

Decaying dark matter: cosmological constraints and implications for the S_8 tension

Guillermo Franco Abellán
Laboratoire Univers et Particules de Montpellier

Based on:

arXiv:2102.12498 (PRD in press)

arXiv:2008.09615 (PRD in press)

with Riccardo Murgia, Vivian Poulin and Julien Lavalle



Tensions in cosmology

With the era of precision cosmology, several discrepancies have emerged

- S_8 with weak-lensing data
[KiDS-1000 2007.15632](#)
- H_0 with local measurements
[Riess++ 2012.08534](#)

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Unaccounted systematics?

- Less exotic explanation
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Physics beyond Λ CDM?

- Reveal properties about the dark sector
- Very challenging

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- S_8 with weak-lensing data
[KiDS-1000 2007.15632](#)  **This talk**
- H_0 with local measurements
[Riess++ 2012.08534](#)  **My previous works:**
[Murgia, GFA, Poulin 2009.10733](#)
[Schöneberg, GFA++ 2107.10291](#)

Unaccounted systematics?

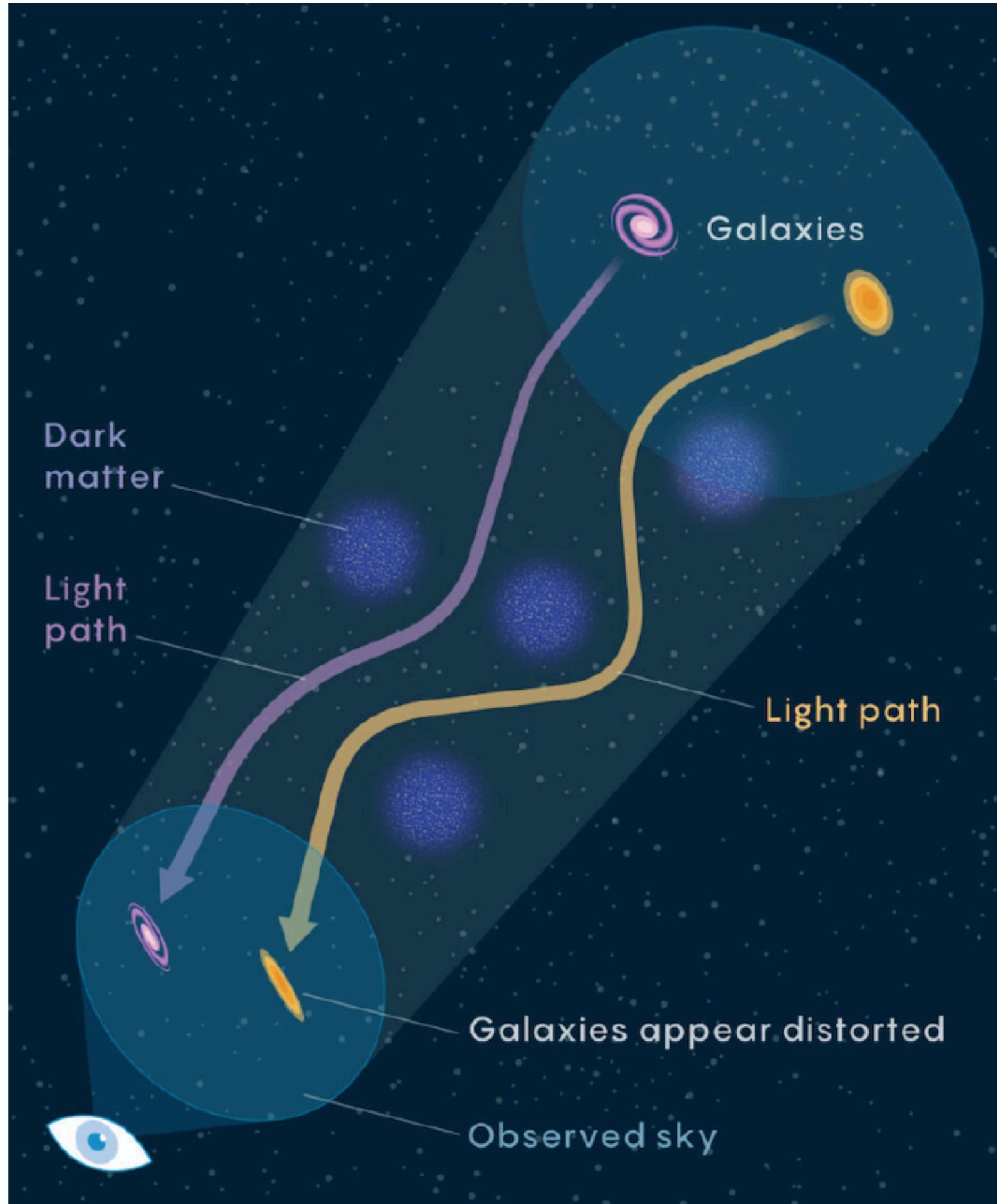
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The S_8 tension

Weak-lensing surveys are mainly sensible to $S_8 \equiv \sigma_8 \sqrt{\Omega_m/0.3}$



KiDS+BOSS+2dfLenS*:

$$S_8 = 0.766^{+0.020}_{-0.014}$$

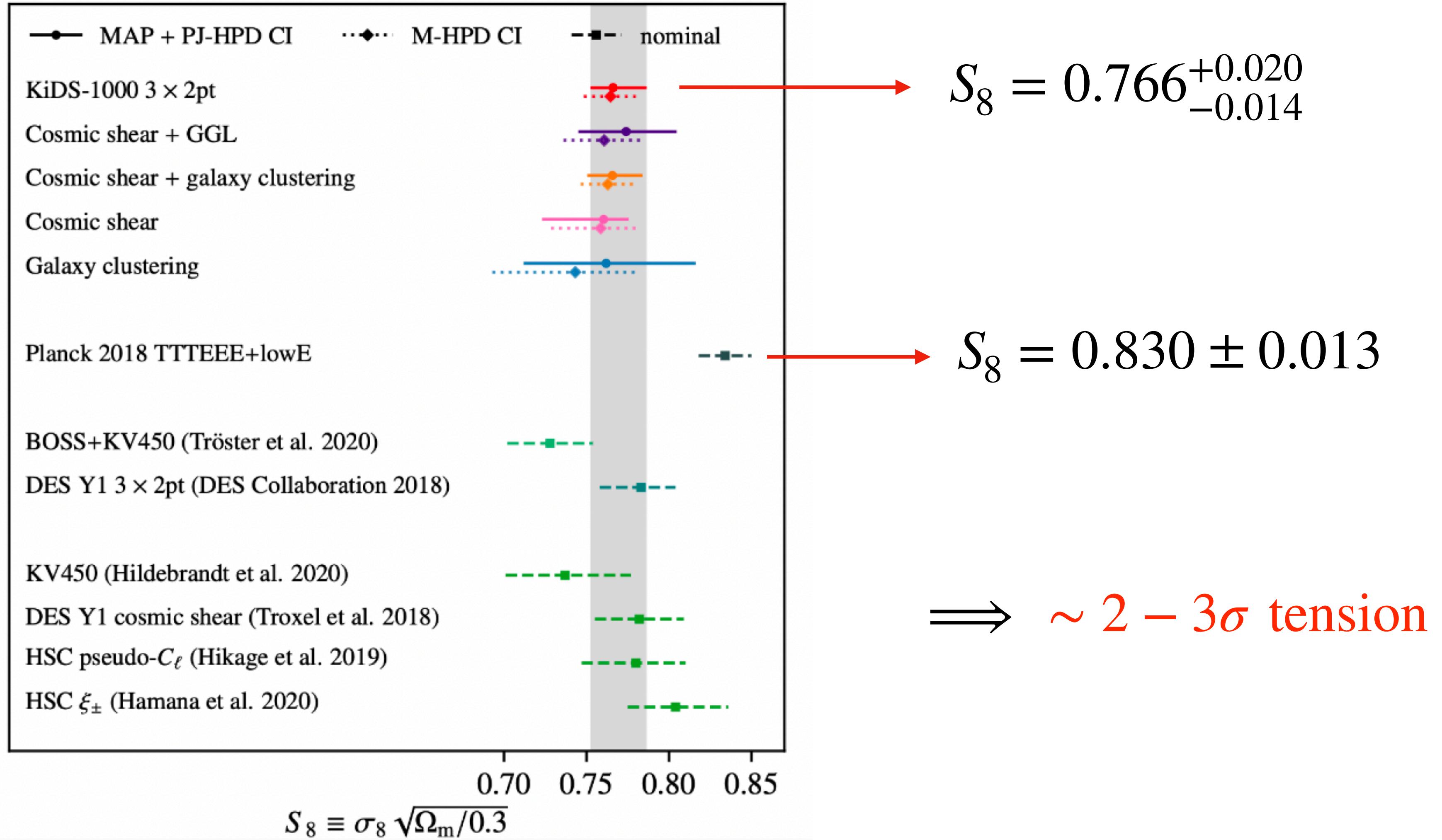
Planck (*under Λ CDM*):

$$S_8 = 0.830 \pm 0.013$$

→ $\sim 2 - 3\sigma$ tension

*Other surveys such as DES, CFHTLens or HSC yield similar results

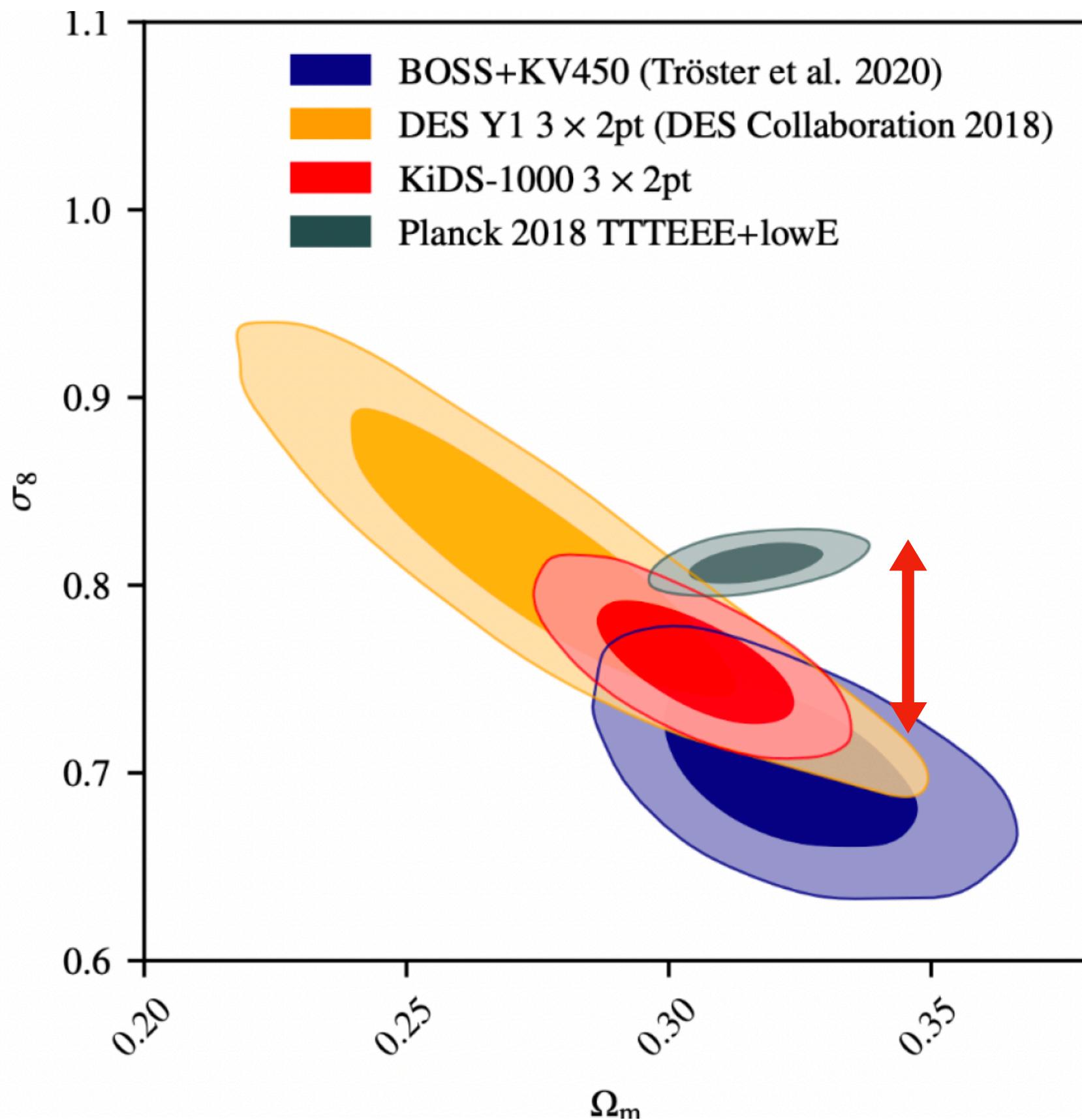
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What is needed to resolve the S_8 tension?

Di Valentino++ 2008.11285

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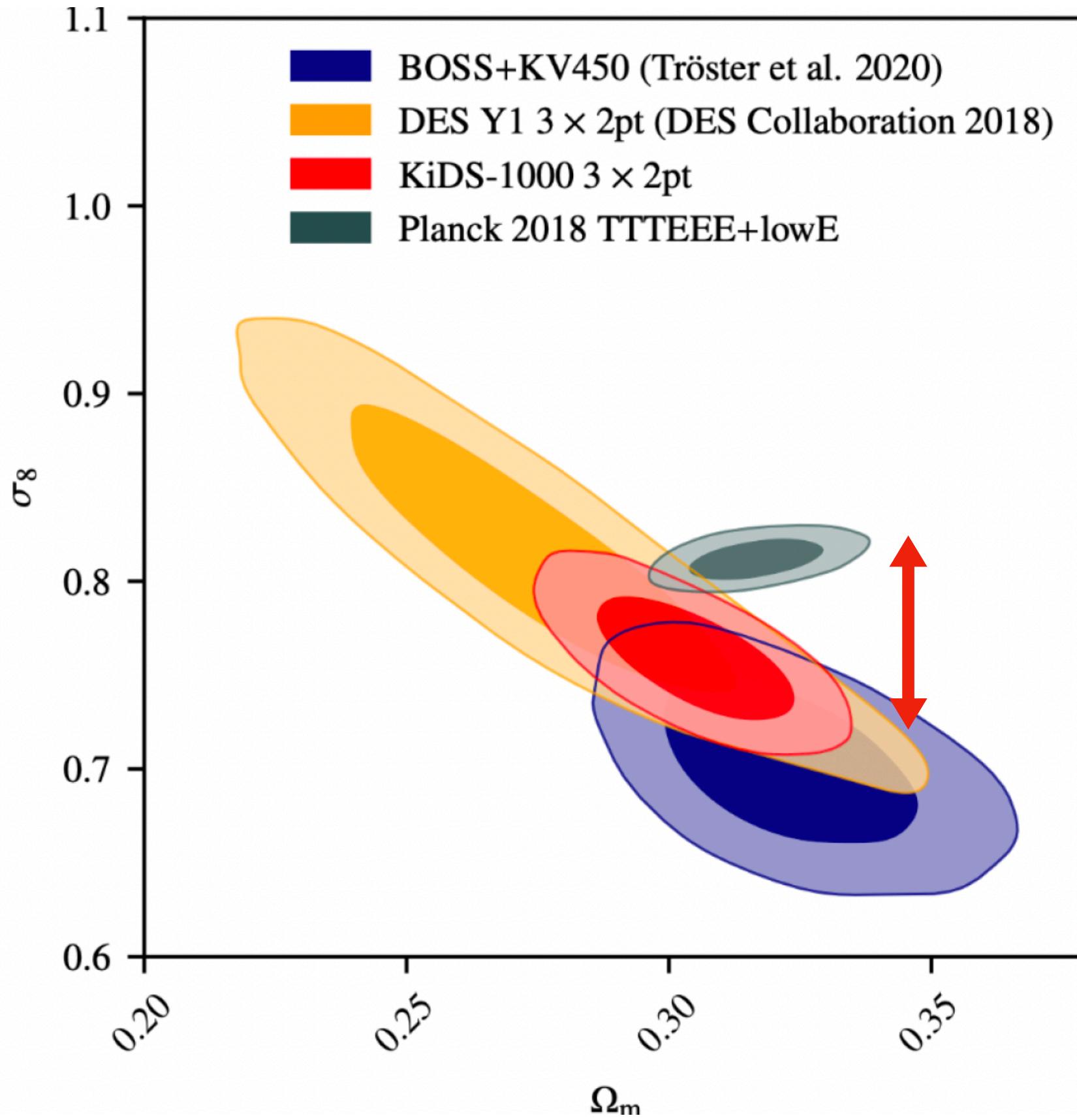
Ω_m should be left unchanged

$$\sigma_8 = \int P_m(k, z=0) W_R^2(k) d\ln k$$

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Ex: Warm Dark Matter

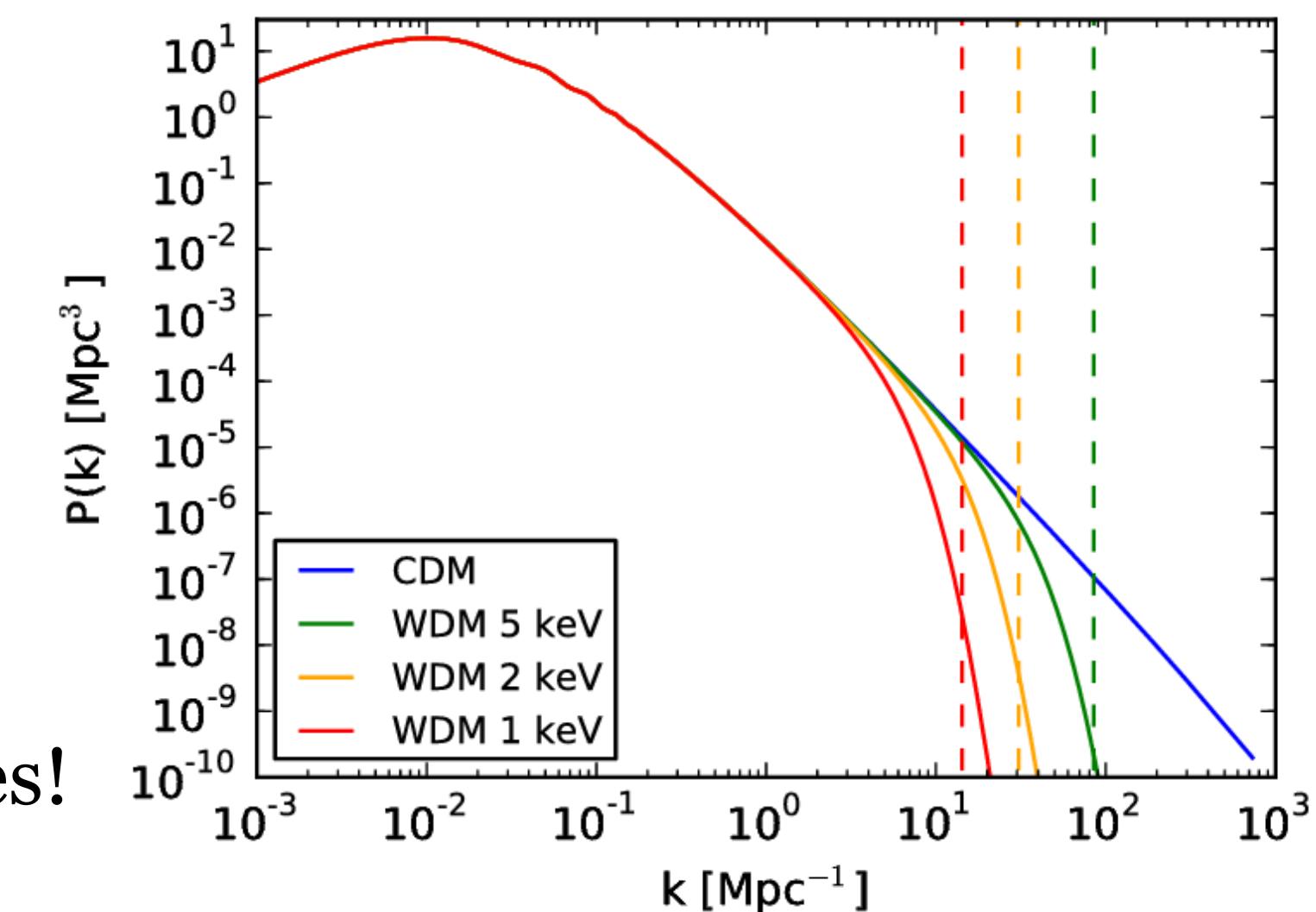
Very constrained by many probes!

Ω_m should be left unchanged

$$\sigma_8 = \int P_m(k, z=0) W_R^2(k) dk$$



Need to **suppress power** at scales $k \sim 0.1 - 1 h/\text{Mpc}$



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- DM decays to **massless** Dark Radiation \longrightarrow **less constrained**,
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From **grav. impact** on CMB : $\Gamma^{-1} \gtrsim 10^2$ Gyr [Audren++ 1407.2418](#)
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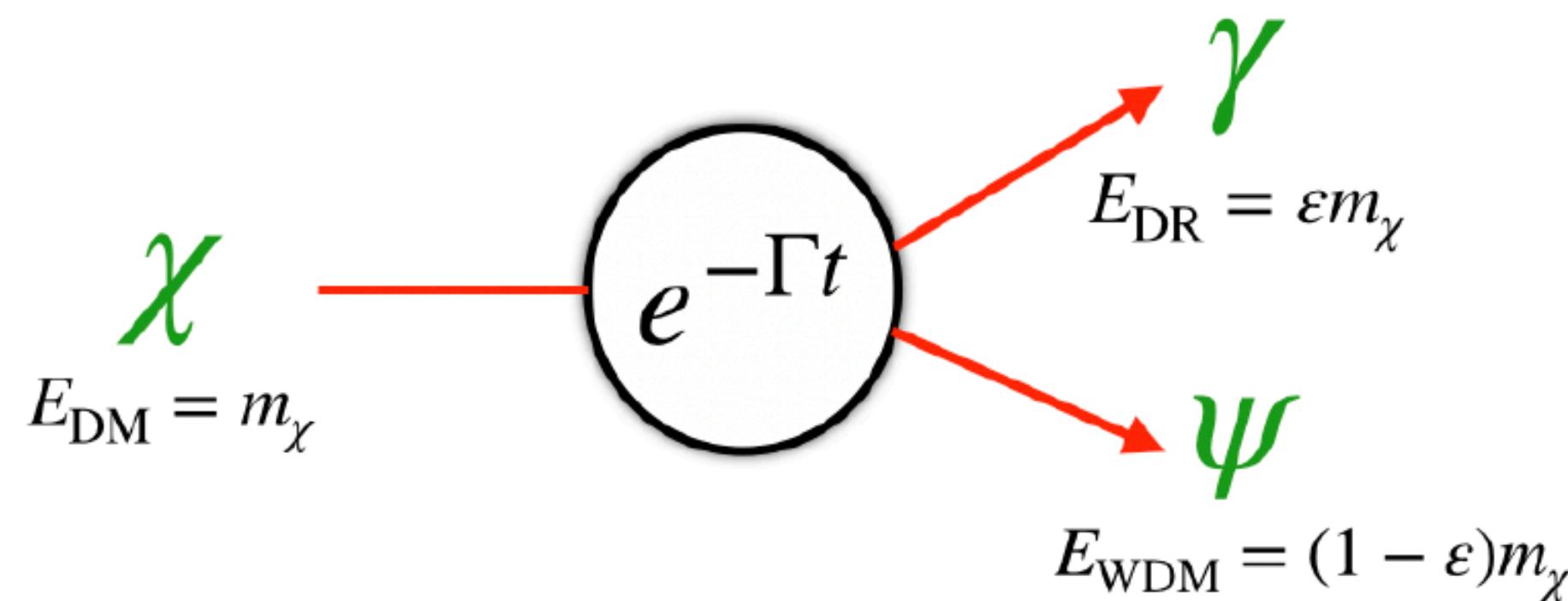
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- What about massive products?

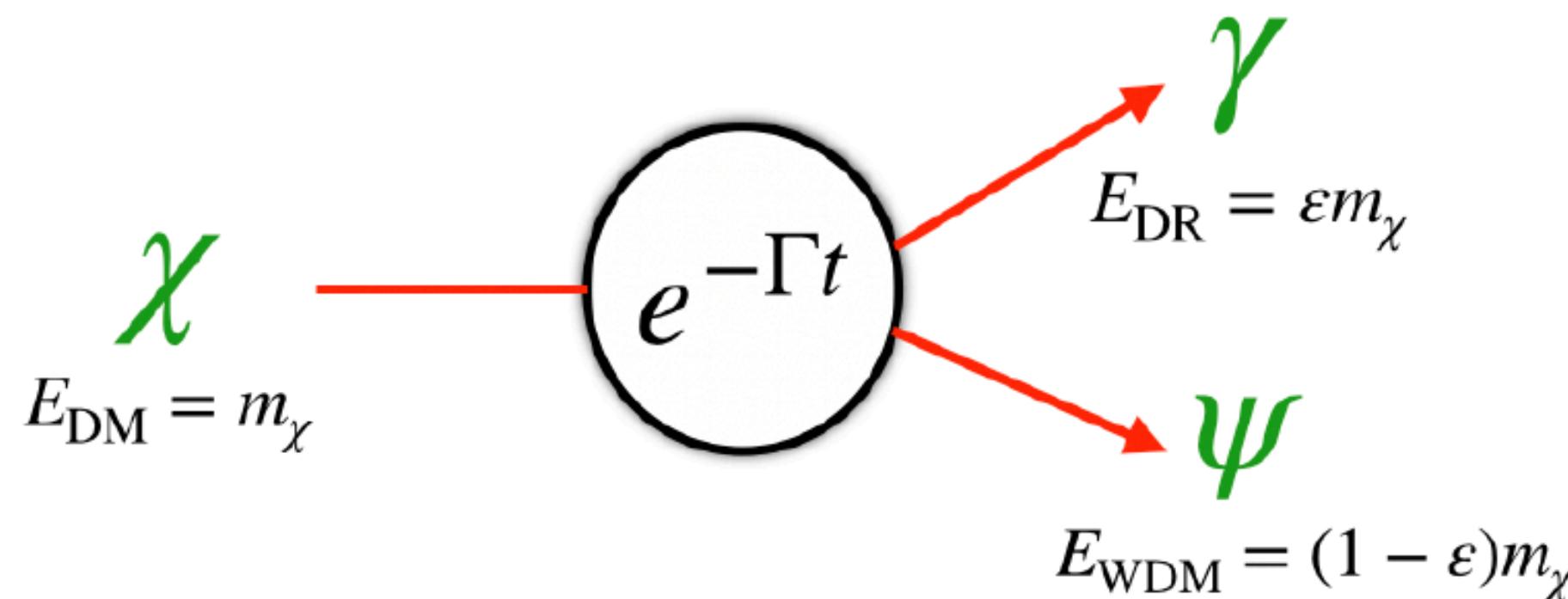
2-body Dark Matter decay

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The model is fully specified by:

$$\{\Gamma, \varepsilon\} \text{ where } \varepsilon = \frac{1}{2} \left(1 - \frac{\mathbf{m}_\psi^2}{\mathbf{m}_\chi^2} \right) \begin{cases} = 0 \text{ for } \Lambda\text{CDM} \\ = 1/2 \text{ for } \text{DM} \rightarrow \text{DR} \end{cases}$$

2-body Dark Matter decay

Aoyama++ 1402.2972	→	Full treatment of perts.	No parameter scan
Vattis++ 1903.06220	→	Resolution to H_0 tension ?	
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Our goal: Perform parameter scan by including full treatment of linear perts,
in order to assess the impact on the S_8 tension

Evolution of perturbations: full treatment

- Effects on $P_m(k)$ and C_ℓ ? Track **linear perts.** for the particles species involved in the decay: δ_i , θ_i and σ_i for $i = dm, dr, wdm$

Evolution of perturbations: full treatment

- Effects on $P_m(k)$ and C_ℓ ? Track **linear perts.** for the particles species involved in the decay: δ_i , θ_i and σ_i for $i = dm, dr, wdm$
- Boltzmann hierarchy of eqs. Dictate the evolution of the **p.s.d. multipoles** $\Delta f_\ell(q, k, \tau)$
 - ◆ DM and DR treatments are **easy**, momentum d.o.f. are integrated out
 - ◆ For WDM, one needs to follow the evolution of the full p.s.d. Computationally expensive $\longrightarrow \mathcal{O}(10^8)$ ODEs to solve!

Evolution of perturbations: fluid equations

New fluid eqs.*[†], based on previous approximation for massive neutrinos

Lesgourgues & Tram, 1104.2935

$$\dot{\delta}_{\text{wdm}} = -3aH(c_{\text{syn}}^2 - w)\delta_{\text{wdm}} - (1 + w) \left(\theta_{\text{wdm}} + \frac{\dot{h}}{2} \right) + a\Gamma(1 - \varepsilon) \frac{\bar{\rho}_{\text{dm}}}{\bar{\rho}_{\text{wdm}}} (\delta_{\text{dm}} - \delta_{\text{wdm}})$$

$$\dot{\theta}_{\text{wdm}} = -aH(1 - 3c_a^2)\theta_{\text{wdm}} + \frac{c_{\text{syn}}^2}{1 + w} k^2 \delta_{\text{wdm}} - k^2 \sigma_{\text{wdm}} - a\Gamma(1 - \varepsilon) \frac{\bar{\rho}_{\text{dm}}}{\bar{\rho}_{\text{wdm}}} \frac{1 + c_a^2}{1 + w} \theta_{\text{wdm}}$$

*Implemented in modified version of public Boltzmann solver CLASS

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where

$$c_a^2(\tau) = w\left(5 - \frac{p_{\text{wdm}}}{P_{\text{wdm}}} - \frac{\bar{\rho}_{\text{dm}}}{\bar{\rho}_{\text{wdm}}}\frac{\Gamma}{3wH}\frac{\varepsilon^2}{1-\varepsilon}\right)\left[3(1+w) - \frac{\bar{\rho}_{\text{dm}}}{\bar{\rho}_{\text{wdm}}}\frac{\Gamma}{H}(1-\varepsilon)\right]^{-1}$$

and

$$c_{\text{syn}}^2(k, \tau) = c_a^2(\tau)[1 + (1-2\varepsilon)T(k/k_{\text{fs}})]$$

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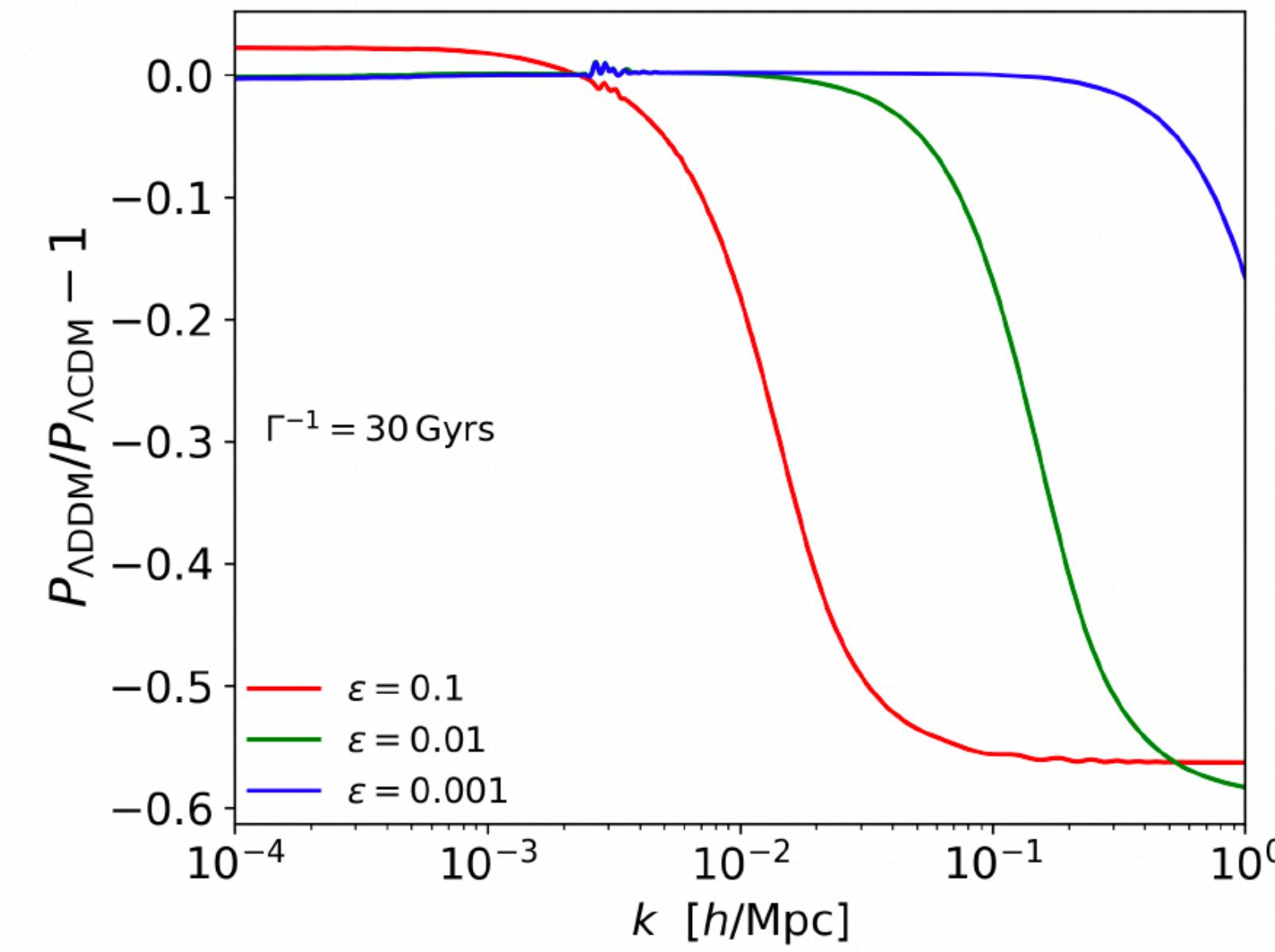
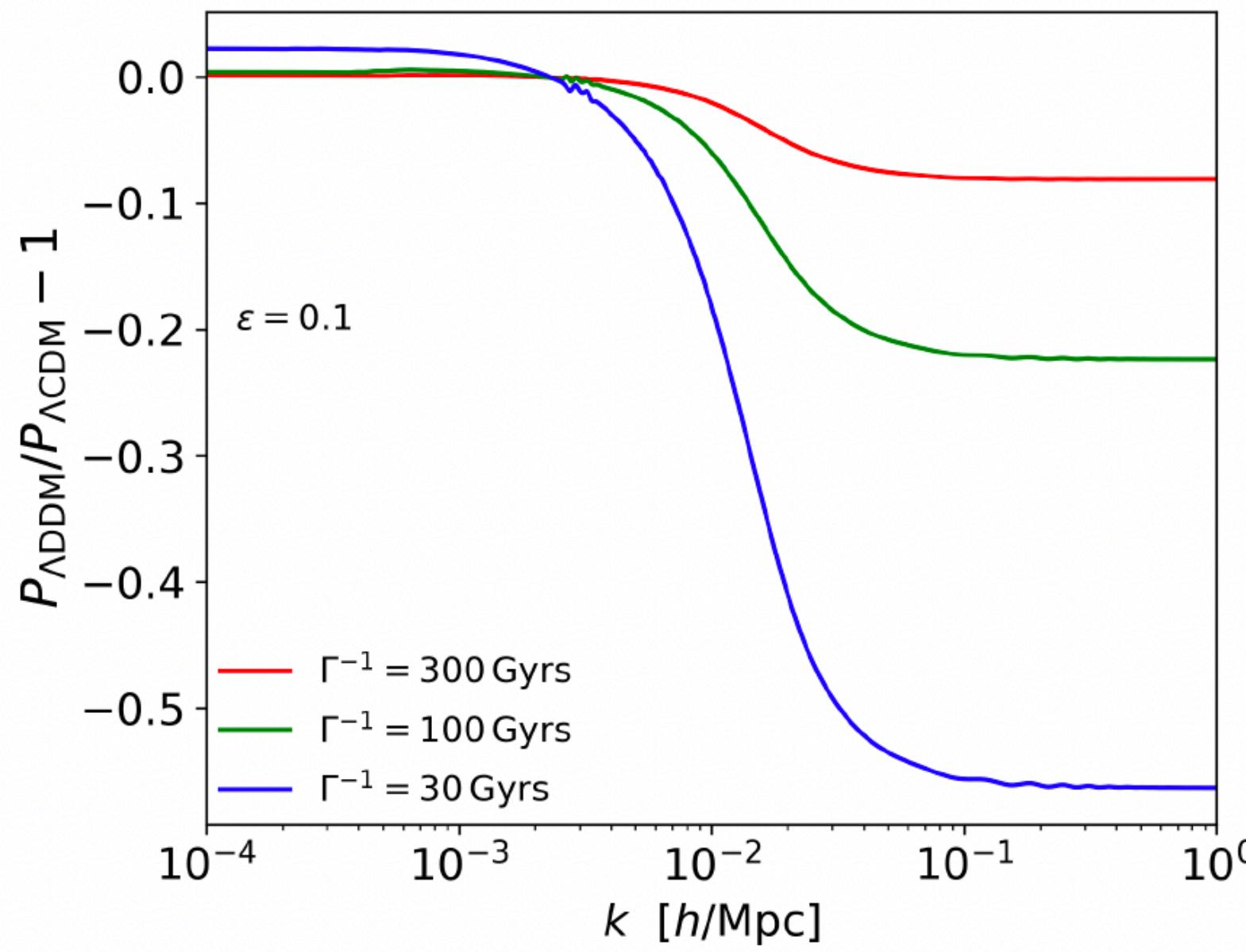
$$c_{\text{syn}}^2(k, \tau) = c_a^2(\tau)[1 + (1-2\varepsilon)T(k/k_{\text{fs}})]$$

CPU time reduced from ~ 1 day to ~ 1 minute!

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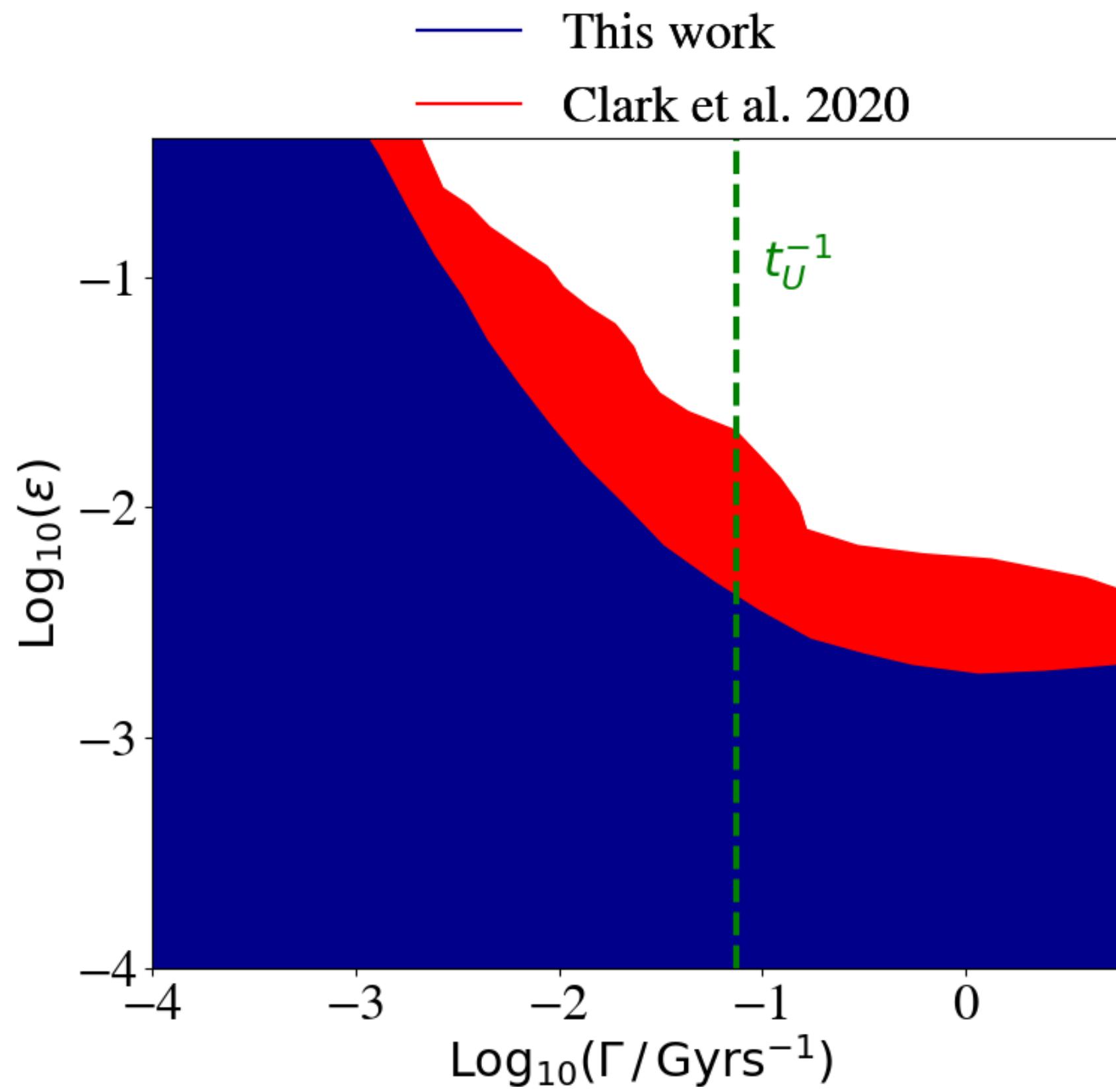
Impact of decaying DM on the matter spectrum

The WDM daughter leads to a power suppression in $P_m(k)$ at small scales $k > k_{fs}$, where $k_{fs} \sim aH/c_a$



General constraints on the 2-body DM decay

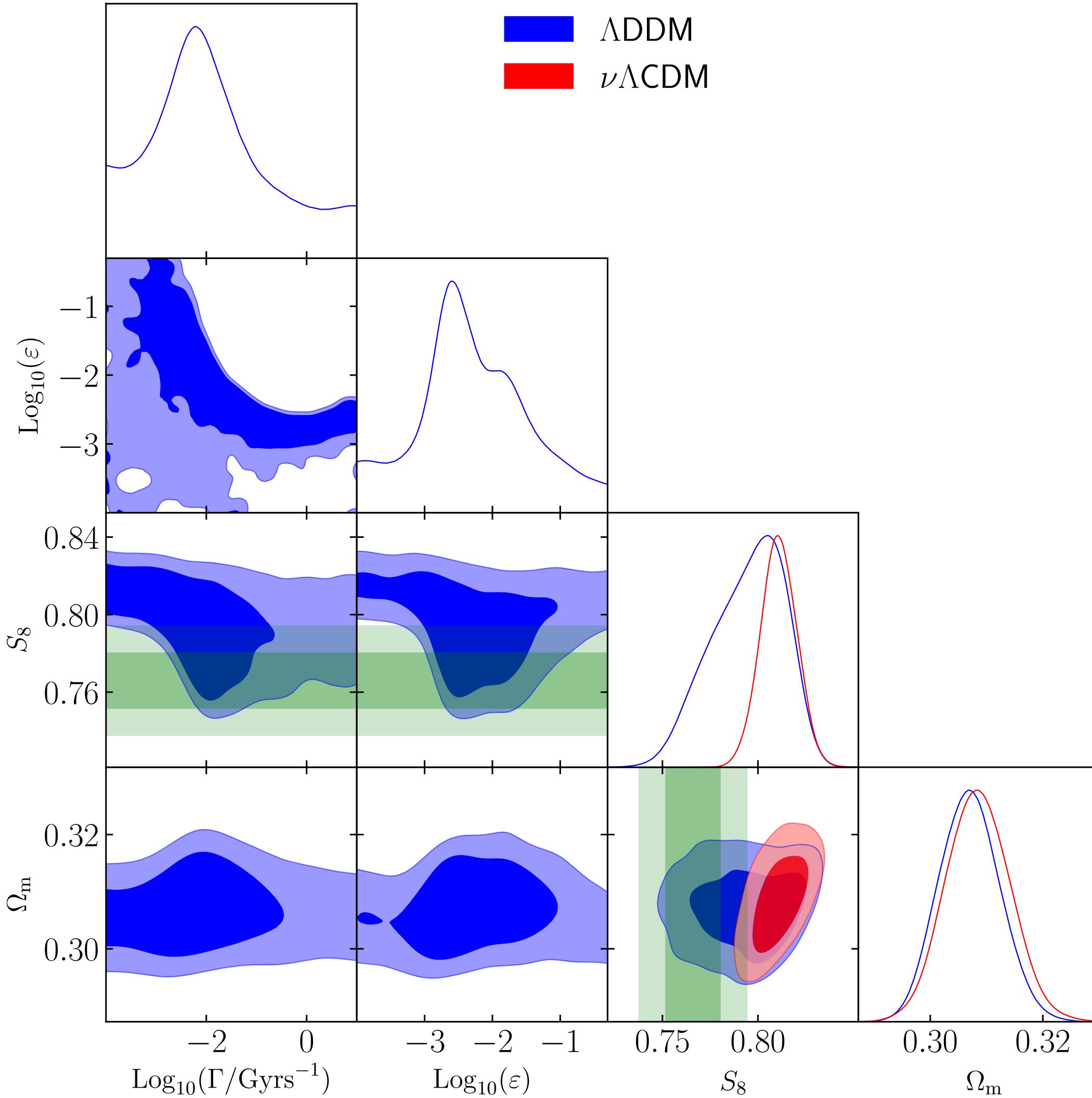
Planck+BAO+SNIa analysis



Strong negative correlation
between ϵ and Γ

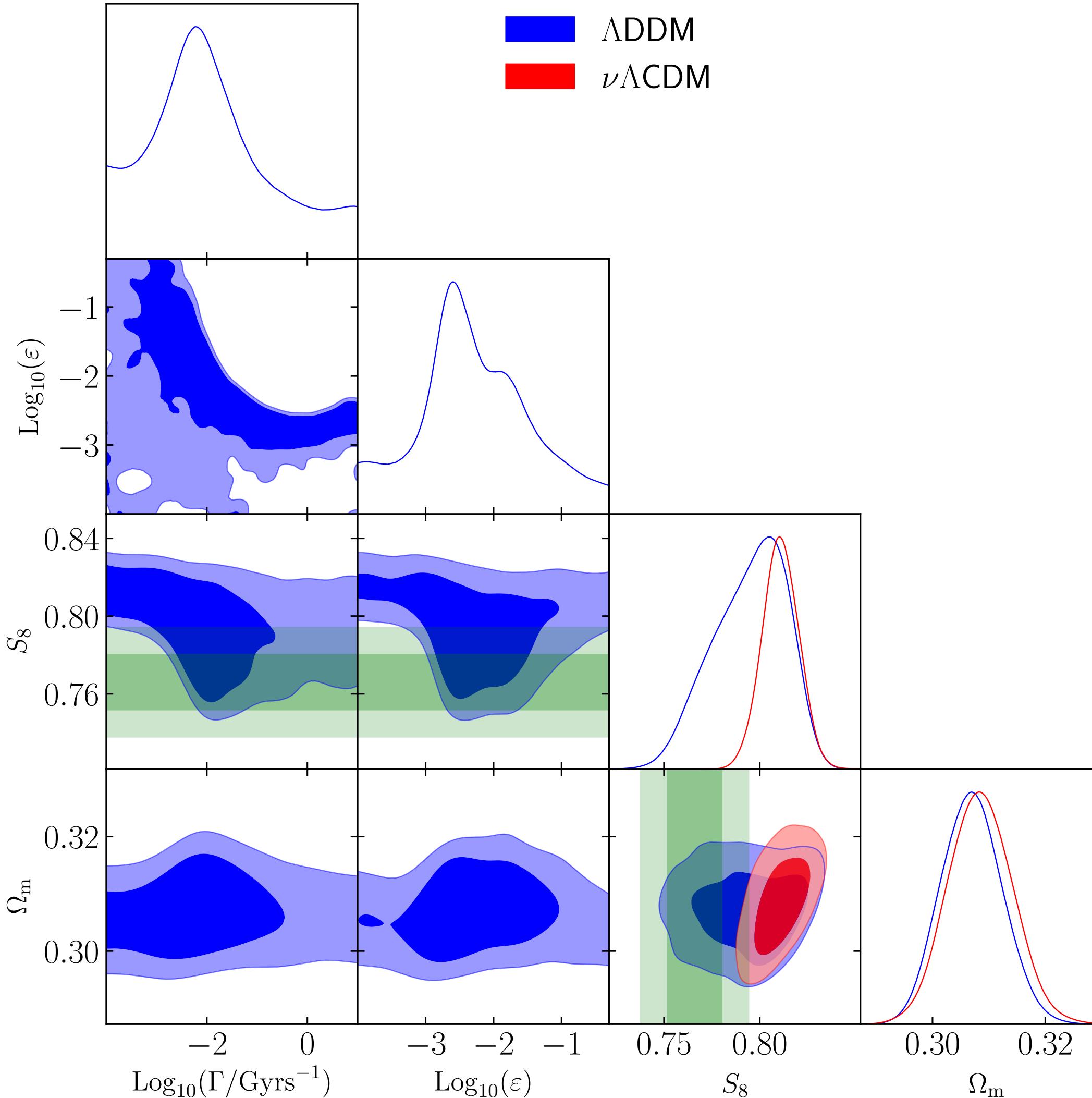
Constraints up to 1 order of
magnitude stronger than
previous literature

Explaining the S_8 tension



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- MCMC analysis using Planck+BAO+SNIa+prior on S_8 from KIDS+BOSS+2dfLenS
- Reconstructed S_8 values are in excellent agreement with WL data!

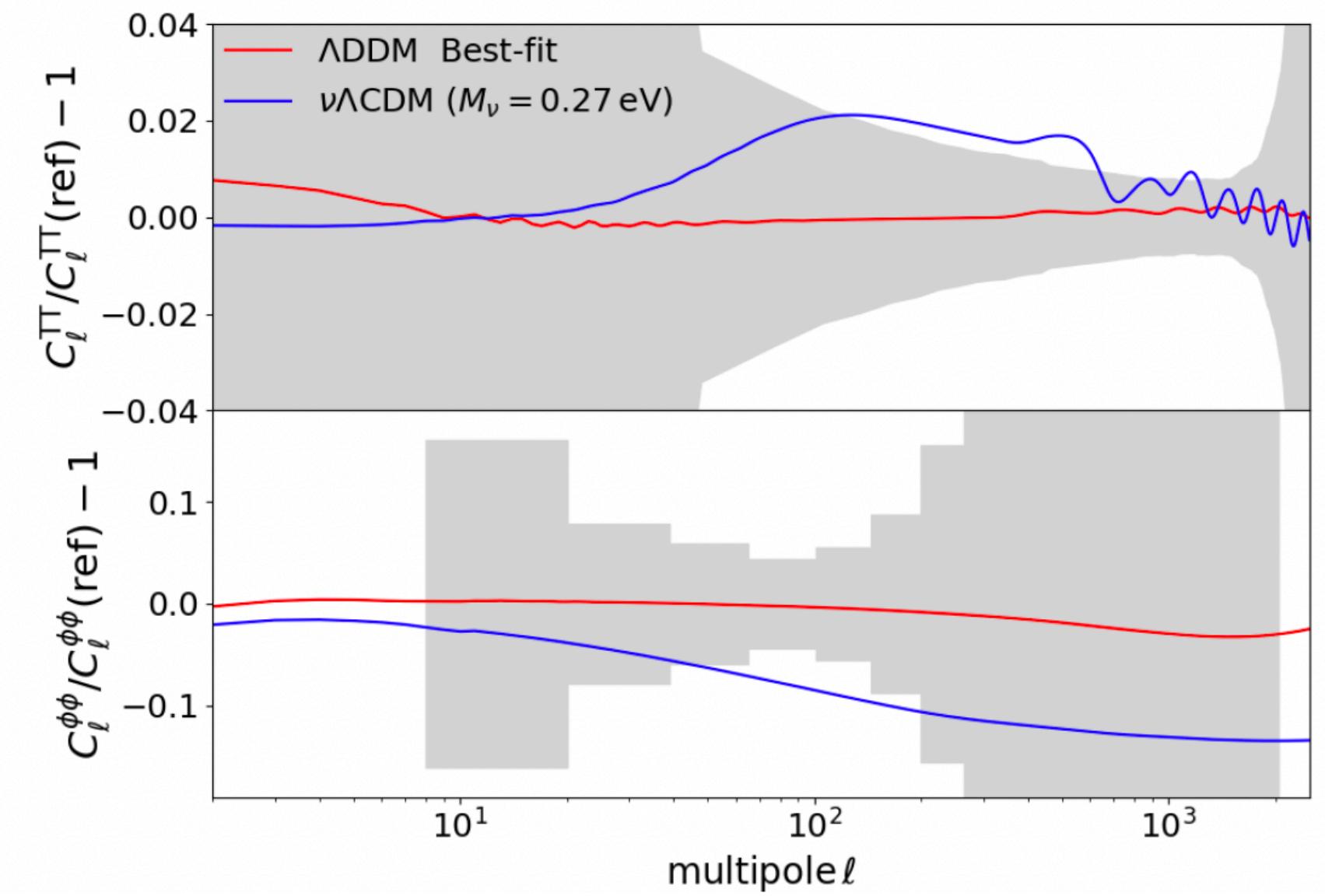
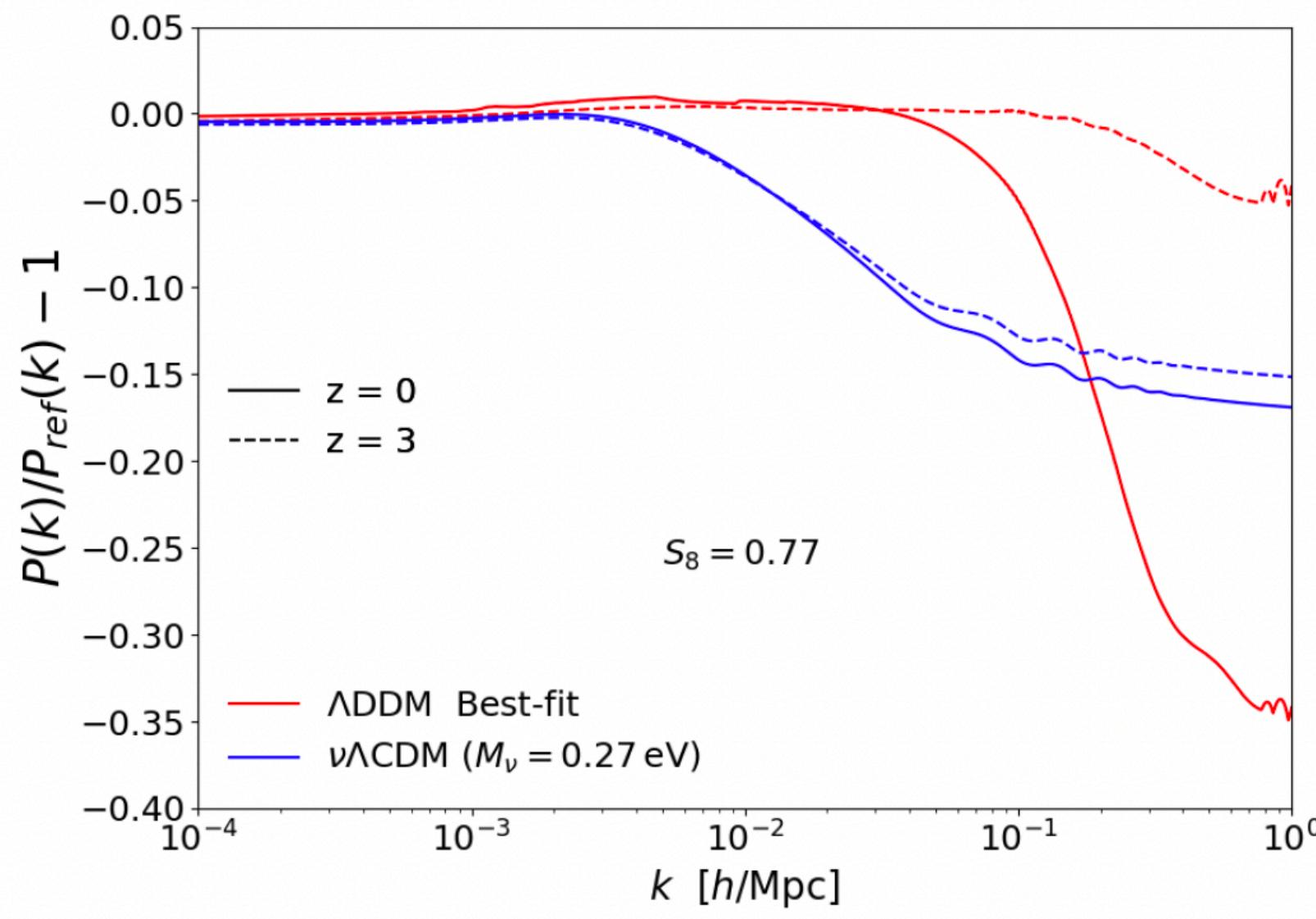
	$\nu\Lambda$ CDM	Λ DDM
χ^2_{CMB}	1015.9	1015.2
$\chi^2_{S_8}$	5.64	0.002

$$\rightarrow \Delta\chi^2_{\min} \simeq -5.5$$

$$\Gamma^{-1} \simeq 55 (\varepsilon/0.007)^{1.4} \text{ Gyr}$$

Why does the 2-body DM decay work better than massive neutrinos?

The 2-body decay gives a better fit thanks to the **time-dependence of the power suppression** and the cut-off scale



Interesting implications

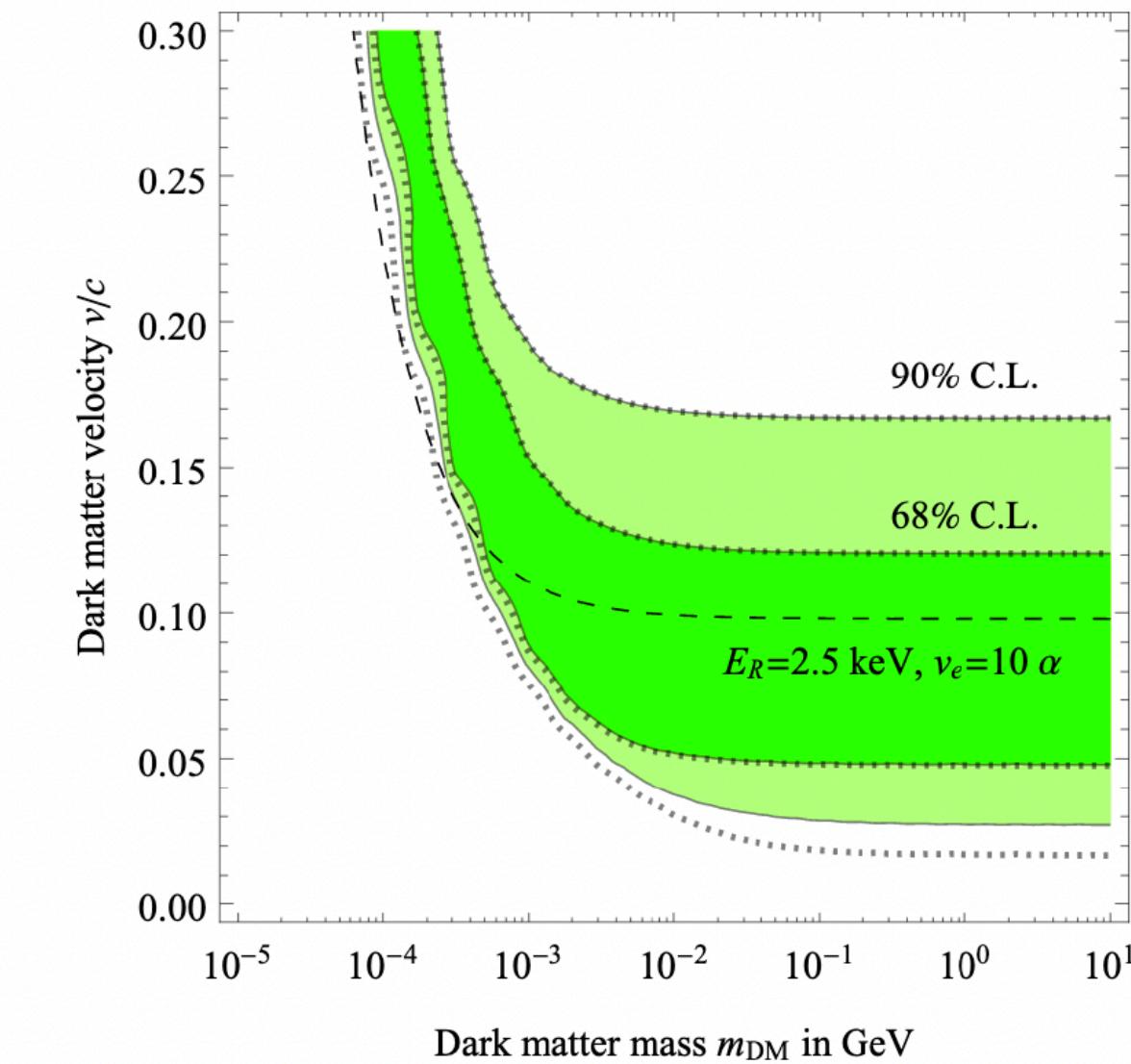
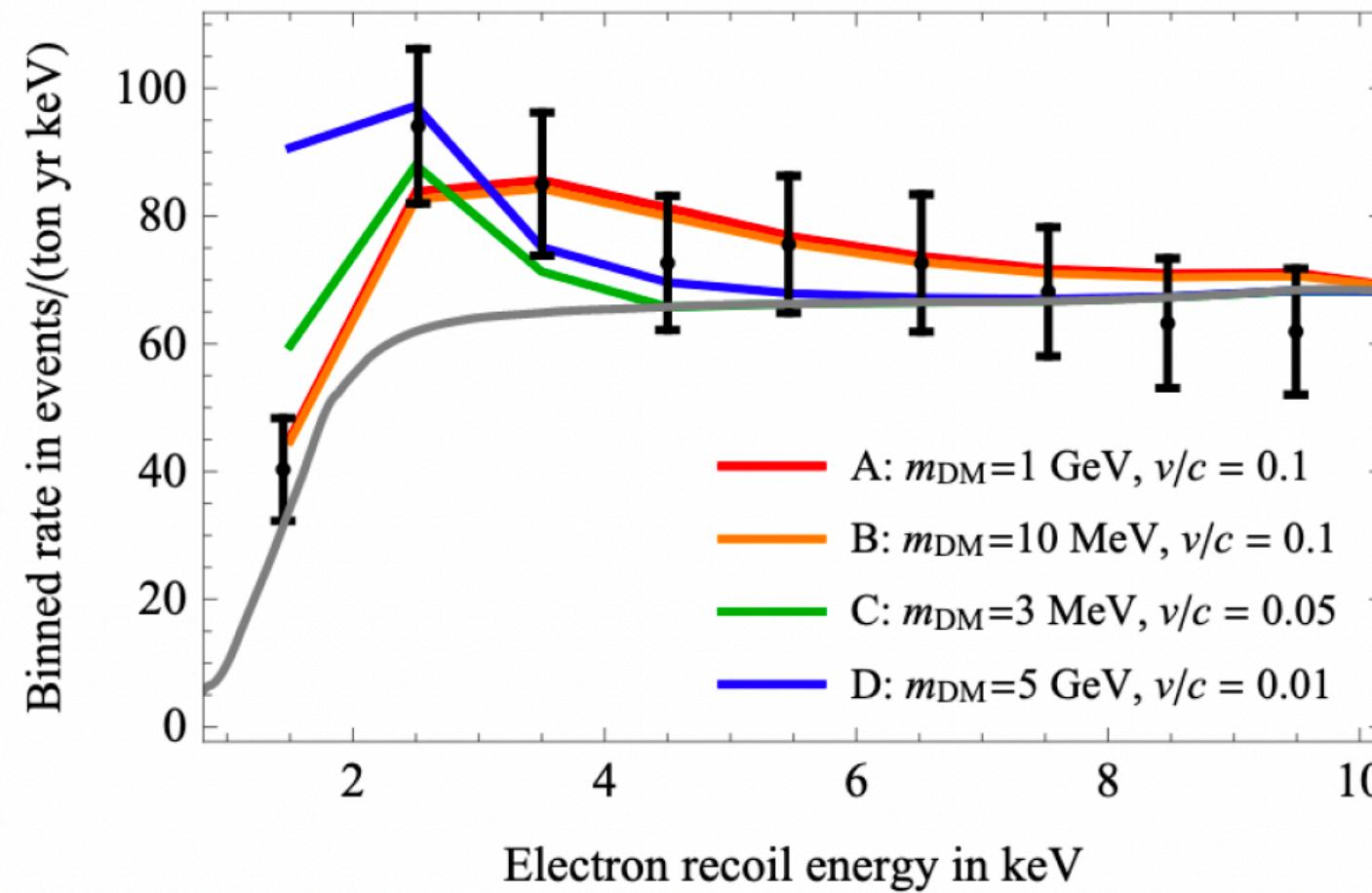
- Model building: Why $\varepsilon \ll 1/2$, i.e. $m_{\text{wdm}} \sim m_{\text{dm}}$?
Ex : Supergravity [Choi&Yanagida 2104.02958](#)

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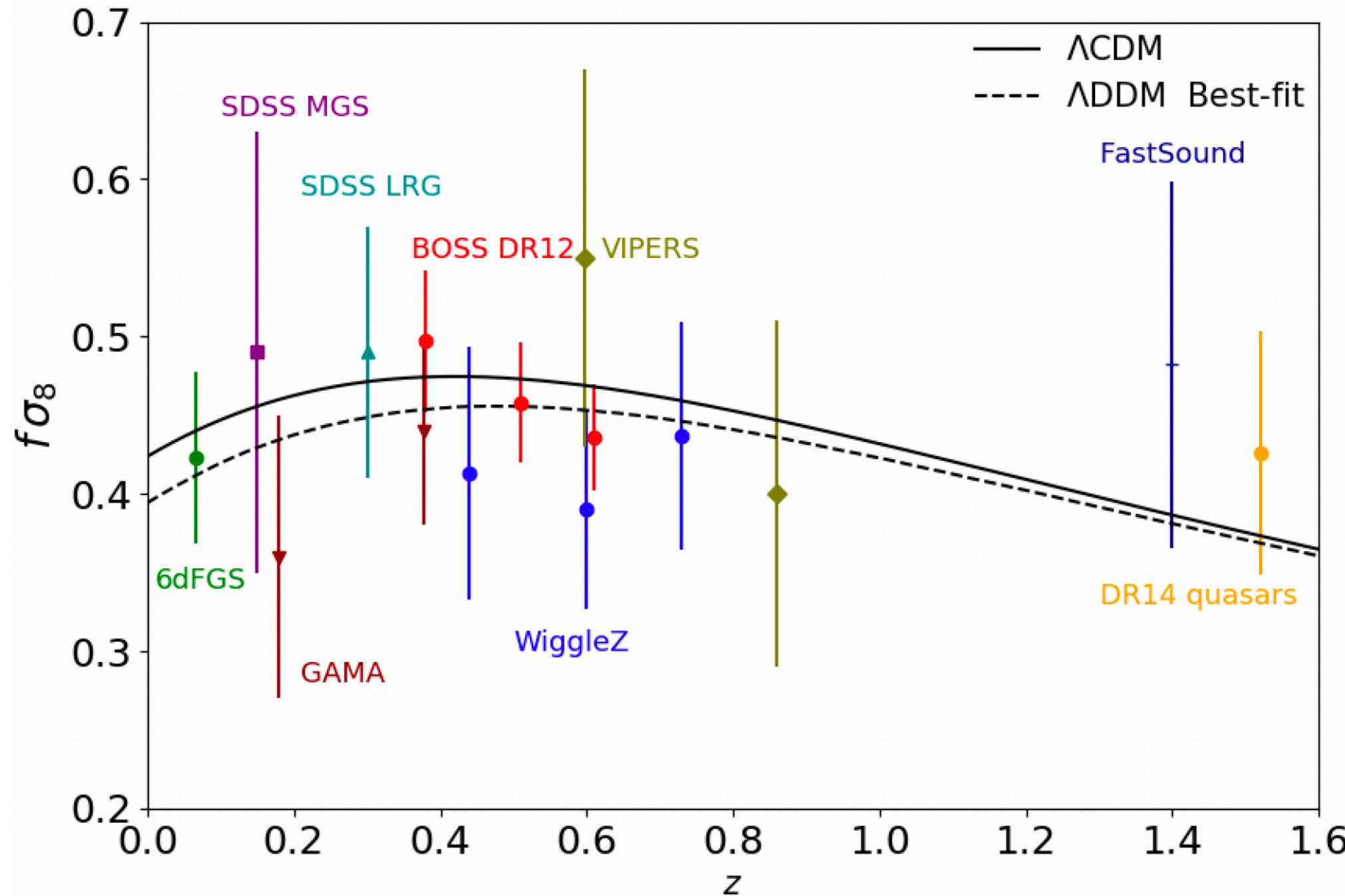
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- **Small-scale crisis of Λ CDM:** Reduction in the abundance of subhalos and their concentrations Wang++ [1406.0527](#)
- **Xenon-1T excess:** It could be explained by a fast DM component, such as the WDM, with $v/c \simeq \varepsilon$ Kannike++ [2006.10735](#)



Prospects for the 2-body DM decay



Accurate measurements of $f\sigma_8$ at $0 \lesssim z \lesssim 1$ will further test the 2-body decay

Next goal: Predict non-linear matter power spectrum
(using either N-body simulations or EFT of LSS)

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- It could have interesting implications for model building, the small-scale crisis, and the recent **Xenon-1T excess**
- Future growth factor measurements can further test this scenario

BACK-UP SLIDES

The full Boltzmann hierarchy

$$f(q, k, \mu, \tau) = \bar{f}(q, \tau) + \delta f(q, k, \mu, \tau)$$

Expand δf in multipoles. The Boltzmann eq. leads to the following **hierarchy**
(in synchronous gauge comoving with the mother)

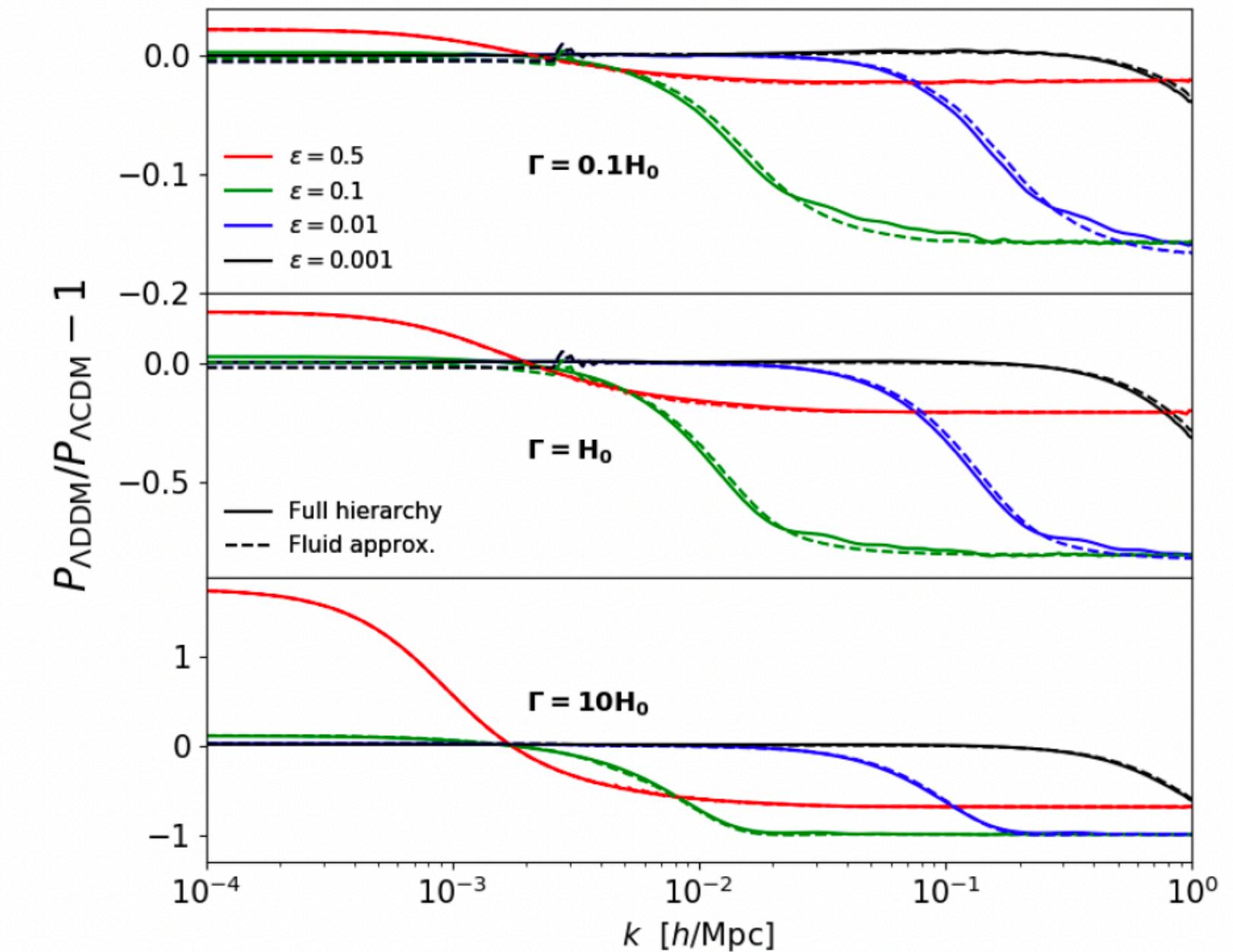
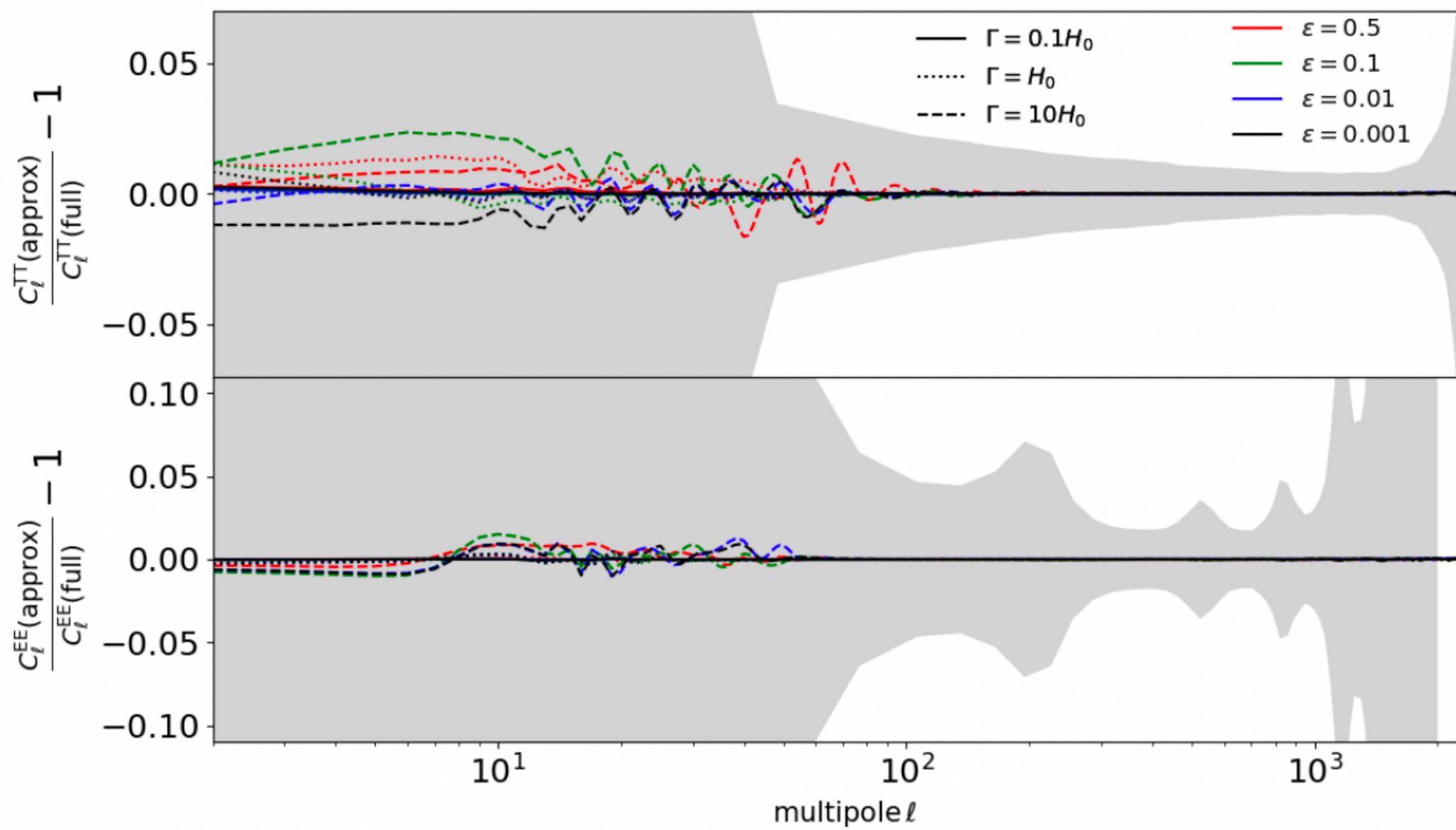
$$\begin{aligned} \frac{\partial}{\partial \tau} (\delta f_0) &= -\frac{\mathbf{q}k}{\mathbf{a}\mathbf{E}} \delta f_1 + q \frac{\partial \bar{f}}{\partial q} \frac{\dot{h}}{6} + \frac{\Gamma \bar{N}_{\text{dm}}(\tau)}{4\pi q^3 H} \delta(\tau - \tau_q) \delta_{\text{dm}}, \\ \frac{\partial}{\partial \tau} (\delta f_1) &= \frac{\mathbf{q}k}{3\mathbf{a}\mathbf{E}} [\delta f_0 - 2\delta f_2], \\ \frac{\partial}{\partial \tau} (\delta f_2) &= \frac{\mathbf{q}k}{5\mathbf{a}\mathbf{E}} [2\delta f_1 - 3\delta f_3] - q \frac{\partial \bar{f}}{\partial q} \frac{(\dot{h} + 6\dot{\eta})}{15}, \\ \frac{\partial}{\partial \tau} (\delta f_\ell) &= \frac{\mathbf{q}k}{(2\ell + 1)\mathbf{a}\mathbf{E}} [\ell \delta f_{\ell-1} - (\ell + 1) \delta f_{\ell+1}] \quad (\text{for } \ell \geq 3). \end{aligned}$$

where $q = a(\tau_q) p_{\max}$. In the relat. limit $\mathbf{q}/\mathbf{a}\mathbf{E} = 1$, so one can take

$F_\ell \equiv \frac{4\pi}{\rho_c} \int dq q^3 \delta f_\ell$ and **integrate out the dependency on \mathbf{q}**

Checking the accuracy of the WDM fluid approx.

We compare the full Boltzmann hierarchy calculation with the WDM fluid approx.

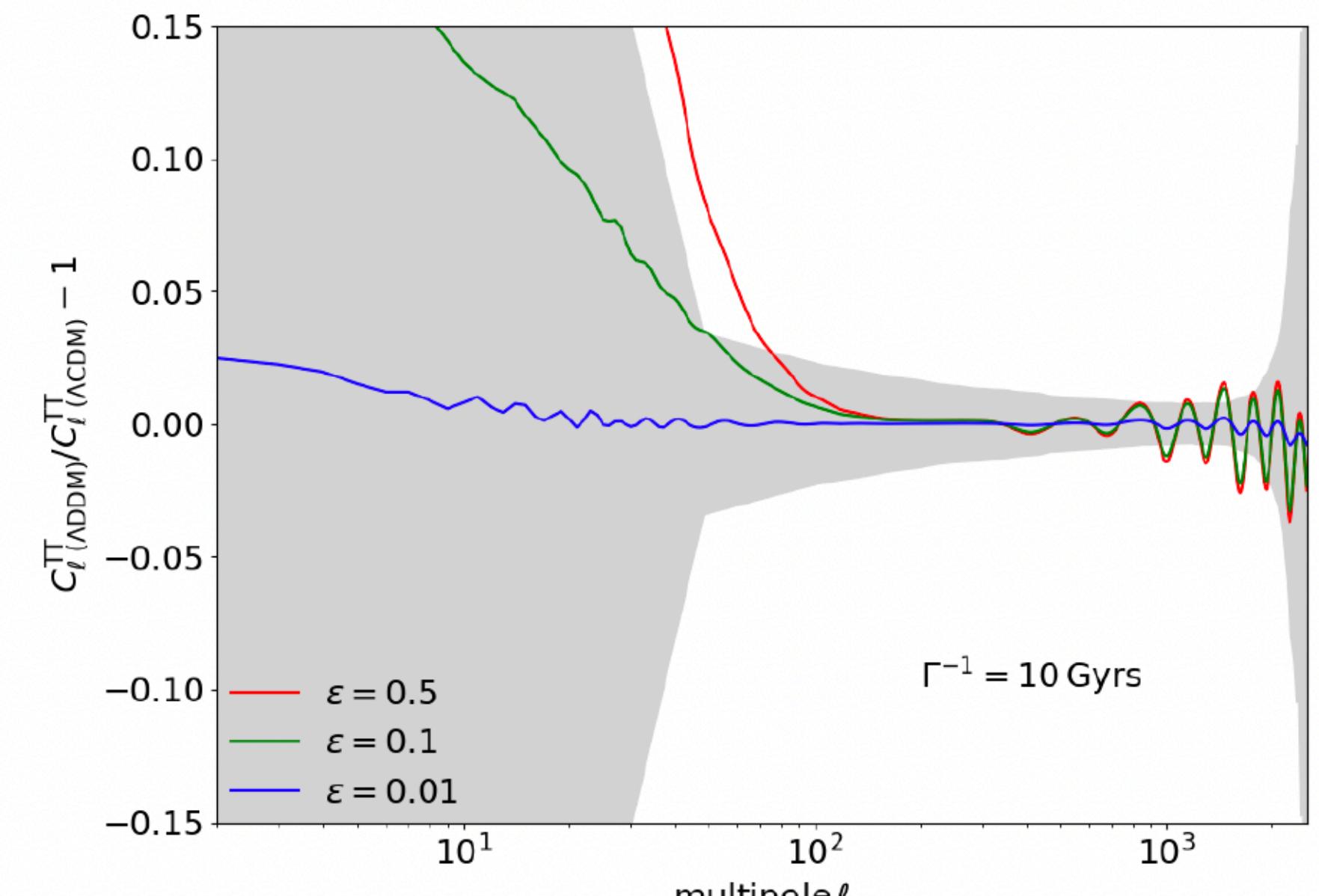
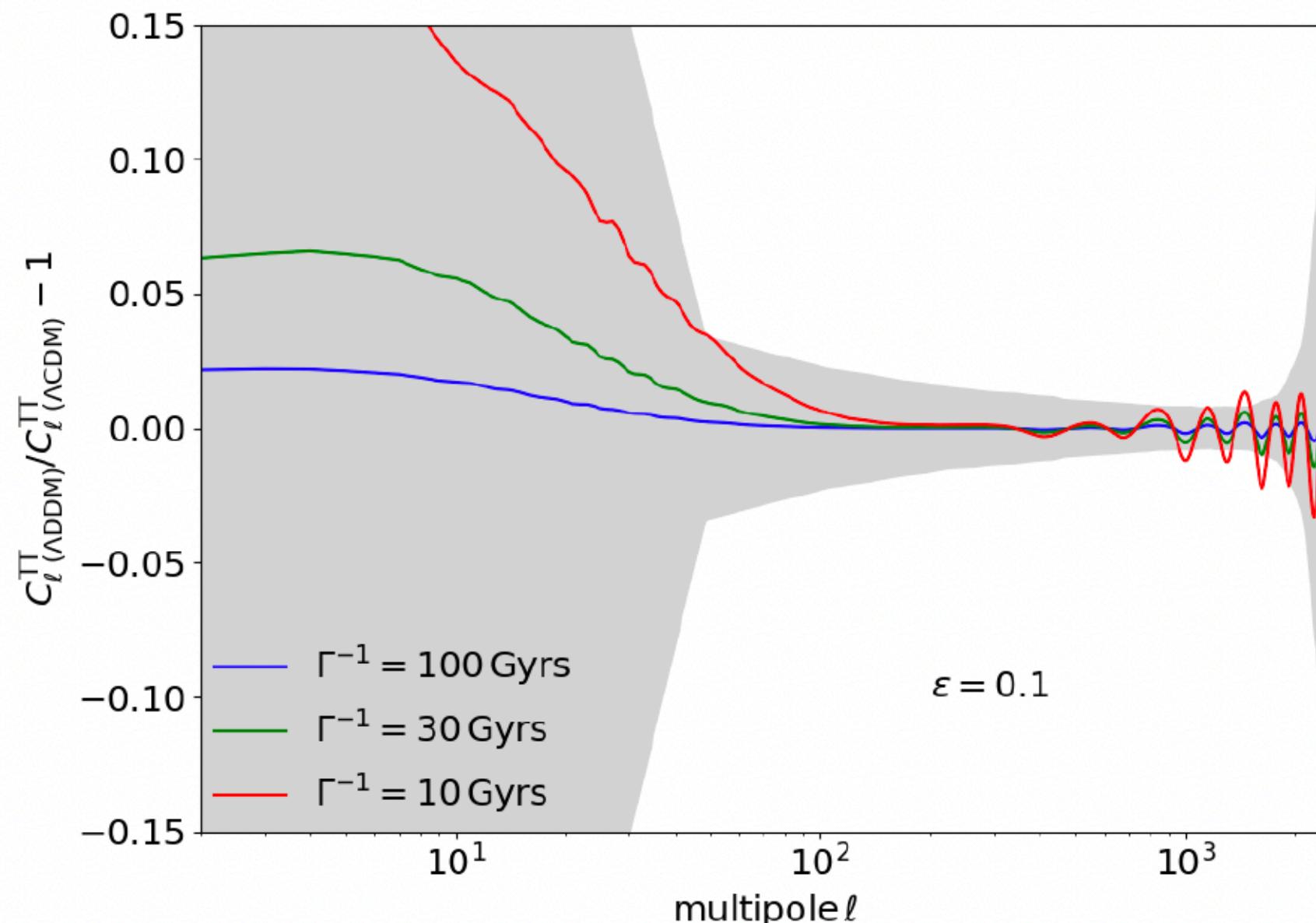


The max. error on S_8 is $\sim 0.65\%$, smaller than the $\sim 1.8\%$ error of the measurement from BOSS+KiDS+2dfLenS

Impact on the CMB temperature spectrum

Low- ℓ : **enhanced** Late Integrated Sachs Wolfe (**LISW**) effect

High- ℓ : **suppressed** lensing (higher contrast between peaks)

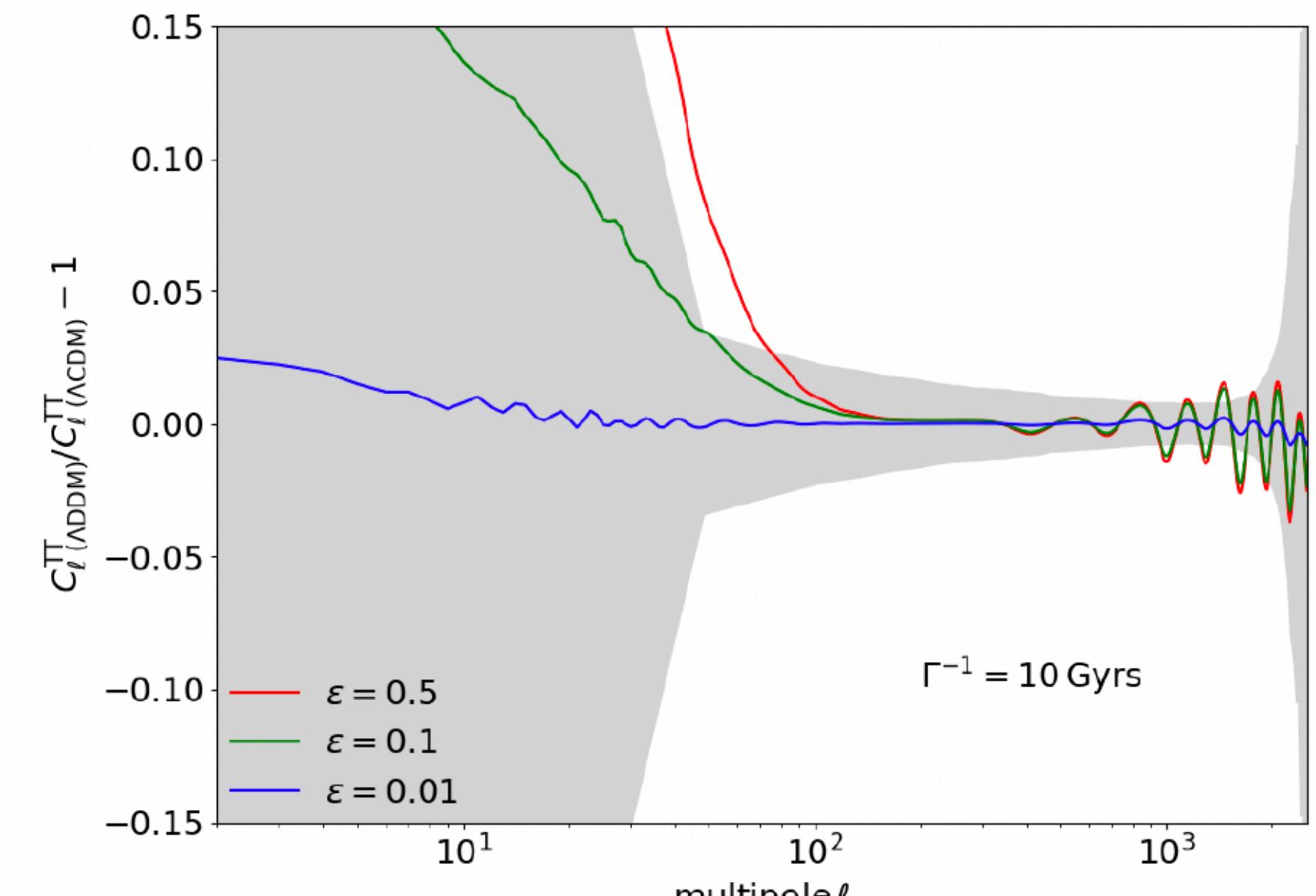
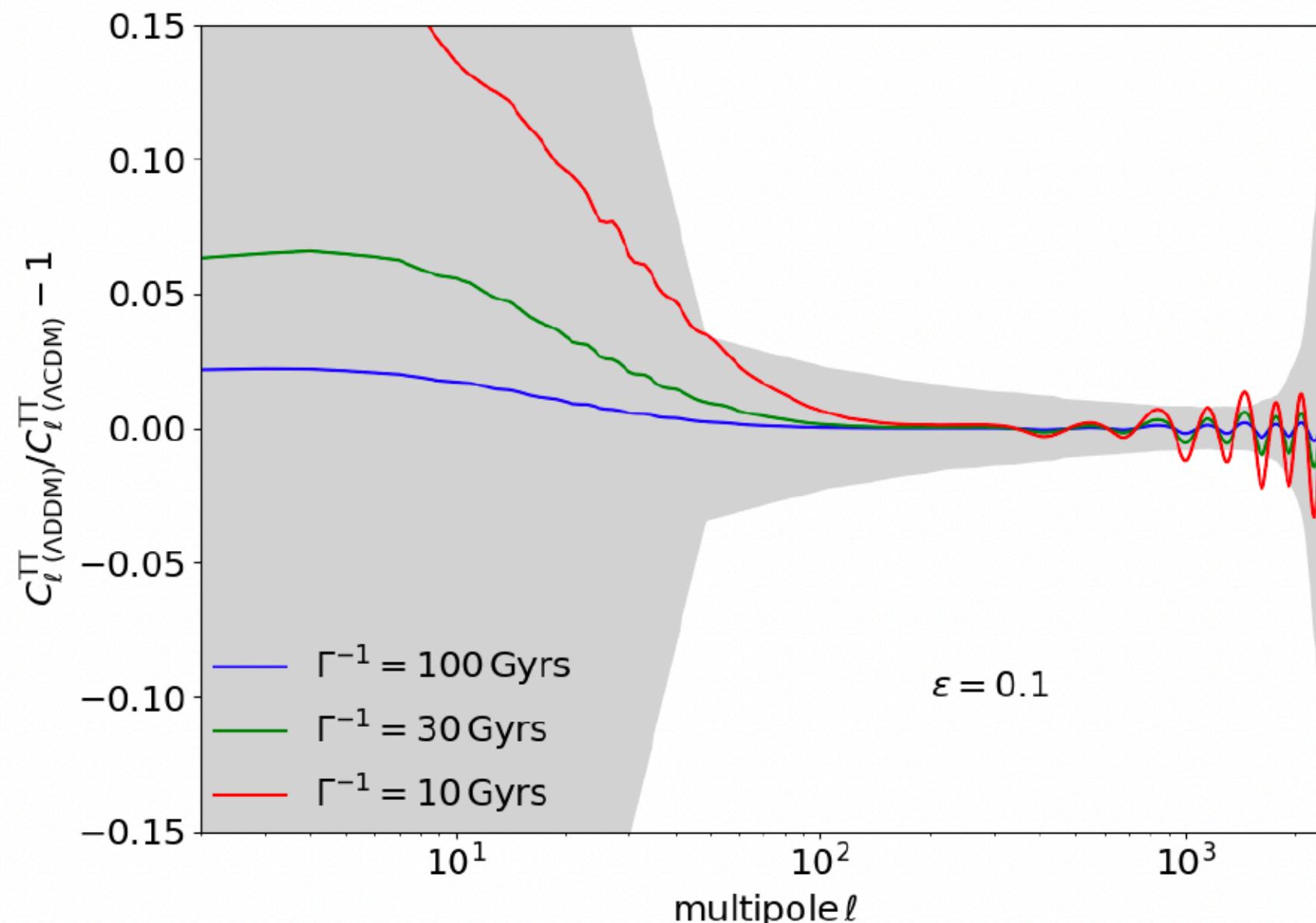


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Resolution to the S_8 tension

The level of detection depends on the level of tension with Λ CDM

