

5 Conclusion & Future work

Conclusion

In conclusion, we've simulated the RSJJ Hamiltonian using the diagrammatic approximation method to determine its criticality. To consider the temperature dependency of the RSJJ system, we developed our idea in Matsubara formalism in steady-state conditions and calculated the partition function for consecutive temperatures. Chapter 2 deals with the representation methodology of field theory and how it expanded into a many-body problem with topological structure. A detailed description of the simulation model is in sec 2.1. We showed how the RSJJ Hamiltonian model is described in a macroscopic state of view. Operator \hat{N} was used for conveying the bosonic interaction effect to the target system, and a detailed matrix form was derived in subsections. Sec 2.2 to Sec 2.3 introduced the basic method for solving many-body Hamiltonian. Specifically, we calculate the partition function for the thermal equilibrium state propagating in imaginary time. We describe the mapping procedure of our RSJJ model in the present pseudo-particle solver method in Sec 2.3. In Chapter 3, the program code for implementing the simulation in a computational environment was analyzed. We evaluate the accuracy of our simulation approach and validate the approximate result at the beginning of Chapter 4, which is that calculated values converged toward the result from Exact diagonalization and higher order expansion method. The detailed procedure is explained in Sec 4.2 and Sec 4.3. As a result of the simulation, criticality existed in the order parameter calculated in discrete temperature intervals, which can be deduced that the system feature goes on an insulating phase when temperature decreases, Whence in the previous study phase, a transition occurred in the fixed interaction parameter. We assumed this different phenomenon was revealed due to the finite temperature condition of our study. We also evaluated the approximation result, which was calculated from the correlation function, and found that it indicates different trends through temperature. Yet it was hard to decide the exact critical point of our simulation progress, so more research should be done to get concrete results for examination.

Future work

First, as the most straightforward task, we can modify the range conditions for the β and u values and proceed with the calculations. By performing calculations for higher u values, we can get more accurate data for the expected changes in criticality predicted by the order parameter. The second task is to introduce higher-order approximation calculation methods. We have a calculation code capable of performing TOA, and using this code to perform calculations within the same parameter range is expected to yield more accurate results. However, it should be considered that this case requires significant time to proceed with the calculations. Third, considering to extend the current framework to the Keldysh contour. The formalism employed in this study, the Matsubara formalism, transforms the time-dependent energy changes into temperature-dependent changes for the calculation. This formalism is vulnerable to considering the effects of non-equilibrium phenomena such as dissipation. We anticipate that more accurate calculations will be feasible under real-time analysis, incorporating non-equilibrium phenomena.

A Circuit Hamiltonian

From the circuit in Figure 3, we can establish the circuit equation for the Josephson junction.

A.1 Josephson junction

A Josephson junction consists of two conductors separated by a thin insulator (S-I-S structure), where the insulator acts as a potential barrier. Two factors determine the energy of a Josephson junction: the energy derived from Cooper pairs tunneling through the potential barrier, analogous to the inductor energy in an LC circuit, and the effective capacitance due to the S-I-S structure of the junction. The effective inductance is determined by the critical current I_c and flux $\Phi = \int_0^t V(t)dt$, the potential difference V across the junction, and the energy E_J of the Cooper pairs that have crossed the potential barrier of the junction. This can be expressed as the following equation:

$$\begin{cases} I = I_c \sin \Phi \\ \frac{\partial \Phi}{\partial t} = \frac{2e}{\hbar} V \end{cases} \quad (\text{A.1})$$

Combining the two equations, we obtain the following relation for Cooper pairs:

$$\begin{aligned} \frac{\partial I}{\partial t} &= \frac{\partial I}{\partial \Phi} \frac{\partial \Phi}{\partial t} = I_c V \frac{2e}{\hbar} \cos \Phi \\ \mathcal{L}_{J_J} &= E_J \cos \frac{2e}{\hbar} \Phi \end{aligned} \quad (\text{A.2})$$

And the energy of the junction due to the effective capacitance is written as follows:

$$\mathcal{L}_{J_c} = \frac{C_J}{2} \dot{\Phi}^2 \quad (\text{A.3})$$

From this, the Lagrangian of the entire Josephson junction can be written as follows:

$$\begin{aligned} \mathcal{L}_J &= \mathcal{L}_{J_c} + \mathcal{L}_{J_J} \\ &= \frac{C_J}{2} \dot{\Phi}^2 + E_J \cos \frac{2e}{\hbar} \Phi \end{aligned} \quad (\text{A.4})$$

Performing Legendre transformation same as the case of LC oscillator, the Hamiltonian of the Josephson junction can be rewritten as follows:

$$H_{JJ} = \frac{Q_J^2}{2C_J} - E_J \cos \frac{2e}{\hbar} \Phi \quad (\text{A.5})$$

This is equivalent to the Hamiltonian of a simple pendulum suspended in a gravitational field. $\cos \Phi$ term makes the Josephson junction behave like a nonlinear harmonic oscillator, where the Josephson energy E_J controls the degree of nonlinearity.

A.2 Full Hamiltonian

By using the potential energy (voltage) applied to the circuit, the Lagrangian of the entire circuit can be expressed as follows.[12]

$$\mathcal{L} = \frac{C_J}{2} \dot{\Phi}^2 + E_J \cos \frac{2e}{\hbar} + \frac{C_C}{2} (\dot{\Phi}_1 - \dot{\Phi})^2 + \sum_{i=2}^N \frac{cx}{2} \dot{\Phi}_i^2 - \sum_{i=2}^N \frac{(\Phi_i - \Phi_{i-1})^2}{2lx} \quad (\text{A.6})$$

To transform the Lagrangian into the Hamiltonian form, we define the variable $Q_i = \frac{\partial \mathcal{L}}{\partial \dot{\Phi}_i}$ for the LC circuit, along with the previously defined canonical conjugate variables. This can be interpreted as the charge accumulated on the capacitor connected to the i -th node. Considering the variable with respect to which the differentiation is performed, the total Lagrangian can be divided into the following two parts.

$$\mathcal{L} = \mathcal{L}_{loc} + \mathcal{L}_{LC} \quad (\text{A.7})$$

By applying appropriate Legendre transformations to each Lagrangian, the following form of the Hamiltonian is obtained.

$$\hat{H} = \frac{(Q + Q_1)^2}{C_J} + \frac{Q_1^2}{2C_C} - E_J \cos \left(\frac{2e}{\hbar} \Phi \right) + \sum_{i=2} \frac{(\hat{\Phi}_i - \hat{\Phi}_{i-1})^2}{2lx} + \sum_{i=2} \frac{\hat{Q}_i^2}{2cx} \quad (\text{A.8})$$

Transmission line composed by LC circuit

The Lagrangian for a waveguide structure using an LC circuit consisting of inductors and capacitors corresponds to the following part.

$$\mathcal{L}_{LC} = \sum_{i=2} \frac{cx}{2} \dot{\Phi}^2 - \sum_{i=2} \frac{(\Phi_i - \Phi_{i-1})^2}{2lx} \quad (\text{A.9})$$

This originates from the form of energy stored in the capacitor due to the potential difference across it and the energy stored in the inductance per unit length of the inductance. By rewriting the Lagrangian into the Hamiltonian using the canonical variables for the LC circuit, the following equation can be obtained.

$$H_{LC} = \sum_{i=2} \frac{(\hat{\Phi}_i - \hat{\Phi}_{i-1})^2}{2lx} + \sum_{i=2} \frac{\hat{Q}_i^2}{2cx} \quad (\text{A.10})$$

Local system

The Lagrangian for the capacitor connecting the Josephson junction and the LC oscillator can be written as follows.

$$\mathcal{L}_{loc} = \frac{C_J}{2} \dot{\Phi}^2 + E_J \cos \frac{2e}{\hbar} + \frac{C_C}{2} (\dot{\Phi}_1 - \dot{\Phi})^2 \quad (\text{A.11})$$

The Legendre transformation for a multivariable function can be written as follows.

$$\begin{aligned} Q &= \frac{\partial \mathcal{L}_{CC}}{\partial \dot{\Phi}} = C_J \dot{\Phi} + C_C (\dot{\Phi} - \dot{\Phi}_1) \quad , \quad \dot{\Phi} = \frac{Q}{C_J} - \frac{C_C}{C_J} (\dot{\Phi} - \dot{\Phi}_1) \\ Q_1 &= \frac{\partial \mathcal{L}_{CC}}{\partial \dot{\Phi}_1} = \frac{C_C}{2} (2\dot{\Phi}_1 - 2\dot{\Phi}) \end{aligned} \quad (\text{A.12})$$

Using this, rewriting the Lagrangian into the Hamiltonian yields the following:

$$H_{loc} = \frac{(Q + Q_1)^2}{C_J} + \frac{Q_1^2}{2C_C} - E_J \cos\left(\frac{2e}{\hbar} \Phi\right) \quad (\text{A.13})$$

Quantization example - 1st node consideration

For the case where it is the first node (i=1), when considering the variables, Hamiltonian can be written as follows.

$$\begin{aligned} \hat{H}_1 &= \frac{(Q + Q_1)^2}{C_J} + \frac{Q_1^2}{2C_C} - E_J \cos\left(\frac{2e}{\hbar} \Phi\right) + \frac{\hat{\Phi}}{2lx} \\ &= \frac{1}{2} \begin{pmatrix} \hat{P} & \hat{P}_1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & \sqrt{\frac{C}{C_p}} \end{pmatrix}^{-1} \begin{pmatrix} \frac{1}{C_J} & \frac{1}{C_J} \\ \frac{1}{C_J} & \frac{1}{C_P} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & \sqrt{\frac{C}{C_p}} \end{pmatrix} \begin{pmatrix} \hat{P} \\ \hat{P}_1 \end{pmatrix} + \\ &\quad \frac{1}{2lx} \begin{pmatrix} \hat{\Psi} & \hat{\Psi}_1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & \sqrt{\frac{C}{C_p}} \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & \sqrt{\frac{C}{C_p}} \end{pmatrix} \begin{pmatrix} \hat{\Psi} \\ \hat{\Psi}_1 \end{pmatrix} - E_J \cos\left(\frac{2e}{\hbar} \Phi\right) \end{aligned} \quad (\text{A.14})$$

Here, C is a constant. If we expand and write the above Hamiltonian, it can be written as below.

$$\hat{H}_1 = \frac{1}{2} \left(\frac{\hat{P}^2}{C_J} + \frac{2}{C_J} \sqrt{\frac{C_P}{C}} (\hat{P} \hat{P}_1) + \frac{\hat{P}_1^2}{C} \right) + \frac{1}{2lx} \frac{C}{C_P} \hat{\Psi}^2 - E_J \cos \frac{2e}{\hbar} \hat{\Psi} \quad (\text{A.15})$$

The operators $\hat{P}_1, \hat{\Psi}_1$ can be rewritten as bosonic operators using the following relation.

$$\begin{aligned} \hat{P}_1 &= -\sqrt{\frac{\hbar C \omega}{2}} (\hat{b}_1^\dagger + \hat{b}_1) \\ \hat{\Psi}_1 &= i\sqrt{\frac{\hbar}{2C\omega}} (\hat{b}^\dagger - \hat{b}) \end{aligned} \quad (\text{A.16})$$

Here, $\omega = \frac{1}{lx C_p}$. Using the above equation to rewrite the Hamiltonian, it can be rewritten into the following form:

$$\hat{H}_1 \approx \frac{\hat{P}^2}{2C_J} - E_J \cos\left(\frac{2e}{\hbar} \hat{\Psi}\right) - g \hat{P} (\hat{b}_1^\dagger + \hat{b}_1) + \hbar \omega \hat{b}_1^\dagger \hat{b}_1 \quad (\text{A.17})$$

Here, $g = \frac{1}{C_J} \sqrt{\frac{\hbar}{2lx}}$. This can be inferred the capacitor in the structure of the LC oscillator is not considered in the example circuit.

B Brief statement for field theory

In the main text, we calculated the partition function using the action S . Here, we provide a brief explanation of this approach. To begin, we can define the partition function of Minkowski space that restricted by the certain energy condition. The dynamic of variables of given space are governed under the action functional, thus partition function of space dictates the probability distribution of possible paths make system propagates onto. The partition function $Z[J]$ can be calculated in path integral method. For field $\phi(x, t)$, corresponding Lagrangian is :

$$\begin{aligned}\mathcal{L} &= \frac{1}{2}\phi^T A \phi + \phi \cdot J \\ &= \frac{1}{2}(-\phi^T \partial_t^2 \phi + \phi^T \nabla^2 \phi - m^2 \phi) + J \cdot \phi\end{aligned}\tag{B.1}$$

Here, A is a form of diagonal matrix and J represents to the potential of the system we want to describe. Consider only time in 4-vector notation, and use the above notation, we can write partition function in following:

$$Z[J] = \int [D\phi] e^{\int d\phi [\frac{1}{2}\partial_\tau^2 \phi - m^2 \phi^2] + J \cdot \phi}\tag{B.2}$$

Which is, in short way:

$$Z[J] = Z[0]W[J]\tag{B.3}$$

A similar form to the above can be found in equations (2.40), (2.41) which are presented in the main text. Here $W[J] = e^{\frac{1}{2}J_i A_{ij}^{-1} J_j}$, $Z[0] = \left(\frac{(2\pi i)^N}{\det A}\right)^{\frac{1}{2}}$. The sub indices i, j means i th and j th element of $J(\phi)$. Note the indices refer to Minkowski space, thus encompassing not only spatial coordinates but also temporal indices. we can recognize that :

$$\langle \phi^{2n} \rangle = \frac{\int [D\phi] \phi^{2n} e^{\frac{1}{2}\phi^i A_{ij} \phi^j + J \cdot \phi}}{Z[0]} = \frac{\Pi_i \frac{\delta}{\delta J_i} \frac{\delta}{\delta J_j} \int [D\phi] e^{-\frac{1}{2}(J_{ij} A_{ij} J_j)}}{Z[0]} = \Pi_i \frac{\delta}{\delta J_i} \frac{\delta}{\delta J_j} \ln Z[J]\tag{B.4}$$

This implies that if we perform the above calculation for a small time variation, we can re-express the equation as an expansion in time. Upon rotating time by i in the complex plane, this result matches the form of the thermal partition function calculated in the main text. Similar, if there are some interaction happened between the field and the environment we can treat its effect as a weak perturbation. Let $-\frac{g}{4}\phi^4$ represents the interaction effect. Then the form of partition function turns out to be:

$$Z[J] = \int [D\phi] e^{\int d\phi [\frac{1}{2}\partial_\tau^2 \phi - m^2 \phi^2] + J \cdot \phi - \frac{g}{4}\phi^4}\tag{B.5}$$

Which can be calculated as the same procedure above, in a more complex structure.

C Wick's theorem

As mentioned in Sec.2.2, Wick's theorem states that when there are an even number of pair creation and annihilation operators for a quantum state, the product of the expectation values of each pair is equal to the product of the total expectation value. A well-known example of pair interaction is the Coulomb interaction which can be written in the form:

$$H(\tau) = \frac{1}{2} \sum_{l, l', \kappa} v(\kappa) (\hat{a}_l^\dagger(\tau) \hat{a}_{l'}^\dagger(\tau) \hat{a}_{l'+\kappa}(\tau) \hat{a}_{l-\kappa}(\tau)) \quad (\text{C.1})$$

Here indices l, l' indicate the lattice point of each electron, and the term $v(\kappa)$ is the scalar value (or function) that represents the interaction effect. Written in the form of the interaction operator introduced in the main text, it can be expanded in time as follows:

$$\begin{aligned} \langle U(\beta) \rangle_0 &= 1 - \frac{1}{2} \sum_{l, l', \kappa} v(\kappa) \int_0^\beta d\tau_1 \langle \hat{a}_l^\dagger(\tau_1) \hat{a}_{l'}^\dagger(\tau_1) \hat{a}_{l'+\kappa}(\tau_1) \hat{a}_{l-\kappa}(\tau_1) \rangle_0 \\ &+ \frac{1}{2(2!)} \sum_{\text{lower indices}} v(\kappa_1) v(\kappa_2) \int_0^\beta d\tau_1 \int_0^\beta d\tau_2 \\ &\underbrace{\langle \mathcal{T} \hat{a}_{l_1}^\dagger(\tau_1) \hat{a}_{l'_1}^\dagger(\tau_1) \hat{a}_{l'_1+\kappa_1}(\tau_1) \hat{a}_{l_1-\kappa_1}(\tau_1) \hat{a}_{l'_2}^\dagger(\tau_2) \hat{a}_{l'_2+\kappa_2}(\tau_2) \hat{a}_{l_2-\kappa_2}(\tau_2) \rangle_0}_{\text{The term that we want to split into a multiplication of pairs of operators}} \end{aligned} \quad (\text{C.2})$$

As mentioned in the equation, our interest lies in replacing the product of expectation values of multiple operators with the product of expectation values of operators that belong to pairs. The general form can be written as:

$$\begin{aligned} W_n &= \langle \mathcal{T} \tilde{b}_1 \tilde{b}_2 \tilde{b}_3 \tilde{b}_4 \cdots \tilde{b}_{n-1} \tilde{b}_n \rangle_0 = \langle \mathcal{T} \tilde{b}_1 \tilde{b}_2 \rangle_0 \langle \mathcal{T} \tilde{b}_3 \tilde{b}_4 \rangle_0 \cdots \langle \mathcal{T} \tilde{b}_{n-1} \tilde{b}_n \rangle_0 \\ &+ \cdots + \langle \mathcal{T} \tilde{b}_n \tilde{b}_{n-1} \rangle_0 \langle \mathcal{T} \tilde{b}_{n-2} \tilde{b}_{n-3} \rangle_0 \cdots \langle \mathcal{T} \tilde{b}_2 \tilde{b}_1 \rangle_0 \\ &= \sum_{P_i} P_i \langle \mathcal{T} \tilde{b}_{i_1} \tilde{b}_{i_2} \rangle_0 \langle \mathcal{T} \tilde{b}_{i_3} \tilde{b}_{i_4} \rangle_0 \cdots \langle \mathcal{T} \tilde{b}_{i_{n-1}} \tilde{b}_{i_n} \rangle_0 \end{aligned} \quad (\text{C.3})$$

Here, the subscripts represent time variables (e.g. $\tilde{b}_1 = \tilde{b}(\tau_1)$), P_i represents the number of all possible permutations, and 0 indicates that the expectation value has been calculated with respect to the local Hamiltonian of the system under consideration. Since the matching of the number of creation and annihilation operators is important, to represent total operators for an arbitrary order, we have used the notation \tilde{b} instead of $\hat{}$ or \dagger . In the case of fermions, the sign changes depending on the order is crucial during permutation, the operators dealt with in the main text are all bosonic cases, so we can neglect the sign problem.