

## Lecture 9: The Big Bang, the horizon and space-time diagrams of the expanding Universe.

**L:** The Friedmann-Robertson-Walker metric is:

$$-ds^2 = c^2 dt^2 - a^2(t) [dr^2 + S_k^2(r)(d\theta^2 + \sin^2 \theta d\phi^2)] . \quad (1)$$

**Q:** What is the history of the expansion  $a(t)$ ? Draw the function  $a(t)$  versus  $H_0 t$ , in a diagram where we fix at the present time,  $a(t_0) = 1$  and  $\dot{a}(t_0) = H_0$ .

**L:** Simple examples:  $\Omega_m = 1$ ,  $a(t) = (3H_0 t/2)^{2/3}$ . Empty Universe,  $\Omega = 0$ :  $a(t) = H_0 t$ . Radiation dominated Universe:  $a(t) = (2H_0 t)^{1/2}$ . Draw also the cases for the open and closed Universe with no radiation or dark energy.

**Q:** For the same value of  $H_0$ , which Universe is older and younger, the open or the closed one? The matter or the radiation dominated one?

**Q:** What happens for the Benchmark model with  $\Omega_{m0} = 0.3$ ,  $\Omega_{\Lambda 0} = 0.7$ ? Is this Universe older or younger than without the cosmological constant?

**L:** Now, we want to draw a space-time diagram of the Universe, which can be conveniently reduced to two dimensions if we eliminate the  $\theta, \phi$  coordinates, so we only look at radial trajectories. Every point in this  $r - t$  diagram represents a whole spherical surface. We would like light rays to be straight lines moving at an angle of 45 degrees, and for that we use units like second and light-second.

**L:** Using comoving coordinates in this diagram is very useful, then all galaxies following the Hubble flow are represented as world lines moving vertically up in cosmic time. But we have a problem: the time-varying scale factor  $a(t)$  means that light-rays no longer move at 45 degrees because, in a light second, they move over 1 comoving light year divided by  $a(t)$ .

**L:** To solve this, we define conformal time, a kind of “comoving time”:  $d\eta = dt/a(t)$ . So now, let us draw the whole history of the Universe in a  $c\eta - r$  diagram. Light rays always move at 45 degrees in this diagram.

Note: these are not exactly like Penrose diagrams, for that you need to define new coordinates to bring infinity into a finite space-time map.

**Q:** What is  $\eta(t)$  in the matter-only Universe?

$$d\eta = (t/t_0)^{-2/3} dt; \quad \eta = 3t_0 (t/t_0)^{1/3} = \frac{2}{H_0} a^{1/2} . \quad (2)$$

**L:** In general,  $d\eta = da/(Ha^2)$ .

**Q:** What is  $\eta$  for the radiation-dominated epoch of the Universe? We use  $H \simeq H_0 \Omega_{r0}^{1/2}/a^2$  (assuming as usual  $a_0 = 1$  at present, and so  $a = 1/(1+z)$ ), and  $t = 1/(2H)$ , to find

$$\eta = \frac{a}{H_0 \Omega_{r0}^{1/2}} = \left( \frac{2t}{H_0} \right)^{1/2} \Omega_{r0}^{1/4} . \quad (3)$$

**Q:** What happens in an open model with only matter and curvature?

$$d\eta = \frac{1}{H_0} \frac{da}{a(\Omega_0/a + 1 - \Omega_0)^{1/2}} . \quad (4)$$

For very large  $a$  (in the long-term future, when matter becomes highly diluted and negligible compared to curvature), we can approximate the integral to:

$$\eta \simeq \frac{1}{H_0 \sqrt{\Omega_0}} \log a + \text{constant} . \quad (5)$$

Note that for large times,  $a \sim t$ .

In all these cases, as  $t$  goes to zero (Big Bang),  $\eta$  goes to zero. As  $t$  goes to infinity,  $\eta$  goes to infinity. When looking to the past, we can only see out to the horizon; to see the Universe further, we have to wait a long time, although we'll be able to see as far as we want if we just wait long enough.

**Q:** The Big Bang is not an event. It is a beginning that took place simultaneously in every point of space. This brings up the horizon problem: why was the Universe so close to homogeneous, without causal communication? And yet, not exactly homogeneous, in such a way that present galaxies formed? In the Big Bang model without anything else, the only answer is that God simply made the Universe like that.

**L:** Now we look at the Benchmark model, which contains dark energy.

$$d\eta = \frac{1}{H_0} \frac{da}{a^2(\Omega_0/a^3 + 1 - \Omega_0)^{1/2}} . \quad (6)$$

For large  $a$ , the result of integrating is:

$$\eta \simeq \eta_{\max} - \frac{1}{H_0 \sqrt{1 - \Omega_0} a} . \quad (7)$$

As a function of time, we can reexpress this as:

$$dt \simeq \frac{da}{H_0 \sqrt{1 - \Omega_0} a} ; \quad t \simeq \frac{1}{H_0 \sqrt{1 - \Omega_0}} (\log a + \text{constant}) ; \quad (8)$$

$$\eta \simeq \eta_{\max} - \frac{\exp[-H_0 \sqrt{1 - \Omega_0} (t - t_1)]}{H_0 \sqrt{1 - \Omega_0}} . \quad (9)$$

This diagram looks quite different. The infinitely far future corresponds to a finite conformal time  $\eta_{\max}$ . Without any matter (only cosmological constant), as  $a$  becomes small  $\eta$  goes to negative infinity (there is no Big Bang). With finite matter, there is a Big Bang but then also a maximum  $\eta_{\max}$  which is never exceeded.