Assignment #10 Selected topics in LGT

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The expression for which we want to integrate out the fermionic terms is:

$$I = \int \left\{ \prod_{k=0}^{N} d\bar{\psi}_k d\psi_k \right\} \exp(-\bar{\Psi}M\Psi + \bar{\eta}\Psi + \bar{\Psi}\eta)$$
 (1)

Let us take for now as a given (and to-prove later on) the following expression:

$$\psi_{i}' = \sum_{j=0}^{N} M_{ij} \psi_{j} \Rightarrow \Psi' = M \Psi \Rightarrow \prod_{k=0}^{N} d\psi_{k} = \det(M) \prod_{k=0}^{N} d\psi_{k}'$$
 (2)

Then, we can transform equation 1 into the following form:

$$I = \det(M) \int \left\{ \prod_{k=0}^{N} d\bar{\psi}_k d\psi_{k'} \right\} \exp(-\bar{\Psi}\Psi' + \bar{\eta}M^{-1}\Psi' + \bar{\Psi}\eta)$$
 (3)

To decompose the exponential in the integrand, note that (taking $a = \bar{\psi}_i \psi_i'$ and $b = \bar{\psi}_j \psi_j'$):

$$ab = \bar{\psi}_i \psi_i' \bar{\psi}_j \psi_j' = \bar{\psi}_j \psi_j' \bar{\psi}_i \psi_i' = ba$$
(4)

and the same applies for the other two terms being added in the argument of the exponential in equation 3 (i.e. $\bar{\eta}_i(M^{-1})_{ij}\psi_j'$ and $\bar{\psi}_i\eta_i$), which means that all the terms being added in the argument of that exponential commute. Therefore, we can separate the exponential in equation 3 in the following way (using the fact that $e^Ae^B=e^{A+B}$ if [A,B]=0):

$$\exp(-\bar{\Psi}\Psi' + \bar{\eta}M^{-1}\Psi' + \bar{\Psi}\eta) = \exp(-\bar{\Psi}\Psi')\exp(\bar{\eta}M^{-1}\Psi')\exp(\bar{\Psi}\eta) = \left(\prod_{k=0}^{N} \exp(-\bar{\psi}_{k}\psi'_{k})\right) \left(\prod_{k=0}^{N} \exp(\sum_{j} \bar{\eta}_{k}(M^{-1})_{kj}\psi'_{j})\right) \left(\prod_{k=0}^{N} \exp(\bar{\psi}_{k}\eta_{k})\right) = \left(\prod_{k=0}^{N} \exp(-\bar{\psi}_{k}\psi'_{k})\right) \left(\prod_{j=0}^{N} \exp(\sum_{k} \bar{\eta}_{k}(M^{-1})_{kj}\psi'_{j})\right) \left(\prod_{k=0}^{N} \exp(\bar{\psi}_{k}\eta_{k})\right) = \left(\prod_{k=0}^{N} (1 - \bar{\psi}_{k}\psi'_{k})\right) \left(\prod_{j=0}^{N} (1 + \sum_{k} \bar{\eta}_{k}(M^{-1})_{kj}\psi'_{j})\right) \left(\prod_{k=0}^{N} (1 + \bar{\psi}_{k}\eta_{k})\right)$$

$$(5)$$

Now, to avoid cumbersome expressions, let us focus on integrals of the form (note that the exponentials can be safely moved around without carrying any -1 sign, as we saw above; the only sign change is coming from moving the $\mathrm{d}\bar{\psi}_k$ to the right to perform each integration):

$$I_{k} = \int d\bar{\psi}_{k} d\psi_{k}' f_{k} = \int d\bar{\psi}_{k} d\psi_{k}' (1 - \bar{\psi}_{k} \psi_{k}') \left(1 + \sum_{j} \bar{\eta}_{j} (M^{-1})_{jk} \psi_{k}' \right) (1 + \bar{\psi}_{k} \eta_{k})$$

$$(6)$$

Expanding the integrand (i.e. f_k) in equation 6 (and displaying the terms relevant to the integration only):

$$f_k = \dots = -\bar{\psi}_k \psi_k' + \left(\sum_m \bar{\eta}_m (M^{-1})_{mk} \psi_k'\right) \bar{\psi}_k \eta_k + \dots$$
 (7)

and then:

$$I_k = 1 + \left(\sum_m \bar{\eta}_m(M^{-1})_{mk}\right) \eta_k = \exp\left(\sum_m \bar{\eta}_m(M^{-1})_{mk} \eta_k\right)$$
 (8)

With all of these reductions, the integral I takes the form:

$$I = S \det(M) \prod_{k=0}^{N} I_k = S \det(M) \prod_{k=0}^{N} \exp\left(\sum_{m} \bar{\eta}_m (M^{-1})_{mk} \eta_k\right)$$
(9)

where S is the sign factor.

Reducing equation 9 some more we finally obtain:

$$I = S \det(M) \exp\left(\sum_{m,k} \bar{\eta}_m(M^{-1})_{mk} \eta_k\right) = S \det(M) \exp\left(\bar{\eta} M^{-1} \eta\right)$$
 (10)

1 Proving that $d\Psi = \det(M)d\Psi'$

...

2 Expression for S

Every time we want to form an integral I_k , we have to drag the necessary $d\bar{\psi}_k$ and $d\psi'_k$ differentials to the far right.

Now, the trick to simply things for the overall integration is to start by computing the integral I_k for k = N, and then proceed downwards towards k = 0 (in steps of 1 for each change in k); this is a simple strategy due to the possibility of moving the exponentials in the integrand around and without making any new sign changes pop-up.

Because the differentials $d\bar{\psi}_k$ and $d\psi'_k$ come in pairs, then there is no need to move any of those differential factors around, and therefore:

$$S = 1 \tag{11}$$