

Particle creation by an external field

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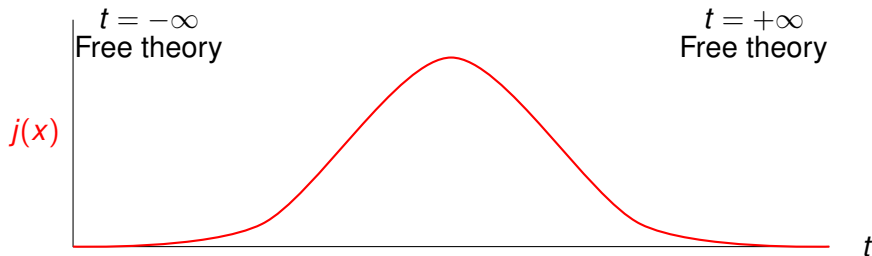
In the presence of an external source field:

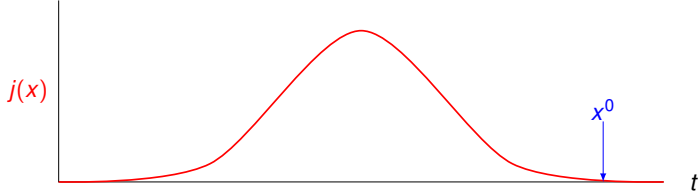
$$(\partial_\mu \partial^\mu + m^2)\phi(x) = j(x) \quad (1)$$

which derives from the Lagrangian:

$$\mathcal{L} = \frac{1}{2}(\partial^\mu \phi \partial_\mu \phi - m^2 \phi^2) + j(x)\phi(x)$$

If $j(x)$ is active for a finite time:





- $t \rightarrow \infty$ and $t \rightarrow -\infty$: free Lagrangian $\mathcal{L}_0 = \frac{1}{2}(\partial^\mu \phi \partial_\mu \phi - m^2 \phi^2)$
 \Rightarrow free-field solution: $\phi_0(x) = \int \frac{d^3 p}{(2\pi)^3 \sqrt{2E_p}} (a_p e^{-ipx} + a_p^\dagger e^{ipx})$
 \Rightarrow usual particle interpretation
- Non-homogeneous differential equation (1)
 \Rightarrow Use retarded Green's function:

$$\begin{aligned} \phi(x) &= \phi_0(x) + i \int d^4 y D_R(x-y) j(y) \\ &= \phi_0(x) + i \int d^4 y \int \frac{d^3 p}{(2\pi)^3 2E_p} \Theta(x^0 - y^0) (e^{-ip(x-y)} - e^{ip(x-y)}) j(y) \end{aligned}$$

- Wait until a time x^0 in which $j(x^0, \mathbf{x}) = 0$
 $\Rightarrow j(x)$ is all in the past $\Rightarrow x^0 > y^0$: $\Theta(x^0 - y^0) = 1$

- Define: Fourier transform of $j(x)$: $\tilde{j}(p) = \int d^4y e^{ipy} j(y)$; $[p^2 = m^2]$

$$\phi(x) = \int \frac{d^3p}{(2\pi)^3 \sqrt{2E_p}} \left\{ (a_p + \frac{i}{\sqrt{2E_p}} \tilde{j}(p)) e^{-ipx} + (a_p^\dagger - \frac{i}{\sqrt{2E_p}} \tilde{j}^\dagger(p)) e^{ipx} \right\}$$

- Hamiltonian after $j(x)$:

$$H = \int \frac{d^3p}{(2\pi)^3} E_p (a_p^\dagger - \frac{i}{\sqrt{2E_p}} \tilde{j}^\dagger(p)) (a_p + \frac{i}{\sqrt{2E_p}} \tilde{j}(p))$$

- vacuum energy:

$$\langle 0|H|0\rangle = \int \frac{d^3p}{(2\pi)^3} \frac{1}{2} |\tilde{j}(p)|^2$$

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Interpretation

- The source $j(x)$ has created particles with momentum \mathbf{p} with a probability density

$$\frac{|\tilde{j}(\mathbf{p})|^2}{2E_p}$$

- The total number of created particles is:

$$\int dN = \int \frac{d^3p}{(2\pi)^3 2E_p} |\tilde{j}(\mathbf{p})|^2$$