## Particle creation by an external field

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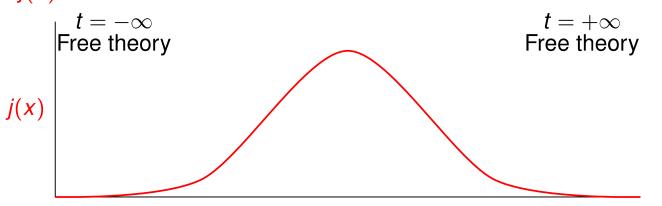
In the presence of an external source field:

$$(\partial_{\mu}\partial^{\mu} + m^2)\phi(x) = j(x) \tag{1}$$

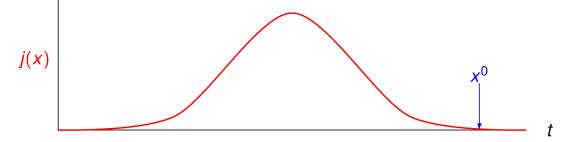
which derives from the Lagrangian:

$$\mathcal{L} = rac{1}{2}(\partial^{\mu}\phi\partial_{\mu}\phi - m^{2}\phi^{2}) + j(x)\phi(x)$$

If j(x) is active for a finite time:



t



- $t \to \infty$  and  $t \to -\infty$ : free Lagrangian  $\mathcal{L}_0 = \frac{1}{2} (\partial^\mu \phi \partial_\mu \phi m^2 \phi^2)$ 
  - $\Rightarrow$  free-field solution:  $\phi_0(x) = \int \frac{\mathrm{d}^3 p}{(2\pi)^3 \sqrt{2E_p}} (a_p e^{-ipx} + a_p^{\dagger} e^{ipx})$
  - ⇒ usual particle interpretation
- Non-homogeneous differential equation (1)
  - ⇒ Use retarded Green's function:

$$\phi(x) = \phi_0(x) + i \int d^4 y D_R(x - y) j(y) 
= \phi_0(x) + i \int d^4 y \int \frac{d^3 p}{(2\pi)^3 2E_p} \Theta(x^0 - y^0) (e^{-ip(x-y)} - e^{ip(x-y)}) j(y)$$

- Wait until a time  $x^0$  in which  $j(x^0, \mathbf{x}) = 0$ 
  - $\Rightarrow$  j(x) is all in the past  $\Rightarrow$   $x^0 > y^0$ :  $\Theta(x^0 y^0) = 1$
- Define: Fourier transform of j(x):  $\tilde{j}(p) = \int d^4y \, e^{ipy} j(y)$ ;  $[p^2 = m^2]$

$$\phi(x) = \int \frac{\mathrm{d}^3 p}{(2\pi)^3 \sqrt{2E_p}} \left\{ (a_p + \frac{i}{\sqrt{2E_p}} \tilde{j}(p)) e^{-ipx} + (a_p^{\dagger} - \frac{i}{\sqrt{2E_p}} \tilde{j}^{\dagger}(p)) e^{ipx} \right\}$$

Hamiltonian after j(x):

$$H = \int rac{\mathrm{d}^3 p}{(2\pi)^3} E_p(a_p^\dagger - rac{i}{\sqrt{2E_p}} \tilde{j}^\dagger(p)) (a_p + rac{i}{\sqrt{2E_p}} \tilde{j}(p))$$

vacuum energy:

$$\langle 0|H|0
angle = \int rac{\mathrm{d}^3 p}{(2\pi)^3} rac{1}{2} | ilde{m{j}}(m{p})|^2$$

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## Interpretation

• The source j(x) has created particles with momentum p with a probability density

$$\frac{|\tilde{j}(p)|^2}{2E_p}$$

The total number of created particles is:

$$\int \mathrm{d}N = \int \frac{\mathrm{d}^3 p}{(2\pi)^3 2 E_p} |\tilde{\boldsymbol{j}}(\boldsymbol{p})|^2$$

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