

# Physics 211C: Solid State Physics

Instructor: Prof. Tarun Grover

Lecture 13

Topic: Mean-field for AFM (large  $S$ ), Schwinger boson approach, Intro to  $\mathbb{Z}_2$  gauge theory

Large -  $S$  for AFM

$$\mathcal{H} = |J| \sum \vec{s}_i \cdot \vec{s}_j$$

(classically)

$\uparrow \downarrow \uparrow \downarrow \uparrow \downarrow \uparrow \downarrow$

Neel state (not an eigenstate of  $\mathcal{H}_{AFM}$ )

$$\mathcal{H}_{FM} |\uparrow\uparrow \dots \uparrow\rangle = E_{gs} |\uparrow \dots \uparrow\rangle$$

$$S^+ = S_x + iS_y = \left( \sqrt{2S - b^\dagger b} \right) b$$

$$S^2 = S - b^\dagger b \quad (b^\dagger b) \text{ small}$$

$$\begin{aligned} \text{define } \tilde{S}^z &= -S^z \text{ on B sublattice} \\ \tilde{S}^x &= S^x, \quad \tilde{S}^y = -S^y \end{aligned} \quad \left. \begin{array}{l} \text{for comm. relations} \\ \text{to stay the same} \end{array} \right\}$$

$$S^z \text{ or } \tilde{S}^z = S - n_b$$

$$S^+ \text{ or } \tilde{S}^+ = \left( \sqrt{2S - n_b} \right) b$$

$$\begin{aligned} \mathcal{H} &= +|J| \sum_{\langle ij \rangle} \left[ S_i^z S_j^z + \frac{1}{2} (S_i^+ S_j^- + \text{h.c.}) \right] \\ &= |J| \sum_{\langle ij \rangle} \left[ -\tilde{S}_i^z \tilde{S}_j^z + \frac{1}{2} (\tilde{S}_i^+ \tilde{S}_j^- + \text{h.c.}) \right] \end{aligned}$$

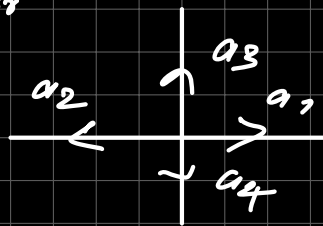
$$\sqrt{2s - \eta_b} \approx (\sqrt{2s}) \left[ 1 - \frac{\eta_b}{4s} + \dots \right]$$

$$\mathcal{H} = \mathcal{H}_0 + \mathcal{H}_1 + \mathcal{H}_2 + \dots$$

$$\mathcal{H}_0 = -s^2 \frac{|J| V_Z}{2}$$

$$\mathcal{H}_1 = |J| s z \sum_{\mathbf{k}} \left[ b_{\mathbf{k}}^\dagger b_{\mathbf{k}} + \frac{\gamma_{\mathbf{k}}}{2} (b_{\mathbf{k}}^\dagger b_{-\mathbf{k}}^\dagger + \text{h.c.}) \right]$$

$$\gamma_{\mathbf{k}} = \frac{1}{Z} \sum_{\vec{a}_i} e^{i \vec{k} \cdot \vec{a}_i}$$



$$= \frac{1}{4} [2 \cos k_x a + 2 \cos k_y a]$$

$$\mathcal{H}_2 = -\frac{|J|}{4} \sum_{\langle i,j \rangle} ( \dots )$$

$$b_{\mathbf{k}} = \cosh(\theta_{\mathbf{k}}) a_{\mathbf{k}} + \sinh(\theta_{\mathbf{k}}) (a_{-\mathbf{k}}^\dagger)$$

$$[a_{\mathbf{k}}, a_{\mathbf{k}'}^\dagger] = \delta_{\mathbf{k}\mathbf{k}'}$$

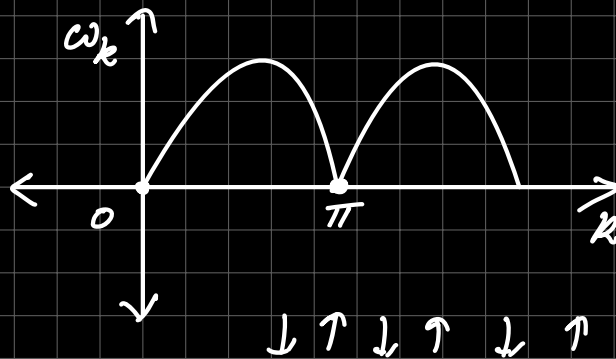
$$\sum_{\mathbf{k}} a_{\mathbf{k}}^\dagger a_{\mathbf{k}} E_{\mathbf{k}}$$

$$\begin{aligned} \mathcal{H}_1 = |J| s z \sum_{\mathbf{k}} & \left[ \cosh(2\theta_{\mathbf{k}}) + \gamma_{\mathbf{k}} \sinh(2\theta_{\mathbf{k}}) \right] a_{\mathbf{k}}^\dagger a_{\mathbf{k}} \\ & + \frac{1}{2} \left[ \sinh(2\theta_{\mathbf{k}}) + \gamma_{\mathbf{k}} \cosh(2\theta_{\mathbf{k}}) \right] [a_{\mathbf{k}}^\dagger a_{-\mathbf{k}}^\dagger + \text{h.c.}] \end{aligned}$$

$$\mathcal{H}_1 = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \left[ a_{\mathbf{k}}^\dagger a_{\mathbf{k}} + \frac{1}{2} \right] - \frac{|J| Z N s}{2}$$

$$\omega_{\mathbf{k}} = |J| s z \sqrt{1 - \gamma_{\mathbf{k}}^2} \rightarrow \text{vanishes at } \vec{k} = 0 \text{ \& } \vec{k} = (\pi, \pi, \dots, \pi) \equiv \vec{Q}$$

$$\vec{s}_k = \frac{\cos k_x + \cos k_y}{2} \quad \left. \begin{array}{l} \rightarrow k=0 \\ \rightarrow k=-1 \end{array} \right\} \vec{s}_k^2 = 1$$



$$s_i^z s_j^z \quad |s|=1$$

$$\approx |s(\mathbf{k})|^2 (\cos k_x + \cos k_y)$$

$$k \approx 0$$

$$\omega_k \approx J s \sqrt{2} |k| = c_s |k|$$

$$k \approx \vec{q}$$

$$\omega_k \approx J s \sqrt{2} |\vec{k} - \vec{q}|$$

energy diff

$$\Delta E = E_{\text{quantum}} - E_{\text{classical}}$$

$$= \frac{1}{2} \sum_k |J| s z \left( -\sqrt{1 - \vec{s}_k^2} - 1 \right) \leq 0$$

Fluctuations / Stability :-

$$\Delta m = \frac{1}{V} \sum_{\vec{r}} e^{i \vec{q} \cdot \vec{r}} \langle s^z(\vec{r}) \rangle - s$$

( $\equiv m_q - m_{cl}$ )

$$= \frac{1}{V} \sum_{\vec{r}} [s - m_b(\vec{r})] - s$$

$$= - \int \frac{d^d k}{(2\pi)^d} \langle b_k^\dagger b_k \rangle$$

$$\text{check} \rightarrow = \frac{1}{2} - \int \frac{d^d k}{(2\pi)^d} \frac{(m_B(k) + \frac{1}{2})}{\sqrt{1 - v_k^2}}$$

$$\approx - \int \frac{d^d k}{k(e^{\beta k} - 1)} + \int \frac{d^d k}{k}$$

diverges in  $d=1$ , contributes at  $T=0$

$\therefore$  At  $T=0$ , no order in  $d=1$

At  $T>0$ ,  $T \int \frac{d^d k}{k^2}$  diverges in  $\underline{d \leq 2}$ .

Some discussions on

mermin-Wagner

theorem

## Large $N$ approximation

$SU(2) \rightarrow SU(N)$  using Schwinger bosons

$$\vec{S} = b^\dagger \frac{\vec{\sigma}}{2} b \quad b = [b_\uparrow, b_\downarrow]$$

$$S^+ = b_\uparrow^\dagger b_\downarrow \quad S^z = \frac{(b_\uparrow^\dagger b_\uparrow - b_\downarrow^\dagger b_\downarrow)}{2}$$

$$[b_\uparrow, b_\uparrow^\dagger] = 1$$

$$[b_\uparrow, b_\downarrow^\dagger] = 0 \quad \vec{S}^2 = S^x^2 + S^y^2 + S^z^2 = S(S+1)$$

$$[b_\downarrow, b_\downarrow^\dagger] = 1$$

$$b_\uparrow^\dagger b_\uparrow + b_\downarrow^\dagger b_\downarrow = 2S$$

$$|s, m\rangle = \frac{(b_\uparrow^\dagger)^{s+m} (b_\downarrow^\dagger)^{s-m} |0\rangle}{\sqrt{(s+m)! (s-m)!}}$$

$Z^\dagger Z$  is invariant  $SU(N)$  transformation.

$$Z \rightarrow U Z \quad U \in SU(N)$$

$$S^x, S^y, S^z \longrightarrow [S_\alpha\beta, S_\gamma\delta] = \delta_{\beta\gamma} S_\alpha\delta - \delta_{\alpha\delta} S_\gamma\beta$$

$$S_\alpha\beta^\dagger = b_\alpha^\dagger b_\beta$$

check that it works out