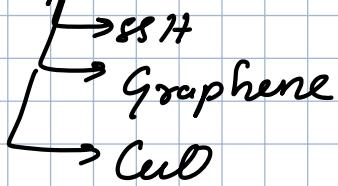
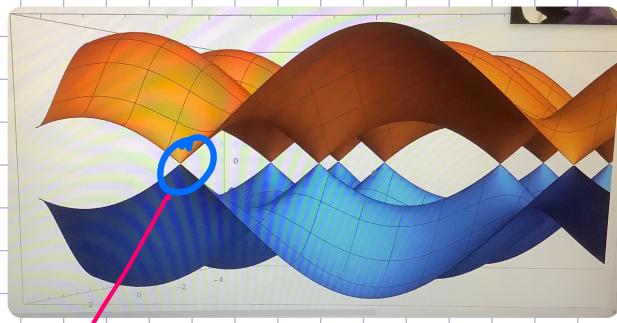
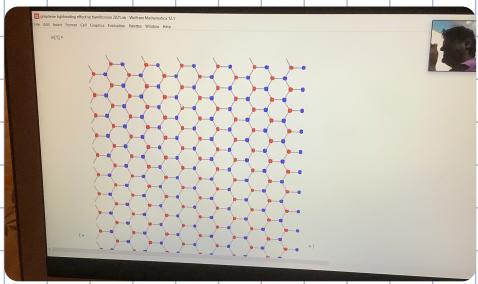


## # Review of tBM from previous class



## # Effective hamiltonian of Graphene near touching points.



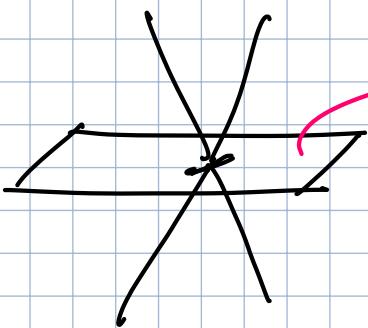
semi-metal  
low densities  
 ↳ gapless but only at  
 precise points  
 metals have more k-space areas  
 where it's gapless.

→ p-orbitals → bands  $1e^\ominus$

$\frac{1}{2}$  filled

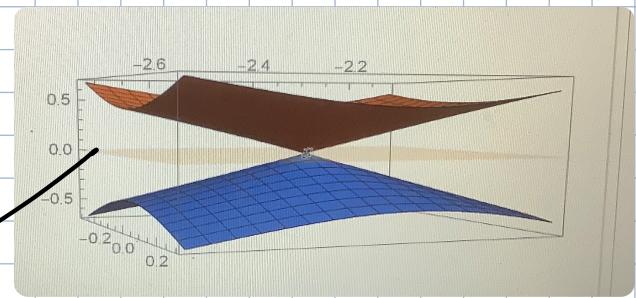
estates arise from

hybridization of these  
orbitals



due to  $1e^\ominus$  of p-orbitals  
 $\Rightarrow \frac{1}{2}$  filling

pristine  
 graphene's  
 fermi energy

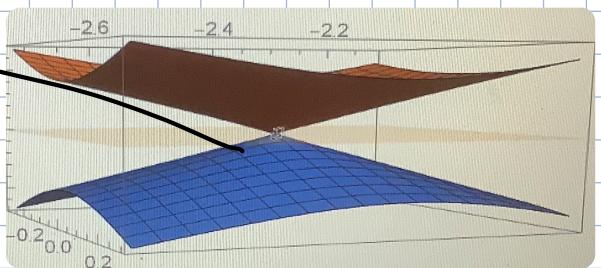


# Taylor series expansion can be done in mathematica.

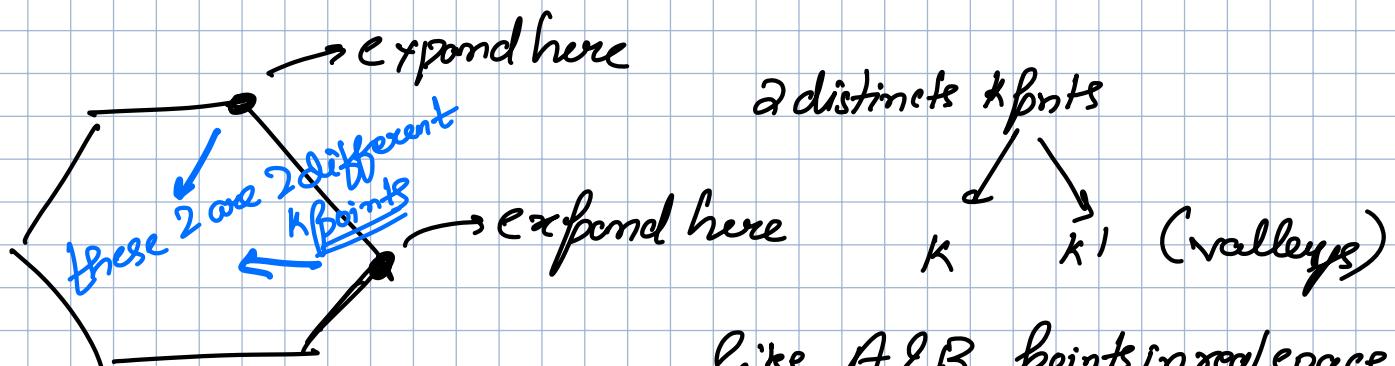
## # Effective Hamiltonian :-

- Linear dispersion for low energy excitations

$$H = \begin{pmatrix} 0 & \frac{-i(2\sqrt{3})}{\hbar} (q_x - iq_y) \\ \frac{i(2\sqrt{3})}{\hbar} (q_x + iq_y) & 0 \end{pmatrix}$$



$H = \vec{\sigma} \cdot \vec{p}$  → Dirac Hamiltonian for massless particles.

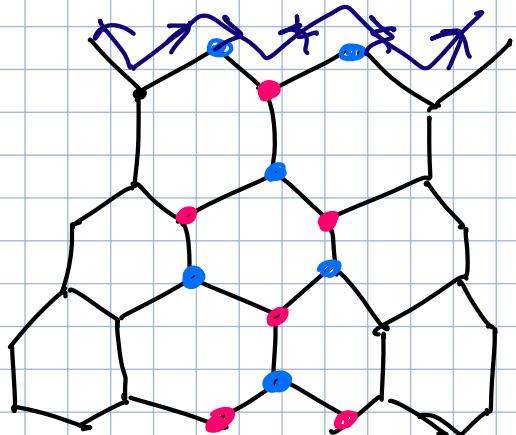


Like A & B points in real space.

→ Structure of Hamiltonian is different at these K points.

## Analogue of Gating for graphene

# Let's consider a GNR.



Q. What are the electronic states present in GNR?

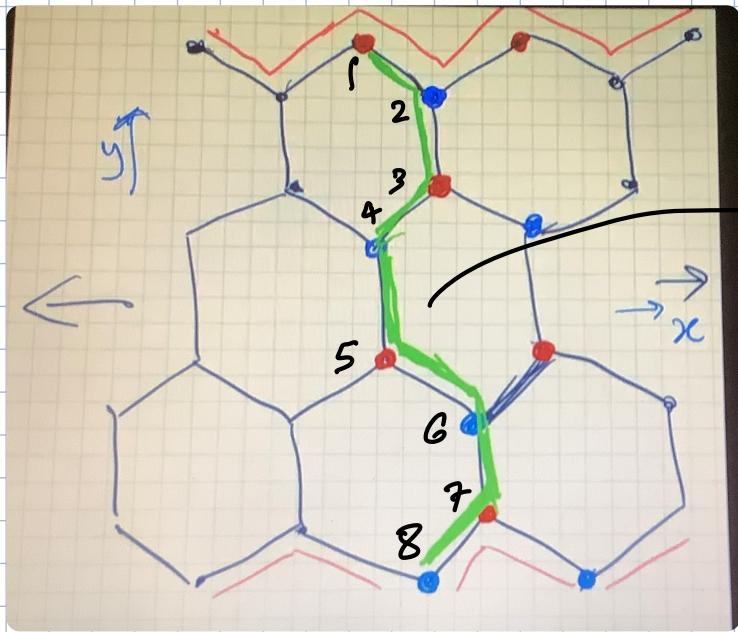
Motivation:- From QM, we know that edge states can be chiral & be ballistic.

We can now see if something similar is

possible in Graphene.

# How do we do it?

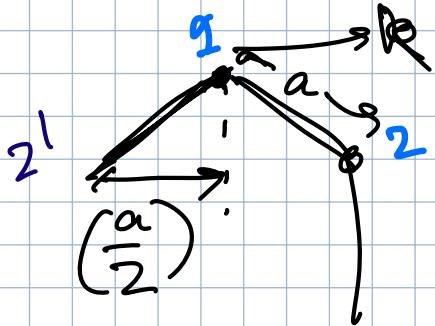
Broute force: Define a Gigantic Unit cell.



Because  $\gamma$ -directional transition symmetry is absent.

Call this a molecule

Now make a hopping matrix.



$$t e^{-ikx_0} + t e^{-ik\sum} \rightarrow \underline{\underline{M}}$$

Disagreement:

$$t e^{-ik\frac{a}{2}} + t e^{+ik\frac{a}{2}} \rightarrow \underline{\underline{M}}$$

he's considering  
intracell  
non-trivial  
contributions  
as well

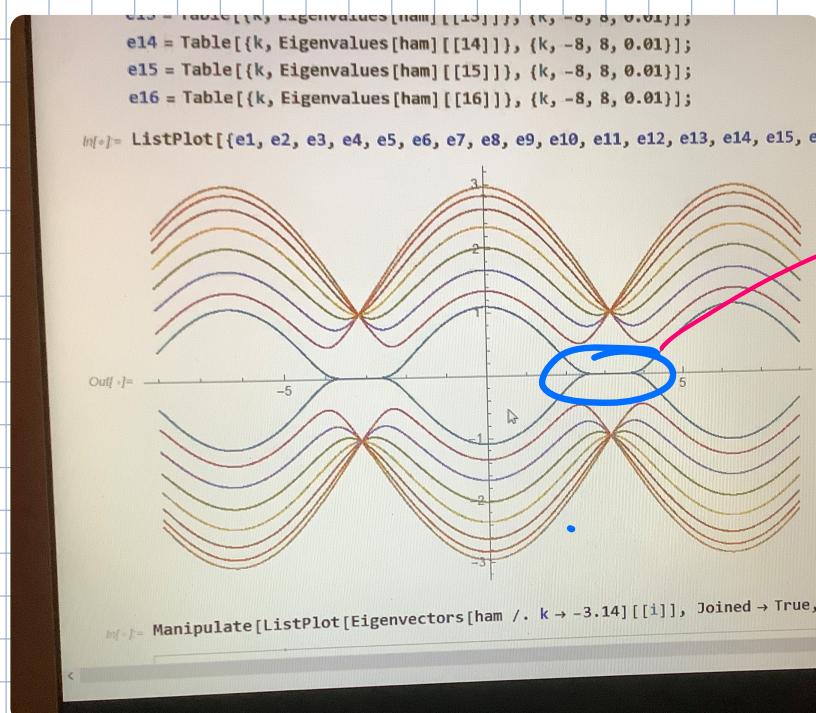
In[1]:= MatrixForm[ham]

MatrixForm=

0	2 Cos[ $\frac{k}{2}$ ]	0	0	0	0	0	0	0	0	0	0	0	0	0	0
2 Cos[ $\frac{k}{2}$ ]	0	1	0	0	0	0	0	0	0	0	0	0	0	0	0
0	1	0	2 Cos[ $\frac{k}{2}$ ]	0	0	0	0	0	0	0	0	0	0	0	0
0	0	2 Cos[ $\frac{k}{2}$ ]	0	1	0	0	0	0	0	0	0	0	0	0	0
0	0	0	1	0	2 Cos[ $\frac{k}{2}$ ]	0	0	0	0	0	0	0	0	0	0
0	0	0	0	2 Cos[ $\frac{k}{2}$ ]	0	1	0	0	0	0	0	0	0	0	0
0	0	0	0	0	1	0	2 Cos[ $\frac{k}{2}$ ]	0	0	0	0	0	0	0	0
0	0	0	0	0	0	2 Cos[ $\frac{k}{2}$ ]	0	1	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	1	0	2 Cos[ $\frac{k}{2}$ ]	0	0	0	0	0
0	0	0	0	0	0	0	0	0	2 Cos[ $\frac{k}{2}$ ]	0	1	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	2 Cos[ $\frac{k}{2}$ ]	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	1	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	2 Cos[ $\frac{1}{2}$ ]	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0

In[2]:= Eigenvalues[ham];

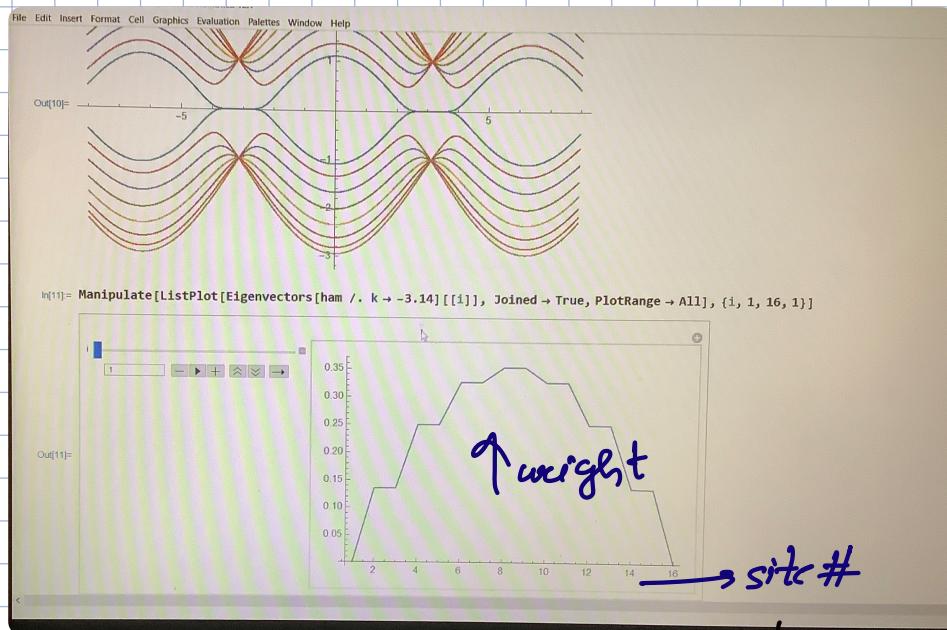
$\downarrow$   
we'll get  
16 bonds



$k = \pi$ , bonds are flat  
looks like  $\sqrt{q} = 0$ , doesn't disperse.

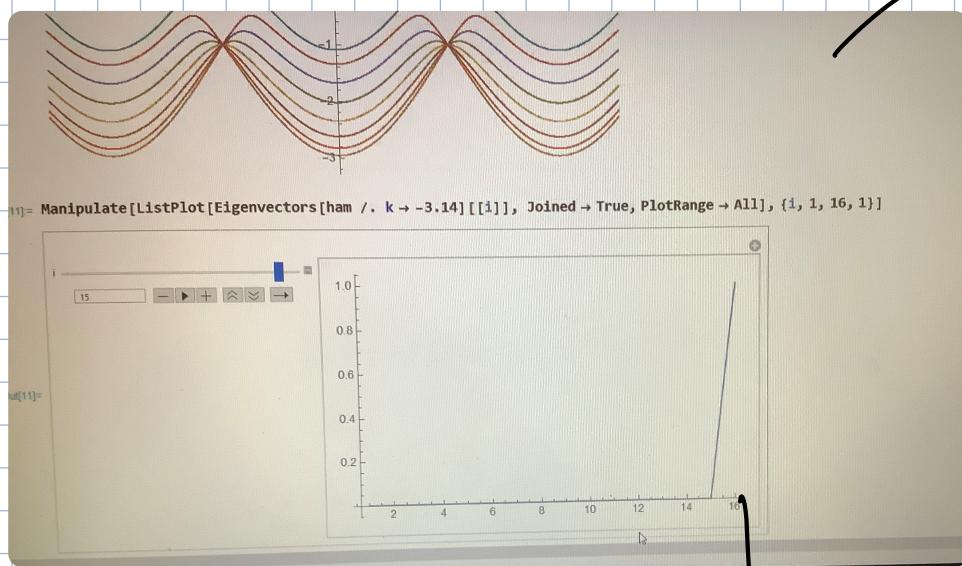
# now look at eigenvectors!

$\begin{pmatrix} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{pmatrix} \rightarrow \text{amplitude}$   
 $\sim \text{weight}$   
 $\text{on a}$   
 $\text{particular size.}$

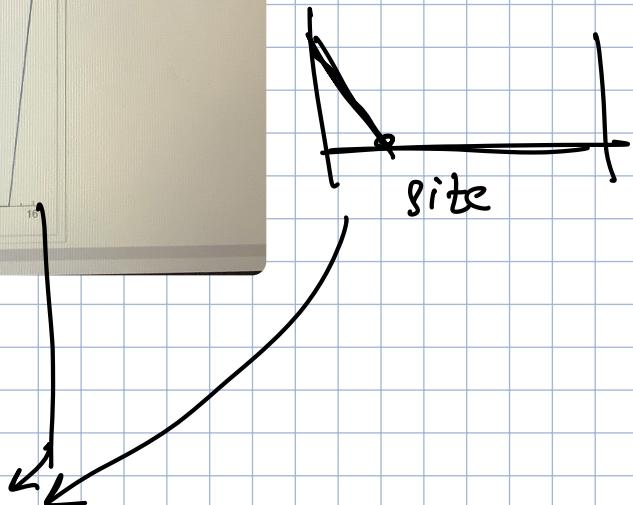


Behold the edge state :-

$\rightarrow$  at  $k = \underline{\underline{-3.14}}$



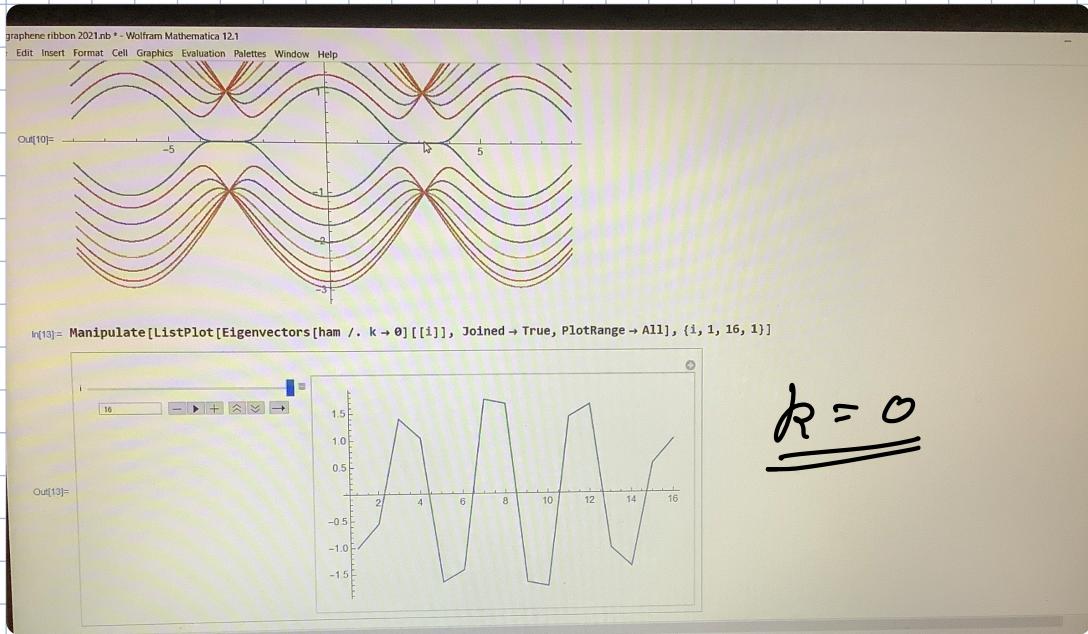
Similarly



$V_g = 0$ , so they're stuck.

- eigenvectors reveal such features  $\rightarrow$  reveal the effect of boundary conditions

At  $k=0$ , it's different:-



H.W.:- Do this for armchair GNR