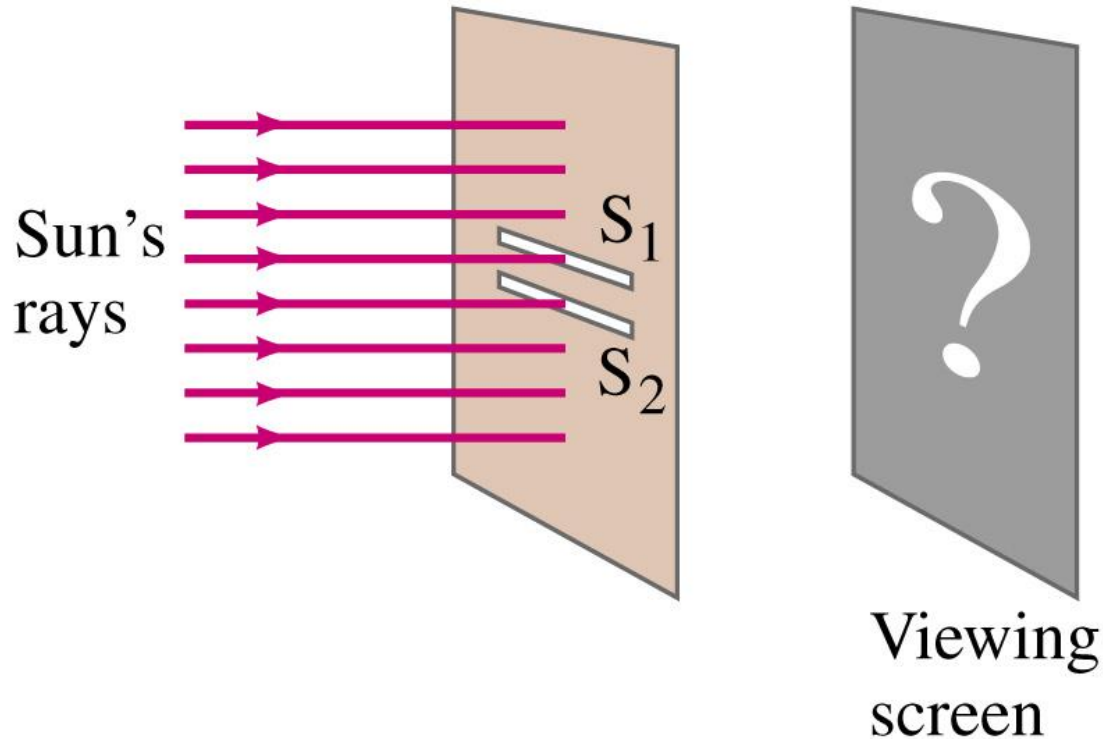


Chapter 30-B Interference



§ 1 Young's Double-Slit Experiment

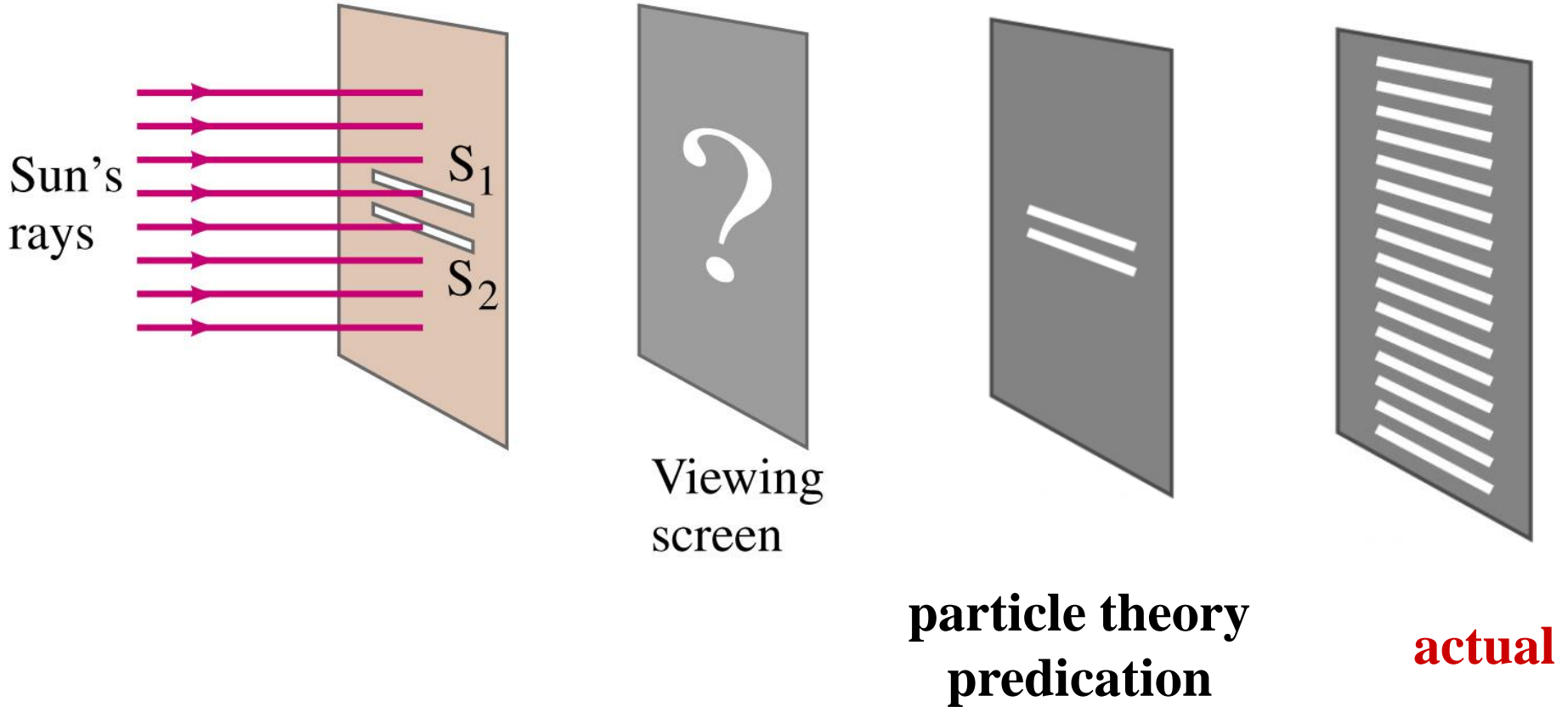


In 1801



**Thomas Young,
English physician
(1773 – 1829)**

Young's Double-Slit Experiment



Review: Two-source interference



- Two identical monochromatic waves from two sources overlap in a region

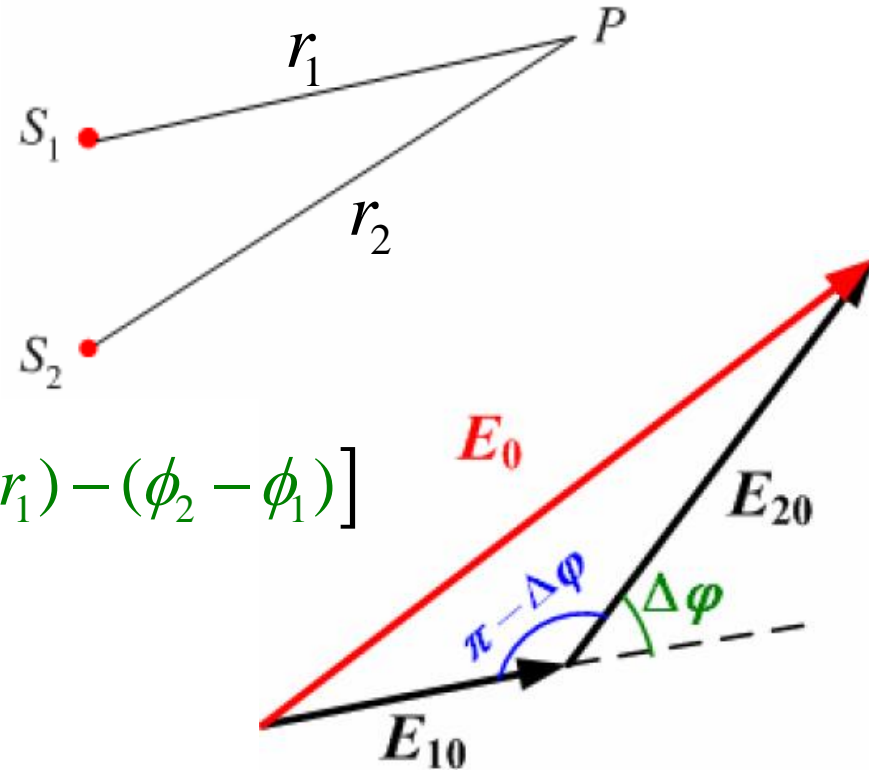
$$E_1 = E_{10} \cos(\omega t - kr_1 + \phi_1)$$

$$E_2 = E_{20} \cos(\omega t - kr_2 + \phi_2)$$

$$E = E_1 + E_2 = E_0 \cos(\omega t + \varphi)$$

$$E_0^2 = E_{10}^2 + E_{20}^2 + 2E_{10}E_{20} \cos[k(r_2 - r_1) - (\phi_2 - \phi_1)]$$

$$I = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos \Delta\varphi$$



Phasor Diagram

Phase difference: $\Delta\varphi = k(r_2 - r_1) - (\phi_2 - \phi_1)$

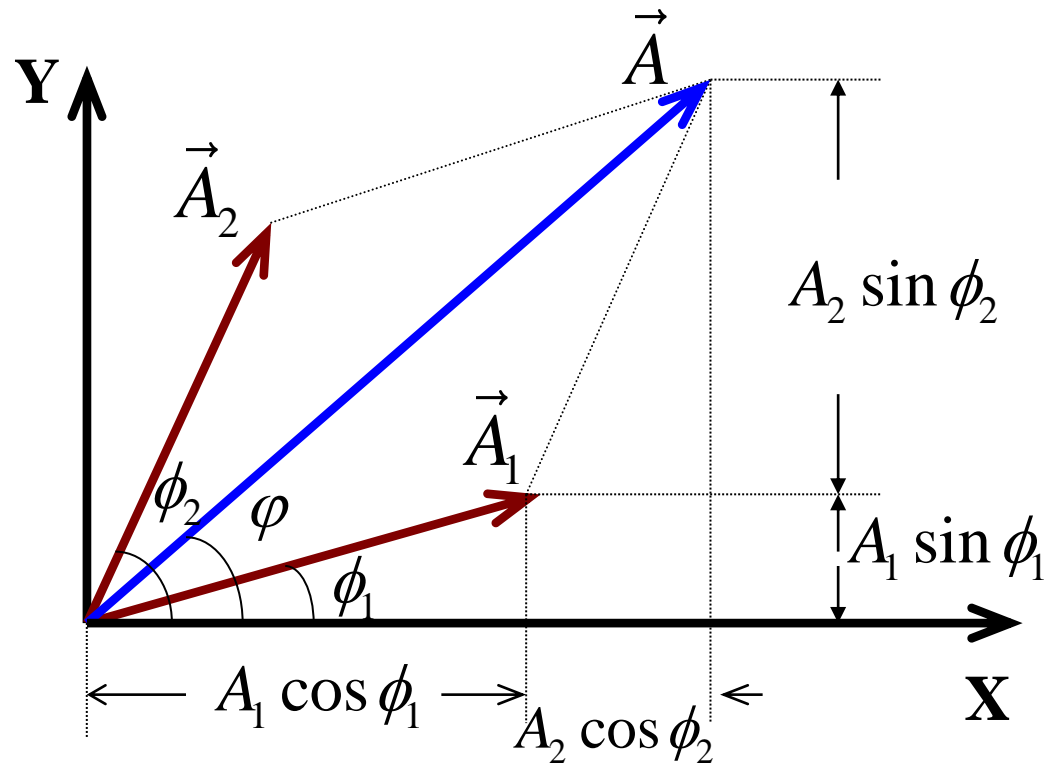
Depending on two factors:

- (1) the location of point P
- (2) the difference of two initial phase angles.

Superposition of SHMs using phasor diagram



Using phasors,



$$A = \sqrt{A_1^2 + A_2^2 + 2A_1A_2 \cos(\phi_2 - \phi_1)}$$

$$\phi = \arctan \frac{A_1 \sin \phi_1 + A_2 \sin \phi_2}{A_1 \cos \phi_1 + A_2 \cos \phi_2}$$

Review: Two-source interference



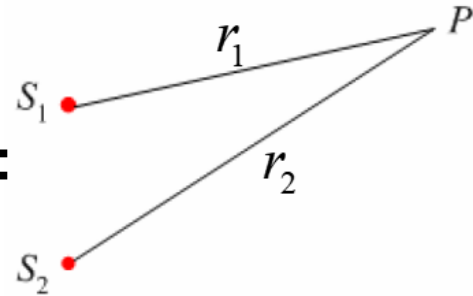
$$I = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos \Delta\varphi, \quad \Delta\varphi = k(r_2 - r_1) - (\phi_2 - \phi_1)$$

➔ **Constructive** and **destructive** interference:

Suppose the two sources are identical $\phi_1 = \phi_2$

For some points, maximum intensity occurs:

$$\Delta\varphi = \frac{2\pi}{\lambda}(r_2 - r_1) = \pm 2m\pi, \quad m = 0, 1, 2, \dots \quad \text{in phase}$$



or path difference: $\delta = r_2 - r_1 = \pm m\lambda, \quad m = 0, 1, 2, \dots$

$$I = I_1 + I_2 + 2\sqrt{I_1 I_2} = \xrightarrow{I_1=I_2} I_{\max} = 4I_1 \quad \text{constructive interference}$$

For some points, minimum intensity occurs :

$$\Delta\varphi = \frac{2\pi}{\lambda}(r_2 - r_1) = \pm (2m+1)\pi, \quad m = 0, 1, 2, \dots \quad \text{out of phase}$$

or path difference: $\delta = r_2 - r_1 = \pm (2m+1)\frac{\lambda}{2}, \quad m = 0, 1, 2, \dots$

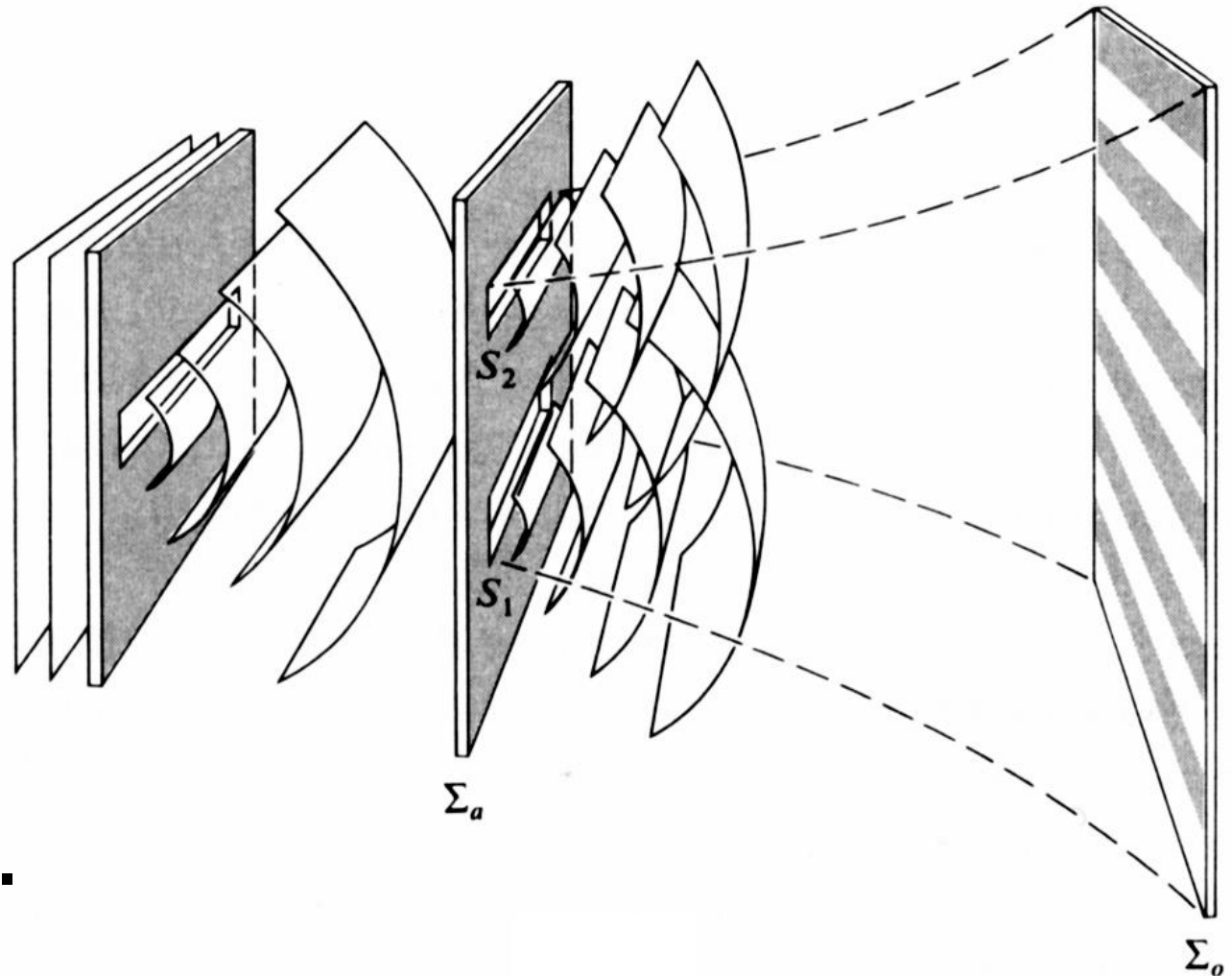
$$I = I_1 + I_2 - 2\sqrt{I_1 I_2} = \xrightarrow{I_1=I_2} I_{\min} = 0 \quad \text{destructive interference}$$

Young's double-slit interference



■ Young's **double-slit** experiment — **wavefront-splitting** interference

- ➡ A light source emits **monochromatic** light. The light is directed at a screen with a narrow **slit** S_0 . The light from slit S_0 falls on a screen with two other narrow **slits** S_1 and S_2 , with distance d apart.



Young's double-slit interference

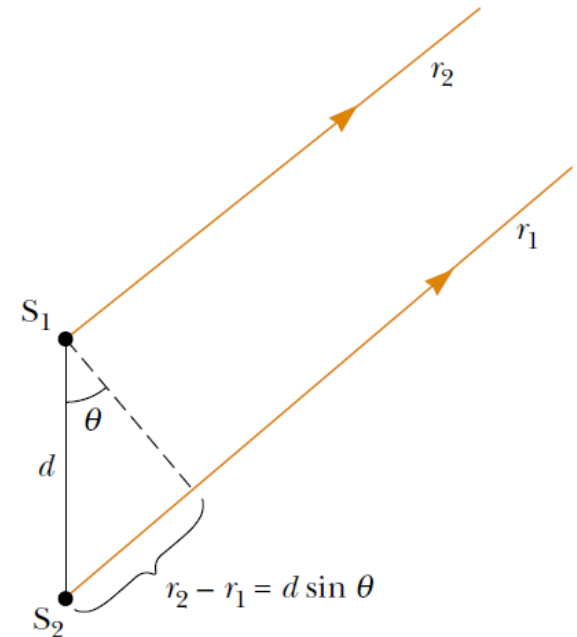
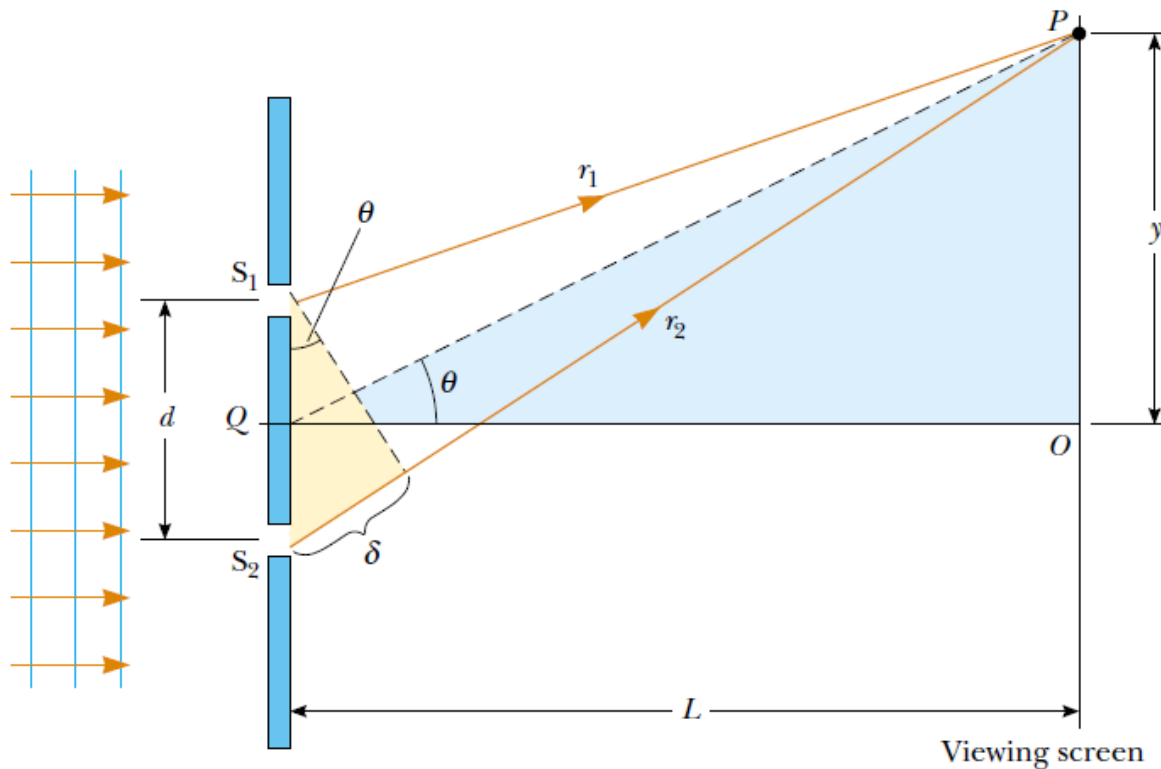


Wave (optical) path length difference:

$$L \gg d, \quad L \gg |y|$$

Generally:
$$\begin{cases} d : 0.1 \sim 1 \text{ mm} \\ L : 1 \sim 10 \text{ m} \\ |y| \leq 1 \sim 10 \text{ cm} \end{cases}$$

$$\delta = r_2 - r_1 \approx d \sin \theta$$

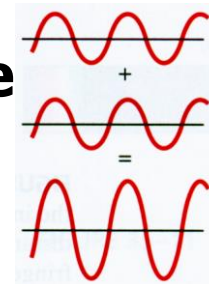


Bright and dark fringes

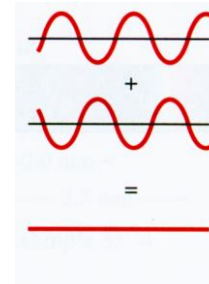


➡ **Bright fringes** — **constructive** interference

$$\delta = d \sin \theta_{\text{bright}} = \pm m \lambda, \quad m = 0, 1, 2, \dots$$



**Construction
interference**



**Destruction
interference**

➡ **Dark fringes** — **destructive** interference:

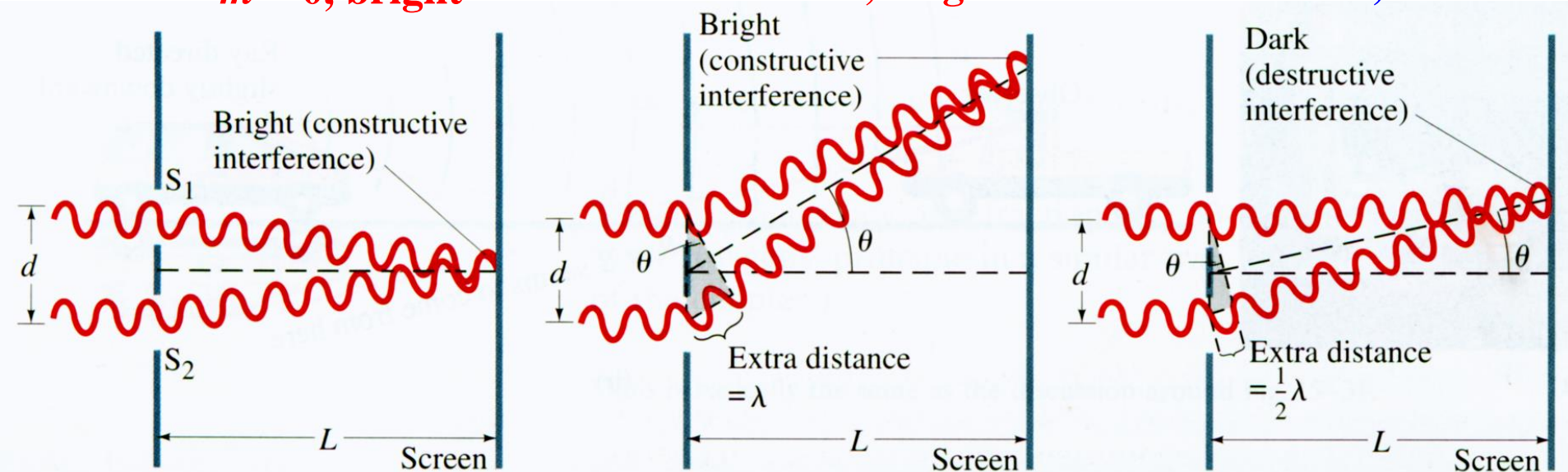
$$\delta = d \sin \theta_{\text{dark}} = \pm \left(m - \frac{1}{2} \right) \lambda, \quad m = 1, 2, \dots$$

m is the order of the fringe

$m = 0$, bright

$m = 1$, bright

$m = 1$, dark



Double-slit interference



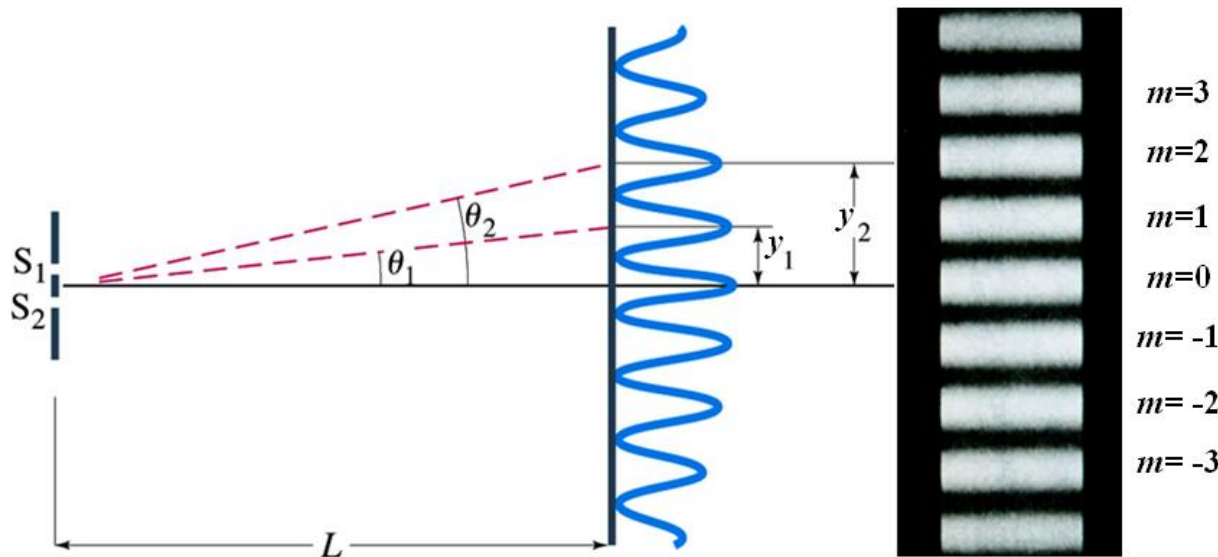
- The positions of bright and dark fringes:

θ is so small that $\sin \theta \approx \tan \theta$, $\delta = d \sin \theta \approx d \tan \theta = d \frac{y}{L}$, $y = \frac{L}{d} \delta$

➤ Bright fringes: $y_{\text{bright}} = \pm m \frac{L}{d} \lambda, \quad m = 0, 1, 2, \dots$

➤ Dark fringes : $y_{\text{dark}} = \pm \left(m - \frac{1}{2} \right) \frac{L}{d} \lambda, \quad m = 1, 2, \dots$

➤ Spacing of the fringes: $\Delta y = y_{m+1} - y_m = (m+1) \frac{L}{d} \lambda - m \frac{L}{d} \lambda = \frac{L}{d} \lambda$



Equal fringe spacing

Intensity in the double-slit interference pattern

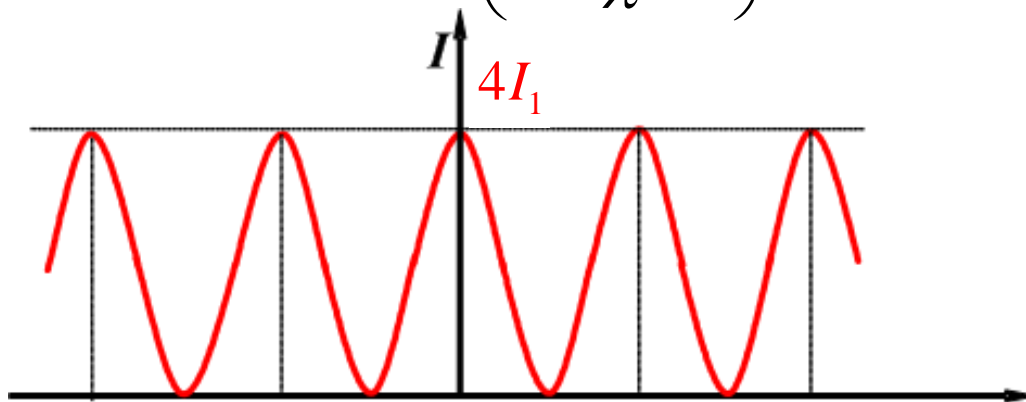


- Intensity distribution: $I = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos \Delta\varphi$, $\Delta\varphi = \frac{2\pi}{\lambda} d \sin \theta$
 ➡ Suppose the width of two slit are the same

$$I_1 = I_2, \quad I = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos \Delta\varphi = 2I_1 (1 + \cos \Delta\varphi) = 4I_1 \cos^2 \left(\frac{\Delta\varphi}{2} \right)$$

$$= 4I_1 \cos^2 \left(\frac{\pi d \sin \theta}{\lambda} \right) \approx 4I_1 \cos^2 \left(\frac{\pi d}{\lambda L} y \right)$$

$$\sin \theta \approx \tan \theta = y / L$$



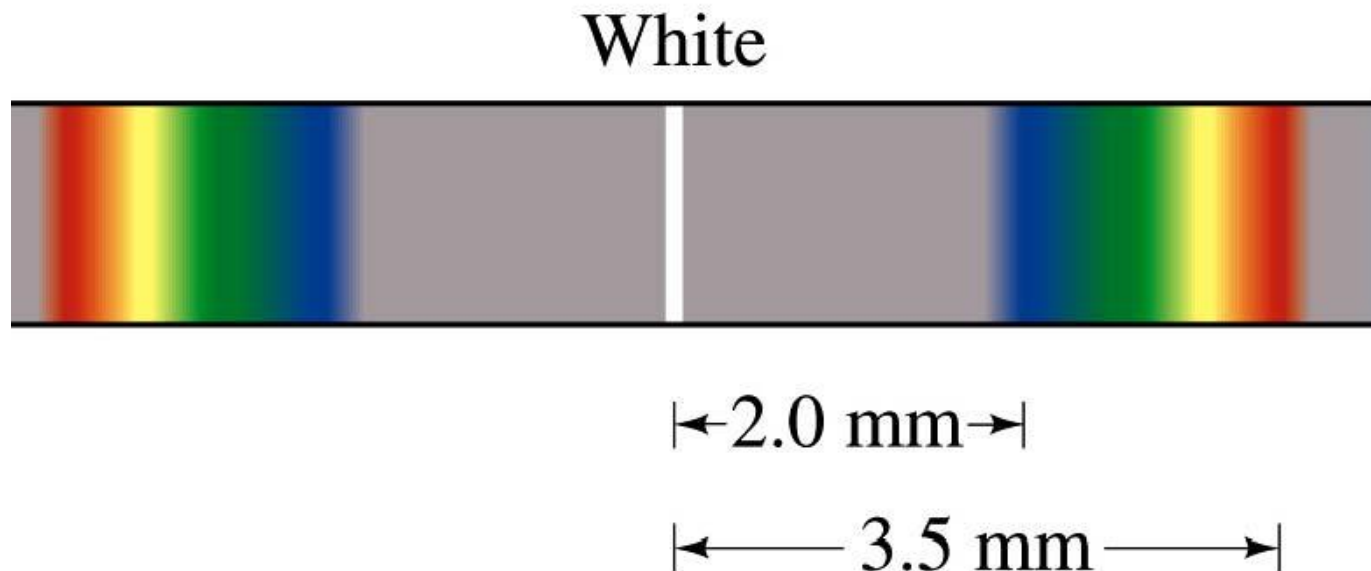
| | | | | | | |
|-----------------|----------------|-----|---------------|----------------|-----------------|------------------|
| -4π | -2π | 0 | 2π | 4π | $\Delta\varphi$ | Phase difference |
| -2 | -1 | 0 | 1 | 2 | m | Order of fringes |
| $-2\lambda/d$ | $-\lambda/d$ | 0 | λ/d | $2\lambda/d$ | $\sin \theta$ | Angular position |
| $-2\lambda L/d$ | $-\lambda L/d$ | 0 | $\lambda L/d$ | $2\lambda L/d$ | y | Position |

Example



Wavelength from double-slit interference

White light passes through two slits 0.50mm apart and an interference pattern is observed on a screen 2.5m away. The first-order fringe resembles a rainbow with violet and red light at either end. The violet light falls about 2.0mm and the red 3.5mm from the center of the central white fringe. Estimate the **wavelengths** of the violet light and the red light.



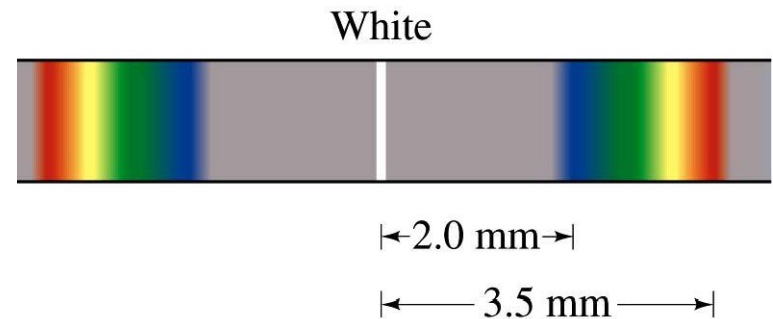
Example



White light passes through two slits **0.50mm** apart and an interference pattern is observed on a screen **2.5m** away. The first-order fringe resembles a rainbow with violet and red light at either end. The violet light falls about **2.0mm** and the red **3.5mm** from the center of the central white fringe. Estimate the **wavelengths** of the violet light and the red light

Solution:

$$y_{\text{bright}} = \pm m \frac{L}{d} \lambda, \quad m = 1$$



For violet light $y = 2.0 \text{ mm}$

$$\lambda_{\text{violet}} = \frac{d}{L} \frac{y}{m} = \frac{(5.0 \times 10^{-4} \text{ m})(2.0 \times 10^{-3} \text{ m})}{2.5 \text{ m}} = 4.0 \times 10^{-7} \text{ m} = 400 \text{ nm}$$

For red light $y = 3.5 \text{ mm}$

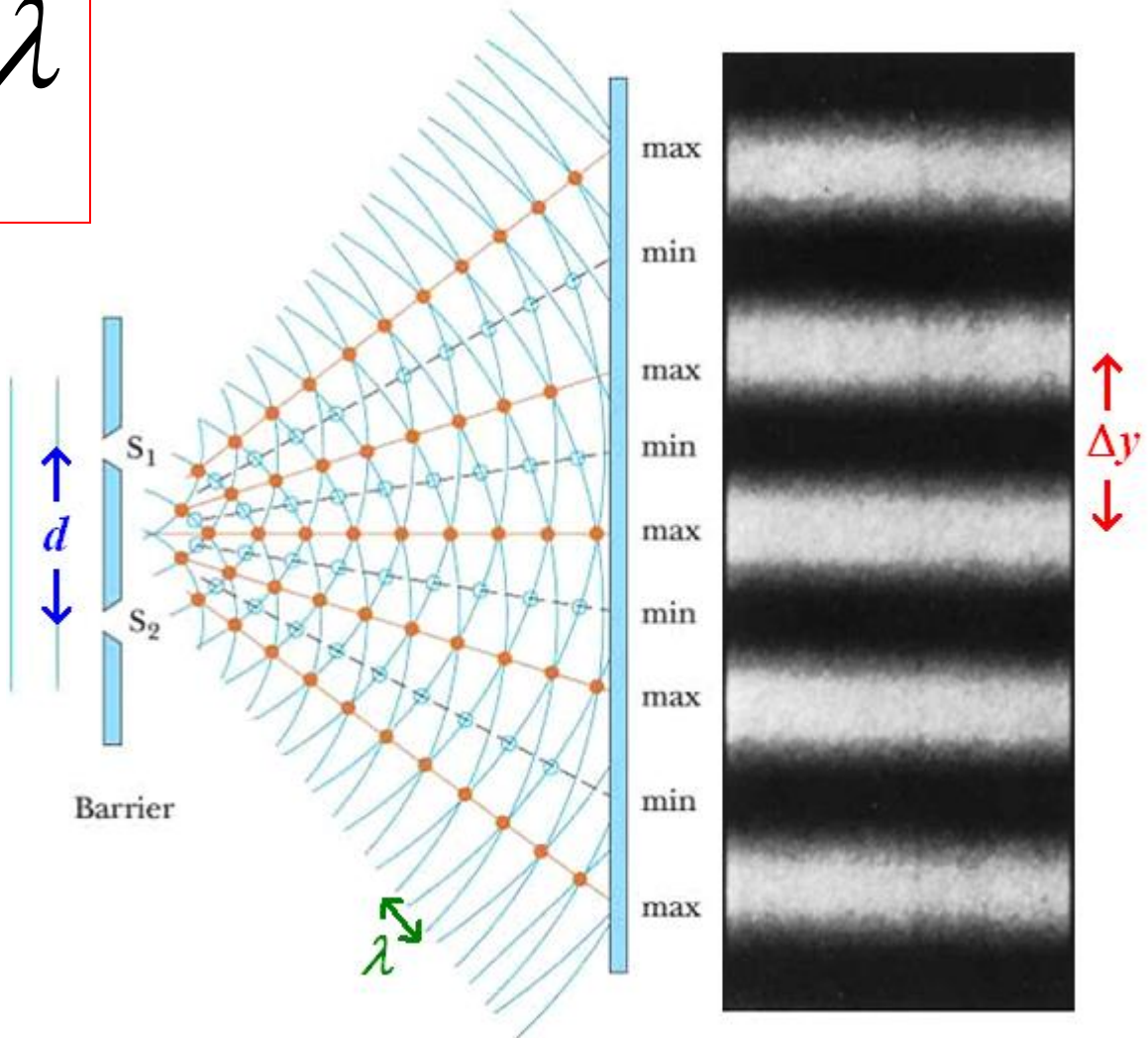
$$\lambda_{\text{red}} = \frac{d}{L} \frac{y}{m} = \frac{(5.0 \times 10^{-4} \text{ m})(3.5 \times 10^{-3} \text{ m})}{2.5 \text{ m}} = 7.0 \times 10^{-7} \text{ m} = 700 \text{ nm}$$

Example (Cont'd)



$$\Delta y = \frac{L}{d} \lambda$$

Fringe spacing is
proportional to λ , L ,
and inverse
proportional to d .



P698, Prob. 6, 7, 9

Prob. 15 (P698)



(a) Consider **three** equally spaced and equal-intensity coherent sources of light (such as adding a third slit to the two slits). Use the **phasor** method to obtain **the intensity** as a function of the phase difference $\Delta\varphi$.

(b) Determine the **positions** of maxima and minima.

Solution:

(a) As shown, the magnitude of $E_{\theta 0}$ is

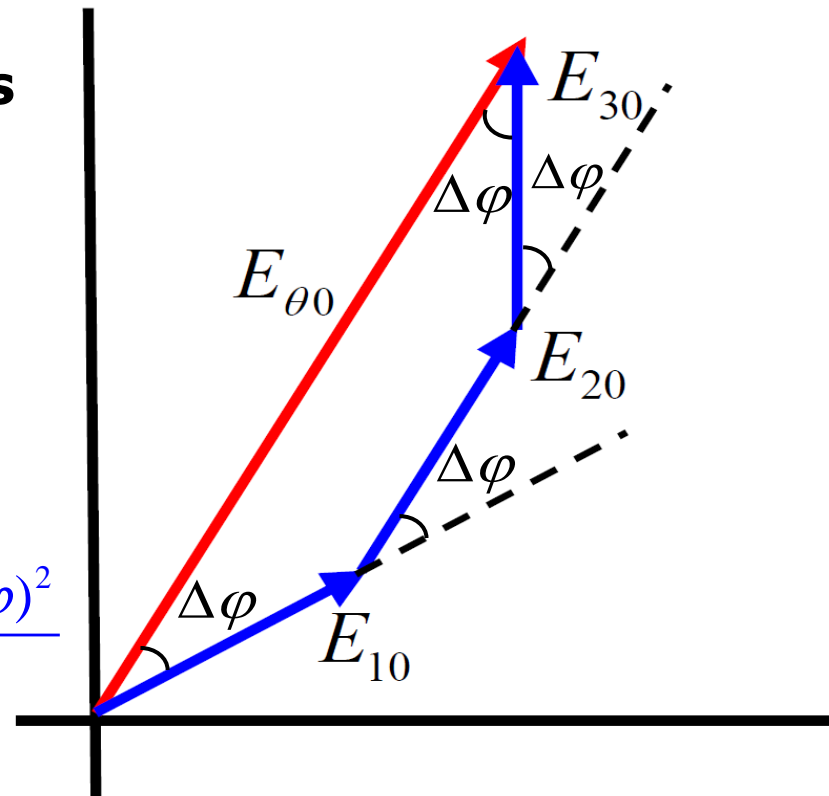
$$\begin{aligned} E_{\theta 0} &= E_{10} \cos \Delta\varphi + E_{20} + E_{30} \cos \Delta\varphi \\ &= E_{10} (1 + 2 \cos \Delta\varphi) \end{aligned}$$

$$\Delta\varphi = \frac{2\pi}{\lambda} d \sin \theta$$

The relative intensity is

$$\frac{I_{\theta}}{I_0} = \frac{E_{\theta 0}^2}{E_{00}^2} = \frac{[E_{10} (1 + 2 \cos \Delta\varphi)]^2}{[E_{10} (1 + 2 \cos 0)]^2} = \frac{(1 + 2 \cos \Delta\varphi)^2}{9}$$

where I_0 is the maximum intensity.



Prob. 15 (P698)



(b) Determine the positions of **maxima** and **minima**.

Solution:

$$I_{\theta} = \frac{(1 + 2 \cos \Delta \varphi)^2}{9} I_0$$

(b) When the three phasors are all in line, the intensity will be at its maximum I_0 .

$$(\Delta \varphi)_{\max} = \frac{2\pi}{\lambda} d \sin \theta = (2\pi)m, \quad m = 0, \pm 1, \pm 2 \dots \quad \sin \theta_{\max} = m \frac{\lambda}{d}$$

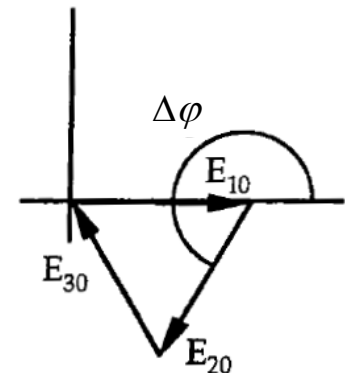
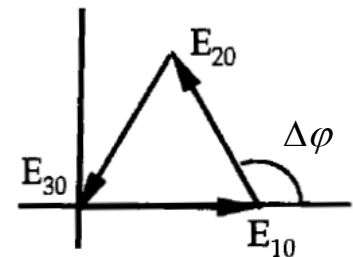
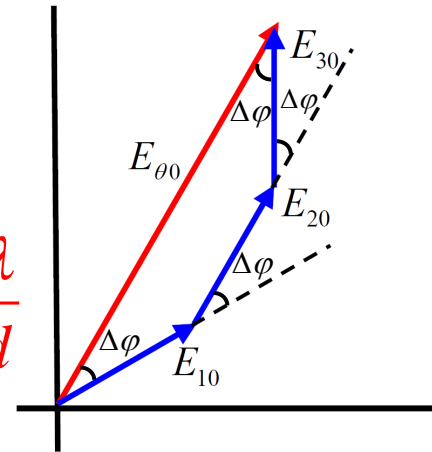
The intensity will be a minimum when

$$1 + 2 \cos \Delta \varphi = 0$$

$$(\Delta \varphi)_{\min} = \arccos\left(-\frac{1}{2}\right) = \begin{cases} \frac{2}{3}\pi + 2m\pi = 2\pi\left(m + \frac{1}{3}\right) \\ \frac{4}{3}\pi + 2m\pi = 2\pi\left(m + \frac{2}{3}\right) \end{cases}, \quad m = 0, \pm 1, \pm 2 \dots$$

$$(\Delta \varphi)_{\min} = 2\pi\left(m + \frac{1}{3}k\right) = \frac{2\pi}{\lambda} d \sin \theta_{\min}, \quad k = 1, 2; \quad m = 0, \pm 1, \pm 2 \dots$$

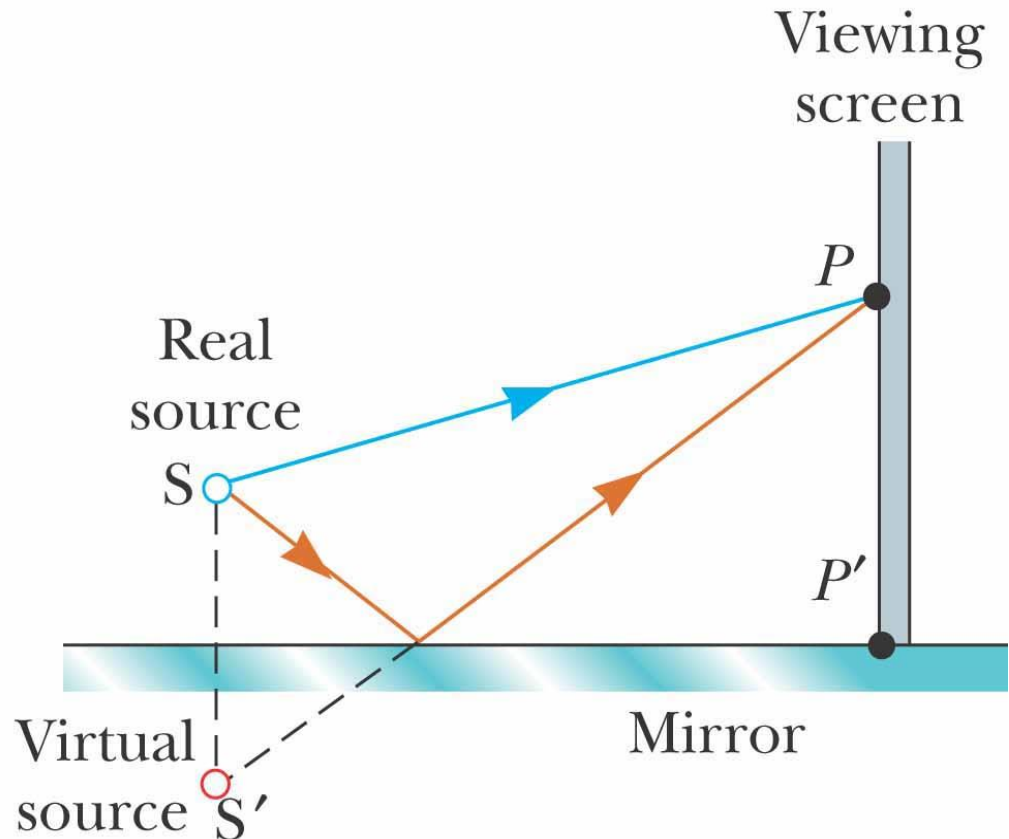
$$\sin \theta_{\min} = \frac{\lambda}{d} \left(m + \frac{1}{3}k\right)$$



Lloyd's mirror (劳埃德镜)



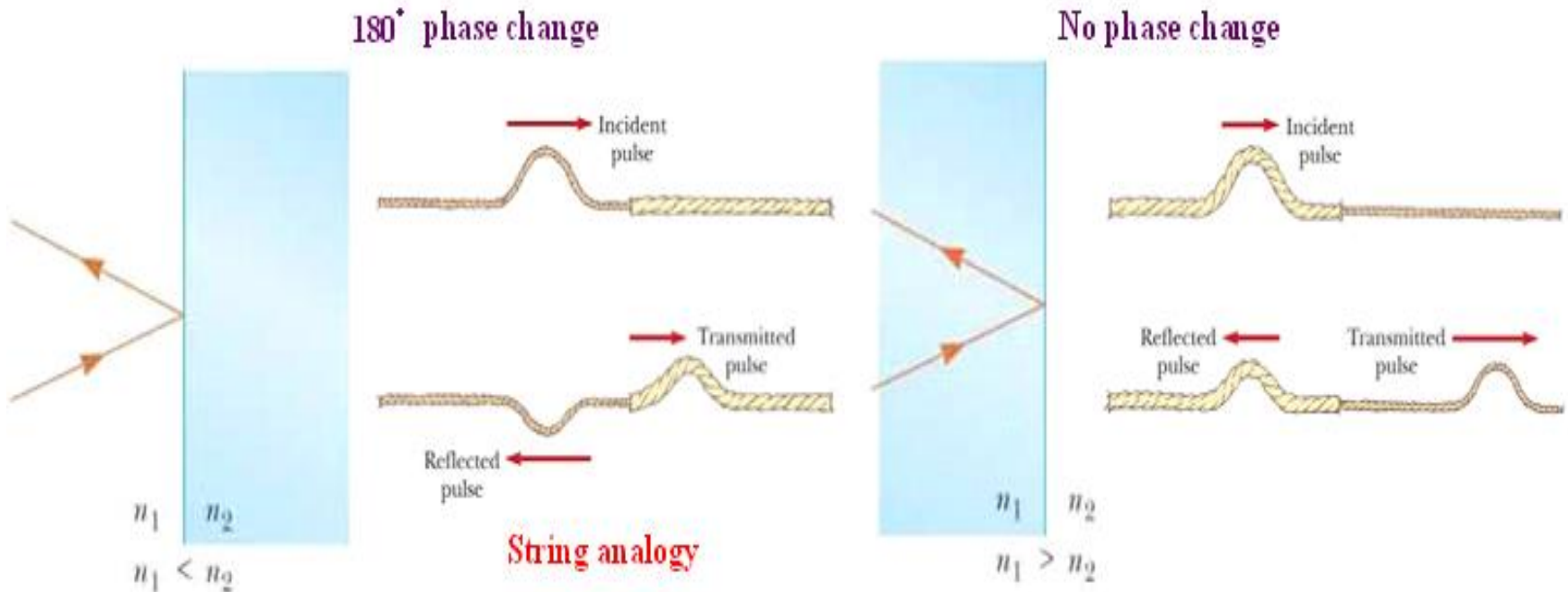
- Lloyd's mirror — another wavefront-splitting interference:
- ➡ The mirror reflects a portion of wavefront coming from slit S . This wave seems to come from the mirror image S' . Another portion of the wavefront proceeds directly from the slit to the screen. The interference pattern in the viewing screen is formed by coherent sources S and its image S' .



Lloyd's mirror



- The phase change at the reflecting boundary:



Lloyd's mirror



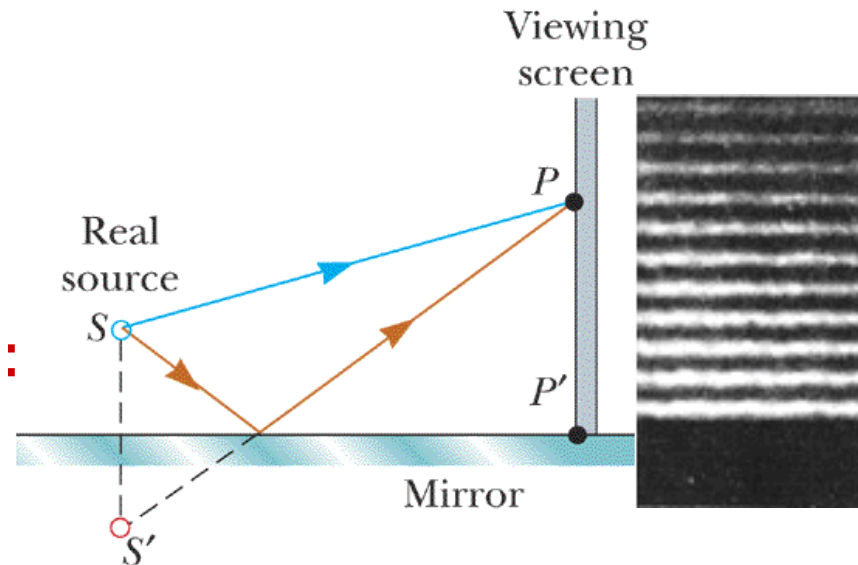
■ The phase difference in Lloyd's mirror interference:

- ➡ The distinguishing feature of Lloyd's mirror is that the reflected beam undergoes a 180° phase shift, corresponds to a **additional optical path length difference $\lambda/2$** .
- ➡ The phase difference:

$$\Delta\varphi = \frac{2\pi}{\lambda} \delta = \frac{2\pi}{\lambda} d \sin\theta \pm \pi$$

- ➡ The optical path length difference:

$$\delta = \underline{S'P} - \underline{SP} \pm \frac{\lambda}{2} = d \sin\theta \pm \frac{\lambda}{2}$$



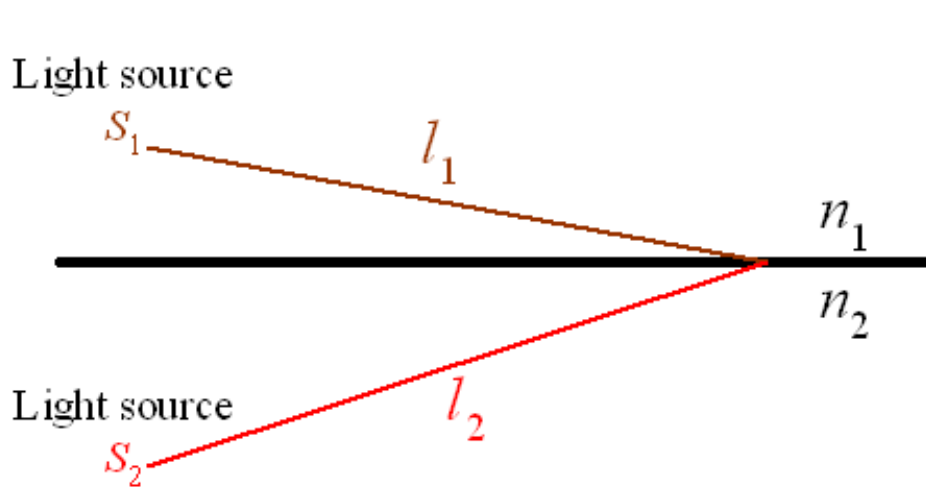
■ The interference pattern for Lloyd's mirror

- ➡ Unlike we will see a bright fringe at P' in Young's double-slit experiment, instead we will observe a **dark** fringe at P' because of 180° phase shift on reflection. The whole pattern is a **reversion** of Young's pattern.

§ 2 Optical Path Length



Comparison of phase retardations when two light waves travel through two **different** routes.



$$\begin{aligned}\Delta\varphi &= 2\pi \frac{l_2}{\lambda_2} - 2\pi \frac{l_1}{\lambda_1} \\ &= 2\pi \frac{l_2}{\frac{\lambda}{n_2}} - 2\pi \frac{l_1}{\frac{\lambda}{n_1}} \\ &= \frac{2\pi}{\lambda} (n_2 l_2 - n_1 l_1)\end{aligned}$$

Optical Path Length



Why introducing the optical path length

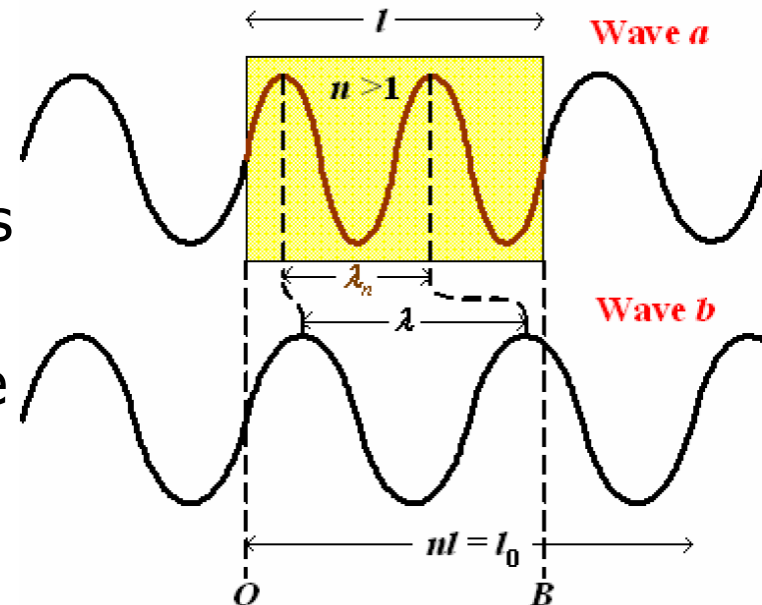
- ➡ Light travels with speed v in **medium** rather than c in **vacuum**.
- ➡ The propagation of wave a undergoes a medium whose index is $n > 1$.
- ➡ When the waves arrive at point B , the **phase retardation** for wave a is

$$\varphi_a = \frac{2\pi}{\lambda_n} l$$

λ_n is the wavelength in the medium. $\lambda_n = \frac{\lambda}{n}$

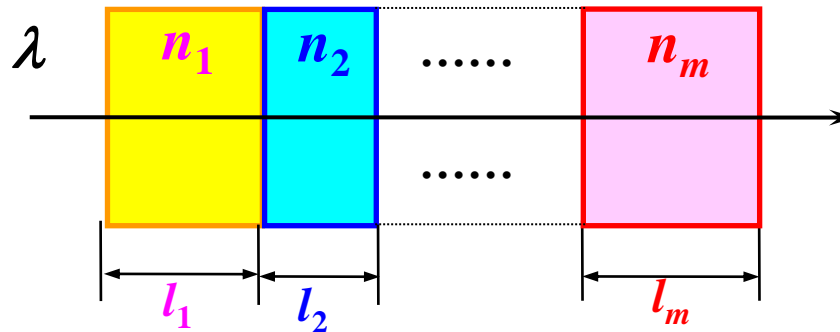
$$\varphi_a = \frac{2\pi}{\lambda} nl = \frac{2\pi}{\lambda} L, \quad L = nl$$

- ➡ The number of wavelength covered by the distance l in the **medium** is equal to the number of wavelength covered by the distance $L = nl$ in the **vacuum**.



■ The optical path

- When a light wave travel across a series of medium composed of m layers, for medium i the index is n_i , the **optical path** traversed by the wave is



$$L = \sum_{i=1}^m n_i l_i$$

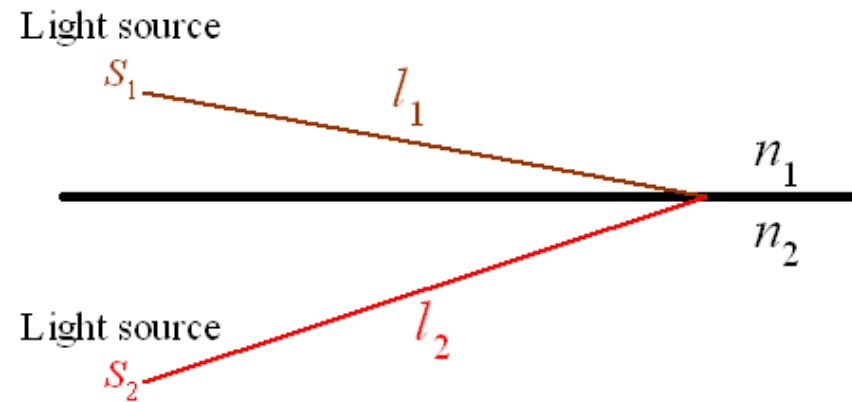
$$\varphi = \frac{2\pi}{\lambda} L$$

Applications of optical path length



- ➡ The relationship between the **optical path length** and **phase retardation**.

Comparison of phase retardations when two light waves travel through two different routes.



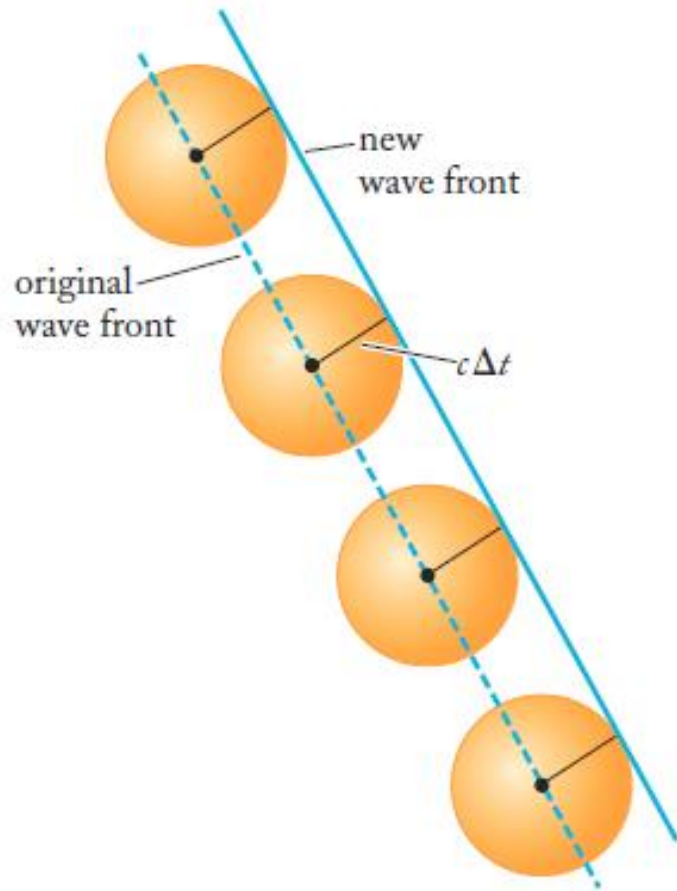
$$\begin{aligned} \text{phase difference} = \Delta\varphi = \varphi_2 - \varphi_1 &= \frac{2\pi}{\lambda_2} l_2 - \frac{2\pi}{\lambda_1} l_1 = \frac{2\pi}{\lambda} (n_2 l_2 - n_1 l_1) = \frac{2\pi}{\lambda} \delta \\ &= \frac{2\pi}{\lambda} \times \text{difference of optical path length} \end{aligned}$$

- ➡ The relationship between the **optical path length** and **traveling time**.

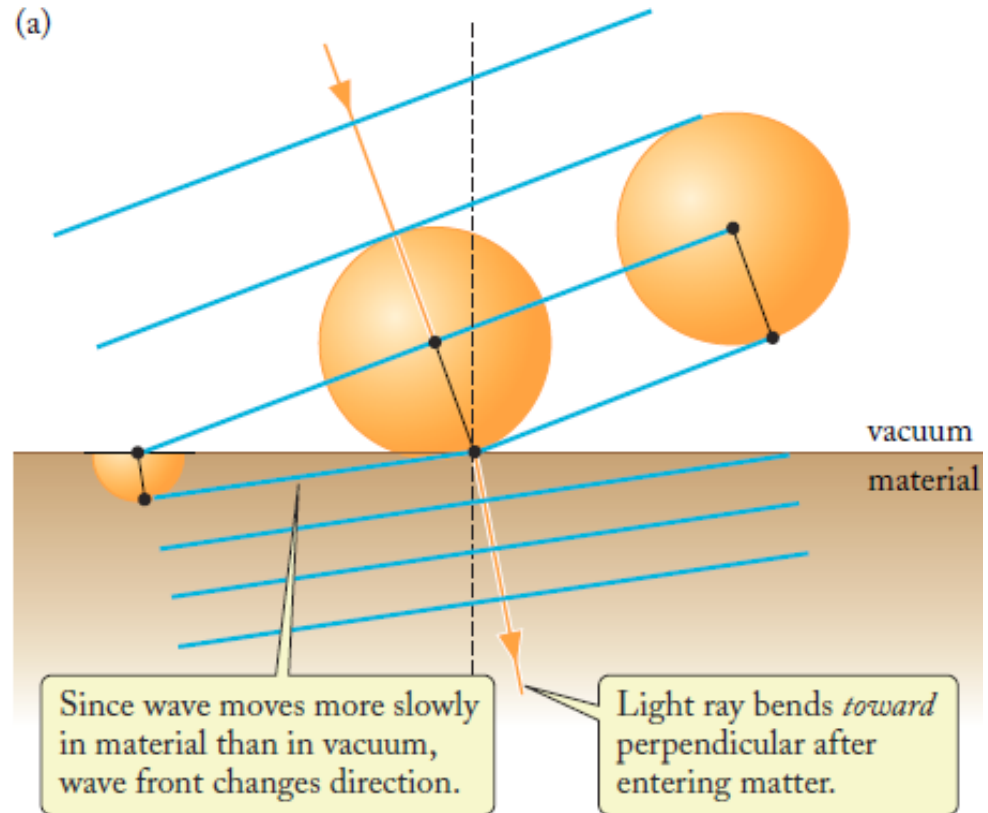
Comparison of time consumed when two light rays traverse through two different routes.

$$\begin{aligned} \text{time difference} = \Delta t = t_2 - t_1 &= \frac{l_2}{v_2} - \frac{l_1}{v_1} = \frac{1}{c} (n_2 l_2 - n_1 l_1) = \frac{\delta}{c} \\ &= \frac{\text{difference of optical path length}}{c} \end{aligned}$$

Applications of optical path length



(a)

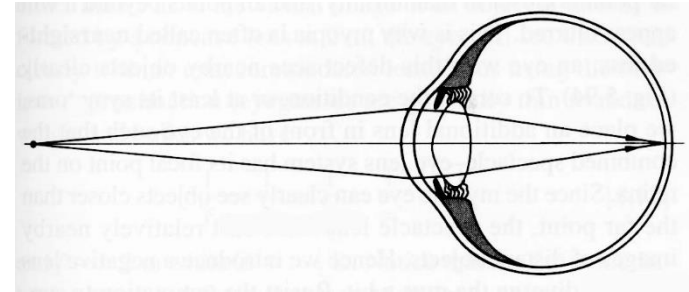
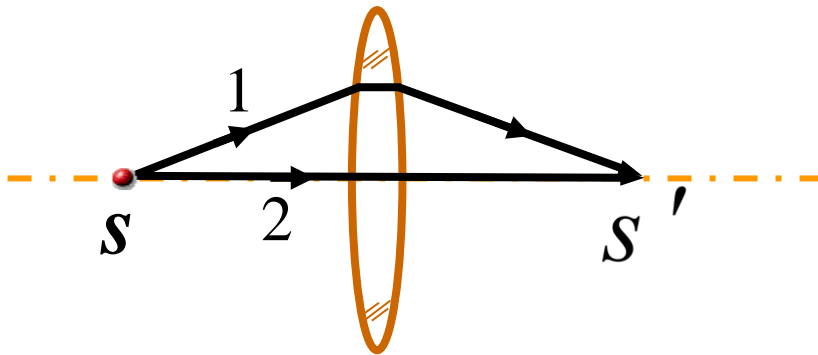


Applications of optical path length



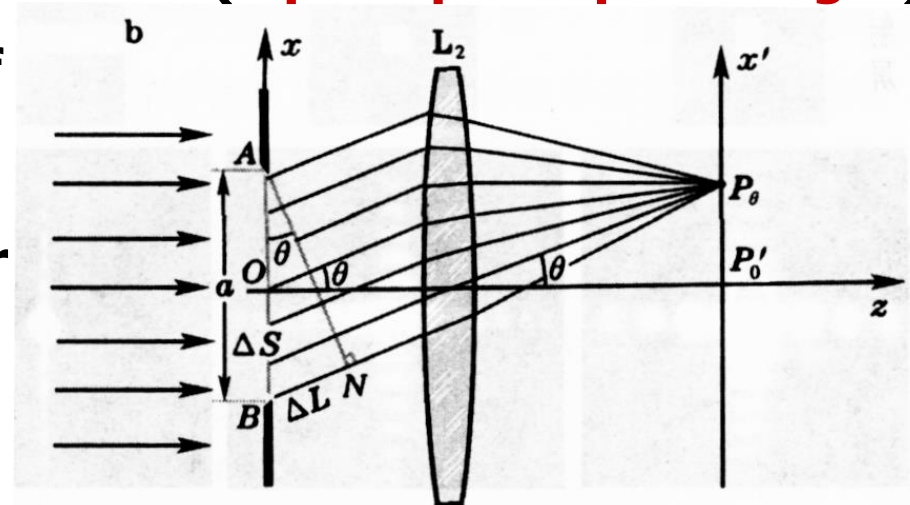
➔ **No** optical path length difference through lens.

The various light pass through the lens would introduce **no** additional optical path difference or phase shift.



A human eye can get stable information of an object because each ray travels with the **same time (equal optical path length)**.

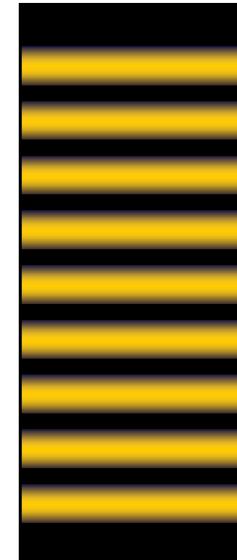
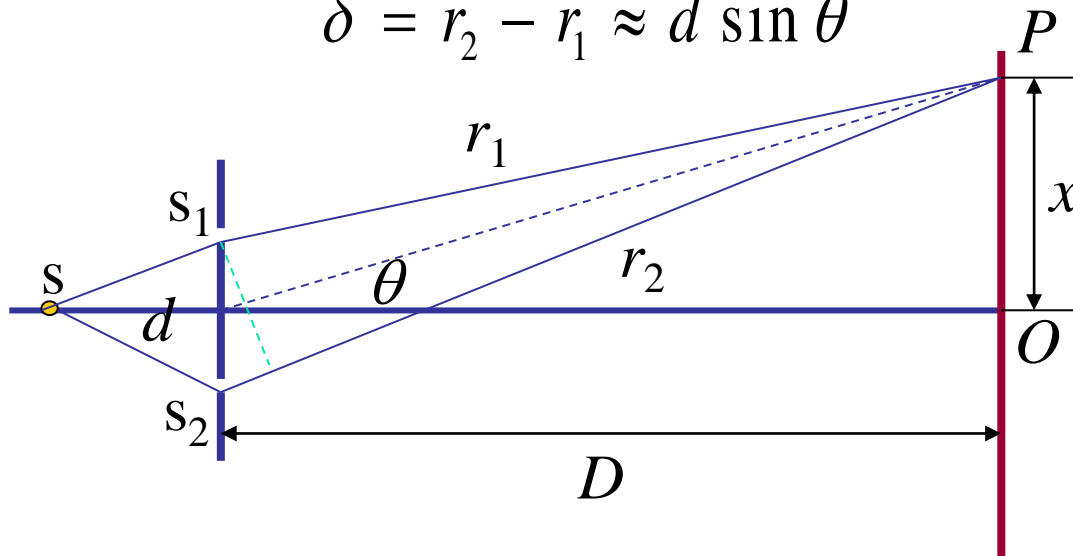
When we make a comparison of optical differences among all rays that will focus on the point P_θ , we can only consider the portion of rays **before** the wavefront AN .



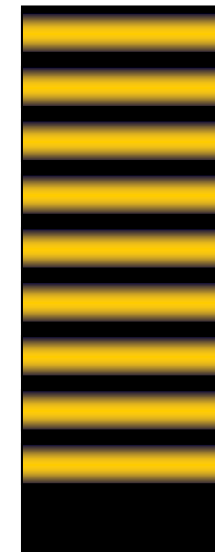
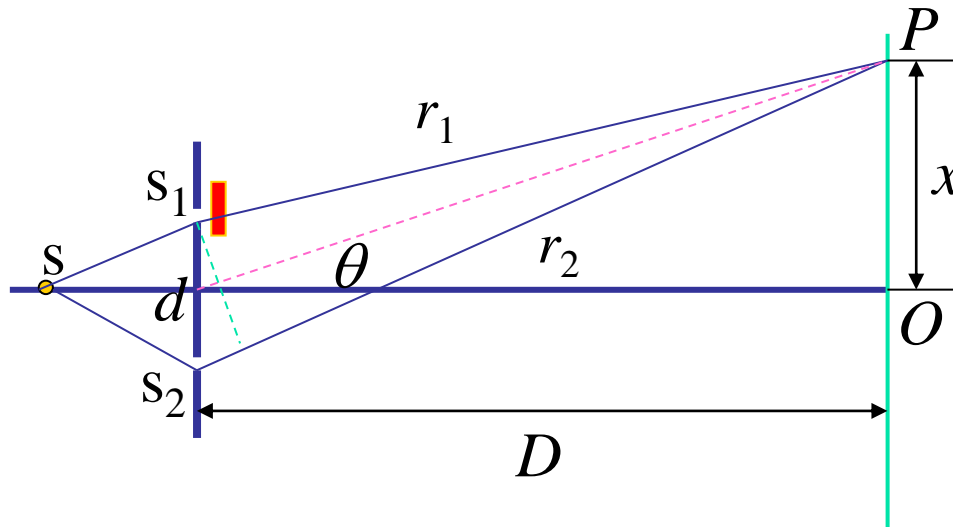
Example



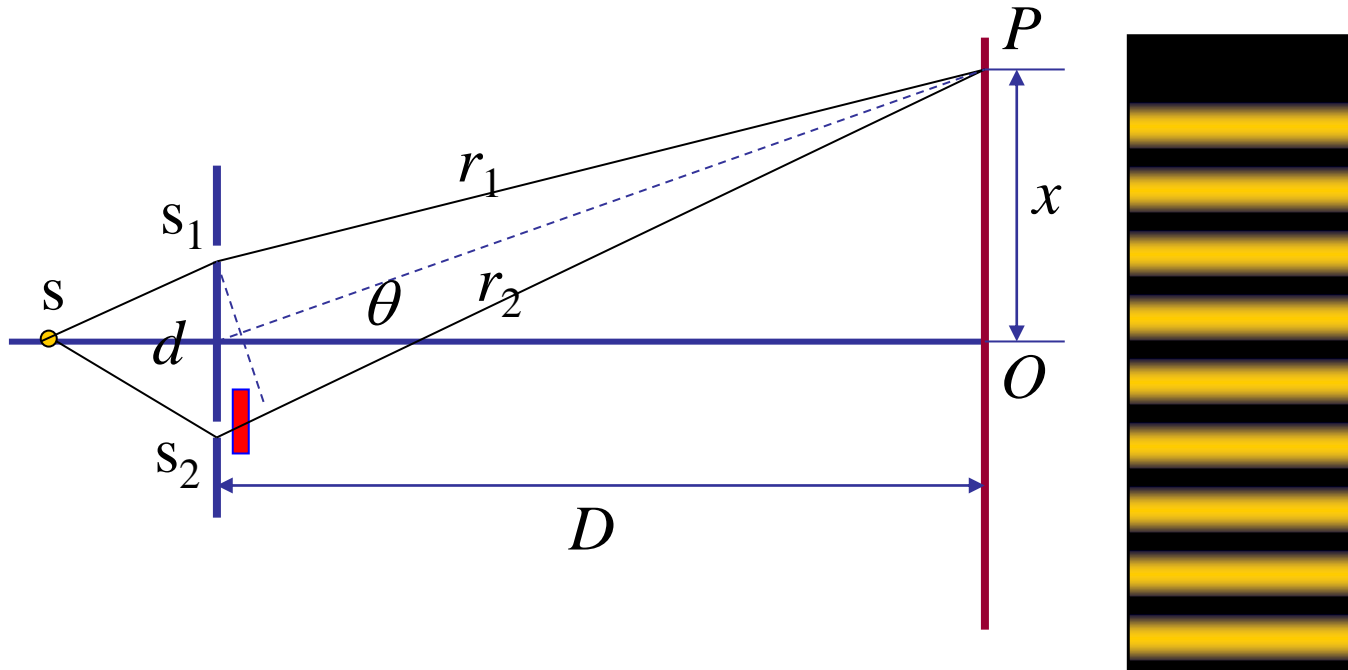
$$\delta = r_2 - r_1 \approx d \sin \theta$$



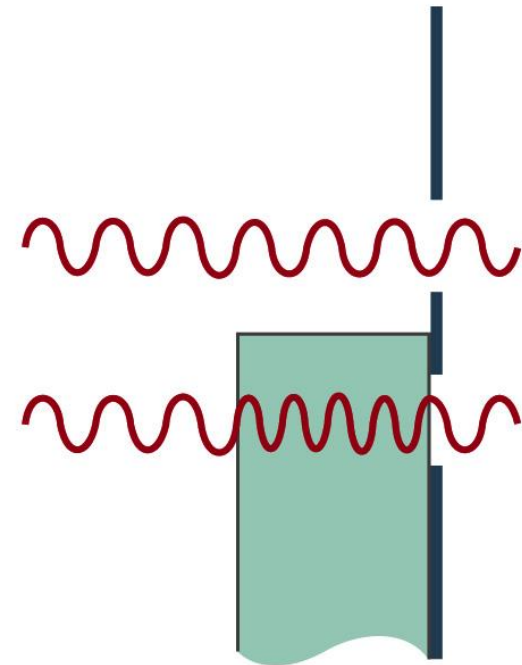
$$\delta' = r_2 - (r_1 - e + ne) < \delta$$



Example



Suppose a thin piece of glass were placed in front of the lower slit so that the two waves enter the slits **180°** out of the phase. Describe in detail the interference pattern on the screen. What is the **minimum thickness** of the glass? (n, λ)



Solution:

$$\delta_{\min} = nt_{\min} - t_{\min} = \frac{\lambda}{2}, \quad t_{\min} = \frac{\lambda}{2(n-1)}$$

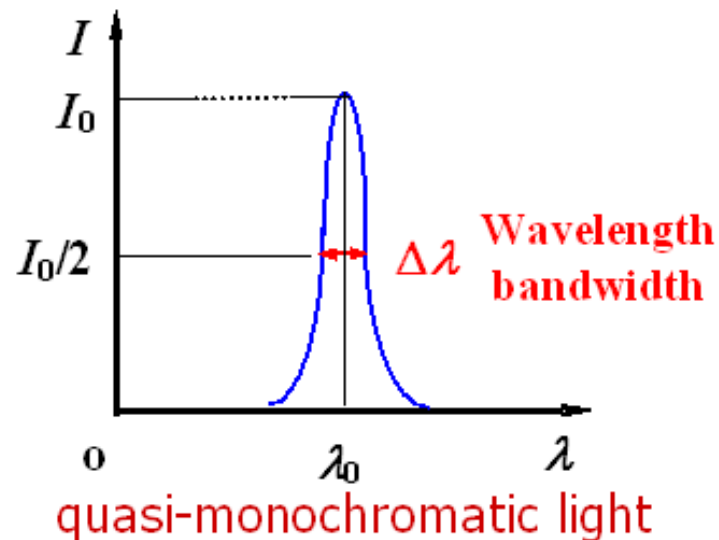


Problems



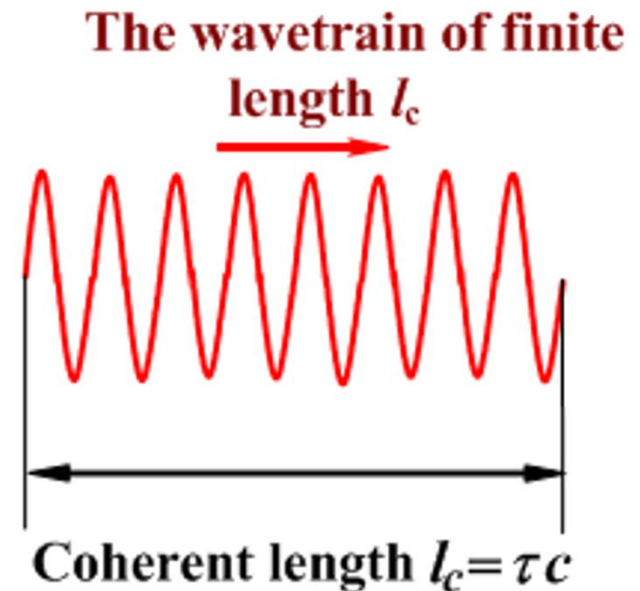
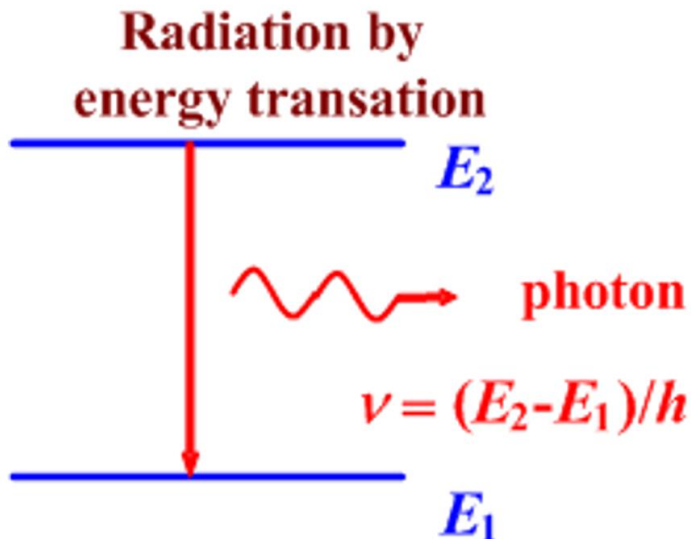
P698, Prob. 10, 11; 14, 16

- **Monochromatic light and quasi-monochromatic light:**
 - ➡ **Monochromatic** light means single-color light having a single wavelength.
 - ➡ Absolutely monochromatic light is an unattainable idealization. Real light have a range of wavelength bandwidth $\Delta\lambda$ around its central wavelength λ_0 . A light with a narrow band $\Delta\lambda \ll \lambda_0$ is called **quasi-monochromatic** light.



■ Wavetrains and coherence length:

- ➡ In the viewpoint of **quantum mechanics**, an atom emits a photon by energy transition from upper energy level to lower level within a short duration time τ .
- ➡ In the viewpoint of **classical physics**, the same thing is described as: a dipole emits a wavetrain in time τ called **coherence time**. The wavetrain has a corresponding **coherence length** $l_c = \tau c$



Coherence and incoherence



Coherence Lengths of Several Sources

| Source | $\lambda_0(\text{nm})$ | $\Delta\lambda(\text{nm})$ | Coherence length l_c |
|---------------------------------|------------------------|----------------------------|----------------------------|
| White Light | 550 | ≈ 300 | $\approx 900 \text{ nm}$ |
| Mercury arc | 546.1 | ≈ 1.0 | $\sim < 0.03 \text{ cm}$ |
| Kr ⁸⁶ discharge lamp | 605.6 | 0.0012 | 0.3 m |
| Stabilized He-Ne Laser | 632.8 | $\approx 10^{-6}$ | $\sim < 400 \text{ m}$ |
| Special He-Ne Laser | 1153 | 8.9×10^{-11} | $15 \times 10^6 \text{ m}$ |

➔ Kr——krypton

➔ He——helium

➔ Ne——neon

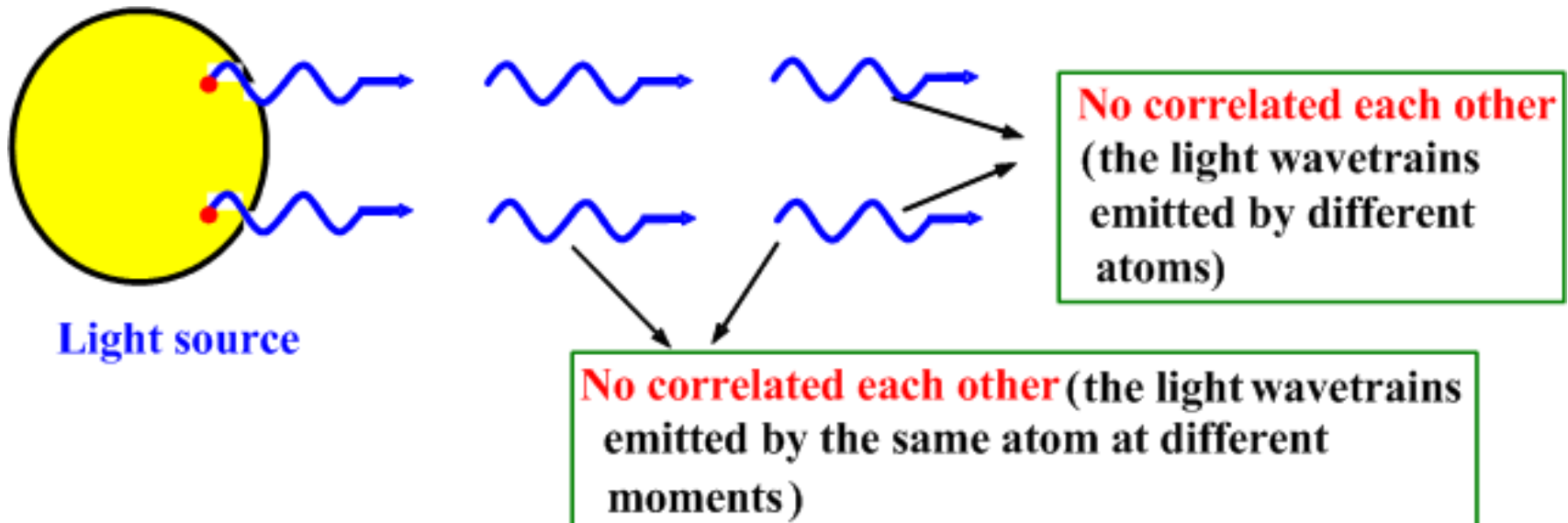
(coherence length) $l_c = \frac{\lambda_0^2}{\Delta\lambda}$

(coherence time) $\tau = \frac{l_c}{c}$

Coherence and incoherence



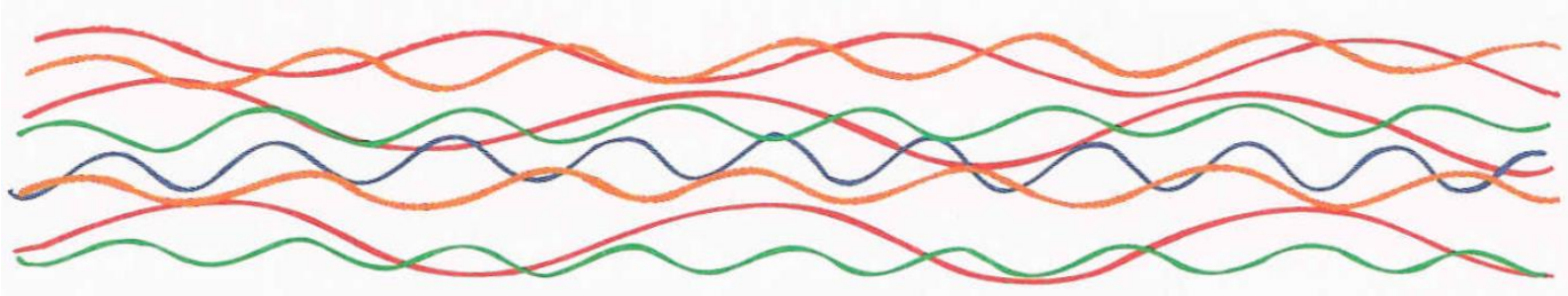
- Atoms emit wavetrains with **random** phase angles at different moments. $\phi = \phi(t)$
- There is no correlation—no definite phase relation—between the wavetrains emitted by **different** atoms.
- There is no correlation—no definite phase relation—between the wavetrains emitted by the same atoms at **different** moments.



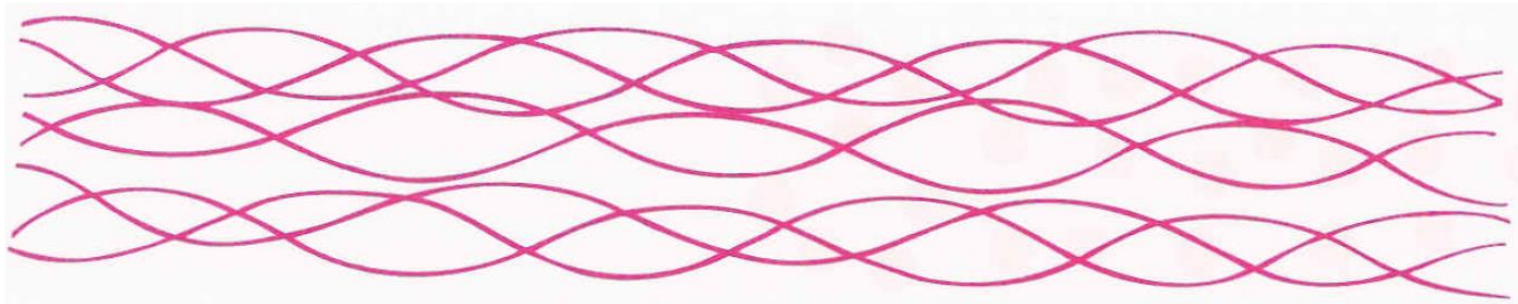
Coherence and incoherence



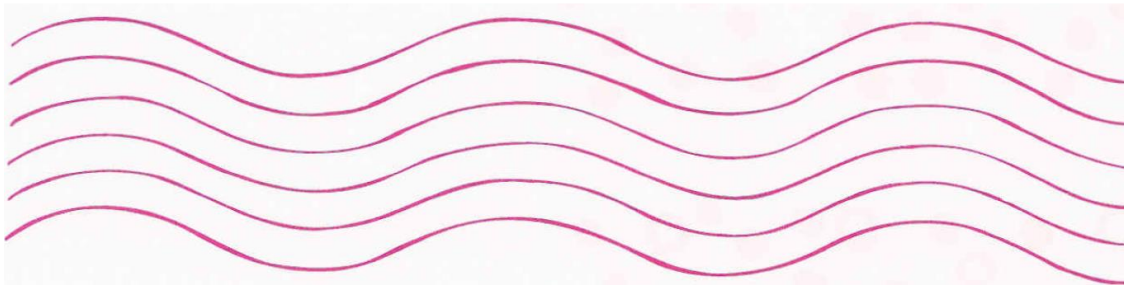
- ➡ **Incoherent white light contains waves of many frequencies (and of many wavelengths) that are out of phase with one another.**



- ➡ **Light of a single frequency and wavelength still contains a mixture of phases.**



- ➡ **Coherent light: All the waves are identical and in phase .**



■ Coherence and incoherence:

- ➡ The intensity of light is the time average of the Poynting vector

$$I = \langle (E_1 + E_2)^2 \rangle = \langle E_1^2 \rangle + \langle E_2^2 \rangle + 2 \langle E_1 E_2 \rangle$$
$$= I_1 + I_2 + 2\sqrt{I_1 I_2} \langle \cos \Delta\varphi \rangle$$

- ➡ The phase difference is generally random due to random phase angles.

$$\Delta\varphi = k(r_2 - r_1) - [\phi_2(t) - \phi_1(t)] = \Delta\varphi(t)$$

- ➡ **Incoherence:** $\langle \cos \Delta\varphi \rangle = 0 \Rightarrow I_{\text{incoh}} = I_1 + I_2$

- ➡ **Coherence:**

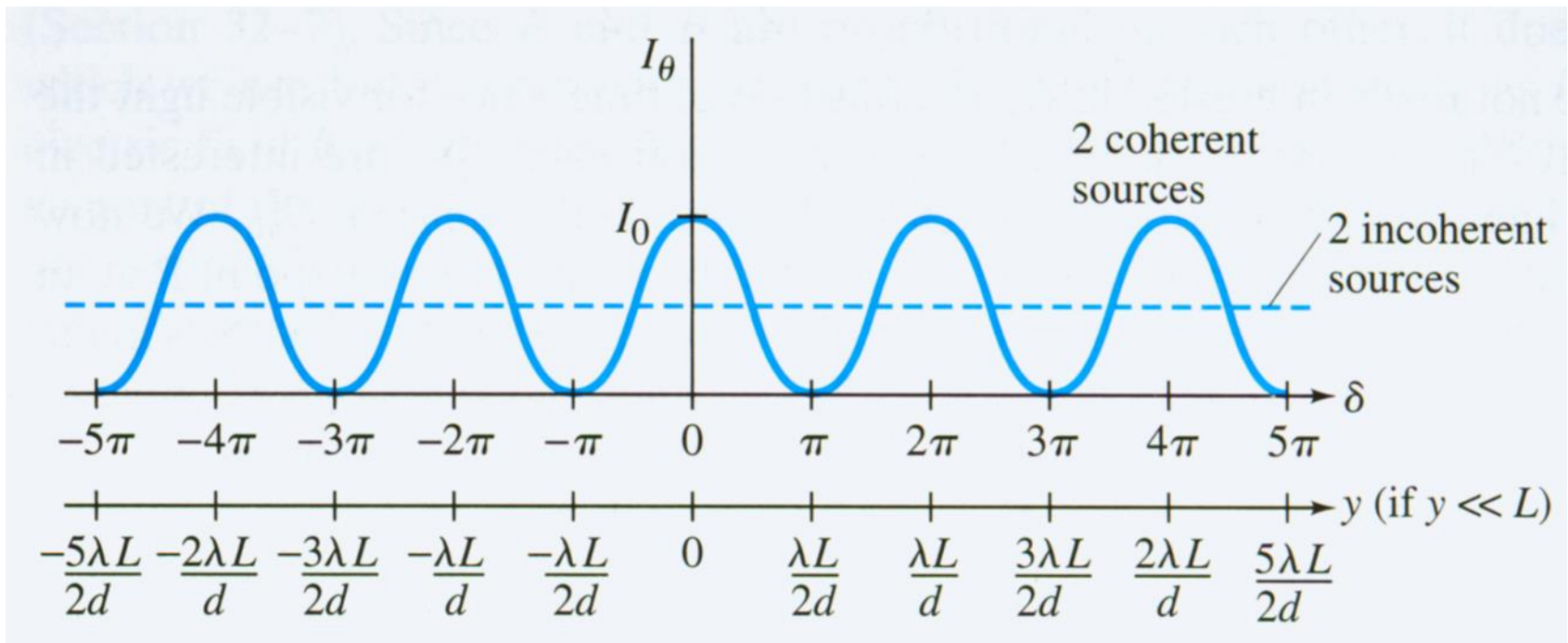
$$I_{\text{coh}} = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos \Delta\varphi = \begin{cases} I_{\text{max}} = I_1 + I_2 + 2\sqrt{I_1 I_2} \\ I_{\text{min}} = I_1 + I_2 - 2\sqrt{I_1 I_2} \end{cases}$$

Coherence and incoherence



- ➡ The energy (intensity) is spread evenly over the screen and there is no interfering maxima and minima, when the sources produce **incoherent** light. $I_{\text{incoh}} = I_1 + I_2$
- ➡ Whereas the energy is distributed in peaks and valleys when the light from two sources is **coherent**.

$$I_{\text{coh}} = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos \Delta\varphi$$



■ The **clever** design of Young's double-slit experiment

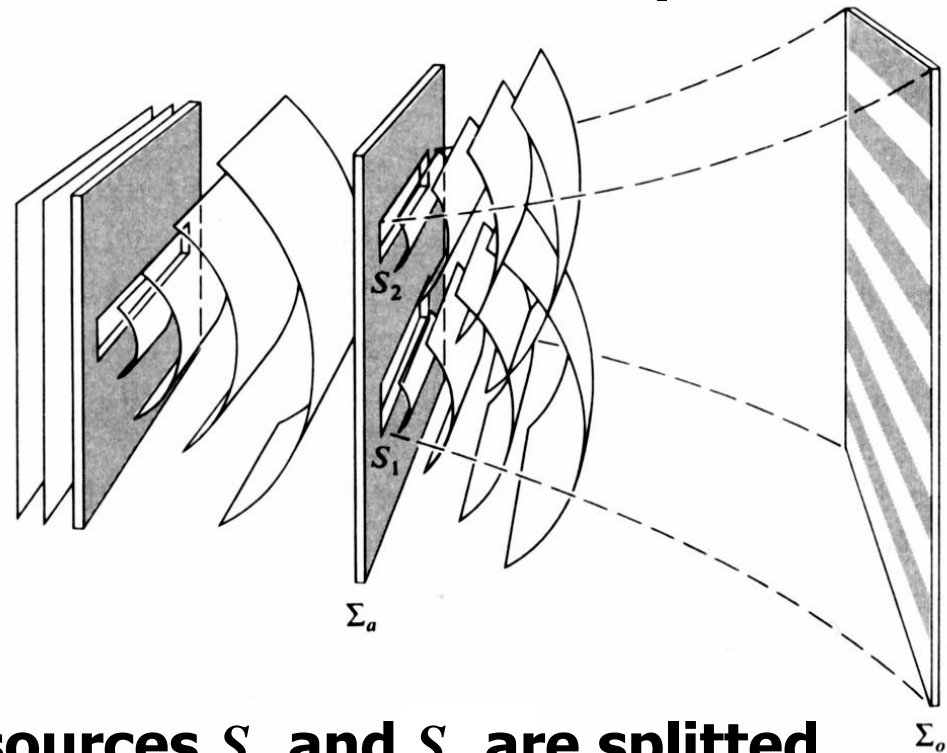
- ➡ The two sources S_1 and S_2 are coherent because they come from the same source S . Therefore their phase angles are the **same**.

$$I = I_1 + I_2 + 2\sqrt{I_1 I_2} \langle \cos \Delta\varphi \rangle$$

$$\phi_1(t) = \phi_2(t) = \phi(t)$$

$$\begin{aligned} \Delta\varphi &= k(r_2 - r_1) - \underbrace{(\phi_2(t) - \phi_1(t))}_{\text{zero}} \\ &= k(r_2 - r_1) \end{aligned}$$

$$\langle \cos \Delta\varphi \rangle \neq 0$$



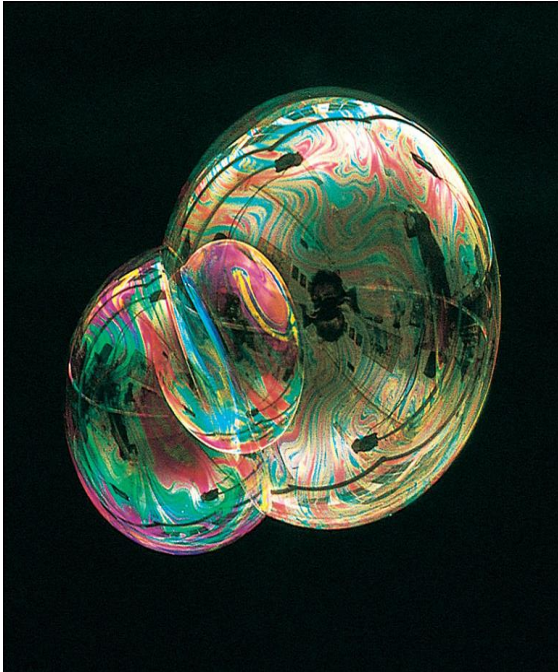
- ➡ Because the two coherent sources S_1 and S_2 are splitted from wavefront of S , so Young's experiment was in the category of **wavefront-splitting** interference.

§ 4 Interference in Thin Films

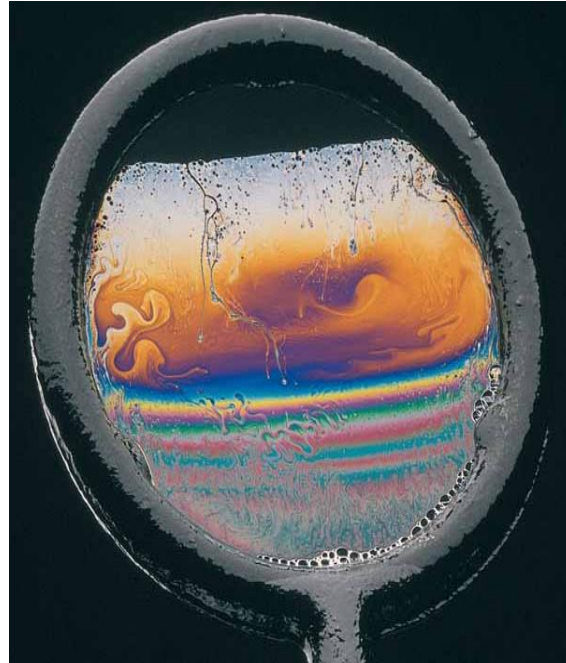


Amplitude-splitting interference

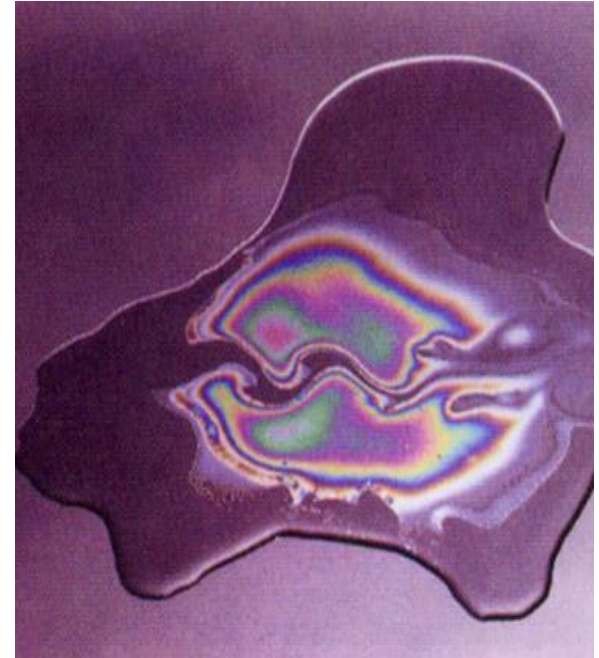
- The everyday phenomena due to the interference in thin films



Soap bubbles



Thin film of soapy water



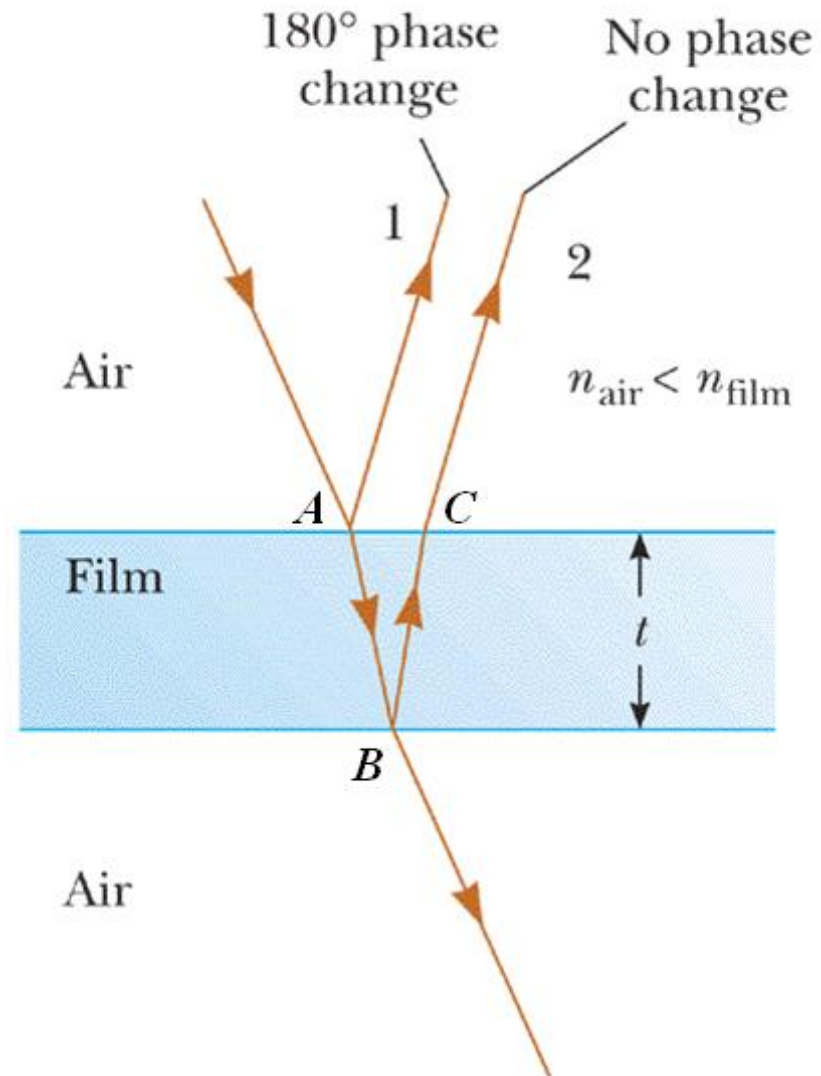
Thin layer of gasoline on water

Amplitude-splitting interference



- **Dielectric films — Double-beam interference**
 - ➔ Part of the incident light is reflected at *A* on the top surface, and part reflected at the bottom surface must travel the **extra** distance *ABC*.
 - ➔ The optical path length **difference** is:
$$\delta \approx 2n_{\text{film}}t + \left(\frac{\lambda}{2}\right)$$

$\lambda/2$ is a additional path depending on index relations at interfaces of the film.



The conditions for constructive and destructive interference

➡ **Bright fringes** —
constructive interference

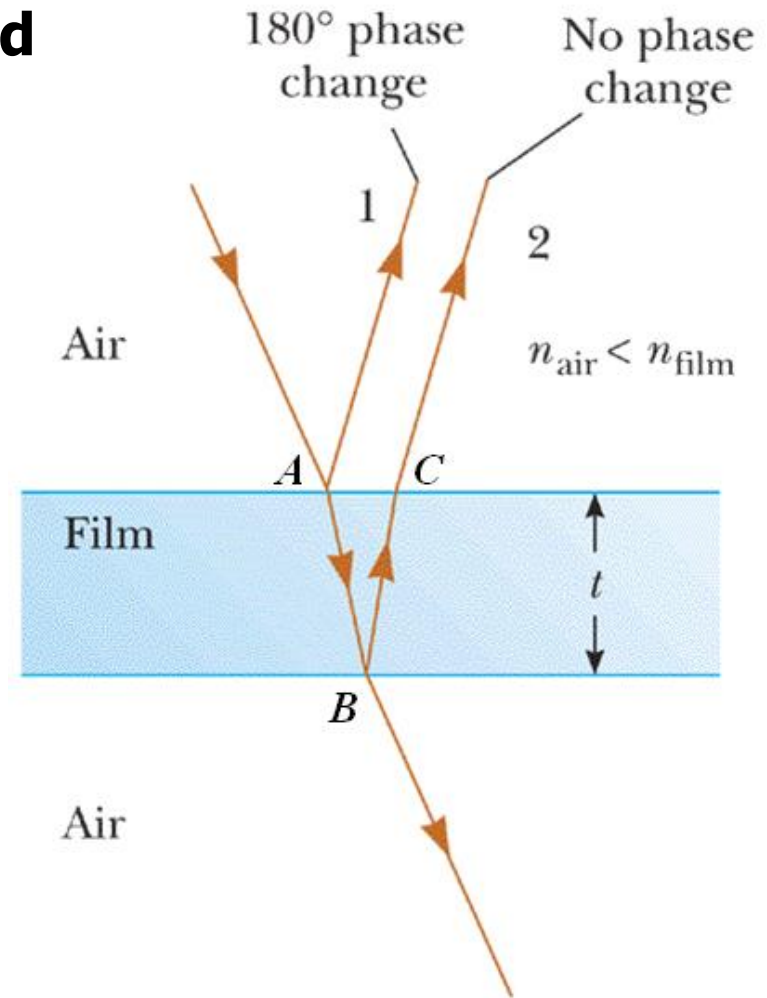
$$\delta = 2nt + \left(\frac{\lambda}{2}\right) = m\lambda$$

$$m = 0, 1, 2, \dots$$

➡ **Dark fringes** —
destructive interference

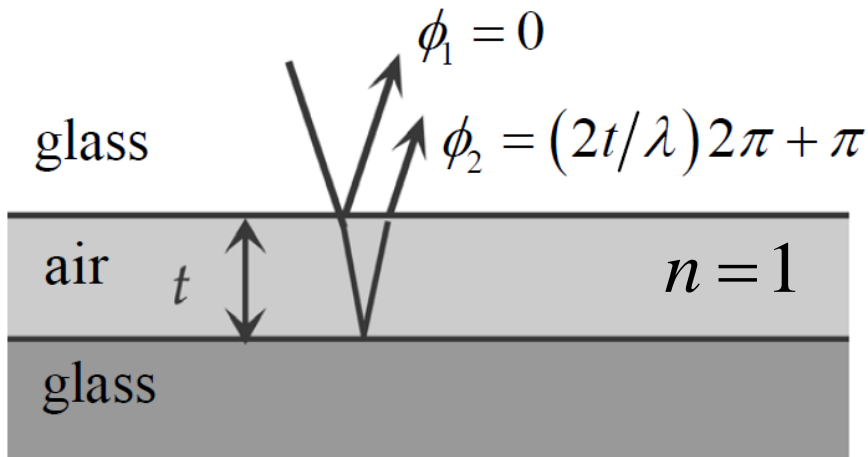
$$\delta = 2nt + \left(\frac{\lambda}{2}\right) = (2m + 1)\frac{\lambda}{2}$$

$$m = 0, 1, 2, \dots$$



What is the **minimum** (non-zero) **thickness** for the air layer between two flat glass surfaces if the glass is to appear **dark** when **640-nm** light is incident normally? What if the glass is to appear **bright**?

Solution:



How to generate fringes?

Dark, $\delta_{\min} = 2nt_{\min} + \frac{\lambda}{2} = \frac{3\lambda}{2}$

$$t_{\min} = \frac{\lambda}{2} = \frac{640}{2} = 320 \text{ nm}$$

Bright, $\delta_{\min} = 2nt_{\min} + \frac{\lambda}{2} = \lambda,$

$$t_{\min} = \frac{\lambda}{4} = \frac{640}{4} = 160 \text{ nm}$$

■ Air wedge

- Air wedge is the wedge of air between the two glass plates.

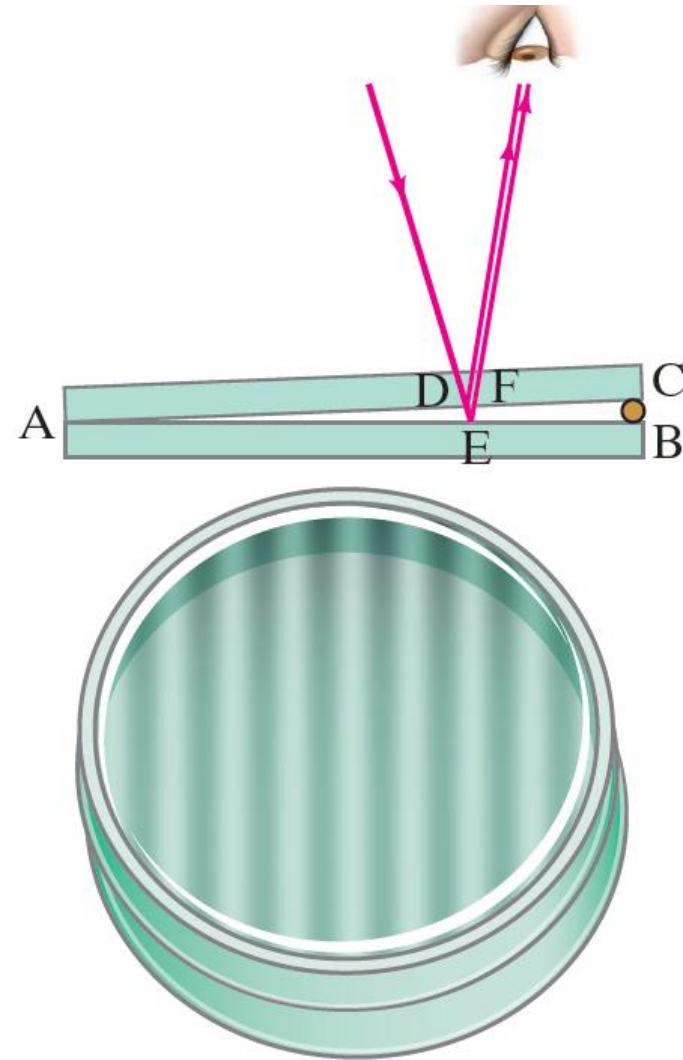
- **Bright** fringes:

$$\delta = 2t + \frac{\lambda}{2} = m\lambda, \quad m = 1, 2, \dots$$

- **Dark** fringes:

$$\delta = 2t + \frac{\lambda}{2} = (2m+1)\frac{\lambda}{2}, \quad m = 0, 1, 2, \dots$$

- Each fringe is the locus (轨迹) of all points in the film for which the optical thickness is constant. For air wedge, $n=1$, so that the fringes correspond to regions of **equal film thickness**.



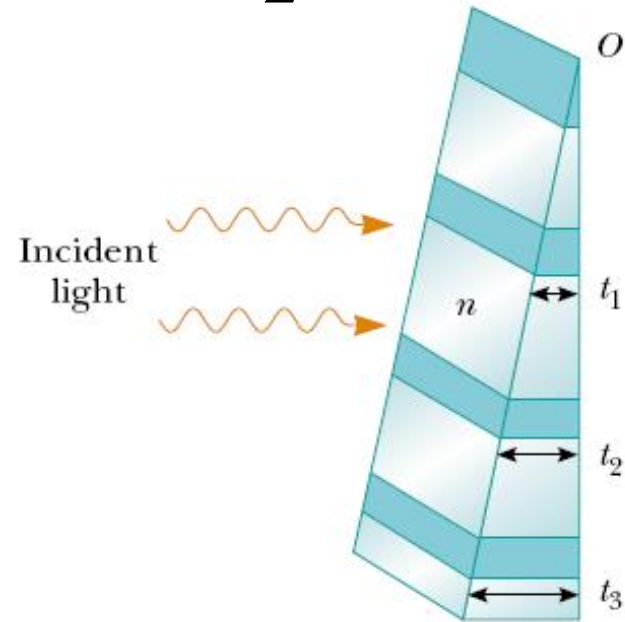
- Why can we see different colors in soap bubbles and other films of varying thickness?

Interference in a vertical film of variable thickness.



The **top** of the film appears darkest where the film is thinnest.

$$\delta = 2nt + \frac{\lambda}{2} = m\lambda \text{ (bright)}$$



If white light is used, bands of different **colors** are observed at different points, corresponding to the different wavelengths of light.

■ The applications of air wedge

➡ At the end where the two plates meet

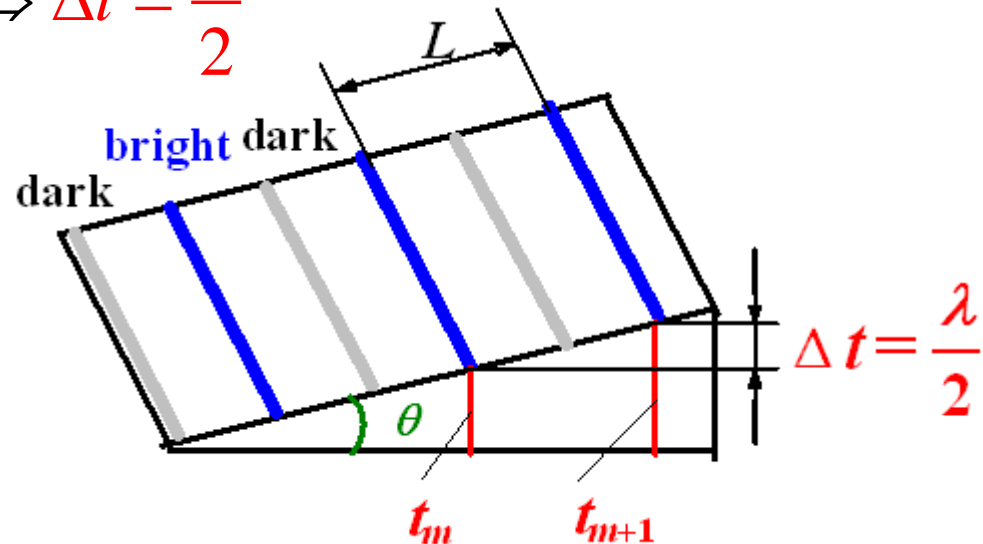
$$\delta = 2t + \frac{\lambda}{2}, \quad t = 0, \quad \delta = \frac{\lambda}{2} \quad \text{(dark fringe)}$$

➡ **The thickness difference** between the adjacent fringes

$$\left. \begin{aligned} \delta_m &= 2t_m + \frac{\lambda}{2} = m\lambda \\ \delta_{m+1} &= 2t_{m+1} + \frac{\lambda}{2} = (m+1)\lambda \end{aligned} \right\} \Rightarrow \Delta t = \frac{\lambda}{2}$$

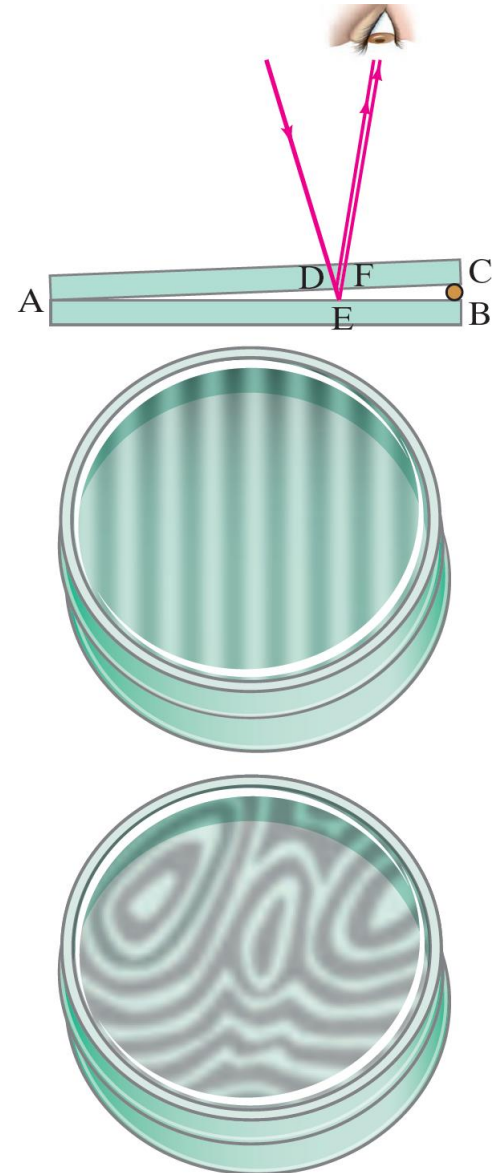
➡ **The spacing of fringes:**

$$L = \frac{\Delta t}{\sin \theta} \approx \frac{\Delta t}{\theta} = \frac{\lambda}{2\theta}$$



The applications of air wedge (cont'd)

- ➡ We can use air wedge to determine the surface features of optical elements (lenses, prisms).
- ➡ If we want examine whether a surface of an optical element is **flat or not**, we put it into contact with an *optical flat*. If the test surface is perfect flat, a series of straight, equal spaced fringes will be appeared near the surface.
- ➡ Now, mirrors that are flat to better than 5 percent of one wavelength, or about $(500 \times 5\% =)$ **25nm**, are available commercially.



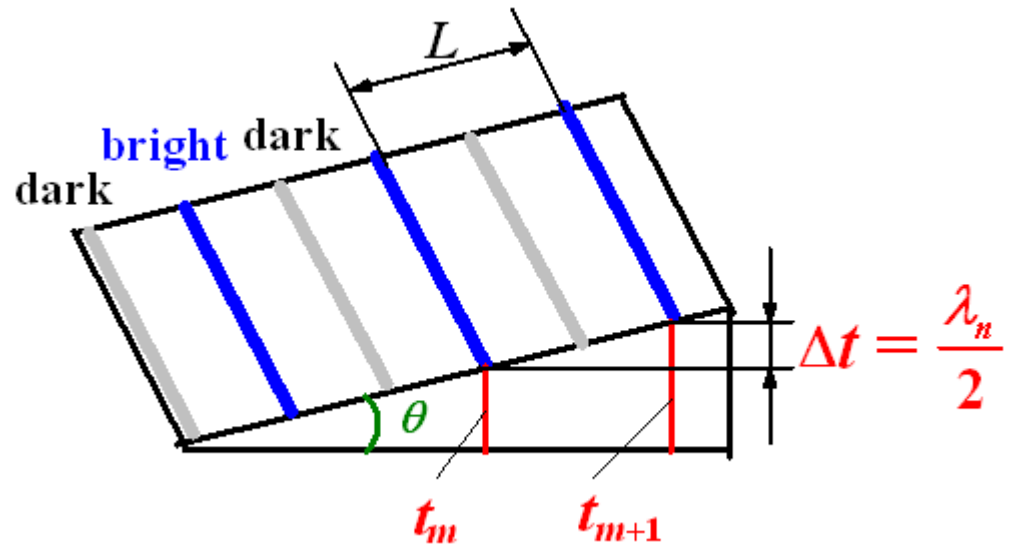
■ The applications of air wedge (cont'd)

- ➡ If the wedge between the two glass plates is filled with some transparent substance other than **air** — say, **water** — the pattern shifts because

$$\lambda_n = \frac{\lambda}{n}$$

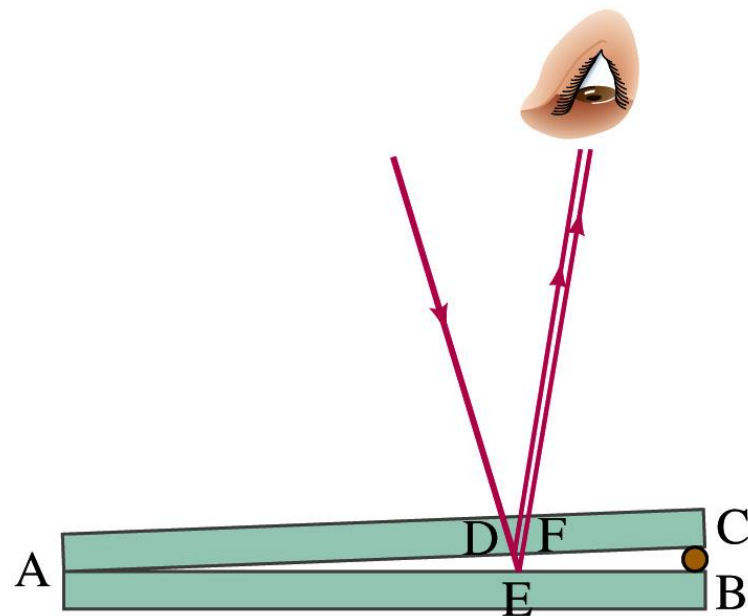
- ➡ The thickness difference between the adjacent fringes now is

$$\Delta t = \frac{\lambda_n}{2} = \frac{\lambda}{2n}$$

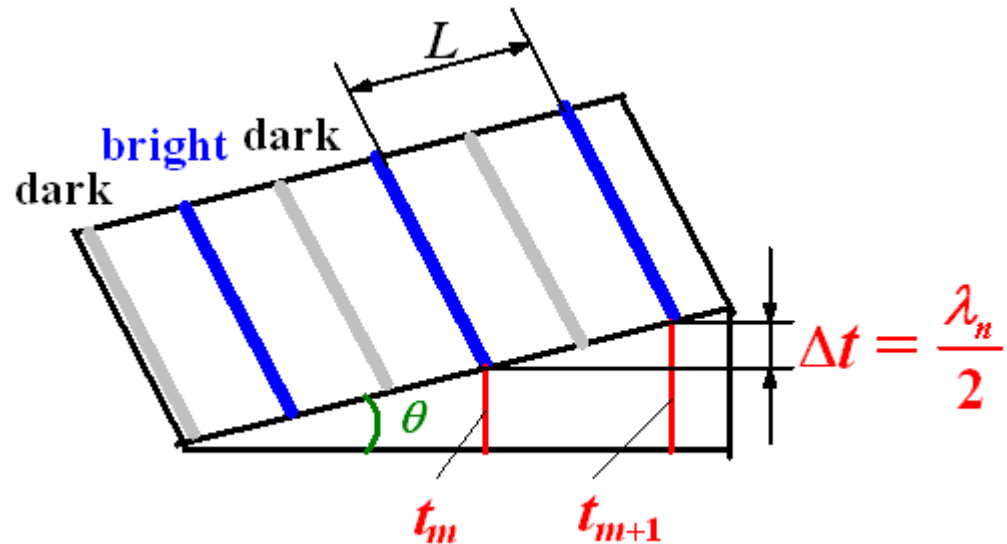
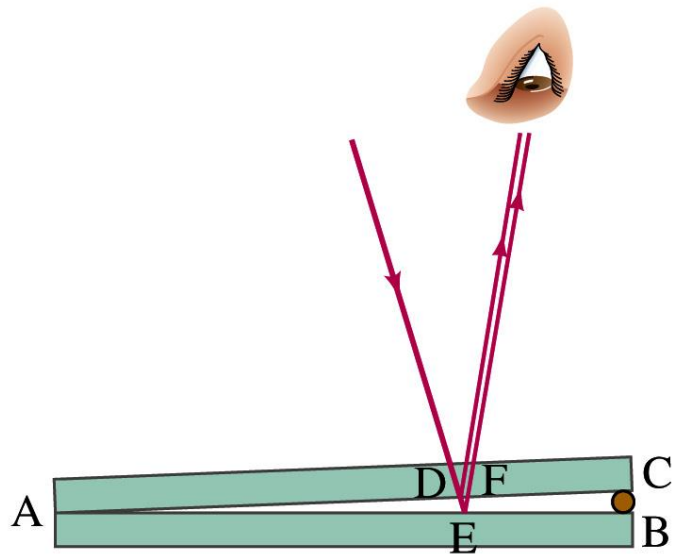


Thin film of air, wedge-shaped.

A very fine wire $7.35 \times 10^{-3} \text{ mm}$ in diameter is placed between two flat glass plates as seen. Light whose wavelength in air is 600 nm falls (and is viewed) **perpendicular** to the plates, and a series of bright and dark bands is seen. **How many** light and dark bands will there be in this case? Will the area next to the wire be bright or dark?



Ex. 30-6 (P692)



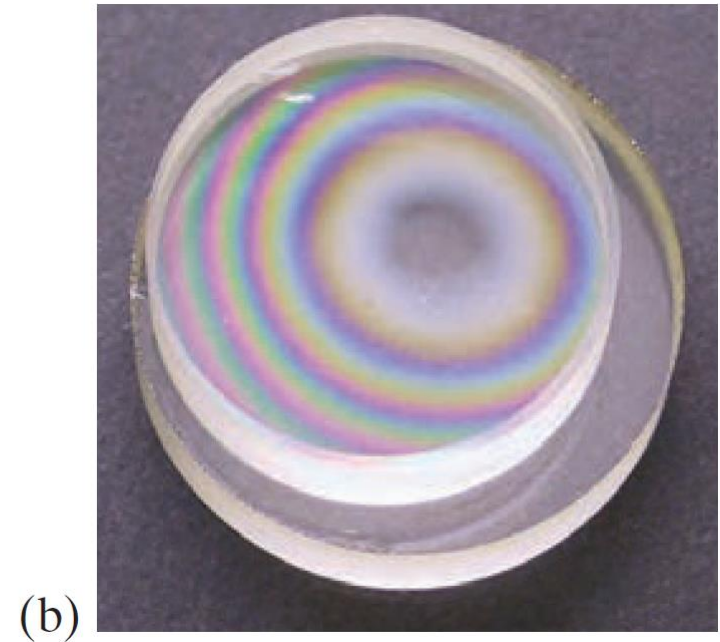
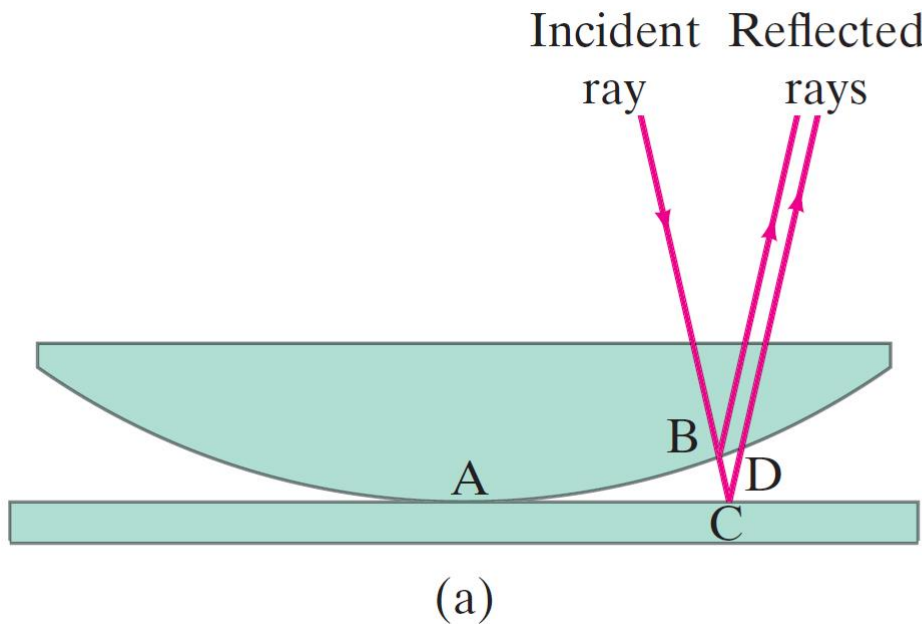
Solution:

$$\frac{t}{\lambda/2} = \frac{7.35 \times 10^{-6}}{600 \times 10^{-9} / 2} = 24.5$$

There will be a **25 dark** lines and **25 bright** lines.

■ Newton's rings

- When a convex (凸起的) surface of a lens is placed in contact with a flat glass surface, a series of **concentric rings** is seen when illuminated from above by monochromatic light.



Newton's rings



➡ The radii of Newton's ring:

$$r = \sqrt{R^2 - (R - t)^2} = \sqrt{2Rt - t^2} \approx \sqrt{2Rt} \quad (t \ll R)$$

➤ For **bright** fringes:

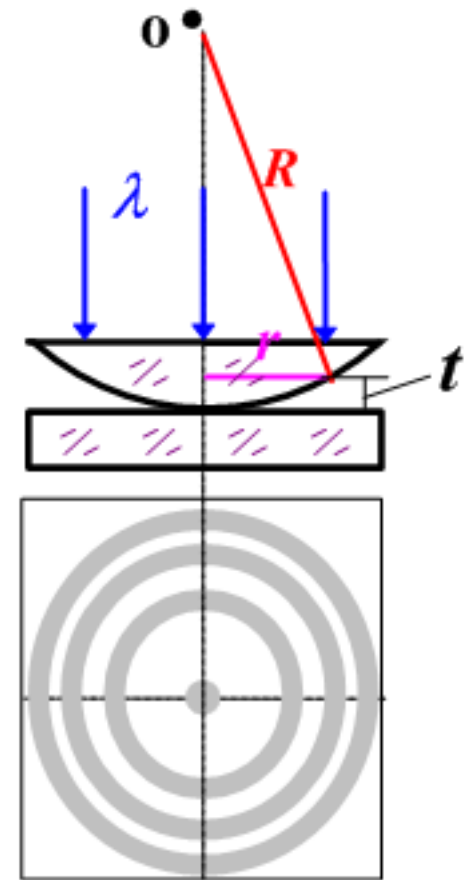
$$2t + \frac{\lambda}{2} = m\lambda, \quad m = 1, 2, 3, \dots$$

$$r_m = \sqrt{\left(m - \frac{1}{2}\right)\lambda R}, \quad m = 1, 2, 3, \dots$$

➤ For **dark** fringes:

$$2t + \frac{\lambda}{2} = (2m + 1)\frac{\lambda}{2}, \quad m = 0, 1, 2, \dots$$

$$r_m = \sqrt{m\lambda R}, \quad m = 0, 1, 2, \dots$$



Example

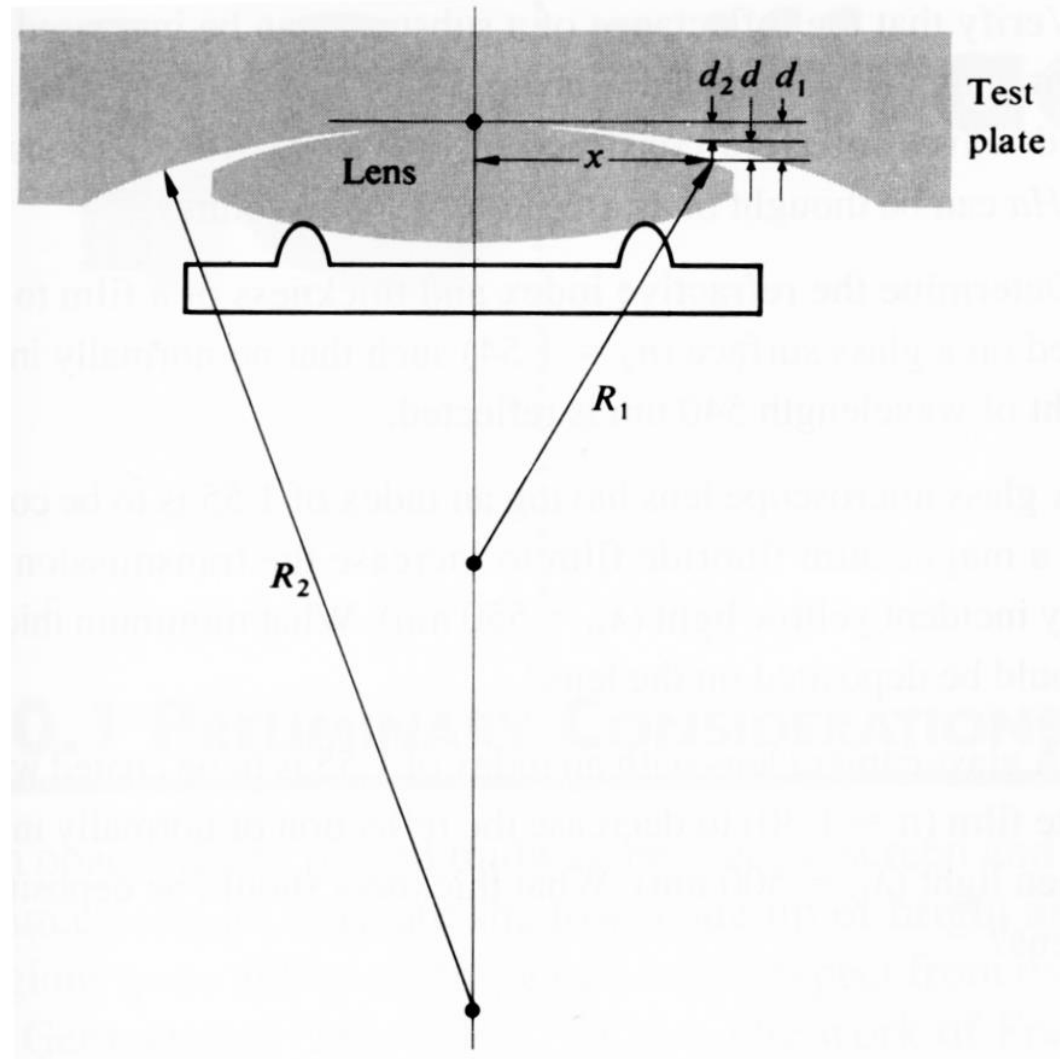


The figure illustrates a setup used for testing lenses.
Show that:

$$d \approx \frac{x^2 (R_2 - R_1)}{2R_1 R_2}$$

Prove that the radius of m -th **dark** fringe is then

$$x_m = \sqrt{\frac{R_1 R_2 m \lambda}{(R_2 - R_1)}}$$



Example



Solution:

$$x \approx \sqrt{2R_1d_1} \approx \sqrt{2R_2d_2}$$

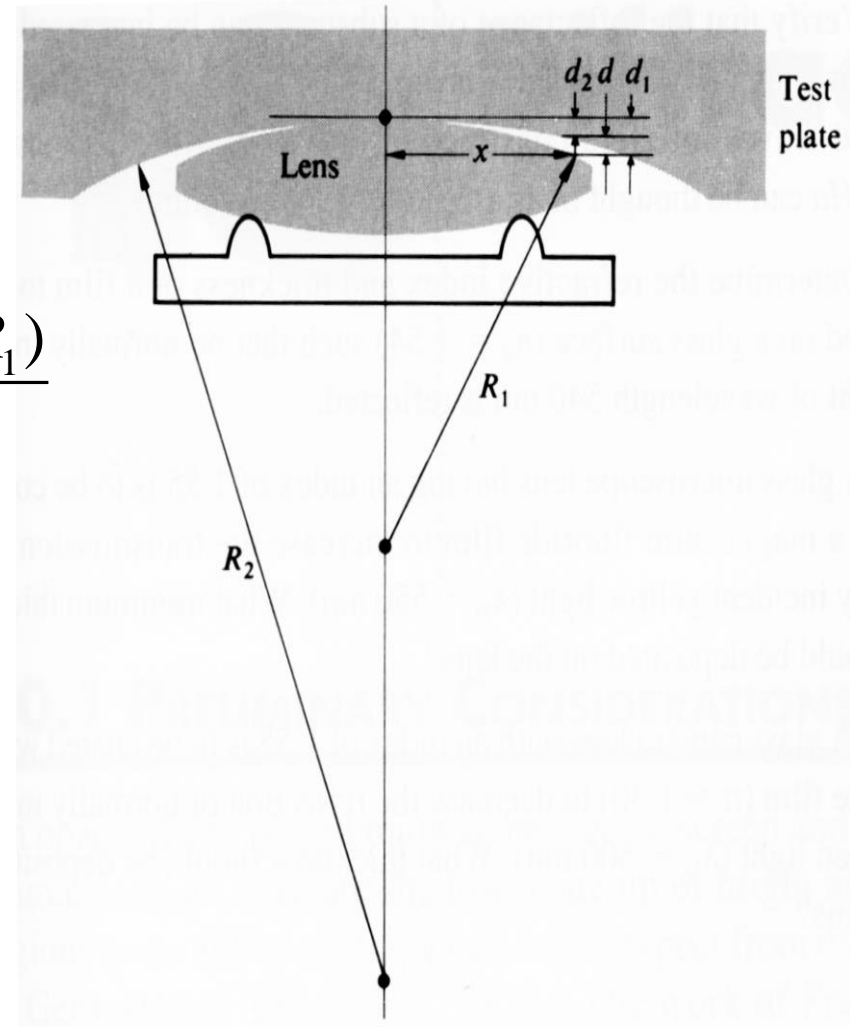
$$d_1 = \frac{x^2}{2R_1}, \quad d_2 = \frac{x^2}{2R_2}$$

$$d = d_1 - d_2 = \frac{x^2}{2} \left(\frac{1}{R_1} - \frac{1}{R_2} \right) = \frac{x^2(R_2 - R_1)}{2R_1R_2}$$

For dark fringes:

$$2d + \frac{\lambda}{2} = (2m+1)\frac{\lambda}{2} \Rightarrow d = m\frac{\lambda}{2}$$

$$x_m = \sqrt{\frac{2R_1R_2d}{(R_2 - R_1)}} = \sqrt{\frac{R_1R_2m\lambda}{(R_2 - R_1)}}$$



Thin non-reflecting coating



■ Thin non-reflecting coating (减反镀膜)

- ➡ A glass surface reflect about 4% of light incident on it. Lenses are often coated with thin films of transparent substance such as MgF_2 ($n_2=1.38$) to reduce the reflection from the glass surface.

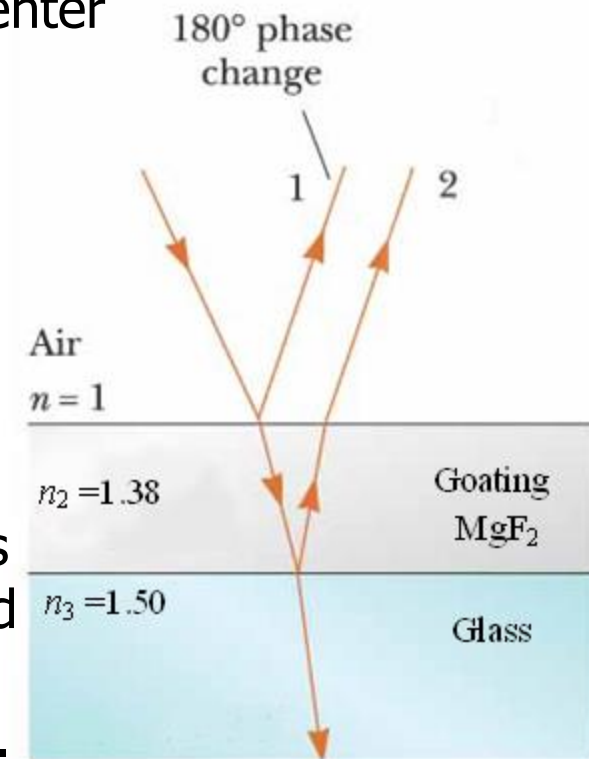
- Often the coating is designed to eliminate the center of reflected spectrum, $\lambda=550\text{nm}$.

- ➡ For MgF_2 , $n_2=1.38$. For glass $n_3=1.50$.

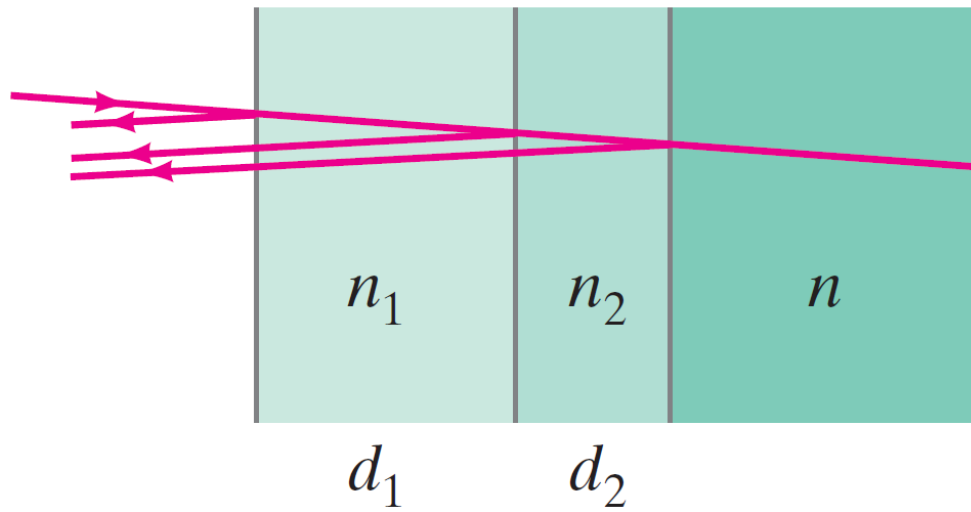
$$\delta = 2n_2t = (2m+1)\frac{\lambda}{2}, \quad m=0$$

- ➡ The thickness of coating is: $t = \frac{\lambda}{4n_2}$

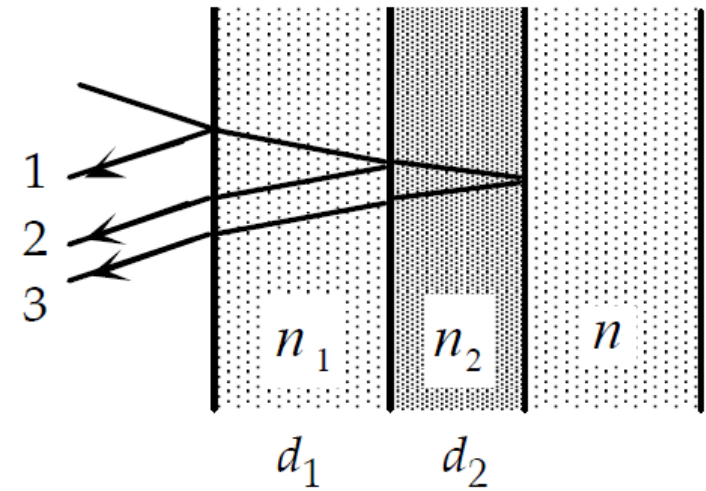
- ➡ In reflected light, yellow light (around 550nm) is reduced, but two extremes of spectrum — red and violet — will not be reduced as much, so such coated lenses is purple (mixture of red and violet).



Very highly reflective mirrors



$$n_1 < n_2 < n$$



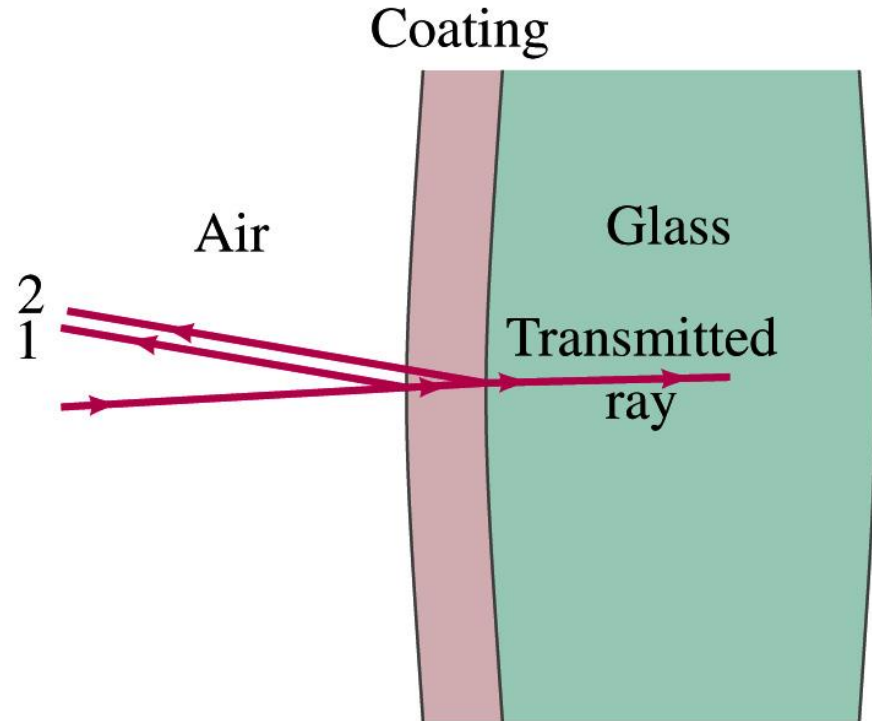
$$d_1 = \frac{\lambda}{2n_1}, d_2 = \frac{\lambda}{2n_2}$$

Non-reflective coating

What is the **smallest thickness** of an optical coating of MgF_2 , whose index of reflection is $n_2=1.38$, which is designed to eliminate reflected light at wavelengths centered at **550 nm** when incident normally on glass for which $n_3=1.50$?

Solution:

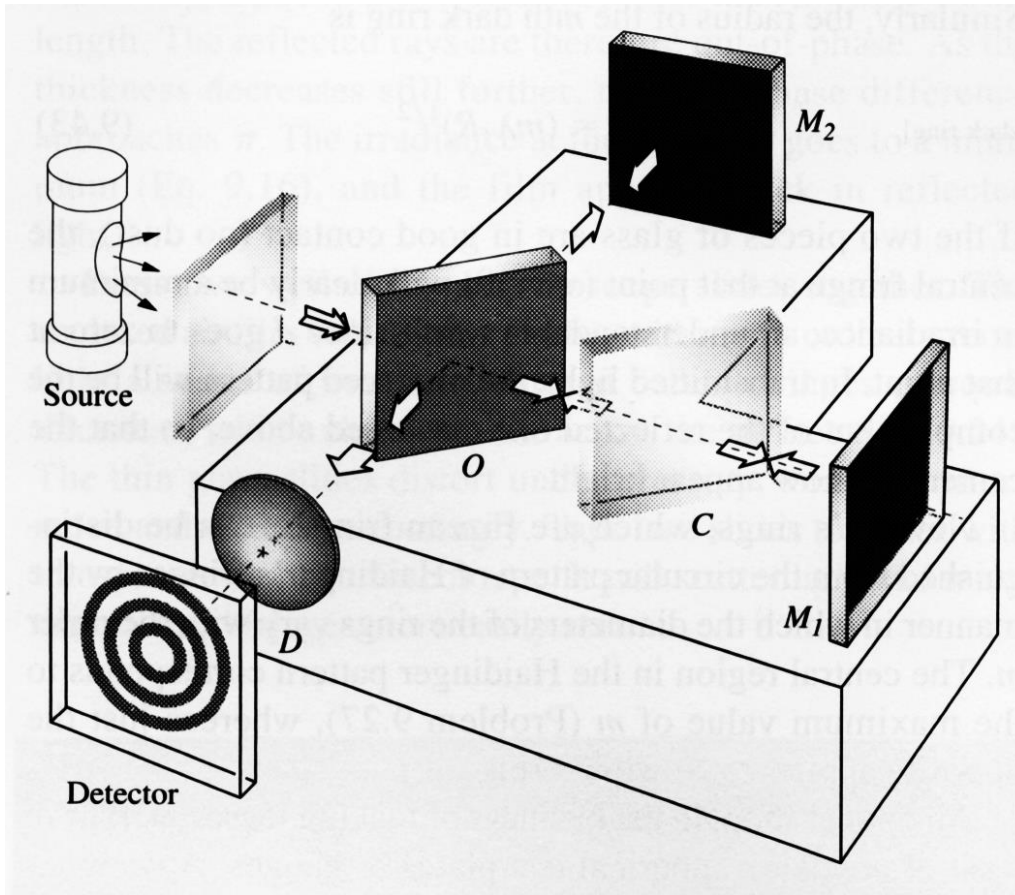
$$t = \frac{\lambda}{4n_2} = \frac{550}{4 \times 1.38} = 99.6 \text{ nm}$$



P699, Prob. 22, 26 (only the radius of
curvature of the lens surface)

P700, Prob. 44, 46

§ 5 Michelson Interferometer



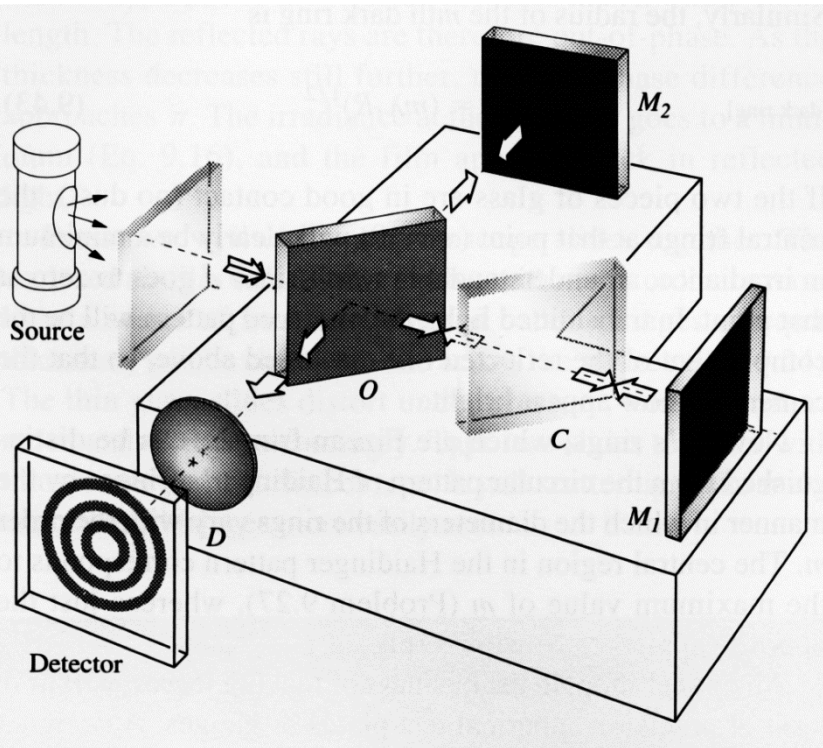
➡ An extended **source** emits a wave. The **beam splitter** with a half-silvered mirror divides the wave into two, one segment traveling to the mirror **M_1** and one to the mirror **M_2** . The two reflected wave are united at the region where the eye locates, and the **interference** can be expected.

➡ The role of **compensator plate**: one beam passes through beam splitter **three** times, whereas the other traverses it only **once**. Each beam will pass through equal thicknesses of glass only when a compensator plate is inserted in the arm 1.

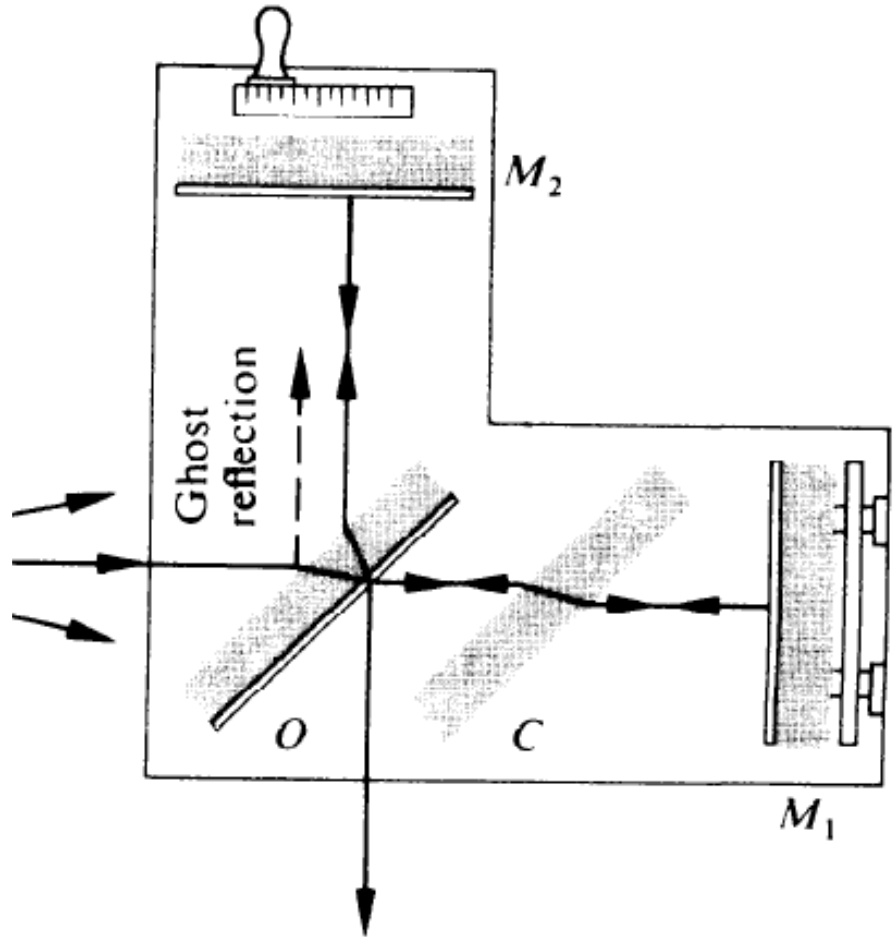
Michelson Interferometer



Top view



Circular fringes are centered on the lens.



Michelson Interferometer



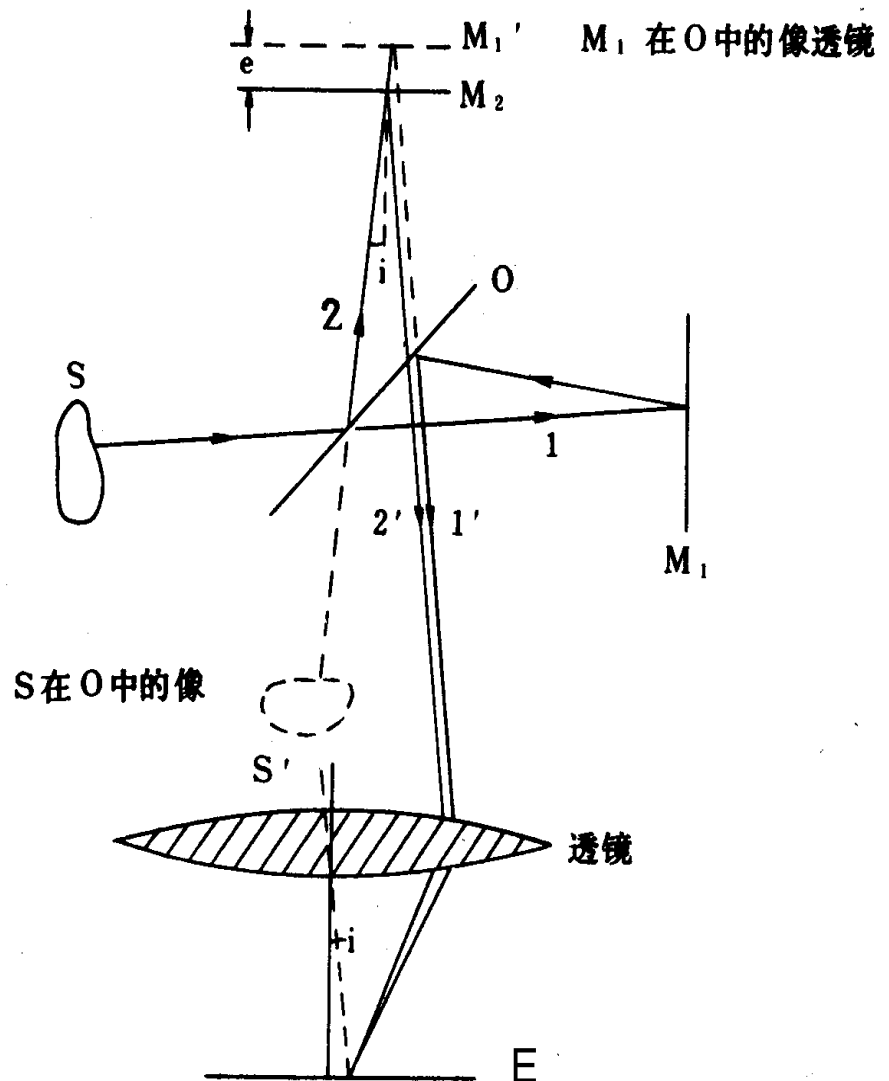
■ To understand how fringes are formed

➡ Re-draw the interferometer as if all the elements were in a straight line. The fringes are formed just like the **thin film** with two surface M_2 and M'_1 (the image of mirror M_1)

➡ The optical path length difference:

$$\delta = 2(L_2 - L_1)$$

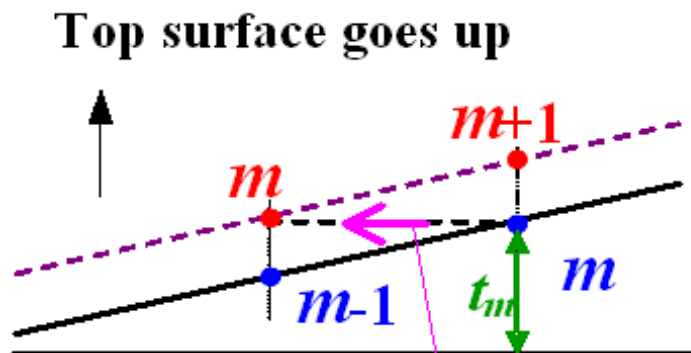
$$i \approx 0$$



The application of Michelson interferometer



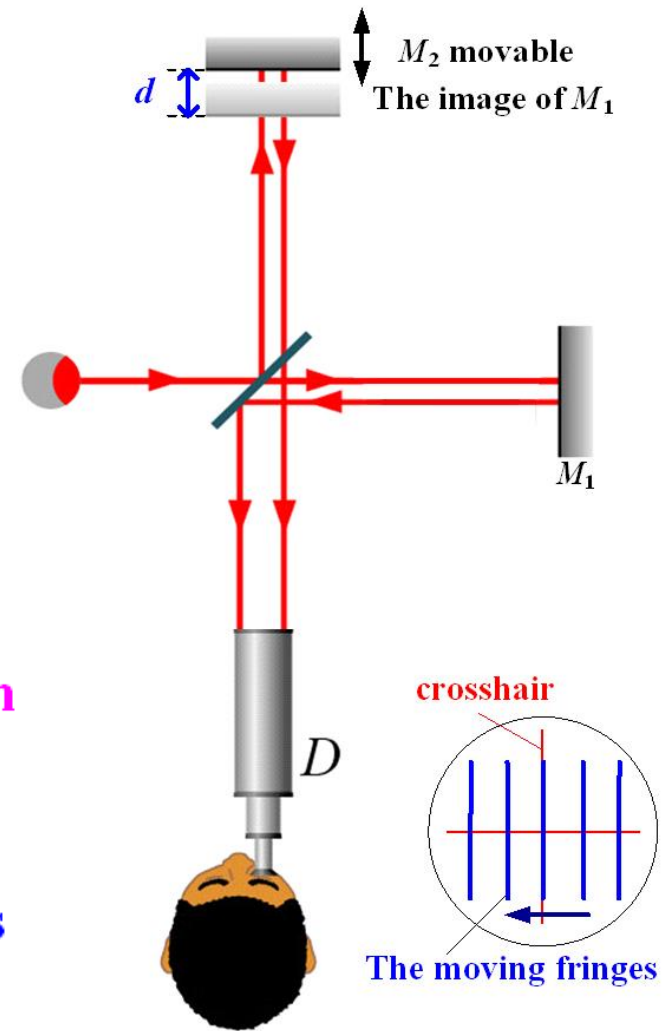
- Very precise length measurement
- ➡ For a thin film, if the top surface shifts up a distance $\lambda/2$, the fringes will move to the left with a distance of one fringe spacing. If we move M_2 slowly either backward or forward a distance $\lambda/2$, each fringe moves to the left or right a distance equal to **one fringe spacing**.



The direction of fringes shifting

The fringe position after shifting

The original position of fringes

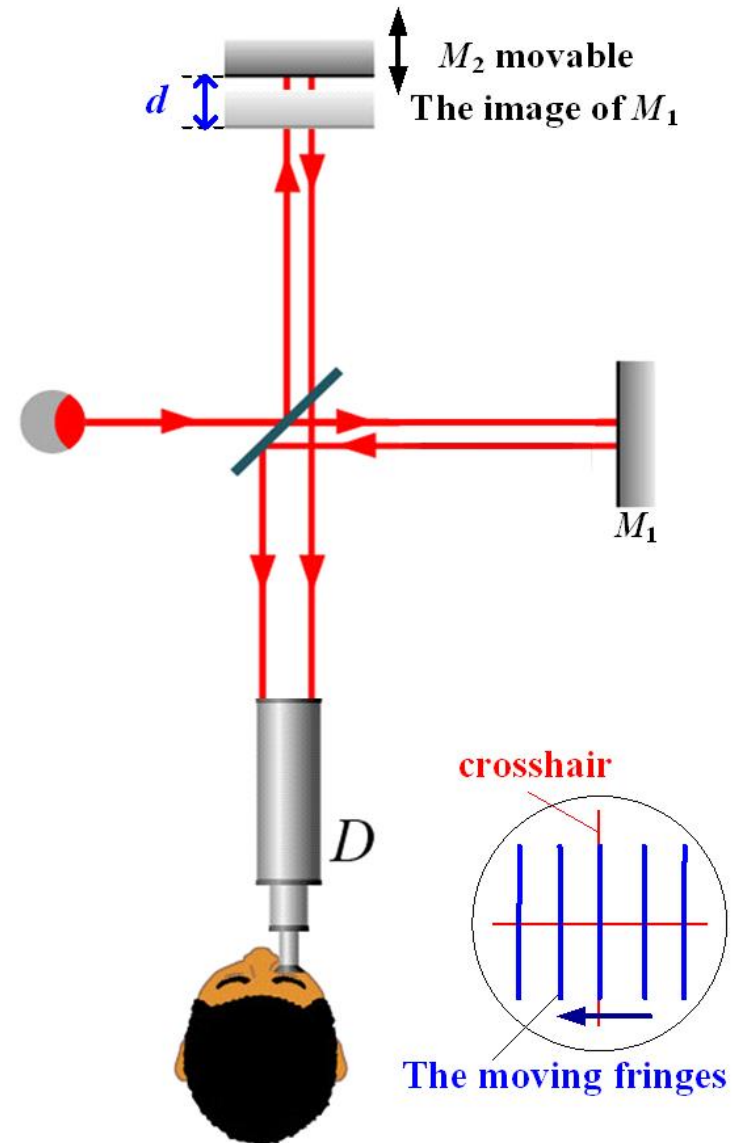


The application of Michelson interferometer



- ➡ If we observe the fringes positions through a telescope with a crosshair eyepieces and N fringes cross the crosshair when we move the mirror M_2 a distance d

$$d = N \frac{\lambda}{2}$$

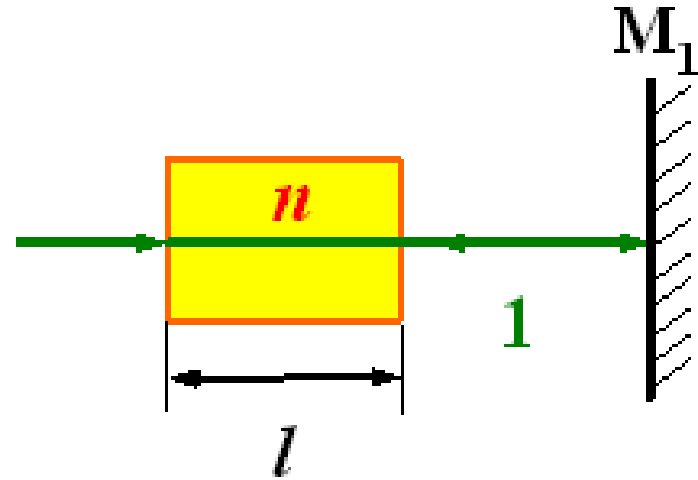


■ Measurement of index of an object

➡ We can insert an transparent object with index of refraction n in one arm, we observe N fringes cross the crosshair.

➡ The difference of optical path length will change

$$\Delta\delta = 2(nl - l) = N\lambda,$$



$$n = 1 + \frac{N\lambda}{2l}$$

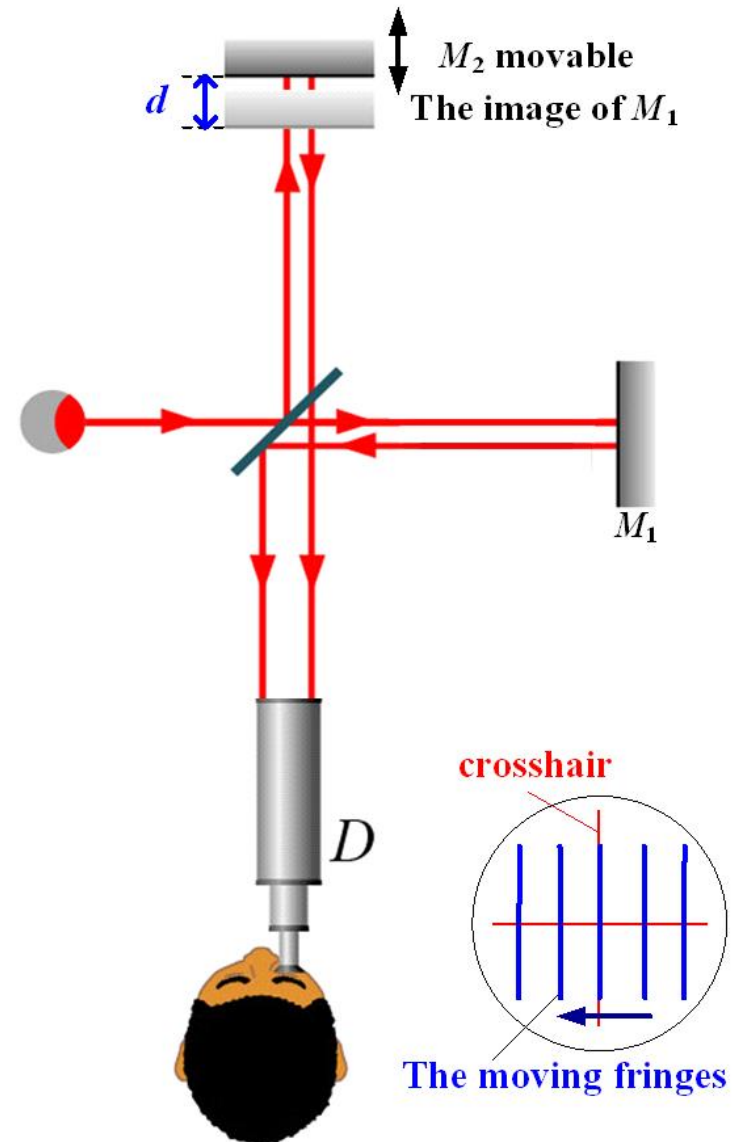
What is the **wavelength** of the light entering an interferometer if **344** bright fringes are counted when the movable mirror moves **0.125 mm**?

Solution:

$$d = N \frac{\lambda}{2}$$

$$\lambda = \frac{2d}{N}$$

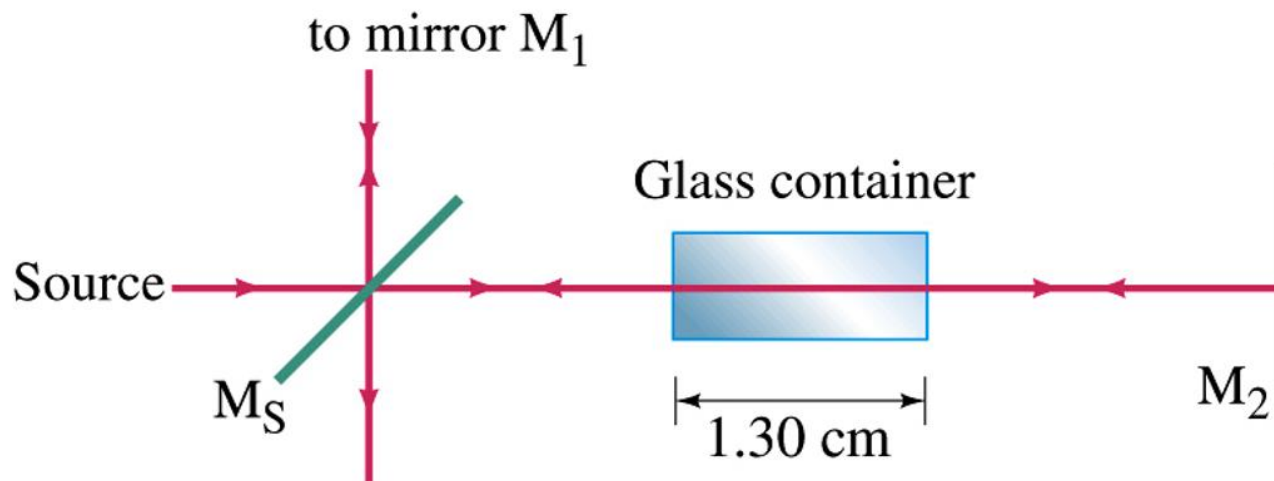
$$= \frac{2 \times 0.125 \times 10^{-3}}{344} = 727 \text{ nm}$$



Prob. 32 (P699)



One of the beams of an interferometer passes through a small glass container containing a cavity **1.30 cm** deep. When a gas is allowed to slowly fill the container, a total of **186** dark fringes are counted to move past a reference line. The light used has a wavelength of **610 nm**. Calculate the **index of refraction** of the gas at its final density, assuming that the interferometer is in vacuum.



P699, Prob.32

P701, Prob.51