Problem Set 2

Math228B Numerical solutions to differential equations

Håkon Olav Torvik

UC Berkeley

(Dated: February 18, 2022)

Problem 1

Part a, See previous problem set

Part b

We have the method

$$U_j^{n+1} = U_j^n + \frac{k\kappa}{2h^2} \left[U_{j-1}^n - 2U_j^n + U_{j+1}^n + U_{j-1}^{n+1} - 2U_j^{n+1} + U_{j+1}^{n+1} \right] - k\gamma \left[(1-\theta)U_j^n + \theta U_j^{n+1} \right],$$
(1)

which models diffusion with decay.

Let

$$U_j^n = e^{ijh\xi}, \quad U_j^{n+1} = g(\xi)U_j^n.$$
 (2)

Von Neumann analysis says that the method is stable provided $|g| \leq 1$.

Inserting (2) into (1), we get the expression for $g(\xi)$ as

$$ge^{ijh\xi} = e^{ijh\xi} + \frac{k\kappa}{2h^2} \left[e^{i(j-1)h\xi} - 2e^{ijh\xi} + e^{i(j+1)h\xi} + ge^{i(j-1)h\xi} - 2ge^{ijh\xi} + ge^{i(j+1)h\xi} \right]$$

$$- k\gamma \left[(1-\theta)e^{ijh\xi} + \theta ge^{ijh\xi} \right],$$

$$g = 1 + \frac{k\kappa}{2h^2} \left[e^{-ih\xi} - 2 + e^{ih\xi} + ge^{-ih\xi} - 2g + ge^{ih\xi} \right] - k\gamma \left[(1-\theta) + \theta g \right],$$

$$g = 1 + \frac{k\kappa}{h^2} \left[\cos(h\xi) - 1 + g\cos(h\xi) - g \right] + k\gamma(\theta - 1) - k\gamma\theta g,$$

$$g - \lambda g + k\gamma\theta g = 1 + \lambda + k\gamma(\theta - 1), \qquad \text{where} \quad \lambda \equiv \frac{k\kappa}{h^2} \left(\cos(h\xi) - 1 \right),$$

$$g = \frac{1 + \lambda + k\gamma(\theta - 1)}{1 - \lambda + k\gamma\theta}.$$

Using the stability condition $|g| \leq 1$, we get

$$1 \ge \left| \frac{1 + \lambda + k\gamma(\theta - 1)}{1 - \lambda + k\gamma\theta} \right|, \tag{3}$$
$$|1 - \lambda + k\gamma\theta| \ge |1 + \lambda + k\gamma(\theta - 1)|.$$

In the second step I used the fact that λ is always negative or equal to 0, such that $1 - \lambda + k\gamma\theta > 1$. WolframAlpha gives that under the conditions $\lambda < 0$ and $k\gamma > 0$, both of which are obviously always true, then

$$\theta \ge \frac{k\gamma - 2}{2k\gamma} = \frac{1}{2} - \frac{1}{k\gamma}.$$

Again, as k and γ are both non-negative, the second term can be dropped, and the method is unconditionally stable for any

$$\theta \geq \frac{1}{2}$$
,

which is what was to be shown. (The conditions set are already true from the formulation of the problem).

 \mathbf{c}

Going back to (3), I now set $\theta = 0$, and do a similar analysis.

$$1 \ge \left| \frac{1 + \lambda - k\gamma}{1 - \lambda} \right|,$$
$$|1 - \lambda| \ge |1 + \lambda - k\gamma|.$$

Again, Wolfram Alpha, under the conditions that the solutions are real and $\lambda < 1$, both of which are again always true, gives the solution as $2\lambda \le k\gamma \le 2$. As $\lambda < 0$, the first inequality is always upheld. This gives unconditional stability for $k \le 2/\gamma$, which is what was to be shown.

Problem 2

To solve the differential equation

$$u_t + au_x = 0,$$

one could use the following scheme

$$U_j^{n+1} = U_{j-2}^{n-1} - \left(\frac{ak}{h} - 1\right) (U_j^n - U_{j-2}^n).$$

Part a

To determine the order of accuracy of the method, I Taylor expand, and subtract the 2 sides from eachother. I Taylor expand to 3rd order.

$$U_{j}^{n} = u(x,t) \equiv u$$

$$u(x,t+k) = u(x-2h,t-k) - \left(\frac{ak}{h}-1\right) (u(x,t)-u(x-2h,t))$$

$$u(x,t+k) = u + ku_{t} + \frac{k^{2}}{2}u_{tt} + \frac{k^{3}}{6}u_{ttt} + \mathcal{O}(k^{4})$$

$$u(x-2h,t) = u - 2hu_{x} + \frac{4h^{2}}{2}u_{xx} - \frac{8h^{3}}{6}u_{xxx} + \mathcal{O}(h^{4})$$

$$u(x-2h,t-k) = u - 2hu_{x} - ku_{t} + \frac{1}{2}\left[4h^{2}u_{xx} + 2 \cdot 2hku_{xt} + k^{2}u_{tt}\right] + \frac{1}{6}\left[-8h^{3}u_{xxx} - 3 \cdot 4h^{2}ku_{xxt} - 3 \cdot 2hk^{2}u_{xtt} - k^{3}u_{ttt}\right] + \mathcal{O}(h^{4} + k^{4})$$

$$\tau = u - 2hu_{x} - ku_{t} + \frac{1}{2}\left[4h^{2}u_{xx} + 2 \cdot 2hku_{xt} + k^{2}u_{tt}\right] + \frac{1}{6}\left[-8h^{3}u_{xxx} - 12h^{2}ku_{xxt} - 6hk^{2}u_{xtt} - k^{3}u_{ttt}\right] - \left(\frac{ak}{h} - 1\right)u + \left(\frac{ak}{h} - 1\right)\left[u - 2hu_{x} + \frac{4h^{2}}{2}u_{xx} - \frac{8h^{3}}{6}u_{xxx}\right] - \left(u + ku_{t} + \frac{k^{2}}{2}u_{tt} + \frac{k^{3}}{6}u_{ttt}\right) + \mathcal{O}(kh^{3} + k^{4})$$

$$= \frac{k}{3}\left(\frac{2h^{2}}{a} + ak^{2} - 3hk\right)u_{xxx} + \mathcal{O}(kh^{3} + k^{4})$$

$$= \mathcal{O}(h^{2}k + k^{3} + hk^{2})$$

where I have used the original equation $u_t = -au_x$ to cancel all frist-derivatives and secondderivatives of u, as well as many of the thrid-derivatives using the following relations.

$$u_{t} = -au_{x}$$

$$u_{tt} = -au_{xt} \quad u_{tx} = -au_{xx}$$

$$u_{ttt} = -au_{xtt} \quad u_{ttx} = -au_{xxt} \quad u_{txx} = -au_{xxx}$$

This seems to be a 3rd order method.

Part b

The true solution to the advection equation for a point is the value of the point upstream at the previous timestep. The skewed leapfrog method with $\nu \equiv ak/h = 1$ will depend on the point two steps back in time, and two upstream, so this is also okay. But for values near 1 it will depend on grid-point outside the true domain of dependence and thus not satisfy the CLF-condition.

Part c

Again, I use the definitions (2).

$$\begin{split} u_j^{n+1} = & u_{j-2}^{n-1} - (\nu - 1)(u_j^n - u_{j-2}^n), \quad \nu \equiv ak/h \\ & \qquad \qquad \Downarrow \\ ge^{i\xi hj} = & e^{i\xi h(j-2)}/g - (\nu - 1)(e^{i\xi hj} - e^{i\xi h(j-2)}) \\ & \qquad \qquad g = & e^{-2i\xi h}/g - \nu + \nu e^{-2i\xi h} + 1 - e^{-2i\xi h} \\ & \qquad \qquad \gamma^2 = & 1 - \nu \gamma e^{i\xi hj} + \nu g + \gamma e^{i\xi h} - g, \quad \gamma \equiv ge^{i\xi h} \\ & \qquad \qquad g\gamma^2 = & g - \nu \gamma^2 + \nu g^2 + \gamma^2 - g^2 \\ & \qquad \qquad \gamma^2(g + \nu - 1) = & g(1 + \nu g - g) \\ & \qquad \qquad \gamma^2 = & g\frac{(1 + \nu g - g)}{(g + \nu - 1)} \end{split}$$

For $|g| = |\gamma| \ge 1$, I get that $\nu = 0 \land \nu = 1 \land |\nu| \ge 2$ will satisfy the sability condition. This was found by plotting $\gamma^2(\nu)$ for $|g| \ge 1$ and finding the regions where it was always less than 1.

This is similar to what I found in the previous part, but with the addition of courant numbers larger than 2.

Problem 3

Part a-e

Code, implemented in the Julia-code uploaded on bCourses. The functions I implemented for part b, c, d and e have different signatures than what was specified in the problem set.

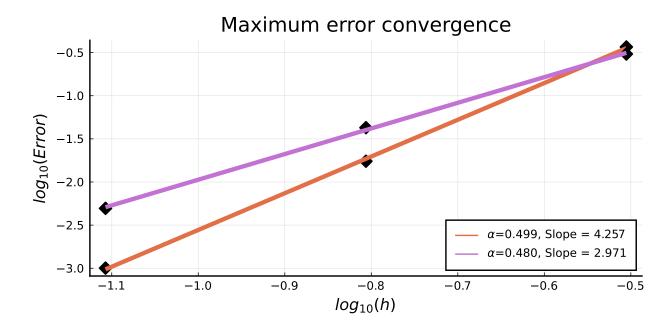


Figure 1: Convergence-plot of the error as the spatial step length gets smaller, for 2 different values of the filtering-parameter α .

I therefore also added some wrapper-functions with the specified signature so autograder would be happy. I have not tested these, but they are very small, so should work alright.

Part f

I implement a solver that sets up the system and iterates to the final solution. I use the max-norm to find the error from the exact solution among all solution components. Figure 1 shows a log-log plot of the error against the step length h for the 2 specified filter-parameters α . For $\alpha = 0.499$ I get that the slope is 4.257, while for $\alpha = 0.480$, it is 2.971.

 $\alpha=0.5$ corresponds to no filtering, so I get better results using less filtering for smaller step sizes. This is as expected, as filtering is not a correct operation when solving such differential equations, but is in this case needed to avoid the solution blowing up. Therefore, as I get a bigger grid, less filtering will be closer to the true solution, while for smaller grids, I need to filter more to maintain stability.

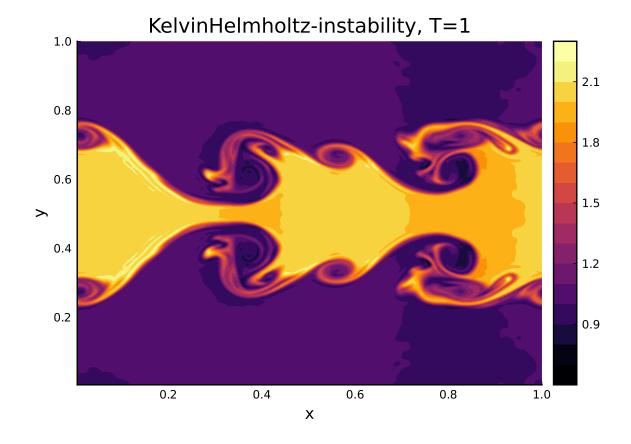


Figure 2: The simulated Kelvin-Helmholtz instability at T=1.

Part g

I implement the specified intital conditions and use my RK4-solver. At T=1, I plot the density of the system in the domain, which is shown in Figure 2. The results looks reasonable.