

General Points (3 Imp points)

(1) Torque in a Stationary Object & Rotating Object:

When a torque is applied to a stationary object it moves in the direction of torque.

When a torque is applied to a rotating / spinning object, & if the value of Torque is much less compared to the angular mom. — then the ang mom. vector rotates. (Just like Centripetal force gives rise to circular motion)

(2) When a Vector rotates that means angular mom rotates then

$$\frac{d\vec{L}}{dt} = \vec{\omega} \times \vec{L} = \vec{\omega}_{\text{ext}} \times \vec{L} \quad \left| \quad \frac{d\vec{r}}{dt} = \vec{\omega} \times \vec{r} \right.$$

• precessional ang vel. written as Ω or ω_p .

• Very Very useful Formula.

(3) While using The formula no 2 — $\vec{\omega}_{\text{ext}}$ or \vec{L} can be calculated about a particular point

- 1) fixed pt on earth
- 2) about CM

(4) Diff between ω_p & ω_s .

ω of precession.

— also written as Ω

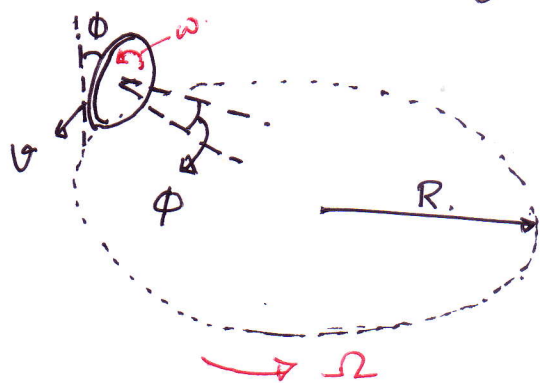
— you may also call this as orbital motion like that of earth

Spin angular velocity

$\omega_s \gg \omega_p$.

It gives rise to large angular momentum L_s

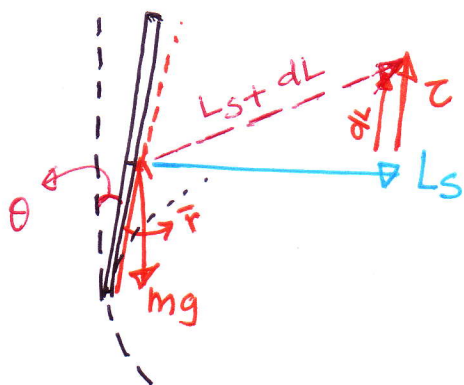
If you start a coin rolling on a table with care, you can make it roll in a circle. As you can see from the drawing, the coin "leans" inward with its axis tilted. The radius of the coin is b , the radius of the circle it follows on the table is R , and its velocity is v . Assume that there is no slipping. Find the angle ϕ that the axis makes with the horizontal. (Kleppner - 7.6)



⇒ It is again a very interesting problem from where we can learn a lot of things. All these problems are based on gyroscopic principle.

If a body is stationary & you apply a torque, we know very well its motion. It will rotate in the direction of torque.

But if the body is in circular motion / rotating about an axis & its angular momentum vector because of its rotation ^{L_s} is quite large compared to torque then the angular momentum vector L_s rotates. This kind of situations are observed in different cases like Gyroscope, bicycle etc. The above example is exactly like that of a bicycle - bicycle does not fall in motion but front wheel either moves towards the left or right due to torque. - dL .



Since the resultant L vector rotates the ω will rotate (changes direction) in a circle.

So because of torque the angular ~~vel~~ momentum vector rotates. So we need to apply.

$$(1) \quad \frac{d\vec{L}}{dt} = \vec{\omega} \times \vec{L} = \vec{\tau}_{ext}.$$

The second important point about these problems. See all the forces acting on the body. We are tempted to take the torque about a point where the ω touches the ground. But that is an accelerating point — a changing point. We need to take the torque about a fixed pt or about the center of mass. Since we don't find any fixed point, so we need to calculate torque about center of mass. That is very important.

$$\vec{\tau}_{ext} = \vec{\omega} \times \vec{L}$$

Taking Torque about CM.

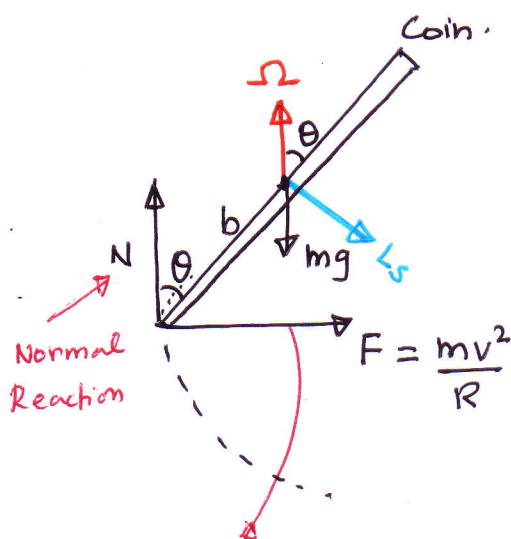
$$Nb \sin \theta - Fb \cos \theta = \omega L_s \sin(90 + \theta)$$

$$mgb \sin \theta - \frac{mv^2}{R} b \cos \theta = \omega L_s \cos \theta.$$

$$L_s = I \omega_s = \frac{mb^2}{2} \omega_s = \frac{mb^2}{2} \frac{v}{b} = \frac{mbv}{2}$$

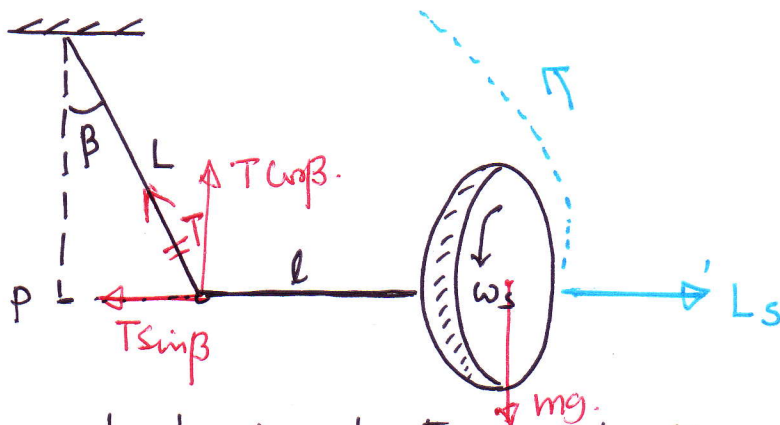
$$\omega = \frac{v}{R}$$

$$\Rightarrow \tan \theta = 3v^2 / 2gR$$



Force of friction giving necessary centripetal force for the body to rotate in a circle.

7.3



We have to find out the angle β . β is very small such that $\sin\beta \approx \beta$.

• Here we can apply the same equation.

$$\frac{d\vec{L}}{dt} = \vec{\omega} \times \vec{L} = \vec{\tau}_{ext}$$

Force $T \cos\beta = Mg$ — (1)

$$T \sin\beta = \frac{Mv^2}{l + L \sin\beta} = M \underline{\Omega^2} (l + L \underline{\sin\beta})$$
 — (2)

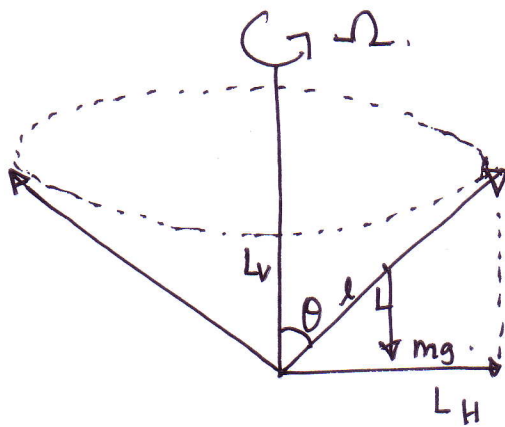
Now we can calculate the Torque about some fixed pt P or about the center of Mass. It will be simpler if we calculate about COM. In the notes given otherwise.

Torque $\frac{T \cos\beta}{\downarrow Mg} l = \Omega L_s = \Omega I_0 \omega_s$
so $Mgl = \underline{\Omega} I_0 \omega_s$ — (3)

Above all the three eqns are given eqns.

$$\begin{aligned} \sin\beta &\approx \beta & \cos\beta &\approx 1 \\ T \sin\beta &= M \Omega^2 (l + L \sin\beta) \\ T \beta &= M \Omega^2 (l + L \beta) \\ Mg \beta &= M \Omega^2 (l + L \beta) \Rightarrow \end{aligned} \left| \begin{aligned} \left(\frac{l + L \beta}{\beta} \right) &= \frac{g}{\Omega^2} = \frac{g (I_0 \omega_s)^2}{(Mgl)^2} \\ \frac{l}{\beta} + L &= () \\ \frac{l}{\beta} &= () - L \\ \beta &= l / () - L \end{aligned} \right.$$

Method-2 (Simpler) Top precessing with Ω



$$\left(\frac{dL}{dt}\right) = \Omega L \sin \theta$$

$$L = I \omega_s$$

$$\Omega I \omega_s \sin \theta = mg l \sin \theta$$

$$\Omega = \frac{mg l}{I \omega_s}$$

→ Earlier in the notes I have solved the problem of Top precessing with angular velocity Ω & calculated its value using cylindrical polar coordinates

→ The above solution is much simpler using the concept of rotating vector that is L_s vector rotates

So
$$\frac{dL_s}{dt} = \Omega \times L_s = \Omega L_s \sin \theta$$

Using this concept the problem can be solved in one step only.