, this may be due to a mistake in my math.
 - ▶ I have a different anisotropic functor, one based exclusively on Mr. Yingst notes, I'll try it to see if the output improves.
 - ▶ At the end of the day, we may consider just using the ground truth solver for extreme transitions.