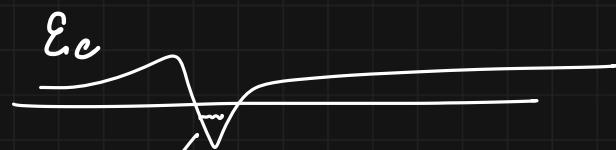


Recap

2D system \rightarrow not only where topo is relevant

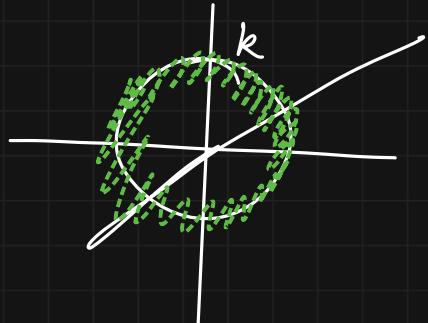


one $\text{f}-\text{semiconductor}$ junction

trapped in a potential well

drift \rightarrow drude picture \rightarrow elastic collisions, momentum gets scrambled

#

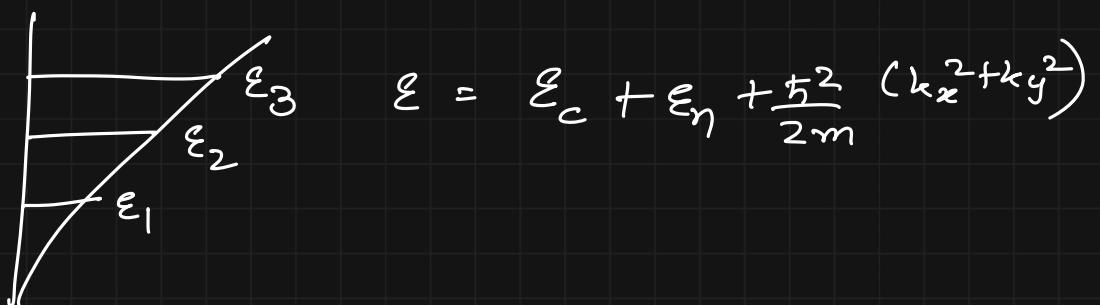


$$\left(\frac{dp}{dt}\right)_{\text{scattering}} = \left(\frac{dp}{dt}\right)_{\text{field}}$$

$$V_d = e \frac{Zm}{m} E \quad \text{or} \quad \mu = \frac{eZm}{m}$$

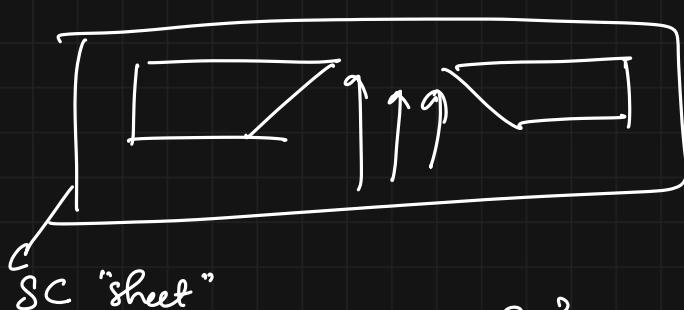
concept of subbands \rightarrow confinement along

$$\psi(z) = \phi_n(z) e^{ik_x x} e^{ik_y y} \quad z\text{-direction}$$



$$\epsilon = \epsilon_c + \epsilon_\eta + \frac{\hbar^2}{2m} (k_x^2 + k_y^2)$$

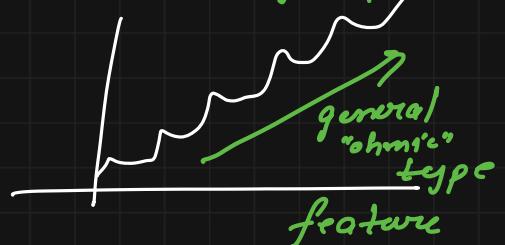
PRL paper \rightarrow quantized conductance



$$\sigma = \frac{2e^2}{h} n$$

Gate voltage

\hookrightarrow # of states that are filled



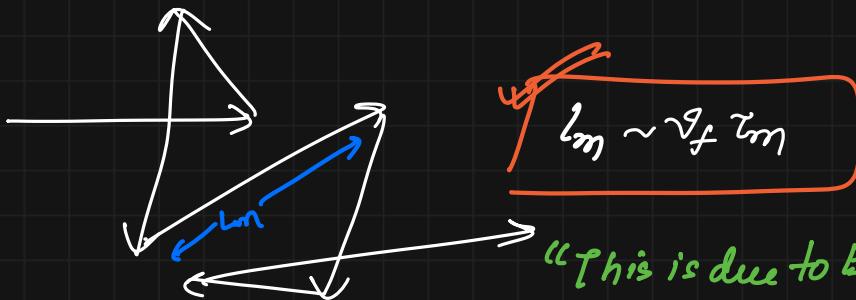
Question 2

diffusive model works? or do we need something extra?

remark

length scales

momentum scattering length.



"This is due to ELASTIC scattering & not INELASTIC scattering"

Problem:-

$$\# \text{ states} \rightarrow \frac{1}{4\pi^2} r^2 \times \frac{d^2 k}{m}$$

$$E = \frac{\hbar^2 k^2}{2m}$$

$$= \frac{1}{2\pi^2} \cdot k dk (2\pi)$$

$$dE = \frac{\hbar^2}{m} k dk$$

$$= \frac{1}{\pi} k dk$$

$$= \frac{1}{\pi} \frac{dE}{\hbar^2} \frac{m}{\hbar^2} \Rightarrow \frac{m}{\pi \hbar^2} = g(E)$$

$$\int \frac{1}{\pi} \times (1) \times dE = \frac{m}{\hbar^2 \pi} E_F = \eta$$

indep of energy

$$\frac{m}{\hbar^2} \times \frac{1}{\pi} \frac{\hbar^2 k_F^2}{2m} = \eta$$

$$k_F = \sqrt{2\pi\eta}$$

periodic BCs

$$\sqrt{\frac{dn}{dE}} = \frac{m}{\pi \hbar^2}$$

$$m V_F = \hbar k_F \Rightarrow V_F = \frac{\hbar k_F}{m} = \frac{\hbar}{m} \sqrt{2\pi\eta}$$

$$L_m = V_F \tau_m \rightarrow \text{momentum scattering}$$

1 more
length
scale

$$\lambda_f = \frac{1}{k_F} \rightarrow \lambda_f = \text{de-Broglie wavelength of electrons}$$

$$\lambda_f \sim (n^{-1/3})$$

$$k_f \sim (n)^{1/3}$$

$$\lambda_f \sim (n)^{-1/3}$$

Problem

What is the typical λ_f for electrons in typical metals?

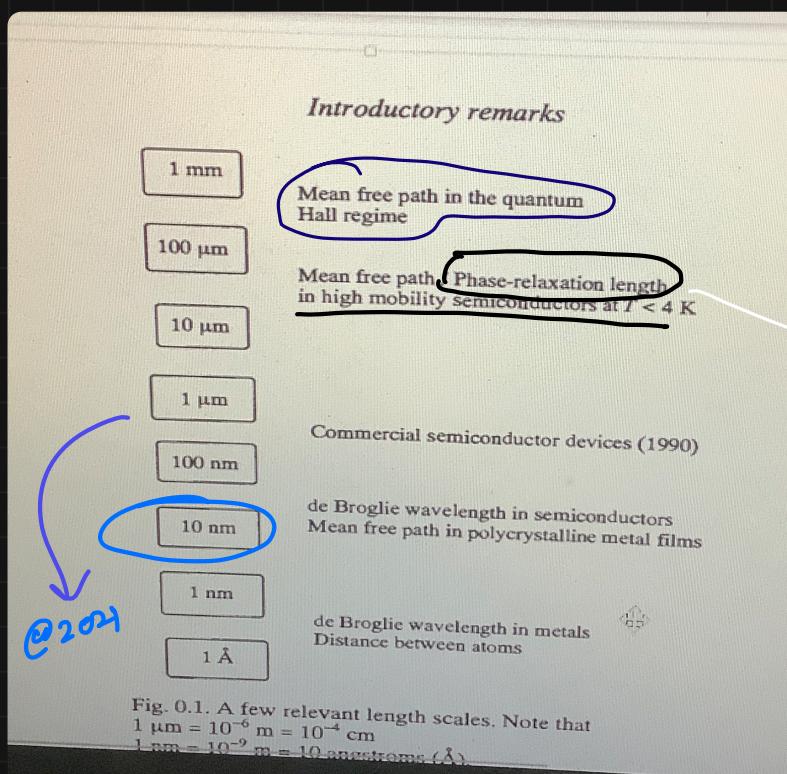
Ans:- $\lambda_f \sim A^0$ (for typical metals) $n \sim 10^{22} \text{ cm}^{-3}$

atomic spacing

$$\lambda_f \sim (n)^{1/3}$$

λ_f semicond ~ larger than metals

↳ essentially as small as n



→ S. Datta

gives cliff
regimes of
transport

Phase relaxatⁿ

If source of collision is inelastic \rightarrow "phase" of electron scrambled

Phase coherence time scale :-

drops an $e\Theta$ into a solid \rightarrow gets scattered due to "inelastic collisions"

(e Θ , phonons, magnetic)

quantum coherence of phase is lost.

(decoherence/dephasing)

elastic scattering \rightarrow no decoherence

\hookrightarrow might add a const. phase

on some timescale $\tau_f \rightarrow e\Theta$ becomes uncorrelated from its

initial phase.

interference effects wash out.

Dist $e\Theta$ moves in this time in a diffusive system

$l_\phi = (D \tau_f)^{1/2}$ = "Coherence length"

Length scales

As inelastic processes freeze out, τ_f is expected to diverge, with power law exponent set by dimensionality and inelastic mechanism:

electron phonon: $\tau_\phi \sim T^{-3}$

electron-electron: $\tau_\phi \sim T^{2(d-4)}$

Phase coherence length $l_\phi = (D \tau_\phi)^{1/2} \uparrow$ as $T \downarrow$

Mean free path $l_e = v_F \tau$ dominated by elastic scattering (impurities, roughness) and is stays constant

mesoscopic physics regime

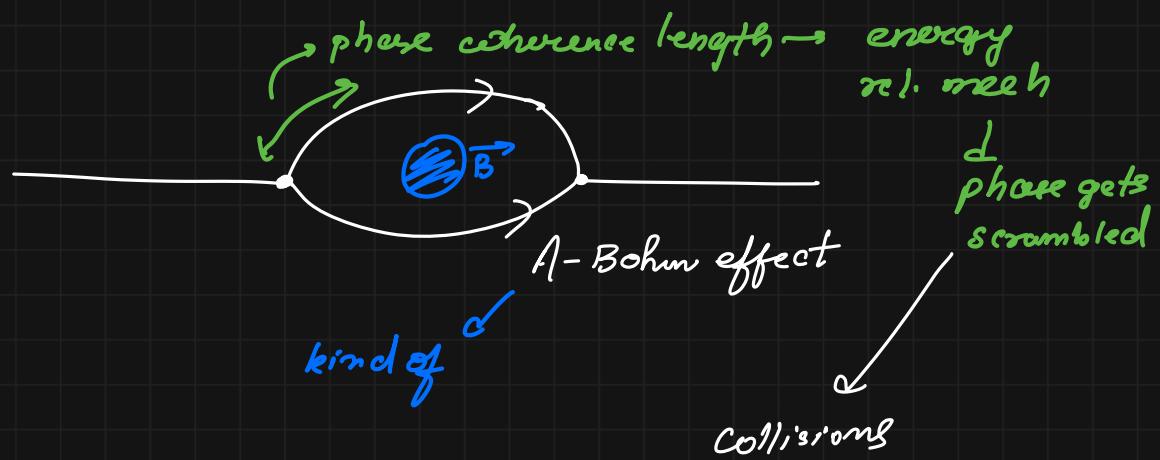
Sample channel length L
 $L < l_\phi$ ballistic,
 $L > l_\phi$ phase coherent,
 $L > l_\phi$ incoherent

$l_\phi > l_e$ at low temperatures, and lots of interesting phenomena observed when $l_e < L < l_\phi$

weak-anti
localizat^n

weak

expt \rightarrow double slit

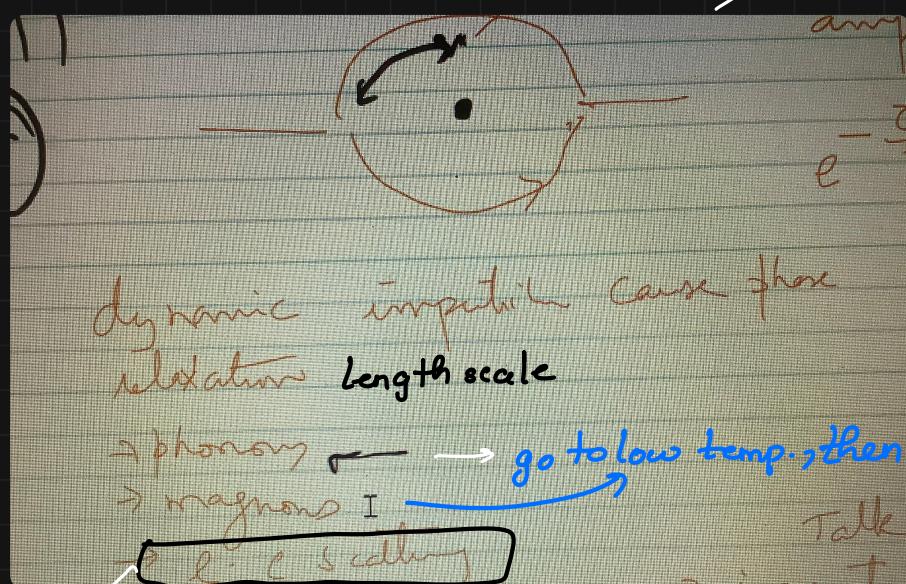
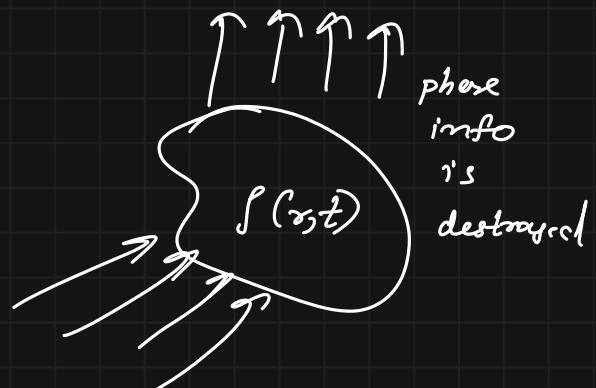
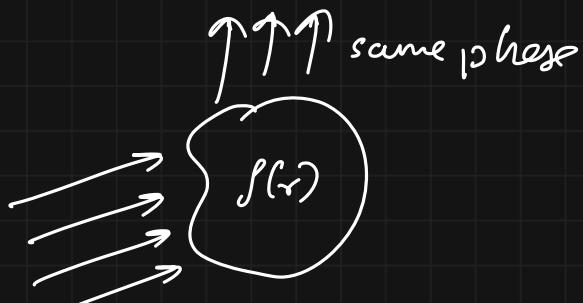


BIG picture

All collisions that have fixed target, then the phase is same in & out.

\hookrightarrow "no time dependence"

\hookrightarrow t dependence \Rightarrow



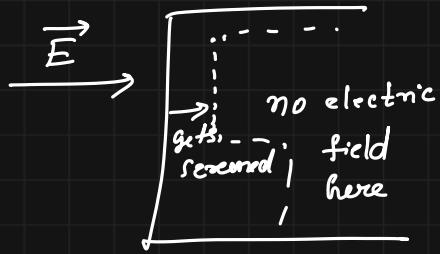
\hookrightarrow at low temp., it's still there.

even in non-correlated system, it's there.

"metals"

lengthscale

Screening length λ_s



diff models
Thomas-fermi
model
 $\lambda_{TF} \sim (\rho_F)^{\text{power}} \propto \text{const.}$
 $\sim \text{few } \text{\AA}$ in metals

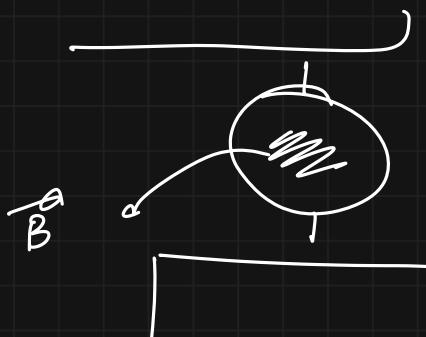
$\lambda_{TF} \sim (\rho_F)^{\text{power}} \propto \text{const.}$
in SC the lengthscale
can be large

e.g. for transistors \rightarrow have a large screening "length"

$R_f < \text{atomic spacing} \rightarrow \text{Particle in a box}$

Paper
measuring
phase
coherence
length

Mesoscopic "physics"



G vs B_{ext}

faster beating oscillations

\hookrightarrow a resetting of phase & oscillation

$d/a < \text{cor. length} \Rightarrow \text{oscillations wash away}$

Φ_0 component $\rightarrow T = \text{knob for phonons}$

amplitude of oscillations
washes away

control inelastic collisions

Datta's Book → phase coherence length

AEM

→ L-lengths → not observed for short magnetic field.

↓
always about the numbers

