## Drafts for Project-Particle Methods for Vortex Problems

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We begin with Poisson's equation in continuous form:

$$\nabla^2 \psi(\mathbf{x}) = -\rho(\mathbf{x})$$

The solution using the Green's function is given by:

$$\psi(\mathbf{x}) = \int_{\mathbb{R}^2} G(\mathbf{x} - \mathbf{x}') \, \rho(\mathbf{x}') \, d\mathbf{x}'$$

Solving Poisson's equation numerically on a discrete grid, means the domain, potential, source, and Green's function are sampled at grid points. To discretize this:

- Let the domain be sampled on a uniform 2D grid with spacing h
- Let  $i = (i_x, i_y)$  index the grid points
- Let  $\mathbf{x}_i = h \cdot i$  be the physical coordinates
- Define  $\rho_j = \rho(\mathbf{x}_j)$  Here, G(i-j) is the Green's function for the Laplacian in 2D (solution to Poisson's equation)

Then, we approximate the integral using a Riemann sum and denote  $\omega_j^g = \rho_j$ .

$$\psi_i = \sum_{i \in \mathbb{Z}^2} G(i-j) \,\omega_j^g$$

where  $\omega_j$  represents the value of  $\rho$  at the grid point j, and the sum runs over all grid points  $j \in \mathbb{Z}^2$ . The function G(i-j) gives the influence of the source at point j on the potential at point i.

The potential on the grid is given by  $\psi_i$ , and the source on the grid is given by  $\omega_i^g$ .

Hockney's algorithm utilizes that The Fourier transform of a convolution equals the product of the Fourier transforms.

$$\mathcal{F}(f * g) = \mathcal{F}(f) \cdot \mathcal{F}(g)$$
$$f * g = \mathcal{F}^{-1}[\mathcal{F}(f) \cdot \mathcal{F}(g)]$$

You will be implementing parts of a particle-in-cell (PIC) method for vortex dynamics, described below. This is primarily an exercise in more elaborate template programming. Generally speaking,

you are integrating an ODE of the form

$$\frac{dX}{dt} = F(t, X) \tag{1}$$

In this problem set our forcing functions will all be independent of time, so you can ignore the a\_time argument, but it is good to have this form available to you when you use RK4 in other projects. We will be using the 4th-order explicit Runge-Kutta integration technique to evolved this system of ODEs. In this case X is the class ParticleSet.

The stages of RK4 all require the calculation of quantities of the form

$$k := \Delta t * F(t + \Delta t, X + k) \tag{2}$$

Your F operator is an evaluation of everything on the right of the equal sign. RK4 is built up by various estimates of what the update to X should be, then recombined to cancel out low order error terms to create a stable method with an error in the solution that is  $O(\Delta t)^4$ .

Specifically, you will implement the class ParticlesVelocities, that has the single member function

The input is the current estimate for k, a\_result, and the output new estimate for k is returned in a\_result:

$$k := \Delta t F(t + \Delta t, X + k)$$

Inputs are the time you are to evaluate the function  $t + \Delta t$ , the timestep to take  $\Delta t$ , the state at the start of the timestep X in this case ParticleSet, and the shift to use to this state in this evaluation of F k, represented by the ParticleShift class. In the case of our particle method, F has no explicit dependence on the first time argument, but we still have implement our class as if it does, in order to conform to the general RK4 interface.

## Specific Instructions

You are to implement in the directory /src/Particles ParticleVelocities::operator() (ParticleShift& a\_k, const double& a\_time, const double& dt, ParticleSet& a\_state) : computes the k's induced on a set of particles by all of the particles in the input ParticleSet displaced by the input k. In addition, you are to implement a driver program that performs the following calculations.

1. A single particle, with strength  $1./h^2$ , placed at (i) (.5,.5), (ii) .4375,5625, (iii) .45,.55. The number of grid points is given by N = 32,  $\Delta t = 1$ .; run for 100 time step. In all of these cases, the displacement of the particle should be roundoff, since the velocity induced by a single particle on itself should vanish. In the case of the initial position of (.5,.5) the displacement should be comparable to roundoff. Output: position of the particle after one step

- 2. Two particles: one with strength  $1/h^2$  located at (.5,.5), the other with strength 0, located at (.5,.25). The number of grid points is given by N = 32. Run for 300 time steps,  $\Delta t = .1$ . The strength 1 particle should not move, while the zero-strength particle should move at constant angular velocity on a circle centered at (.5,.5) of radius .25. Output: graph of the time history of the radius and angle.
- 3. Two particles: located at (.5,.25) and (.5,.75) both with strength  $1/h^2$ . The number of grid points is given by N = 32. Both particles should move at a constant angular velocity on a circle centered at (.5,.5) of radius .25. Output: graph of the time history of the radius and angle for both particles.
- 4. Two-patch problem. For each point  $i \in [0...N_p]$ ,  $N_p = 128, 256$ , place a particle at the point  $ih_p$ ,  $h_p = \frac{1}{N_p}$  provided that

$$||ih_p - (.5, .375)|| \le .12 \text{ or } ||ih_p - (.5, .625)|| \le .12.$$

The strength of each of the particles should be  $h_p^2/h^2$ . This corresponds to a pair of patches of vorticity of constant strength. Take the grid spacing N=64. Integrate the solution to time T=15, plotting the result at least every 1.25 units of time (to make a nifty movie, plot every time step). Set  $\Delta t=.025$ . We will provide a reference solution against which you can compare yours.

By setting ANIMATION = TRUE in your makefile, you can produce a pair of plotfiles every time step (particle locations, vorticity field on the grid). The default is to produce a pair of plotfiles at the end of the calculation for the two-patch case.

## Description of Algorithm for Computing the Velocity Field

1. Depositing the charges in the particles on the grid.

$$\omega_{m{i}}^g = \sum_k \omega^k \Psi(m{i}h - m{x}^k)$$

where the  $x^k$ 's are the positions of the particles in a\_state displaces by the input a\_k's.

$$\omega^g \equiv 0$$

$$\begin{split} \boldsymbol{i}^k &= \left\lfloor \frac{\boldsymbol{x}^k}{h} \right\rfloor \\ \boldsymbol{s}^k &= \frac{\boldsymbol{x}^k - \boldsymbol{i}^k h}{h} \\ \boldsymbol{\omega}^g_{\boldsymbol{i}^k} + &= \boldsymbol{\omega}^k (1 - s_0^k) (1 - s_1^k) \\ \boldsymbol{\omega}^g_{\boldsymbol{i}^k + (1,0)} + &= \boldsymbol{\omega}^k s_0^k (1 - s_1^k) \\ \boldsymbol{\omega}^g_{\boldsymbol{i}^k + (0,1)} + &= \boldsymbol{\omega}^k (1 - s_0^k) s_1^k \\ \boldsymbol{\omega}^g_{\boldsymbol{i}^k + (1,1)} + &= \boldsymbol{\omega}^k s_0^k s_1^k \end{split}$$

2. Convolution with the Green's function to obtain the potential on the grid, using Hockney's algorithm. The Hockney class will be constructed and maintained in ParticleSet - all you have to do is call it at the appropriate time.

$$\psi_{\boldsymbol{i}} = \sum_{\boldsymbol{j} \in \mathbb{Z}^2} G(\boldsymbol{i} - \boldsymbol{j}) \omega_{\boldsymbol{j}}^g$$

3. Compute the fields on the grid using finite differences.

$$\vec{U}_{i}^{g} = \left(\frac{\psi_{i+(0,1)} - \psi_{i-(0,1)}}{2h}, -\frac{\psi_{i+(1,0)} - \psi_{i-(1,0)}}{2h}\right)$$

4. Interpolate the fields from the grid to the particles.

$$ec{U}^k = \sum_{m{i} \in \mathbb{Z}^2} ec{U}_{m{i}} \Psi(m{x}^k - m{i}h)$$

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ight
floor \ oldsymbol{s}^k &= rac{oldsymbol{x}^k - oldsymbol{i}^k h}{h} \end{aligned}$$

$$\begin{split} \vec{U}^k = & \vec{U}^g_{\pmb{i}} (1 - s^k_0) (1 - s^k_1) \\ + & \vec{U}^g_{\pmb{i}+(1,0)} s^k_0 (1 - s^k_1) \\ + & \vec{U}^g_{\pmb{i}+(0,1)} (1 - s^k_0) s^k_1 \\ + & \vec{U}^g_{\pmb{i}+(1,1)} s^k_0 s^k_1 \end{split}$$

Note that the operatorParticleVelocities::operator() requires you to return in a\_k the quantities  $\Delta t \vec{U}^k$ .