## IIIT-H, Monsoon Semester 24, Final Exam: 22<sup>nd</sup> November 2024 Introduction to Quantum Field Theory (SC1.421)

Full Marks 60, Duration: 180 min

There are 6 questions. The question paper consists of two sides of a sheet.

## 1. The Hamiltonian of a real scalar field is

$$H=\intrac{d^3k}{(2\pi)^3}E_kigg(a_k^\dagger a_kigg)$$
 ,

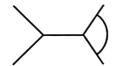
where  $a_k^\dagger$  and  $a_k$  are raising and lowering operators such that  $[a_k, a_{k'}^\dagger] = (2\pi)^3 \delta^3(k-k')$ , and  $[a_k, a_{k'}] = 0 = [a_k^\dagger, a_{k'}^\dagger]$ .

(a) Show that 
$$[H, a_k] = -E_k a_k$$

(a) Show that 
$$H(a_k|E\rangle) = (E - E_k)(a_k|E\rangle)$$
. What is the physical interpretation of the result. [3]

(c) Prove that 
$$\langle E|H|E\rangle \geq 0$$
. [3]

2. Consider the following Feynman diagram of a two-particle to two-particle scattering in a scalar field theory.



- (a) If the interaction term in the Lagrangian (density) that generates the diagram is  $-\frac{\lambda}{n!}\phi^n$  then what is the value of n?
- (b) Write down the amplitude (i.e., matrix element  $i\mathcal{M}$ ) for the diagram using Feynman rules. Momentum and their directions must be labeled. [6]
- 3. Without any specific representation, and only using the properties of  $\gamma$  matrices

(a) Show that 
$$\frac{1}{2}(1-\gamma_5)\frac{1}{2}(1-\gamma_5) = \frac{1}{2}(1-\gamma_5)$$
 and  $\frac{1}{2}(1-\gamma_5)\frac{1}{2}(1+\gamma_5) = 0$  [2]

(b) Explicitly show that 
$$\text{Tr}(\gamma^{\mu}\gamma^{\nu}\gamma^{5}) = 0$$
. Hint: You may use anti-commutation of  $\gamma$  matrices. [3]

(c) Calculate the trace 
$$\text{Tr}(\gamma^{\mu}p_1\gamma^{\nu}(1-\gamma_5)p_2)$$
, where  $p = \gamma_{\mu}p^{\mu}$ . [5]

- 4. The plane-wave solutions to the Dirac equation are:  $u^s(\vec{p}) = \begin{pmatrix} \sqrt{p \cdot \sigma} \xi^s \\ \sqrt{p \cdot \bar{\sigma}} \xi^s \end{pmatrix}$ , where  $\sigma^{\mu} = (1, \vec{\sigma})$  and  $\bar{\sigma}^{\mu} = (1, -\vec{\sigma})$  and  $\xi^s$ , with s = 1, 2, is a basis of orthonormal two-component spinors, satisfying  $(\xi^r)^{\dagger} \cdot \xi^s = \delta^{rs}$ . Show that
  - (a)  $\mathbf{u}^r(\vec{p})^{\dagger} \cdot \mathbf{u}^s(\vec{p}) = 2p_0 \delta^{rs}$
  - (b)  $\bar{\boldsymbol{u}}^r(\vec{\boldsymbol{p}}) \cdot \boldsymbol{u}^s(\vec{\boldsymbol{p}}) = 2m\delta^{rs}$
  - (c)  $\sum_{s=1}^{2} u^{s}(\vec{p}) \bar{u}^{s}(\vec{p}) = \vec{p} + m$

You may need the identity  $(p.\sigma)(p.\bar{\sigma}) = p^2 = m^2$ .

[2+2+2]

5. Under Lorentz Transformation,

$$\psi^{\alpha}(x) \to S[\Lambda]^{\alpha}_{\beta} \psi^{\beta}(\Lambda^{-1}x),$$

where  $\Lambda = \exp\left(\frac{1}{2}\Omega_{\rho\sigma}M^{\rho\sigma}\right)$ , and  $S[\Lambda] = \exp\left(\frac{1}{2}\Omega_{\rho\sigma}S^{\rho\sigma}\right)$ .

The generators of the Lorentz Algebra satisfy,

$$[M^{\rho\sigma}, M^{\tau\nu}] = g^{\sigma\tau} M^{\rho\nu} - g^{\rho\tau} M^{\sigma\nu} + g^{\rho\nu} M^{\sigma\tau} - g^{\sigma\nu} M^{\rho\tau}. \tag{1}$$

Consider  $S^{\rho\sigma} = \frac{1}{4} [\gamma^{\rho}, \gamma^{\sigma}].$ 

(a) Compute the commutation relation,  $[S^{\mu\nu}, \gamma^{\rho}]$ .

- (b) Use the above relations to show that the  $S^{\mu\nu}$  matrices associated with the spinor representation, satisfy the
- Lorentz agents (c) Write down the form of the rotation matrix under spinor representation. Writing the rotation parameters as
- Write down the first of  $\Omega_{12} = -\varphi^3$ ) work out the value of the spinor representation  $S[\Lambda]$  under a rotation  $\Omega_{\infty} = -\epsilon_{\infty} \lambda \varphi^{\Lambda}$  (meaning  $\Omega_{12} = -\varphi^3$ ) by 2x about the z axis.
- 6. The Lagrangian of the spinors along with the interaction of the vector field is:

$$\mathcal{L} = ar{\psi} (i \gamma^{\mu} \partial_{\mu} - m) \psi + e ar{\psi} \gamma^{\mu} \psi A_{\mu}.$$

Consider the process  $e^-(p_1)$   $e^+(p_2) \to e^-(p_3)$   $e^+(p_4)$ , also known as Bhabha Scattering. For simplification, you Consider the process  $(m_e) = (m_e)$  and  $(m_e) = (m_e)$  where the mass of the electron can be neglected  $(m_e = 0)$ .

- (a) Draw the leading order Feynman diagrams for the desired process, with appropriate labeling.
- (b) Use the Feynman rules to write down the total amplitude M. (Be sure, you have the correct relative sign between these diagrams.)
- (c) Square the amplitude, average over the initial particle's spin, and sum over the final particle spins. Use the completeness relations,  $\sum_{s=1}^{2} u^{s}(\vec{p}) \bar{u}^{s}(\vec{p}) = \not p + m_{e}$ ,  $\sum_{s=1}^{2} v^{s}(\vec{p}) \bar{v}^{s}(\vec{p}) = \not p m_{e}$ , and evaluate the traces using
- (d) Introduce the Mandelstam variables, s, t, and u, in terms of  $p_1, p_2, p_3, p_4$ . Show that the final result is

$$rac{1}{4}|\mathcal{M}|^2 = 2e^4\left[rac{u^2+t^2}{s^2} + rac{2u^2}{st} + rac{u^2+s^2}{t^2}
ight]$$

. Express all 4-momentum vectors in the center of mass frame of reference, in terms of a suitably chosen set of variables such as E and  $\theta$ . Considering  $m_e = 0$ , what is the value for s + t + u?

(e) Use crossing symmetry to compute  $\frac{1}{4}|\mathcal{M}|_{e^-e^-\to e^-e^-}^2$  for the process  $e^-(p)$   $e^-(k)\to e^-(p')$   $e^-(k')$  and show

$$\frac{1}{4}|\mathcal{M}|_{e^-e^-\to e^-e^-}^2 = 2e^4\left[\frac{s^2+t^2}{u^2} + \frac{2s^2}{ut} + \frac{u^2+s^2}{t^2}\right]$$

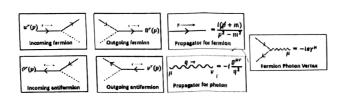
[2+3+5+3+2]

## Some formulas you may need

• The Lagrangian of the spinors along with the interaction of the vector field is:

$$\mathcal{L} = ar{\psi}(i\gamma^{\mu}\partial_{\mu} - m)\psi + ear{\psi}\gamma^{\mu}\psi A_{\mu}.$$

The relevant Feynman rules are listed below:



• In the Weyl representation the  $\gamma$  matrices are

$$\gamma^0 = \begin{pmatrix} 0 & \mathbf{1} \\ \mathbf{1} & 0 \end{pmatrix}$$
 ,  $\gamma^i = \begin{pmatrix} 0 & \sigma^i \\ -\sigma^i & 0 \end{pmatrix}$ 

where, 1 is a  $2 \times 2$  identity matrix and  $\sigma^i$  are  $2 \times 2$  Pauli matrices.

• The Pauli Matrices satisfy  $\{\sigma_i, \sigma_j\} = 2\delta_{ij}$ ,  $\sigma^i \sigma^j = i \epsilon^{ijk} \sigma_k$  and are denoted by:

$$\sigma^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \sigma^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \sigma^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix},$$

Trace formulas

Tr[I] = 4, Tr[any odd number of  $\gamma'$ s] = 0,  $\text{Tr}[\gamma^{\mu}\gamma^{\nu}] = 4g^{\mu\nu}$   $\text{Tr}[\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}] = 4(g^{\mu\nu}g^{\rho\sigma} - g^{\mu\rho}g^{\nu\sigma} + g^{\mu\sigma}g^{\nu\rho}),$   $\text{Tr}[p_{2}^{\nu}\gamma^{\mu}p_{1}^{\nu}\gamma_{\nu}p_{3}^{\nu}\gamma_{\mu}p_{4}^{\nu}\gamma^{\nu}] = -32(p_{1} \cdot p_{4})(p_{2} \cdot p_{3}),$   $\text{Tr}[p_{2}^{\nu}\gamma^{\mu}p_{4}^{\nu}\gamma_{\nu}p_{3}^{\nu}\gamma_{\mu}p_{1}^{\nu}\gamma^{\nu}] = -32(p_{1} \cdot p_{4})(p_{2} \cdot p_{3}),$ 

$$M = \sigma$$
  $V = M$   $\rho = \rho$   $\sigma = v$