

Review

- **Electric charge**
 - **positive, negative**
 - **Quantized:** $e = 1.6 \times 10^{-19} \text{ C}$
- **Unlike** charges attract; **like** charges repel.
- **Law of conservation** of electric charge
 - **The net amount of electric charge produced in any process is zero.**

§ 1 Coulomb's Law

- The electrostatic force on **point charge** q_2 **due to** q_1 , written F_{21} , can be expressed in vector form as (inverse square law):

$$\vec{F}_{21} = k_e \frac{q_1 q_2}{r^2} \hat{r}_{12} = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r^2} \hat{r}_{12}$$

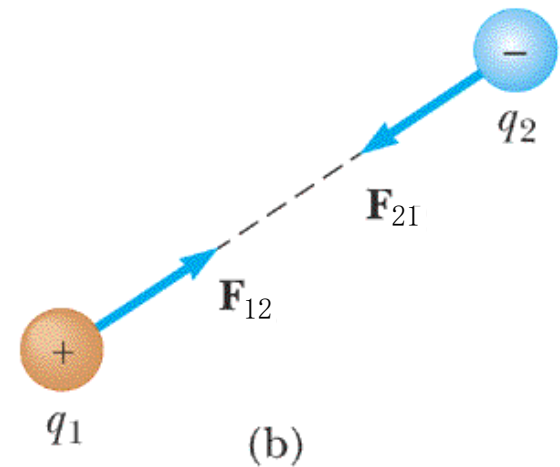
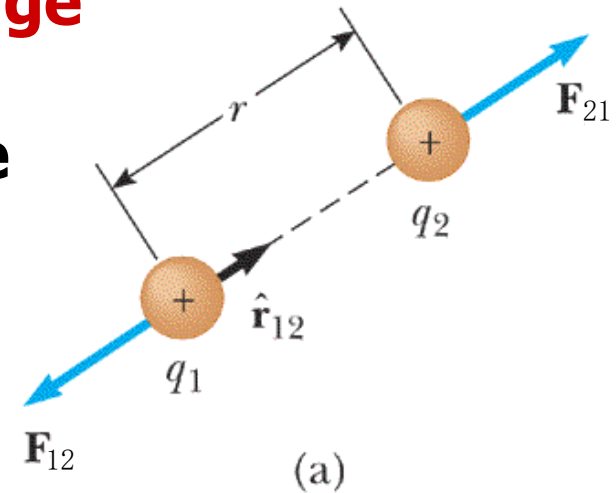
$$k_e = 8.99 \times 10^9 \text{ N} \cdot \text{m}^2 / \text{C}^2$$

—— **Coulomb constant**

$$\epsilon_0 = 8.8542 \times 10^{-12} \text{ C}^2 / \text{N} \cdot \text{m}^2$$

—— **permittivity of free space**
(electric constant)

\hat{r}_{12} is a unit vector directed from q_1 **toward** q_2



Coulomb's Law



$$\vec{F}_{21} = k_e \frac{q_1 q_2}{r^2} \hat{r}_{12} = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r^2} \hat{r}_{12}$$

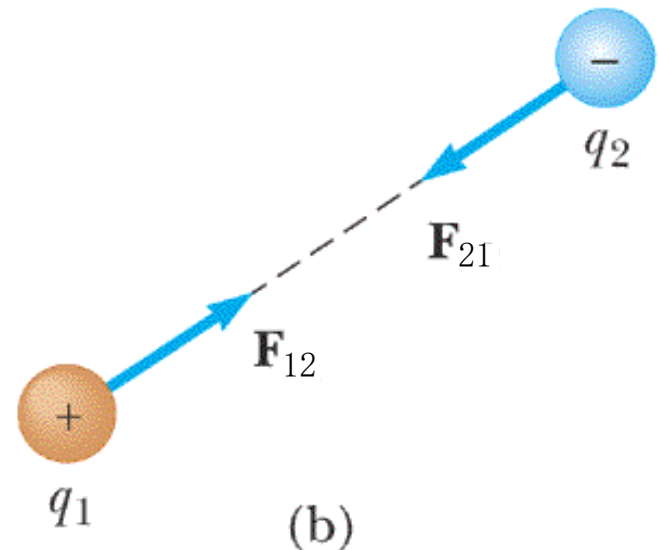
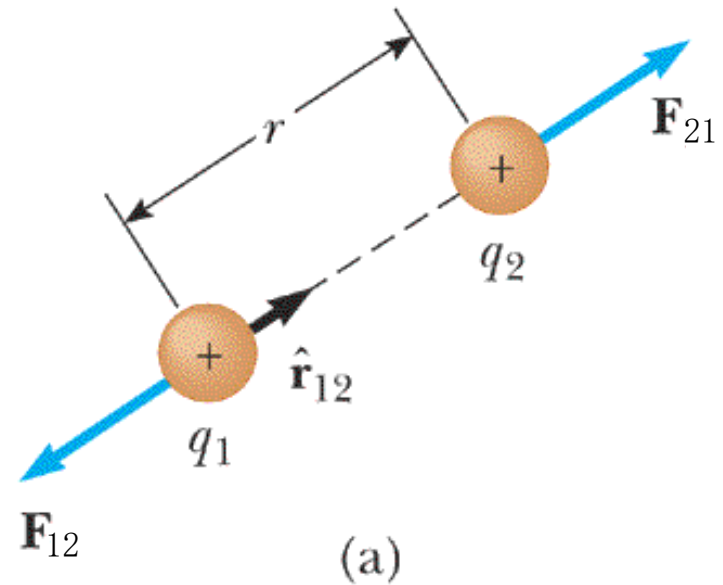
$$k_e = 8.99 \times 10^9 \text{ N} \cdot \text{m}^2 / \text{C}^2$$

$$\epsilon_0 = 8.8542 \times 10^{-12} \text{ C}^2 / \text{N} \cdot \text{m}^2$$

➡ **From Newton's third law**

$$\vec{F}_{21} = -\vec{F}_{12}$$

➡ **The Coulomb forces between the two charges having the same sign are **repulsive**, while the two charges with opposite sign result in **attractive** Coulomb forces.**



Coulomb's Law vs. Newton's law of gravitation



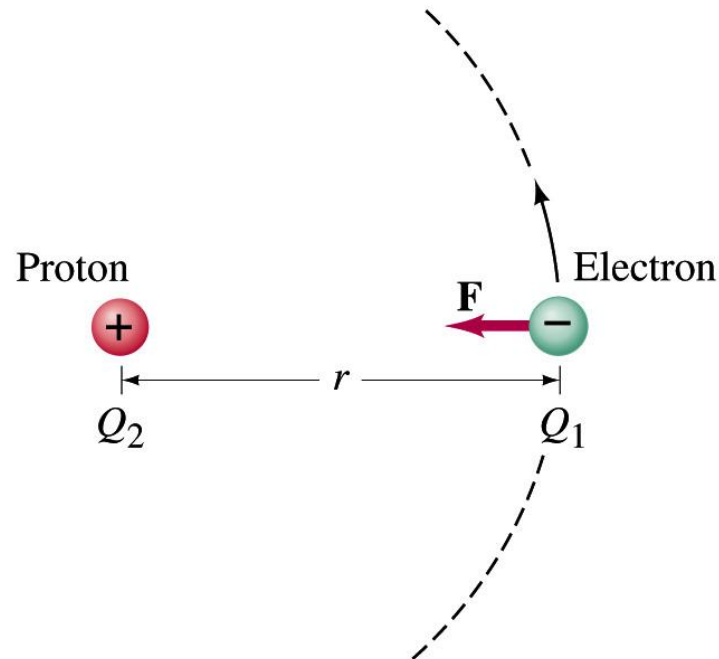
$$\vec{F}_{21} = k_e \frac{q_1 q_2}{r^2} \hat{r}_{12}, \quad \vec{F}_{21} = -G \frac{m_1 m_2}{r^2} \hat{r}_{12}$$

- ➡ Both are **inverse square** laws, and charge q plays the same role in Coulomb's law as that the mass m plays in Newton's law of gravitation.
- ➡ One difference between the two laws is that gravitational forces are always attractive, whereas electrostatic forces can be either **repulsive** or **attractive**.

Coulomb's Law vs. Newton's law of gravitation



Ex.19-1 (P462): the two forces in the hydrogen atom



$$r = 5.3 \times 10^{-11} \text{ m}$$

Particle	Charge (C)	Mass (kg)
Electron (e)	$-1.602\,191\,7 \times 10^{-19}$	$9.109\,5 \times 10^{-31}$
Proton (p)	$+1.602\,191\,7 \times 10^{-19}$	$1.672\,61 \times 10^{-27}$
Neutron (n)	0	$1.674\,92 \times 10^{-27}$

Coulomb's Law vs. Newton's law of gravitation



The electrostatic force:

$$F_e = k_e \frac{e^2}{r^2} = (8.99 \times 10^9 \text{ N} \cdot \text{m}^2 / \text{C}^2) \frac{(1.60 \times 10^{-19} \text{ C})^2}{(5.3 \times 10^{-11} \text{ m})^2} = 8.2 \times 10^{-8} \text{ N}$$

The gravitational force:

$$F_g = G \frac{m_e m_p}{r^2} = (6.67 \times 10^{-11} \text{ N} \cdot \text{m}^2 / \text{kg}^2) \frac{(9.11 \times 10^{-31} \text{ kg})(1.67 \times 10^{-27} \text{ kg})}{(5.3 \times 10^{-11} \text{ m})^2} = 3.6 \times 10^{-47} \text{ N}$$

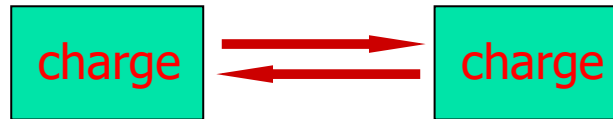
➤ **The gravitational force is weaker than the electrostatic force by factor of about 10^{-39} .**

§ 2 Electric fields

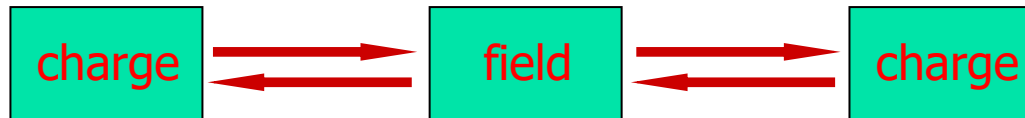


■ Does Coulomb's law mean that the interaction between separated charges is an **action-at-a-distance**? Is the interaction direct and instantaneous?

- ➡ Historical view: The interaction model for charges is the **action-at-a-distance**



- ➡ The real interaction model for charges—The interaction between two charges is realized through the **electric fields** established around the charges.

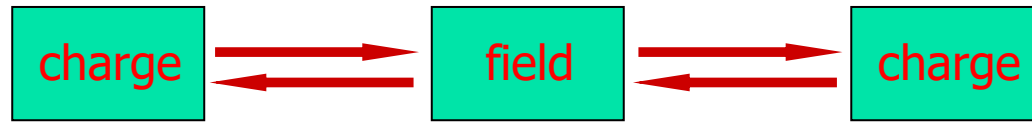


The first charge sets up an electric field, and the second charge interacts with the electric field of the first charge.

The Electric Field

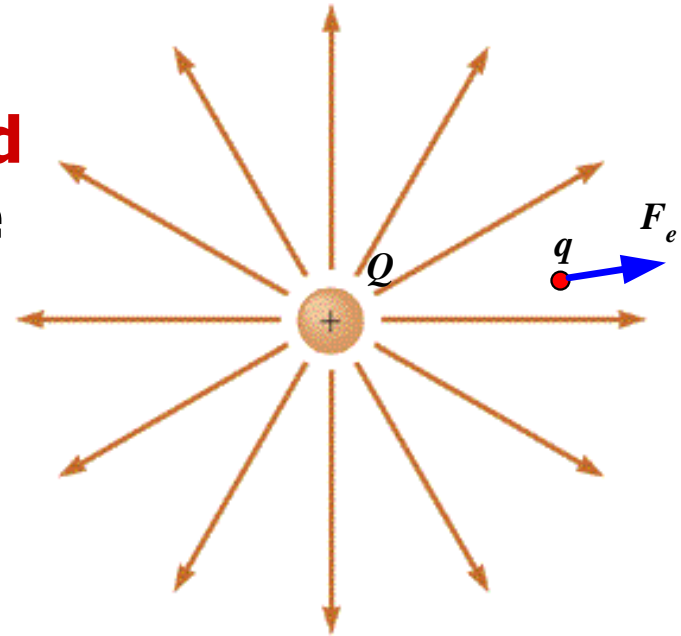


■ The problem of determining the interaction between the charges is therefore reduced to two separate problem



➤ Determine, by measurement or calculation, the **electric field** established by the first charge at every point in space.

➤ Calculate the **force** that the field exerts on the second charge placed at a particular point in space.



The Electric Field and The Electric Force



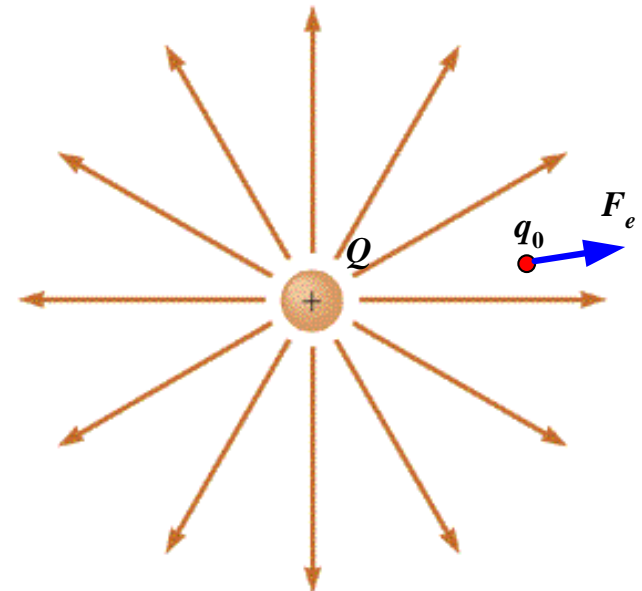
■ The definition of electric field

- ➡ The definition of the electric field \vec{E} in terms of the electric force \vec{F}_e exerted on a **positive** test charge q_0 placed at a particular point.

$$\vec{E} \equiv \frac{\vec{F}_e}{q_0} \quad \text{SI unit: N/C or V/m}$$

The direction of \vec{E} is the same as the direction of \vec{F}_e .

The **test particle** q_0 is used only to detect the existence of the field and evaluate its strength. The existence and strength of the electric field is feature of electric field itself, **not** dependent on the q_0 .



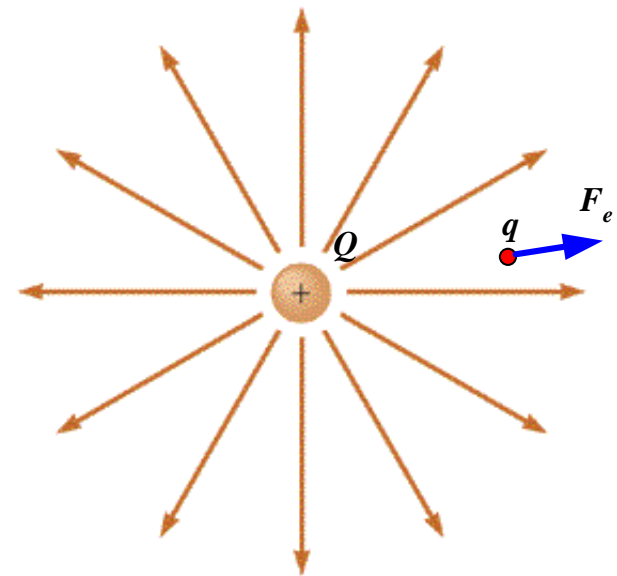
The Electric Field and the Electric Force



- The electric force exerted on a charge.
 - ➔ Once the electric field is known at some point, the electric force on any particle with **charge** q placed at that point can be calculated by

$$\vec{F}_e = q\vec{E}$$

- ➔ Here the electric field \vec{E} is caused by other charges that may be present, **not** by the charge q .



The electric field of point charge



The calculation of electric field due to the **individual** point charges

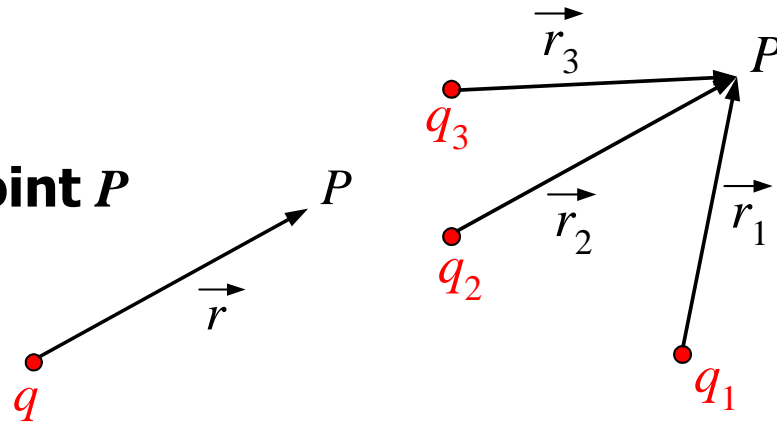
➔ The electric field due to **single** point charge

According to **Coulomb's law**, a test q_0 experience a electric force

$$\vec{F}_e = \frac{1}{4\pi\epsilon_0} \frac{q q_0}{r^2} \hat{r}$$

The electric field created by q at point P

$$\vec{E} \equiv \frac{\vec{F}_e}{q_0} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r}$$



➔ The electric field due to a **series** of point charges distributed in space.

According to the **superposition principle**, the total electric field at point P

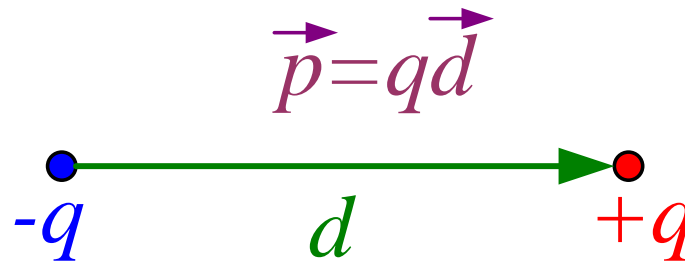
$$\vec{E} = \sum_i \vec{E}_i = \frac{1}{4\pi\epsilon_0} \sum_i \frac{q_i}{r_i^2} \hat{r}_i$$

Example — **The Electric Dipole** (P475 § 19-11)

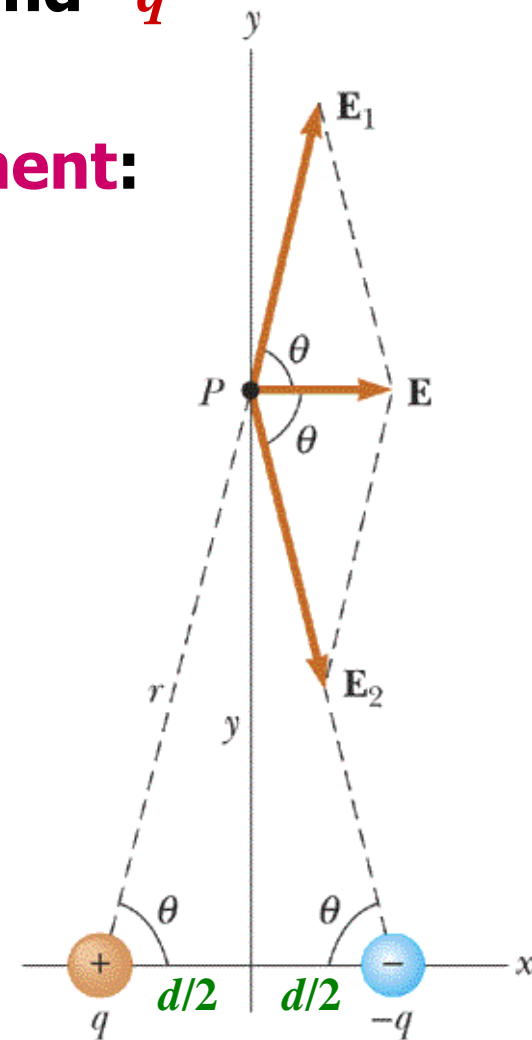


Equal positive and negative charges $+q$ and $-q$ separated by a fixed distance d .

► Definition of the **electric dipole moment**:



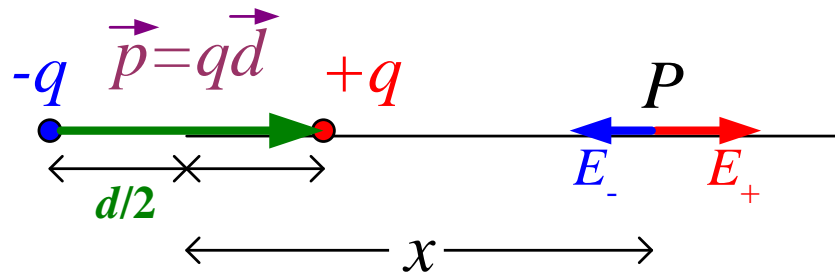
- (a) Find the electric field E due to the dipole along the x axis at the point P ;
- (b) Find the electric field E due to the dipole along the y axis at the point P .



Example Cont'd



Solution: (a)



$$E = E_+ - E_-$$

$$= \frac{q}{4\pi\epsilon_0} \left[\frac{1}{(x - d/2)^2} - \frac{1}{(x + d/2)^2} \right]$$

$$= \frac{q}{4\pi\epsilon_0} \frac{2xd}{(x^2 - d^2/4)^2}$$

$$= \frac{2xp}{4\pi\epsilon_0 (x^2 - d^2/4)^2}, \quad x \gg d/2, \quad E = \frac{1}{2\pi\epsilon_0} \frac{p}{x^3} \propto \frac{p}{x^3}$$

Example Cont'd

Solution: (b)

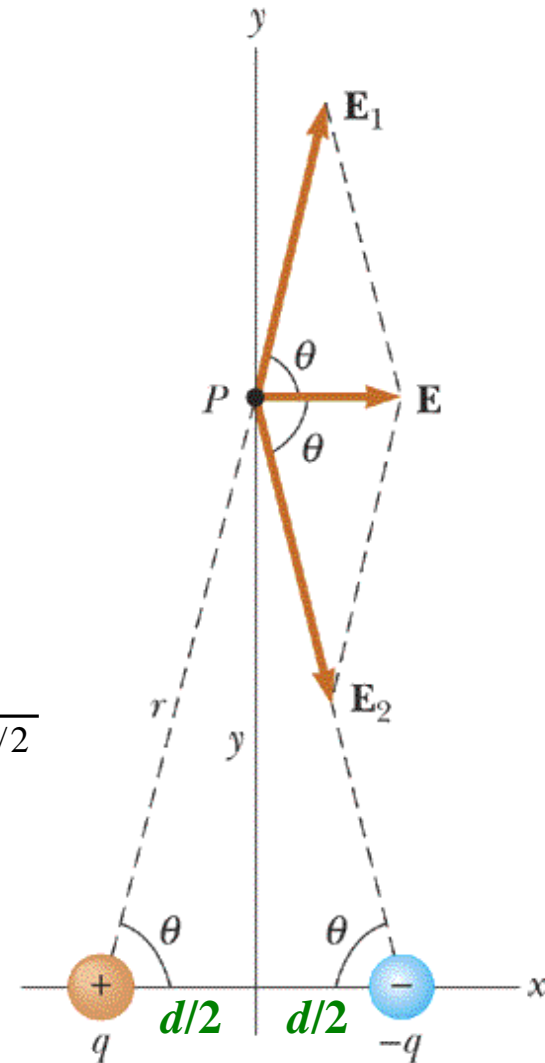
$$\vec{E} = \vec{E}_+ + \vec{E}_-$$

$$E_+ = E_- = \frac{1}{4\pi\epsilon_0} \frac{q}{y^2 + (d/2)^2}$$

$$E = 2E_+ \cos \theta = \frac{2}{4\pi\epsilon_0} \frac{q}{y^2 + (d/2)^2} \frac{(d/2)}{\sqrt{y^2 + (d/2)^2}}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{qd}{[y^2 + (d/2)^2]^{3/2}} = \frac{1}{4\pi\epsilon_0} \frac{p}{[y^2 + (d/2)^2]^{3/2}}$$

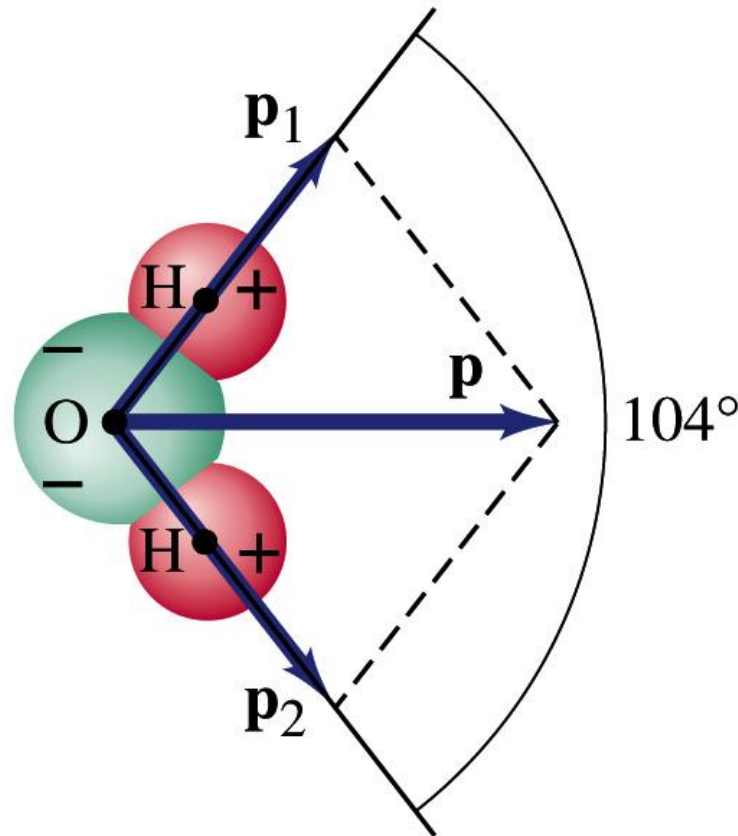
$$y \gg d/2, \quad E = \frac{1}{4\pi\epsilon_0} \frac{p}{y^3} \propto \frac{p}{y^3}$$



Prob. 43 (P521)



The dipole moment of a water molecule



The **water molecule** has a permanent polarization resulting from its nonlinear geometry. We can model the water molecule and other polar molecules as **dipoles**.

The torque on an electric dipole



Example: Find the torque exerted on an electric dipole.

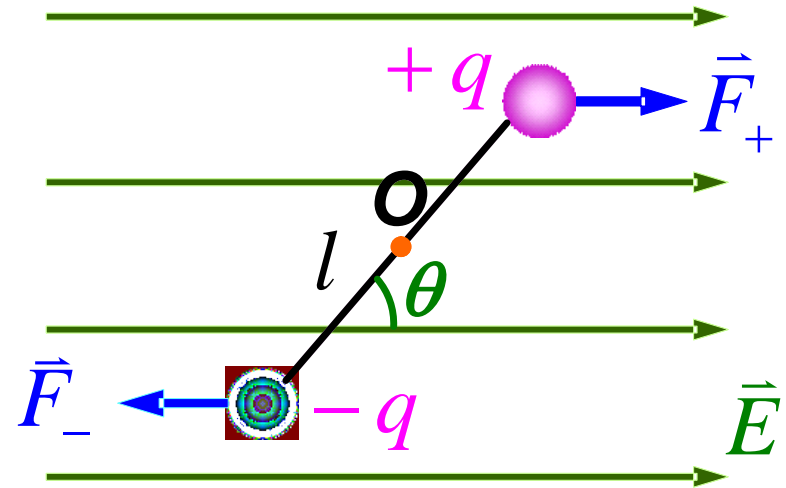
Solution: to the origin point O

The magnitude:

$$\begin{aligned}\tau &= \frac{l}{2} F_+ \sin \theta + \frac{l}{2} F_- \sin \theta \\ &= qEl \sin \theta = pE \sin \theta\end{aligned}$$

The **vector** description:

$$\vec{\tau} = q\vec{l} \times \vec{E} = \vec{p} \times \vec{E}$$



The effect of this torque is to try to turn the dipole so \vec{p} is **parallel** to \vec{E} .

The electric field due to continuous charge distributions (P468 § 19-7)

- The calculation method for electric field due to continuous charge distributions

The procedure:

- Divide the charge distribution into small elements dq ;
- Model each element as a point charge;

$$d\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{dq}{r^2} \hat{r}$$

- Apply the superposition principle to get the total field at P

$$\vec{E} = \int d\vec{E} = \frac{1}{4\pi\epsilon_0} \int \frac{dq}{r^2} \hat{r}$$

- Charge distribution manners

The linear charge density:

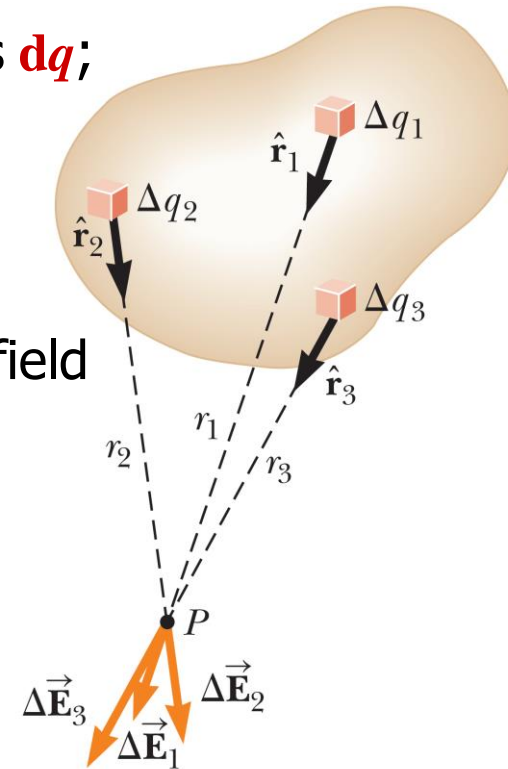
$$dq = \lambda dl$$

The surface charge density:

$$dq = \sigma dA$$

The volume charge density:

$$dq = \rho dV$$

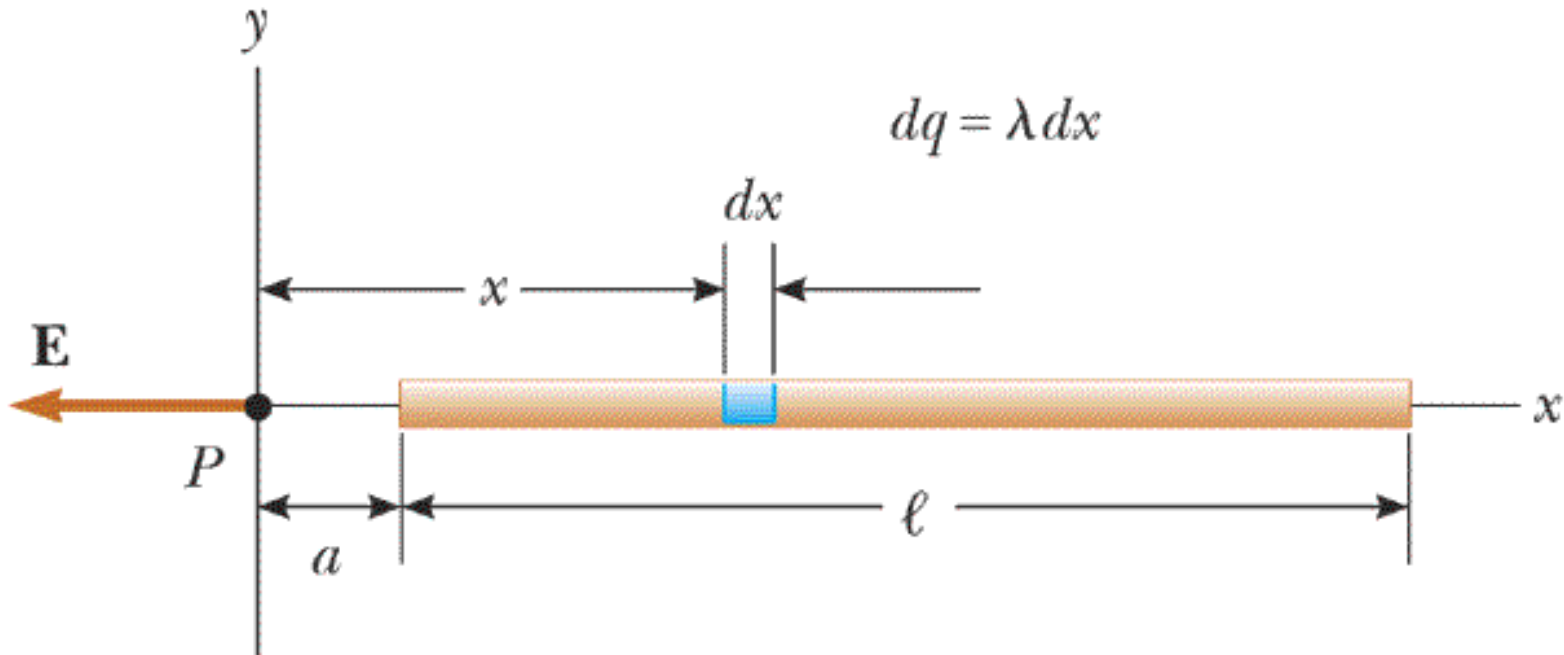


Example

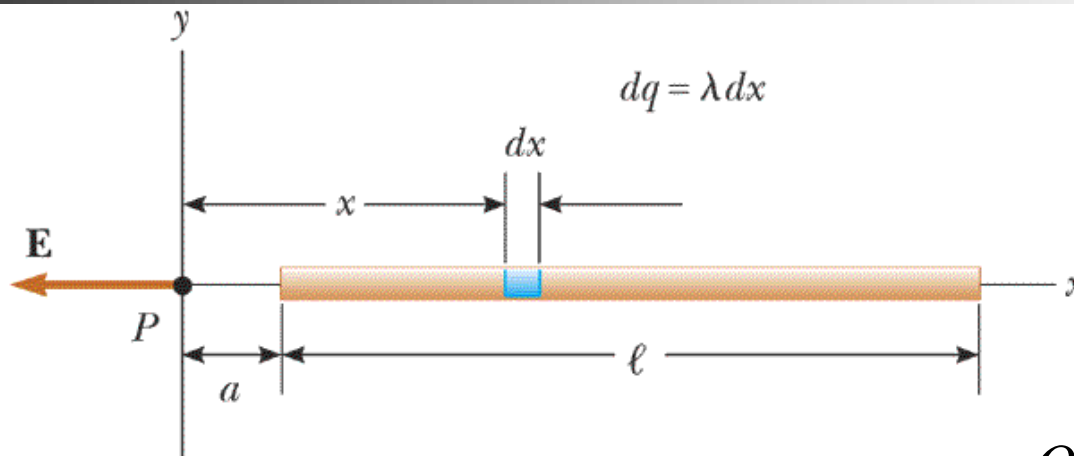


The electric field due to a charged rod

A rod of length ℓ has a uniform linear charge density and a total charge Q . Calculate the **electric field** at a point P along the axis of the rod, a distance a from one end.



Example



Solution:

Step 1: Choose the segment dq . $dq = \lambda dx = \frac{Q}{l} dx$

Step 2: Write the expression of \vec{E} due to dq .

$$dE = -\frac{1}{4\pi\epsilon_0} \frac{dq}{x^2} = -\frac{1}{4\pi\epsilon_0} \frac{\lambda dx}{x^2}$$

Step 3: Obtain the total field \vec{E} by integration.

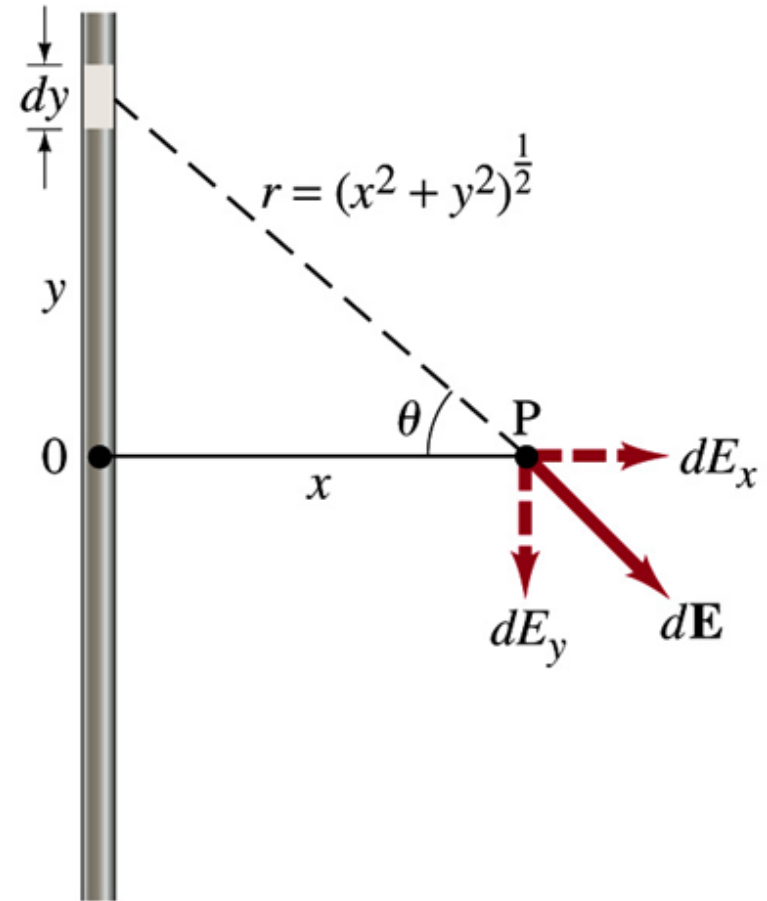
$$\begin{aligned} E &= -\frac{\lambda}{4\pi\epsilon_0} \int_a^{a+l} \frac{dx}{x^2} = \frac{\lambda}{4\pi\epsilon_0} \left[\frac{1}{x} \right]_a^{a+l} = -\frac{\lambda}{4\pi\epsilon_0} \left(\frac{1}{a} - \frac{1}{a+l} \right) \\ &= -\frac{Q}{4\pi\epsilon_0 a(a+l)} \end{aligned}$$

Example (P469 Ex.19-10)



The electric field of long line of charge

A **long** wire has a uniform linear charge density λ . Calculate the **electric field** at a point P a distance x from the wire.



Example

Solution:

Step 1: Choose the segment dq . $dq = \lambda dy$

Step 2: Write the expression of \vec{E} due to dq .

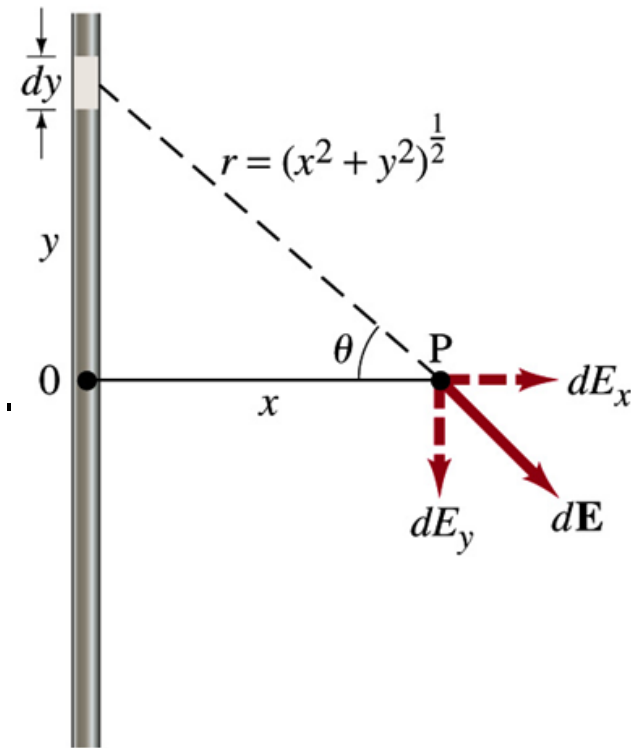
$$dE = \frac{1}{4\pi\epsilon_0} \frac{dq}{r^2} = \frac{1}{4\pi\epsilon_0} \frac{\lambda dy}{(x^2 + y^2)}$$

$$dE_x = dE \cos \theta, \quad dE_y = dE \sin \theta$$

Step 3: Obtain the total field \vec{E} by integration.
for the symmetry:

$$E_y = \int dE \sin \theta = 0$$

$$E_x = \int dE \cos \theta = \frac{\lambda}{4\pi\epsilon_0} \int \frac{\cos \theta dy}{x^2 + y^2}$$

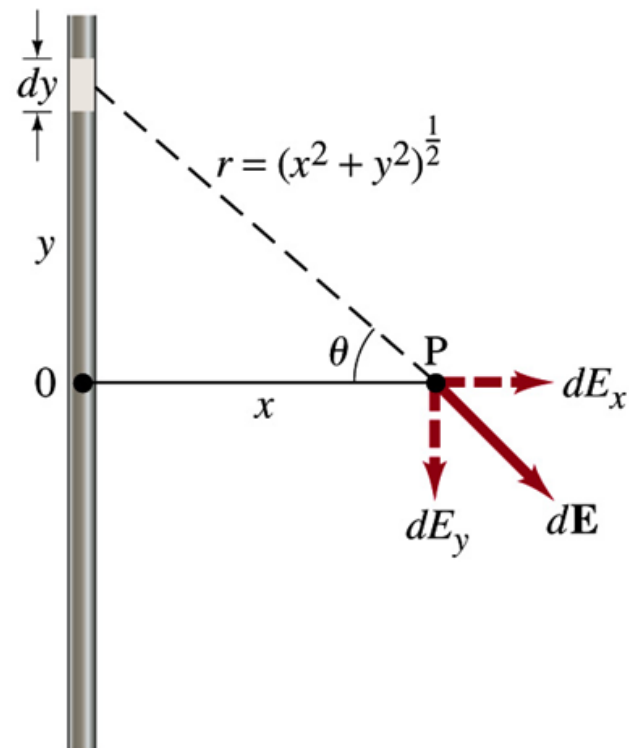


Example cont'd

$$E_x = \int dE \cos \theta = \frac{\lambda}{4\pi\epsilon_0} \int \frac{\cos \theta dy}{x^2 + y^2}$$

$$y = x \tan \theta, \quad dy = \frac{x d\theta}{\cos^2 \theta}, \quad x^2 + y^2 = \frac{x^2}{\cos^2 \theta}$$

$$\begin{aligned} E_x &= \frac{\lambda}{4\pi\epsilon_0} \frac{1}{x} \int_{-\pi/2}^{\pi/2} \cos \theta d\theta \\ &= \frac{\lambda}{4\pi\epsilon_0 x} (\sin \theta) \Big|_{-\pi/2}^{\pi/2} = \frac{1}{2\pi\epsilon_0} \frac{\lambda}{x} \end{aligned}$$

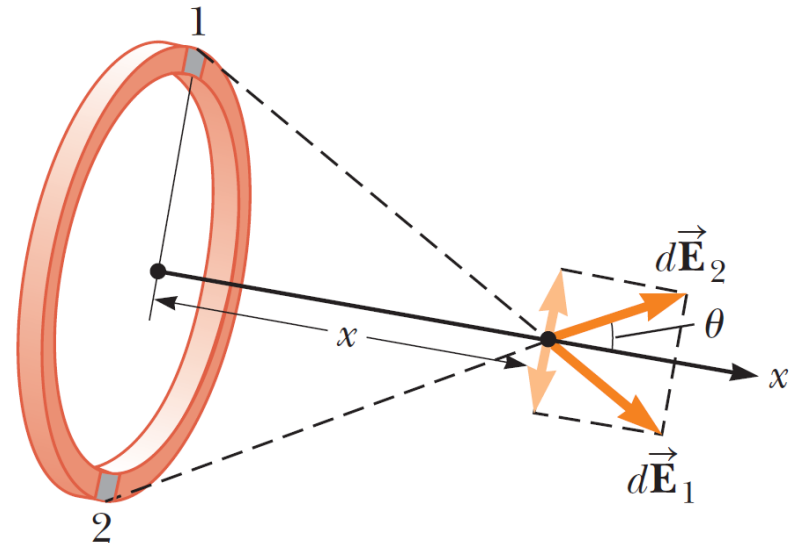
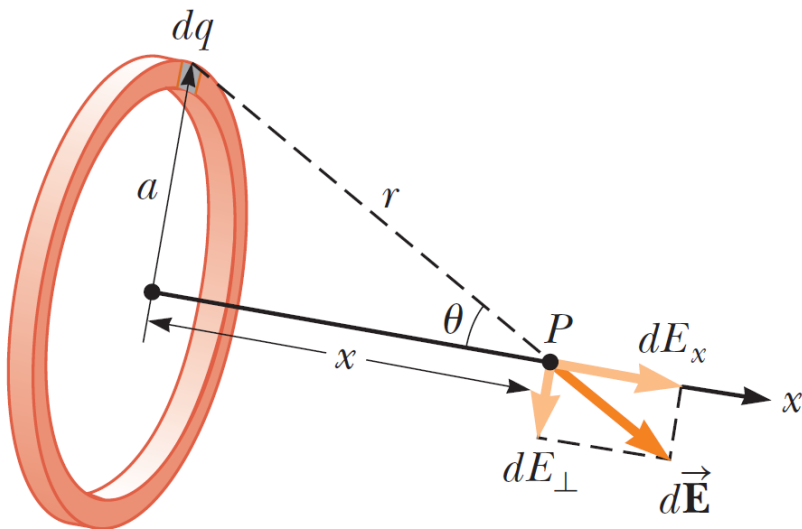


Example (P468 Ex.19-9)



The electric field of a uniform ring of charge

A ring of radius a has a uniform positive charge distribution, with a total charge Q . Calculate the **electric field** at a point P on the axis of the ring, at a distance x from the center of the ring.



Example

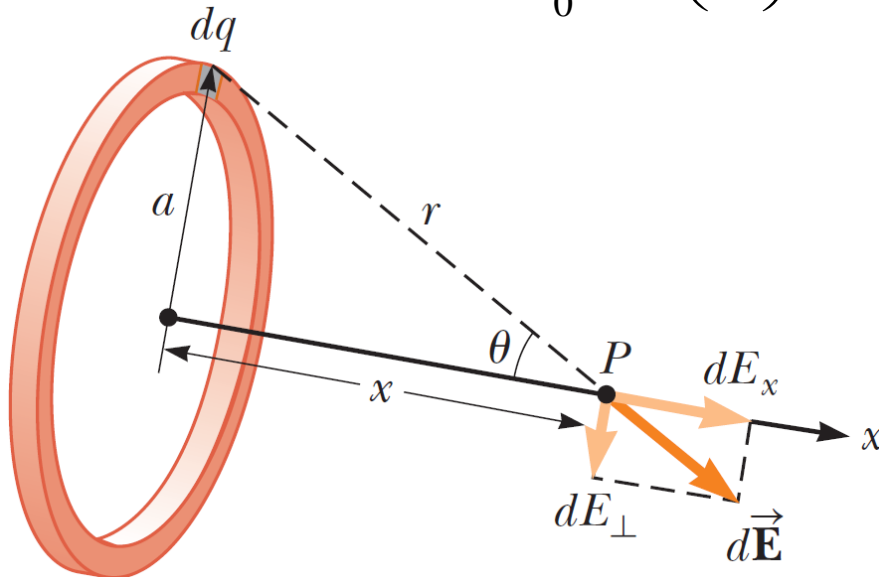


Solution: Choose the segment dq .

Write the expression of $d\vec{E}$ due to dq . $dE = \frac{1}{4\pi\epsilon_0} \frac{dq}{r^2}$

This field include the x component $dE_x = dE \cos \theta$, and perpendicular component dE_\perp , which is canceled by another dE_\perp on the opposite side of the ring.

$$dE_x = dE \cos \theta = \frac{1}{4\pi\epsilon_0} \frac{dq}{r^2} \left(\frac{x}{r} \right), \quad E_x = \frac{1}{4\pi\epsilon_0} \frac{x}{r^3} \int dq = \frac{x}{4\pi\epsilon_0 r^3} Q$$

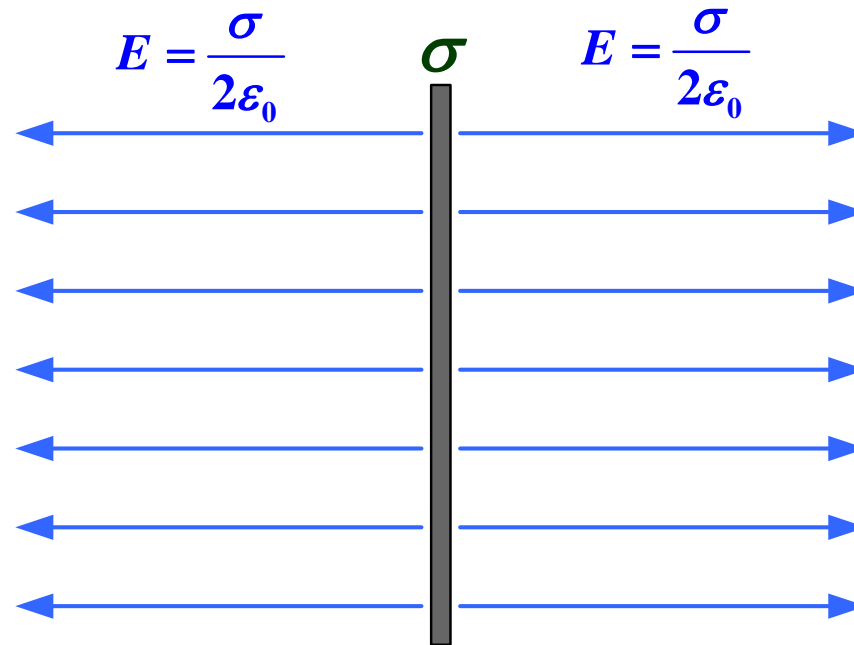


$$= \frac{Q}{4\pi\epsilon_0 r^2} \cos \theta$$

$$= \frac{x}{4\pi\epsilon_0 (x^2 + a^2)^{3/2}} Q$$

Example

The electric field of **an infinite plane** sheet of charge
An **infinite** plane sheet with uniform surface charge density σ . Calculate the **electric field** at a point P on the axis perpendicular to the plane, at a distance x from the plane.



Example



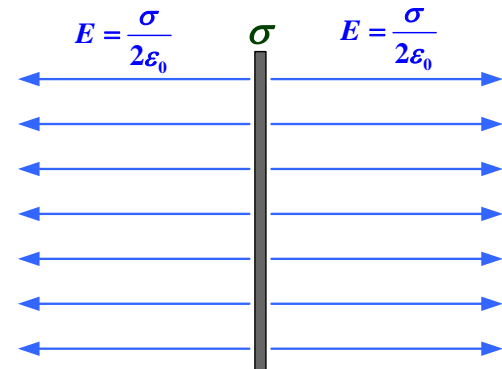
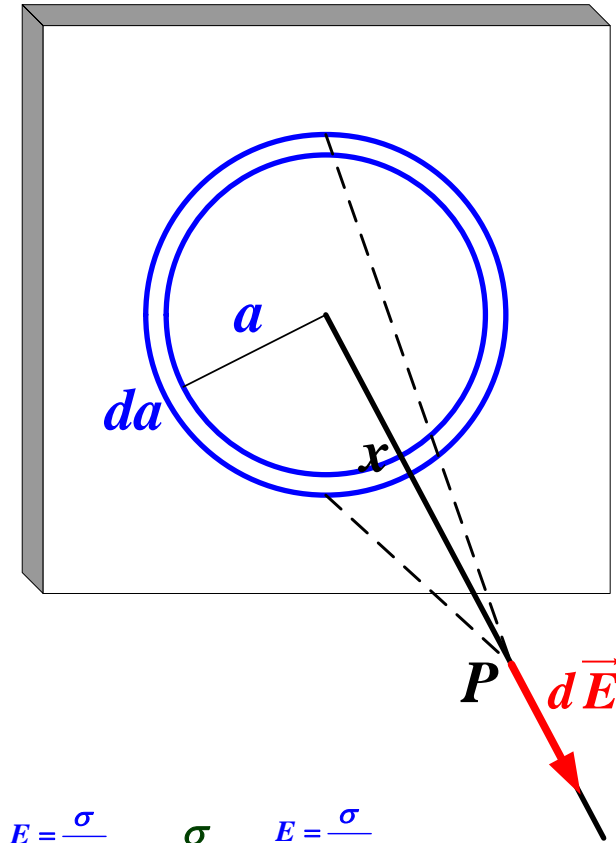
Solution I: Choose the **ring** with radius a and width da as **dq** . Starting from:

$$dE = \frac{x}{4\pi\epsilon_0 (x^2 + a^2)^{3/2}} dq, \quad dq = \sigma dA = \sigma (2\pi a) da$$

$$E = \int_0^\infty \frac{\sigma x a da}{2\epsilon_0 (x^2 + a^2)^{3/2}} = \frac{\sigma}{4\epsilon_0} \int_0^\infty \frac{x da (x^2 + a^2)}{(x^2 + a^2)^{3/2}}$$

$$= \frac{\sigma}{4\epsilon_0} \left[-2 \frac{x}{\sqrt{x^2 + a^2}} \right]_0^\infty = \frac{\sigma}{2\epsilon_0}$$

The electric field keeps constant at any distance from the plane — the field is uniform everywhere.



Example cont'd



Solution II: Choose the infinite lengthy **rod** as **dq** , which is at the distance **a** from the center, and width **da** . Using the conclusion of previous example:

$$E = \frac{1}{2\pi\epsilon_0} \frac{\lambda}{r}$$

where **r** is the distance from the rod along its perpendicular bisector.

$$\lambda = \sigma da$$

$$dE = \frac{\sigma}{2\pi\epsilon_0} \frac{da}{r}$$



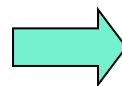
$$dE_x = 2(dE) \cos \theta = \frac{\sigma}{\pi\epsilon_0} \cos \theta \frac{da}{r}$$

$$= \frac{\sigma}{\pi\epsilon_0} d\theta$$

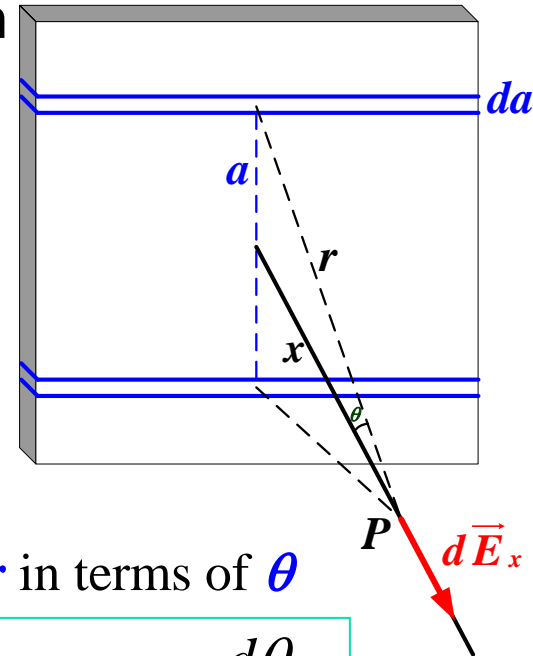
Find the **a** and **r** in terms of **θ**

$$a = x \tan \theta, \quad da = x \frac{d\theta}{\cos^2 \theta}$$

$$\frac{1}{r} = \frac{\cos \theta}{x}$$

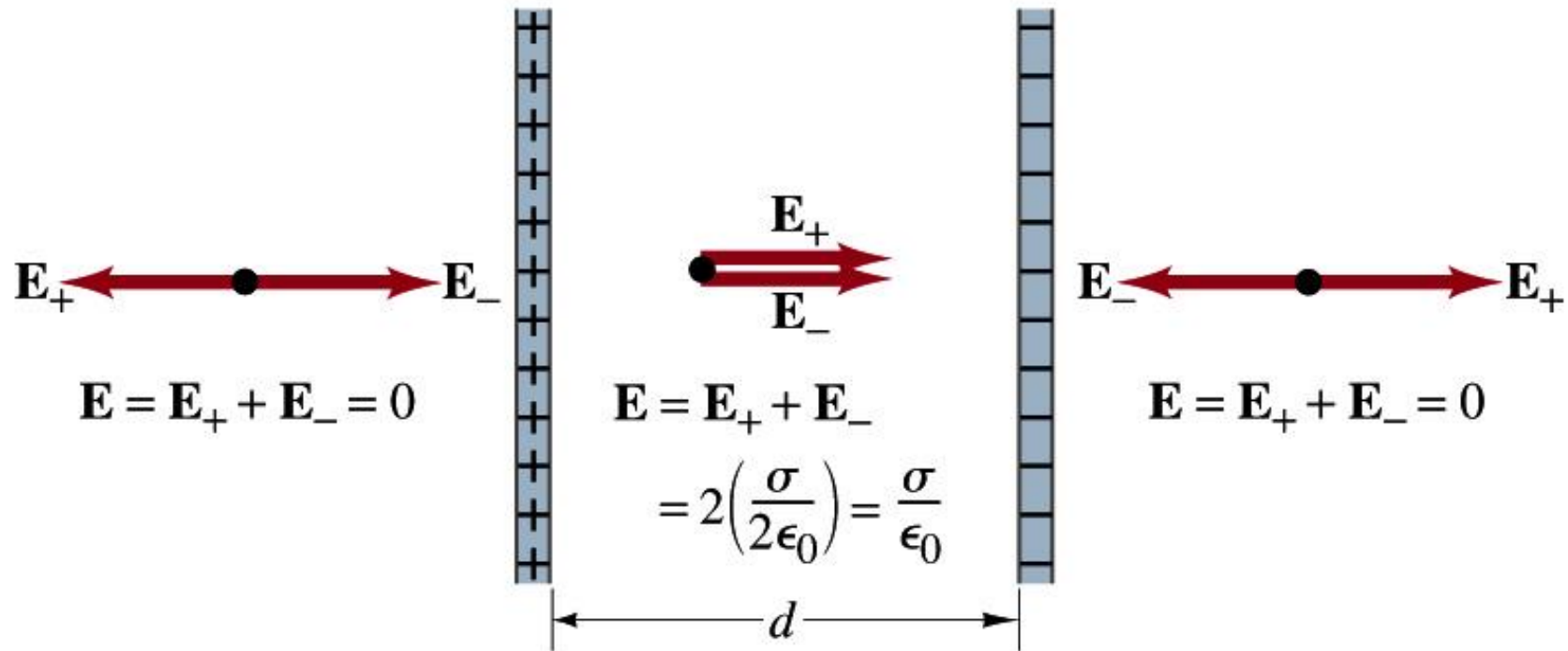


$$E_x = \frac{\sigma}{\pi\epsilon_0} \int_0^{\pi/2} d\theta = \frac{\sigma}{2\epsilon_0}$$



Example (P471 Ex. 19-12)

The electric field of two parallel plates



Problem-Solving Strategy to Calculating Electric field Due to Continuous Charge Distributions



- Divided the charge distribution into small elements dq .
 - ➔ Model each element as a point charge.

$$d\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{dq}{r^2} \hat{r}$$

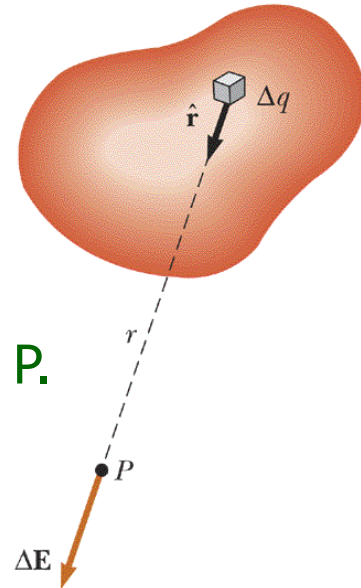
- ➔ Apply the superposition principle to get the total field at P.

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \int \frac{dq}{r^2} \hat{r}$$

- Establish a convenient coordinate system to complete the integral.
 - ➔ Using component representations to solve the vector integral separately.

- Choose appropriate infinitesimal charge element to simplified the integral.

- ➔ Generally choice: $dq = \lambda dl$ for line distribution, $dq = \sigma dA$ for surface distribution, and $dq = \rho dV$ for volume distribution.
- ➔ Using some symmetry of charge distribution to canceling some field components.



Problem-Solving Strategy to Calculating Electric field Due to Continuous Charge Distributions cont'd



- Using the known low-dimensional results for calculating field due to high-dimensional charge distribution.
 - ➡ For example, using the result of line charge distribution in a rod or a ring as the bases for calculation of field in the case of surface charge distribution.
 - ➡ Using the symmetry as possible as you can. For a disk of charge, adopt the ring result as the base. For a plane of charge, choose the rod result as the base.

Ch19 Prob. 49, 50 (P481, 483)

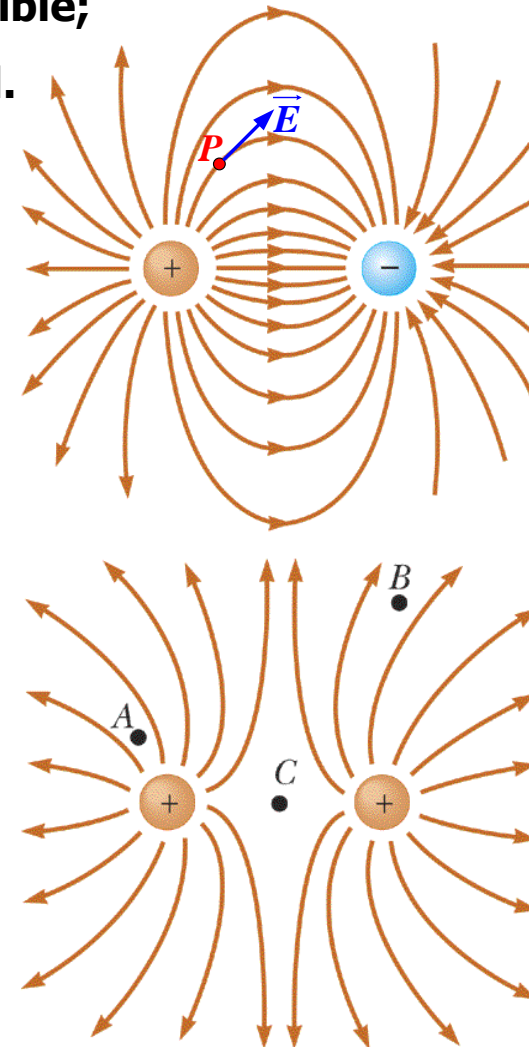
§ 3 Electric Field Lines (P471 § 19-8)



Why introduce electric field lines?

A **graphic** way for description of electric field

Visualize the electric field which is not visible;
Clarified the characteristic of electric field.



■ Electric field lines are related to the electric field in the following manner:

- ➔ **Direction** — is **tangent** to the electric field line at that point.
- ➔ **Magnitude** — is **proportional** to the number of electric field lines per unit area through the cross-sectional surface in that region. E is larger where the field lines are close together and smaller where they are far apart.

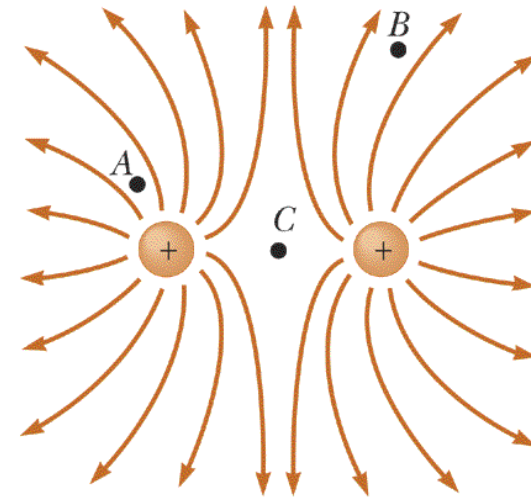
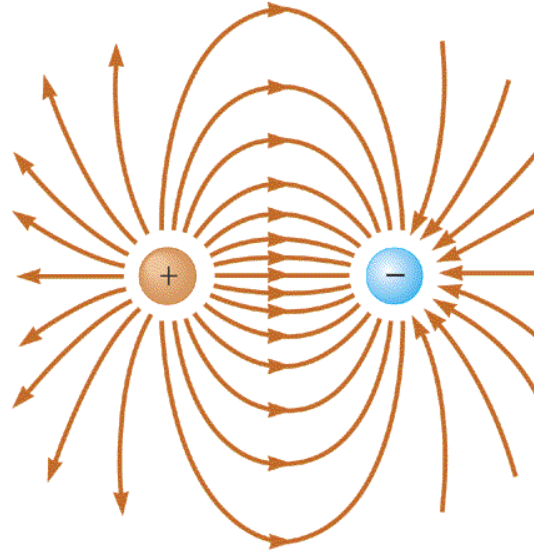
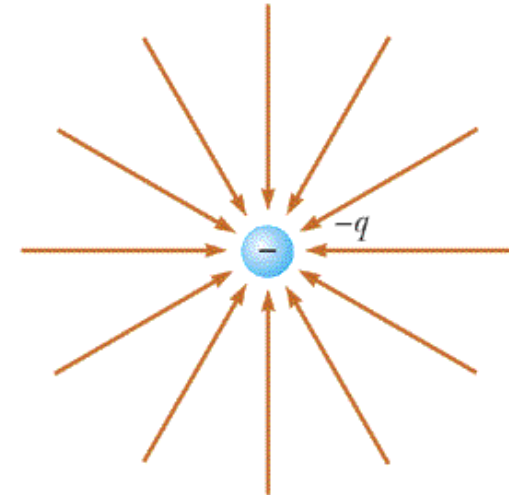
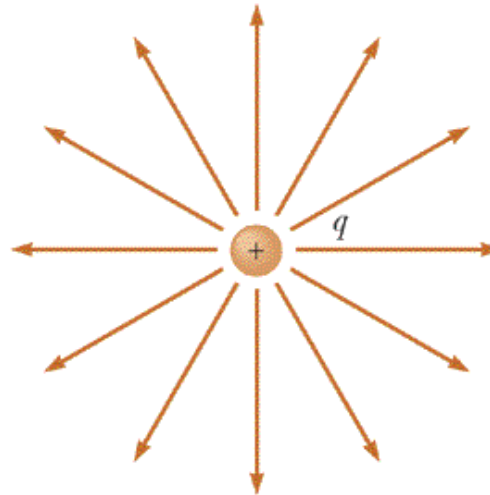
$$E \propto \frac{dN}{dA_{\perp}}$$

The fundamental **properties** for electric field lines



➔ **No** two field lines can cross each other.

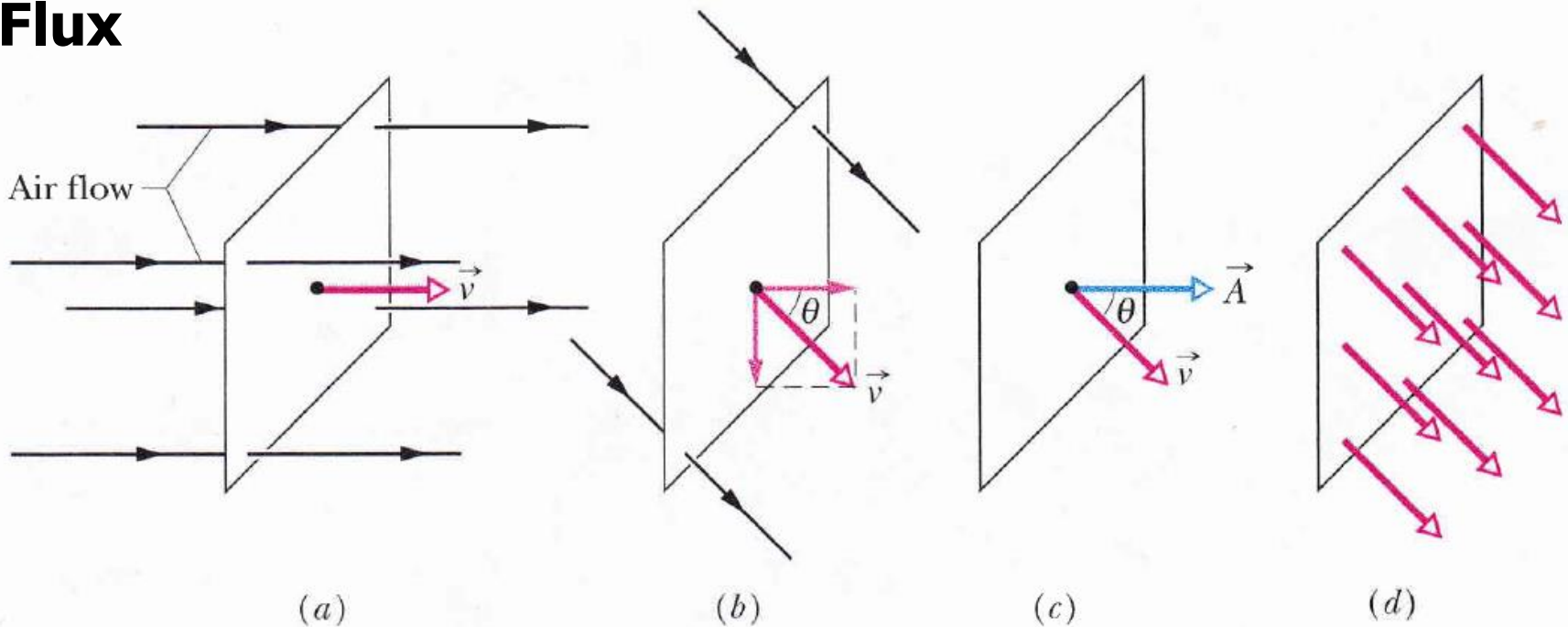
➔ **Begin** on positive charges (or infinite far away) and **end** on negative charges (or infinite far away). In the case of an excess of one type of charge, some lines will begin or end infinitely far away.



§ 4 Electric Flux (P487 § 20-1)



■ Flux



Imagine a airstream of uniform velocity \vec{v} at a small square surface of area A .
The flow of air volume (incompressible) $\Phi = vA \cos \theta = \vec{v} \cdot \vec{A}$
through the surface per unit time:

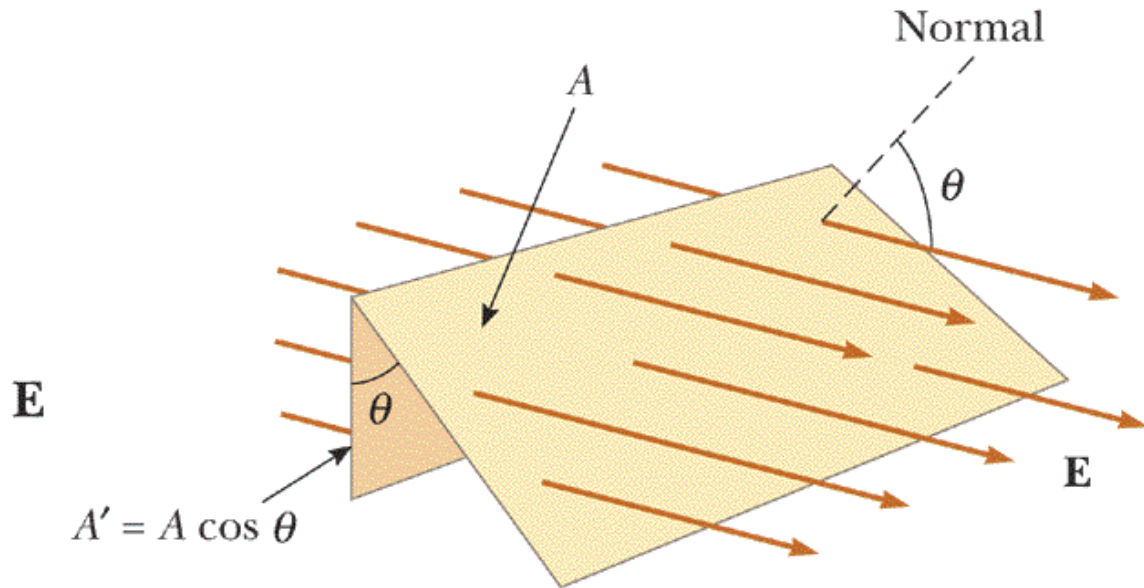
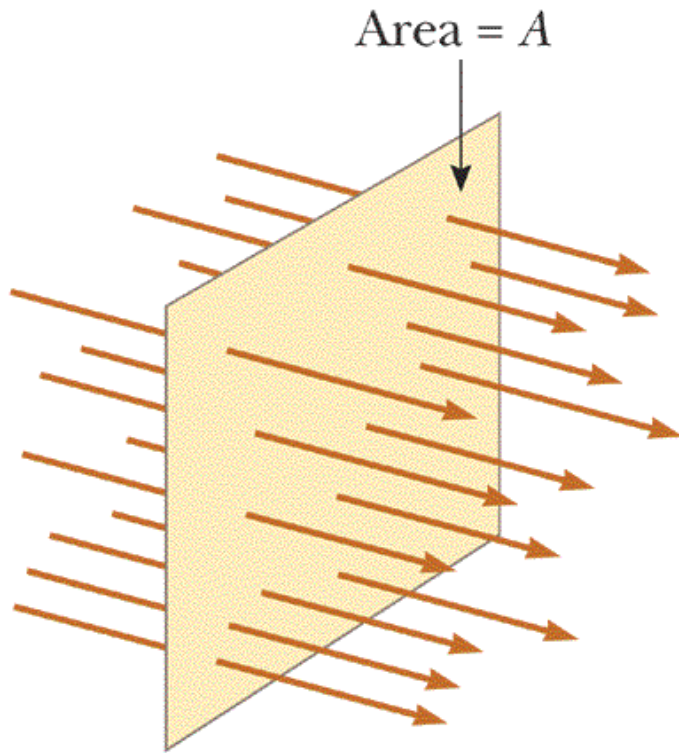
Velocity field: $\vec{v}(\vec{r})$ (a **volume** flux)

Φ — Flux of the **velocity field** through the surface

Electric Flux Φ_E



For **uniform** electric field



$$\Phi_E = EA \text{ (perpendicular area),} \quad \Phi_E = EA' = EA \cos \theta = \vec{E} \cdot \vec{A}$$

Electric Flux



- ➔ For **general** electric field that may vary in both magnitude and direction, **curved** surface — be divided into a large number of a small element of area: ΔA_i .

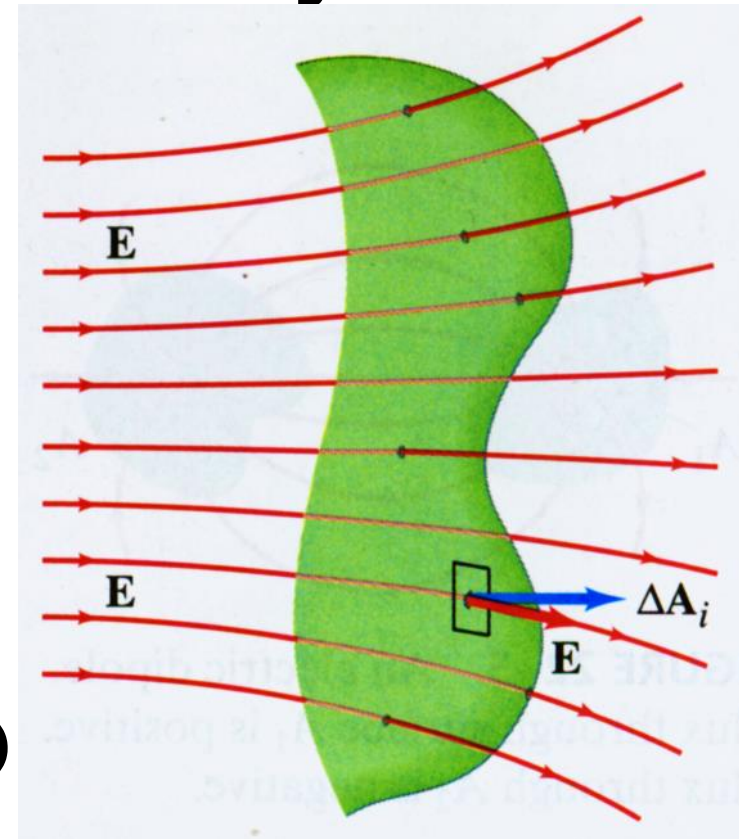
The flux through a small **element**:

$$\Delta\Phi_E = E_i \Delta A_i \cos \theta_i = \vec{E}_i \cdot \Delta \vec{A}_i$$

The total electric flux:

$$\Phi_E = \lim_{\Delta A_i \rightarrow 0} \sum \vec{E}_i \cdot \Delta \vec{A}_i = \iint_{\text{surface}} \vec{E} \cdot d\vec{A}$$

(surface integral)



$$E \propto \frac{dN}{dA_{\perp}}, \quad d\Phi_E = \vec{E} \cdot d\vec{A} = E dA_{\perp} \propto dN$$

the **number** of field lines

Electric flux through a **closed** surface



For a closed surface, **outward** direction is defined to be **positive**.

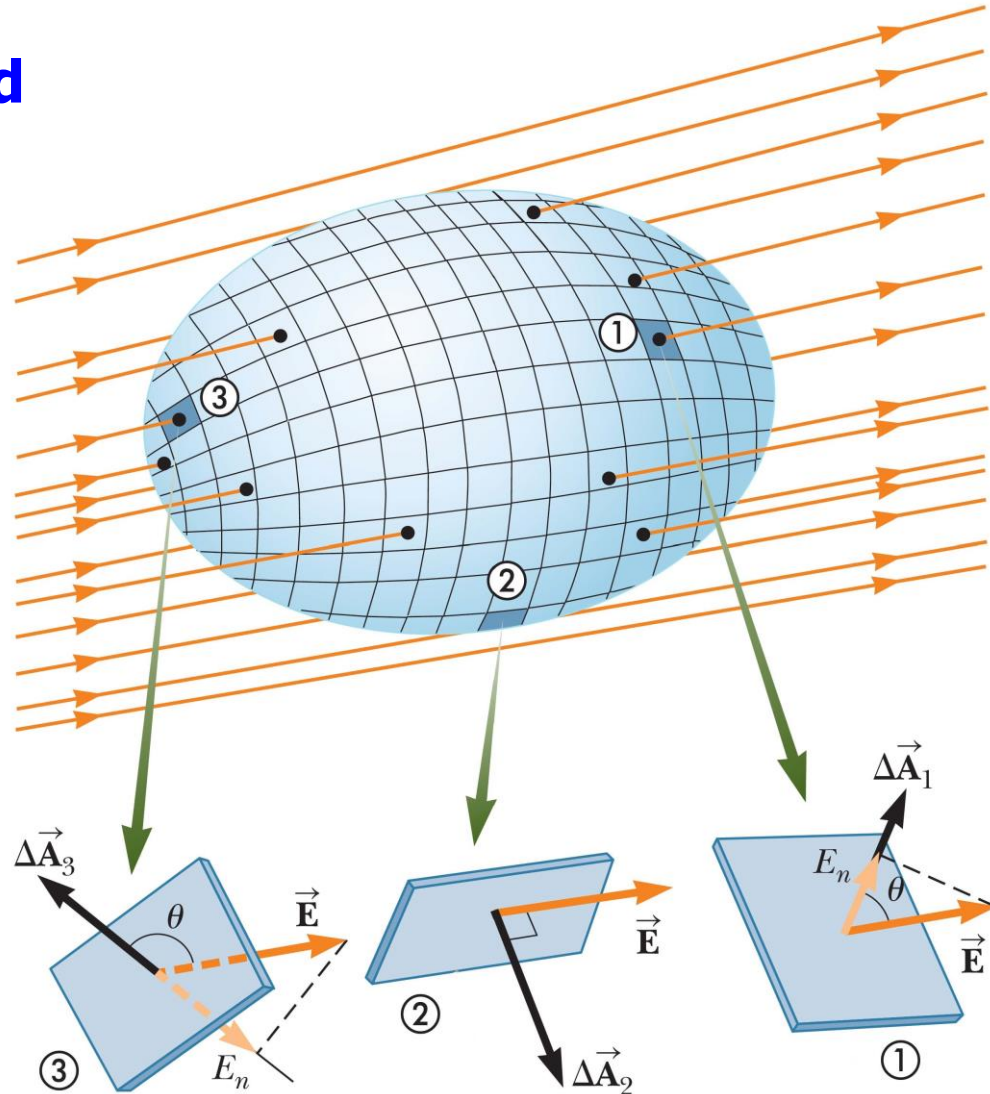
At point ①, $\theta < 90^\circ$, $\Phi_E > 0$.

At point ②, $\theta = 90^\circ$, $\Phi_E = 0$.

At point ③, $\theta > 90^\circ$, $\Phi_E < 0$.

➔ The net electric flux through a **closed** surface:

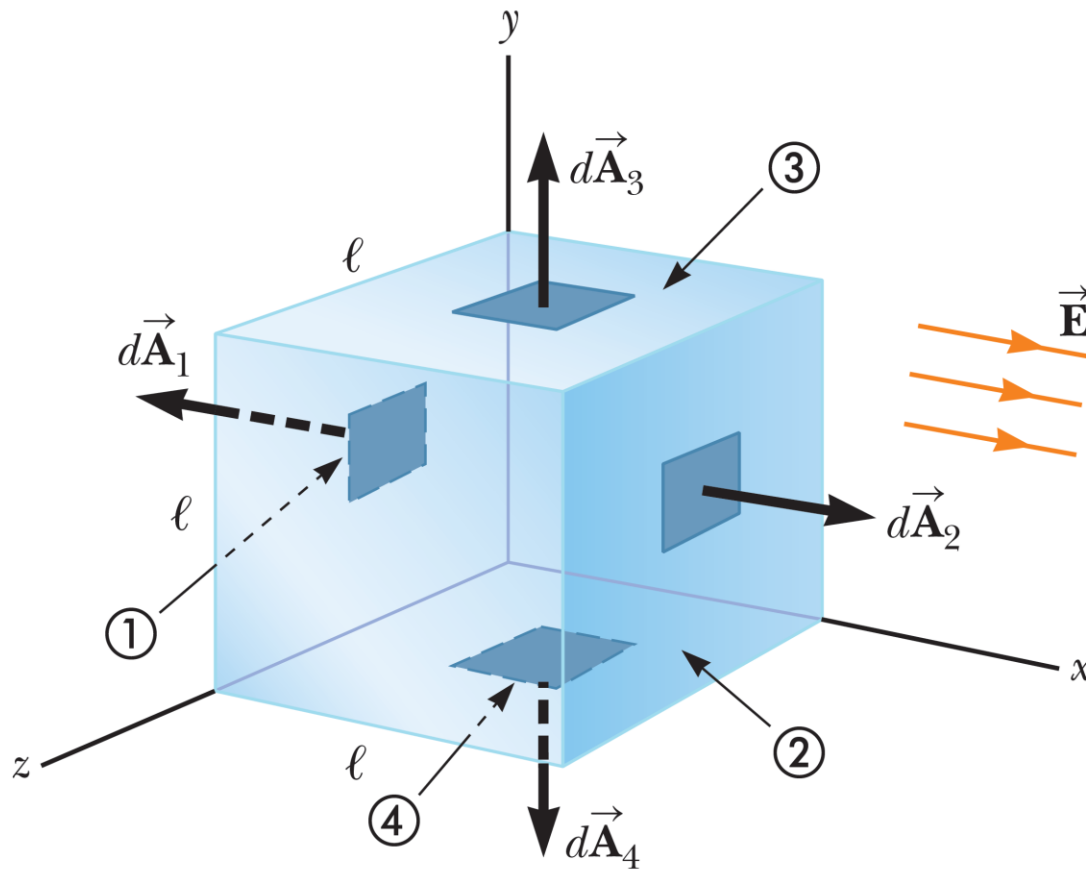
$$\Phi_E = \oiint_S \vec{E} \cdot d\vec{A}$$



Example



Consider a uniform electric field \vec{E} directed along the $+x$ axis. Find the **net electric flux** through the surface of a cube of edges ℓ shown in the figure.



Example



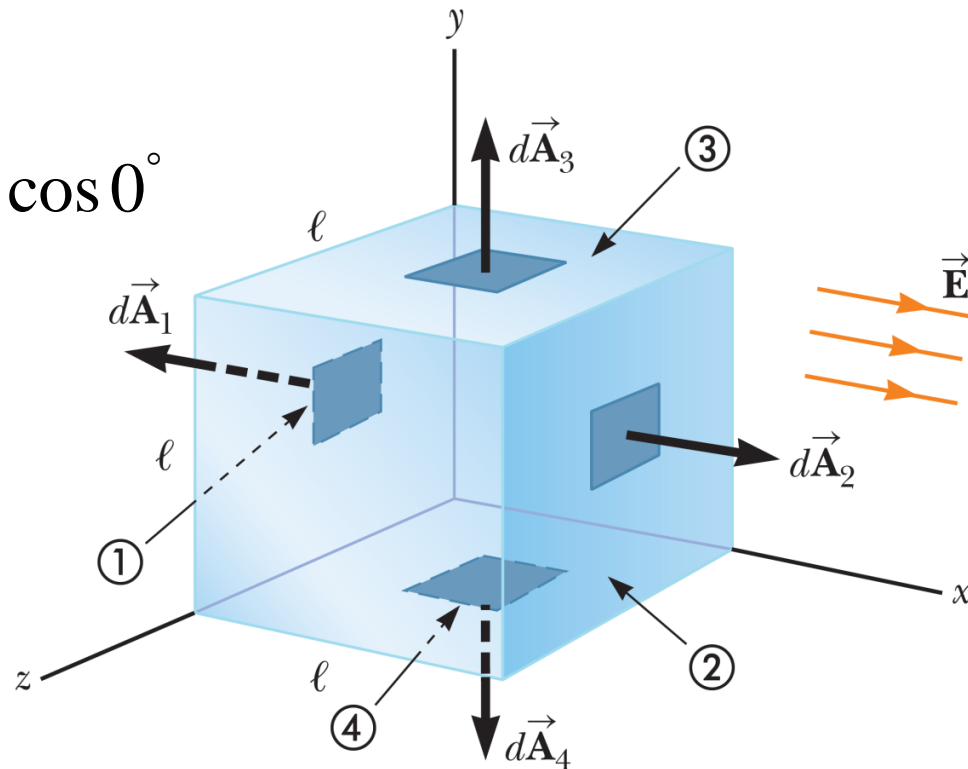
Solution: For the faces labeled ③ and ④ (⑤ and ⑥), the orientation of $d\vec{A}$ is perpendicular to \vec{E} .

The net flux through the surface of cube.

$$\Phi_E = \iint_1 \vec{E} \cdot d\vec{A} + \iint_2 \vec{E} \cdot d\vec{A}$$

$$= \iint_1 EdA \cos 180^\circ + \iint_2 EdA \cos 0^\circ$$

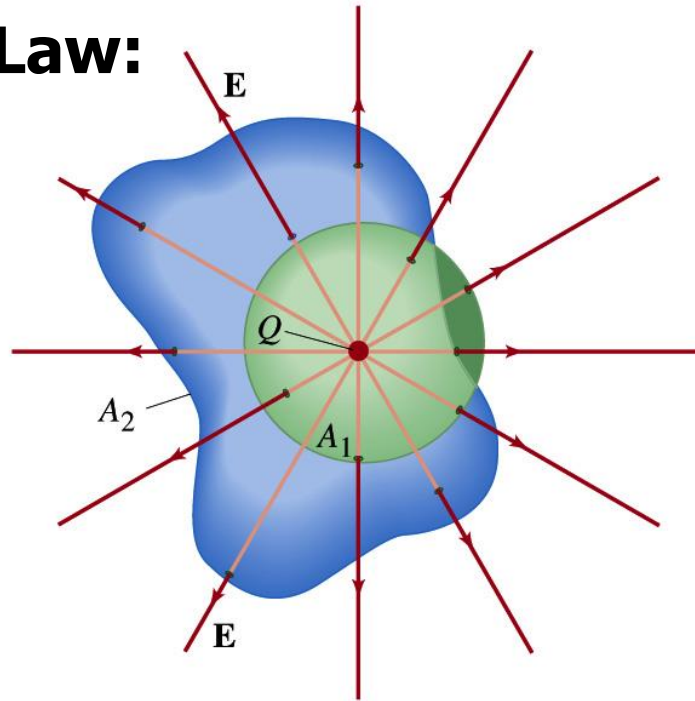
$$= -EA + EA = 0$$



§ 5 Gauss's Law



■ Gauss's Law:



$$\Phi_E = \oiint_S \vec{E} \cdot d\vec{A} = \frac{q_{inside}}{\epsilon_0}$$



Karl Friedrich Gauss
German mathematician
and astronomer
(1777–1855)

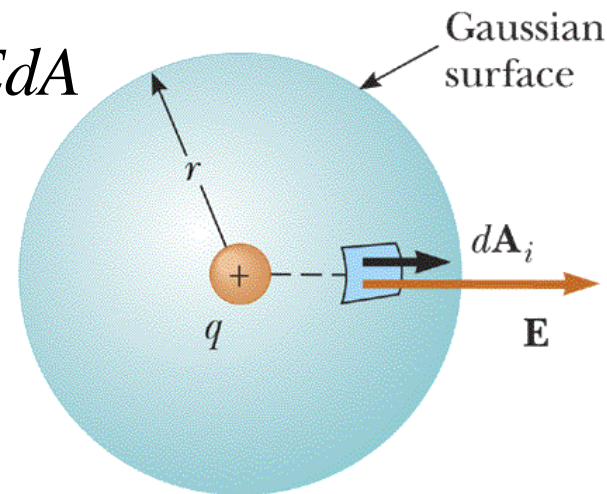
➡ The **net** electric flux through any closed surface is equal to the net charge **inside** the surface divided by ϵ_0 .

The Gauss's Law



- A point charge q locates at the center of a **spherical** surface.

$$\begin{aligned}\Phi_E &= \oint\limits_{\text{spherical surface}} \vec{E} \cdot d\vec{A} = \oint\limits_{\text{spherical surface}} E(dA) \cos \theta = \oint\limits_{\text{spherical surface}} E dA \\ &= E \oint\limits_{\text{spherical surface}} dA = \left(\frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \right) (4\pi r^2) = \frac{q}{\epsilon_0}\end{aligned}$$



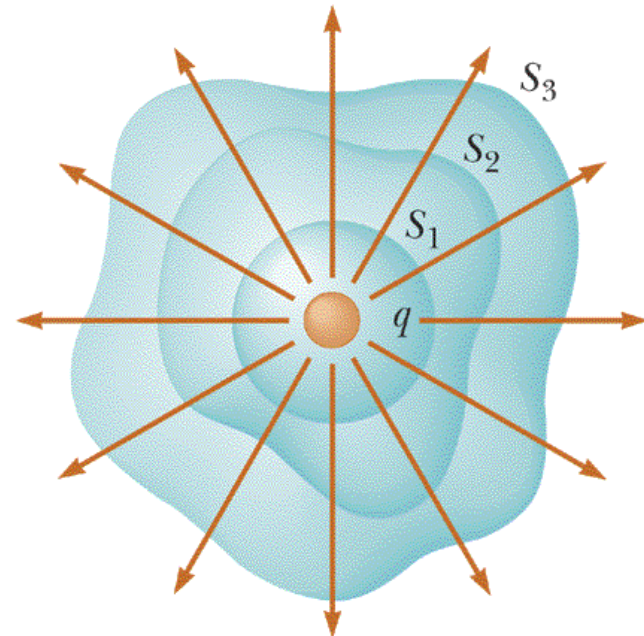
- ➡ The net flux is proportional to the **charge** inside the surface;
- ➡ The net flux is **independent** of the radius r — every field line from the charge must pass through the surface
- ➡ The fact that the net flux is independent of the radius is consequence of **inverse-square** dependence of the electric field according to Coulomb's law.

The Gauss's Law



- The charge q inside, the closed surface **not spherical**.
 - The flux that passes through spherical surface S_1 has the value q / ϵ_0 .
 - The number of electric field lines through the spherical surface S_1 is equal to the number of electric field lines through the non-spherical surfaces S_2 and S_3 .

$$\oiint_{S_1} \vec{E} \cdot d\vec{A} = \oiint_{S_2} \vec{E} \cdot d\vec{A} = \oiint_{S_3} \vec{E} \cdot d\vec{A} = \frac{q}{\epsilon_0}$$

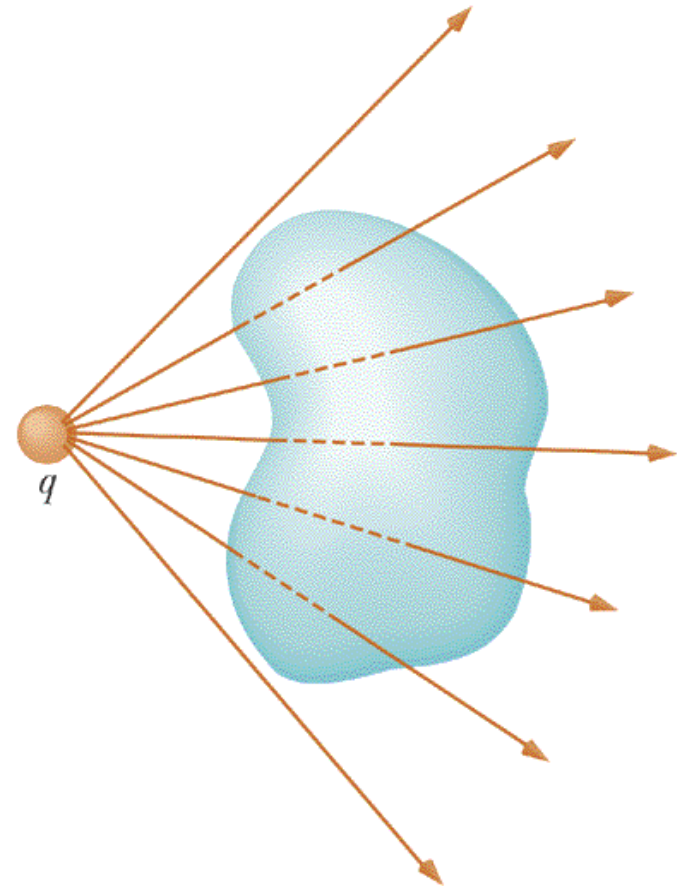


The Gauss's Law



- A point charge locates **outside** a closed surface of arbitrary shape.
 - ➔ The number of electric field lines entering the surface equals the number leaving the surface.

$$\Phi_E = \Phi_E^{in} + \Phi_E^{out} = 0$$



The Gauss's Law



- The **series** of charges, some inside, some outside the closed surface.

➡ The total electric field at any point:

$$\vec{E} = \sum_i \vec{E}_i$$

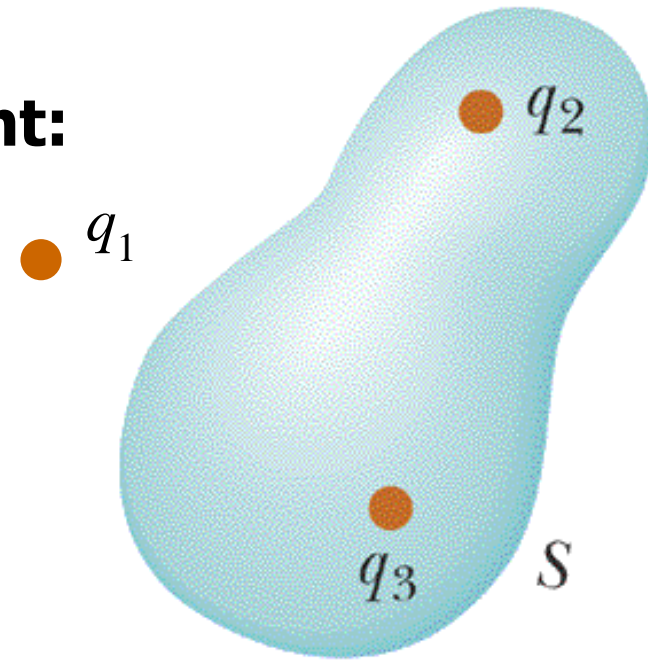
$$\Phi_E = \oiint_S \vec{E} \cdot d\vec{A} = \oiint_S \sum_i \vec{E}_i \cdot d\vec{A}$$

$$= \sum_i \oiint_S \vec{E}_i \cdot d\vec{A} = \sum_i \Phi_{Ei}$$

$$\Phi_{Ei} = \begin{cases} \frac{q_i}{\epsilon_0} & \text{if } q_i \text{ inside the } S \\ 0 & \text{if } q_i \text{ outside the } S \end{cases}$$



$$\Phi_E = \sum_i \Phi_{Ei} = \frac{q_{\text{inside}}}{\epsilon_0}$$



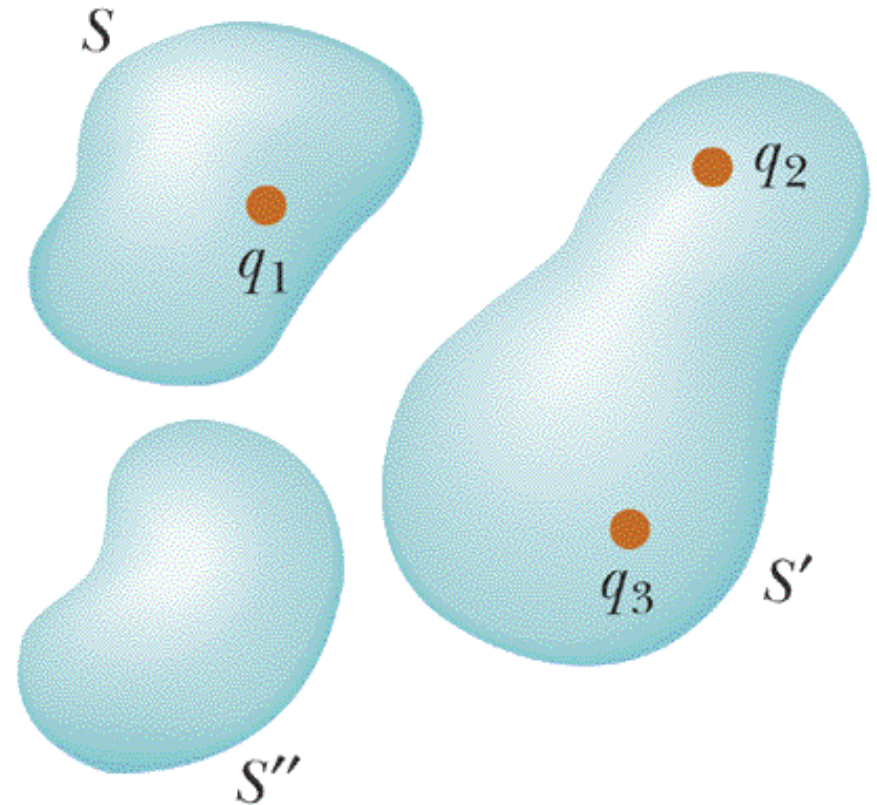
Some comments on Gauss's Law

$$\Phi_E = \oiint_S \vec{E} \cdot d\vec{A} = \frac{q_{\text{inside}}}{\epsilon_0}$$

$$\oiint_S \vec{E} \cdot d\vec{A} = \frac{q_1}{\epsilon_0}$$

$$\oiint_{S'} \vec{E} \cdot d\vec{A} = \frac{q_2 + q_3}{\epsilon_0}$$

$$\oiint_{S''} \vec{E} \cdot d\vec{A} = 0$$



- ➡ The flux only depends on the charges **inside** (enclosed).
- ➡ E on the left side of Gauss's law is the E in Gaussian surface and it is not necessarily due to the charge inside the surface. It is produced by **all** the charges in the space.
- ➡ **Zero** flux doesn't mean the **zero** field.

Gauss's Law and Coulomb's Law

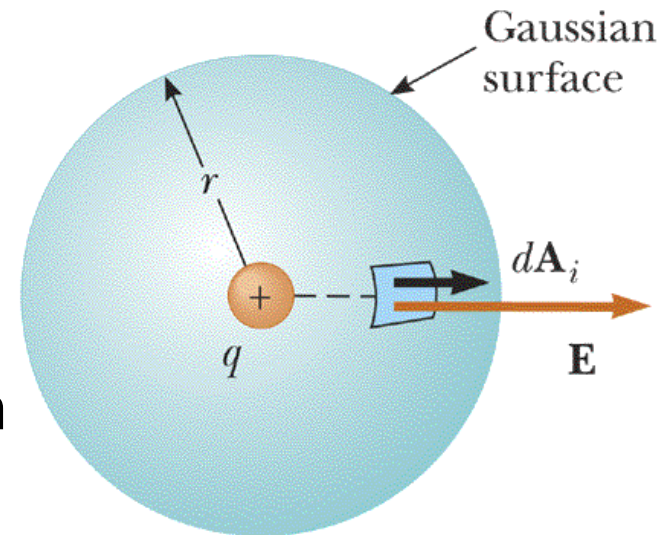


- ➡ Gauss's law is deduced from Coulomb's law. Coulomb's law can **also** be deduced from Gauss's law.

For an isolated charge q locates inside a spherical surface

$$\oiint_S \vec{E} \cdot d\vec{A} = \oiint_S E dA = E \oiint_S dA = E(4\pi r^2) = \frac{q}{\epsilon_0} \quad \Rightarrow \quad E = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2}$$

- ➡ These two laws can regarded as **equivalent** in the situation of **electrostatics**, Gauss's law is found to hold also for electric fields generated by changing magnetic field.
- ➡ Gauss's law is a more general law than Coulomb's law, and so is regarded as a more **fundamental** equation of electromagnetisms.



§ 6 Application of Gauss's Law to Symmetric Charge Distributions

Generally

The charge distribution is known

Coulomb's Law



The electric field

The electric field is known

Gauss's Law



The charge distribution

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \int \frac{dq}{r^2} \hat{r}$$

$$\Phi_E = \oiint_S \vec{E} \cdot d\vec{A} = \frac{q_{inside}}{\epsilon_0}$$

- ➡ For **special case** where the charge distribution possesses a high degree of symmetry, Gauss's law can be used to **calculate the electric field**.

Example

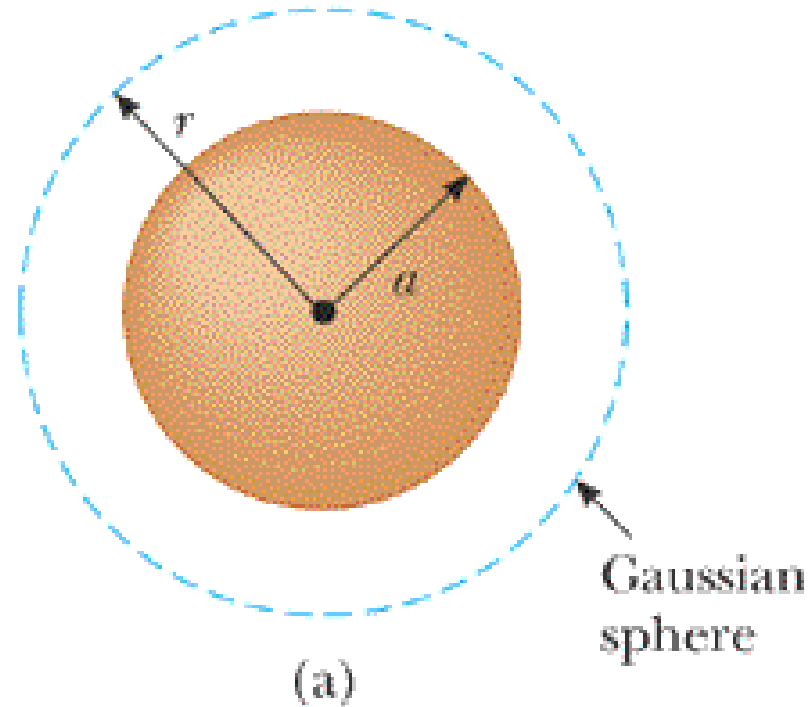


A spherical symmetric charge distribution

An insulating solid sphere of radius a has a uniform volume charge density ρ and carries a total positive charge Q .

(1) Calculate the magnitude of the electric field at a point **outside** the sphere.

(2) Find the magnitude of the electric field at a point **inside** the sphere.



A spherical symmetric charge distribution



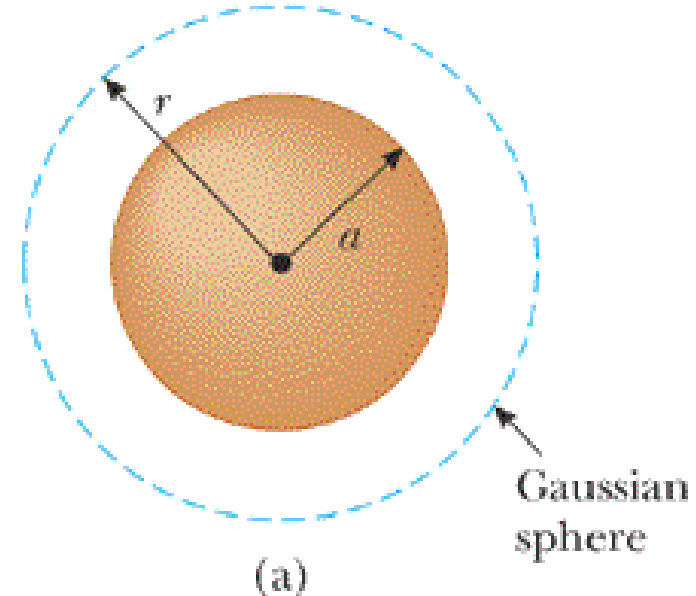
Solution:

(1) Select a spherical gaussian surface of radius $r > a$

$$\begin{aligned}\oiint_S \vec{E} \cdot d\vec{A} &= \oiint_S E(dA) = E \oiint_S dA = E(4\pi r^2) \\ &= \frac{q_{in}}{\epsilon_0} = \frac{Q}{\epsilon_0}\end{aligned}$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \quad (\text{for } r > a)$$

This result is identical to that obtained for a **point charge.**



A spherical symmetric charge distribution

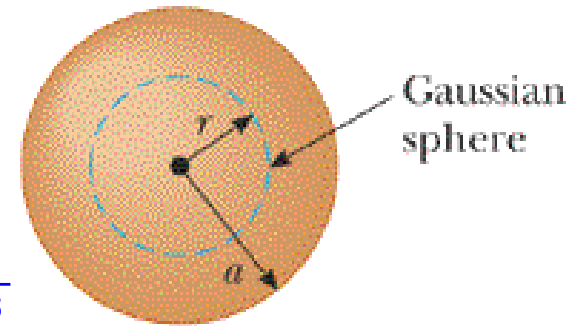


Solution: (2) Select a spherical gaussian surface of radius $r < a$

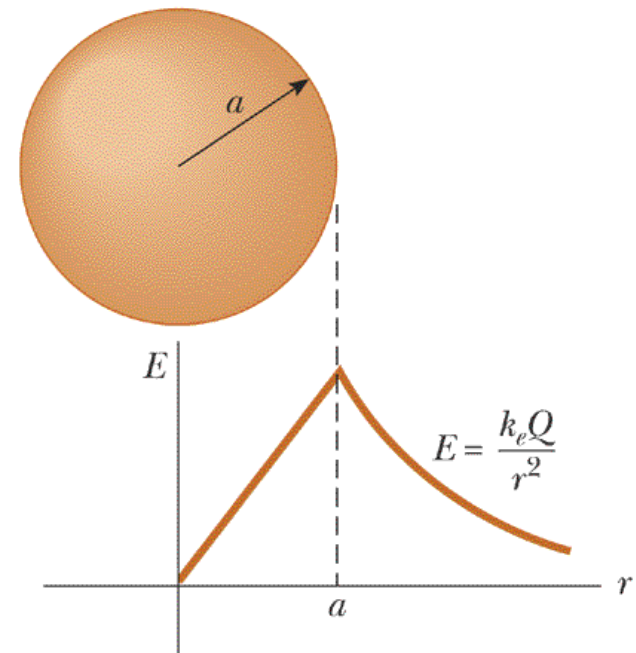
$$\oiint_S \vec{E} \cdot d\vec{A} = \oiint_S E(dA) = E \oiint_S dA = E(4\pi r^2)$$
$$= \frac{q_{in}}{\epsilon_0} = \frac{1}{\epsilon_0} \rho \left(\frac{4}{3} \pi r^3 \right) = \frac{1}{\epsilon_0} \left(\frac{Q}{\frac{4}{3} \pi a^3} \right) \left(\frac{4}{3} \pi r^3 \right) = \frac{Q}{\epsilon_0} \frac{r^3}{a^3}$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{Q}{a^3} r \quad (\text{for } r < a)$$

The expressions of electric fields inside and outside the sphere match when $r=a$.



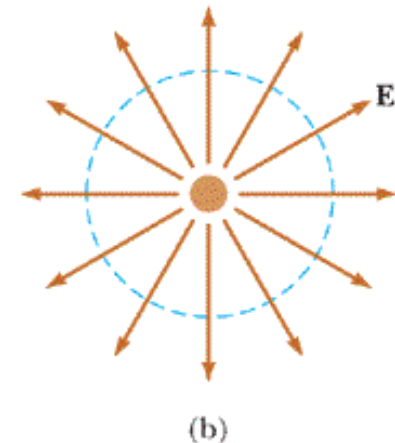
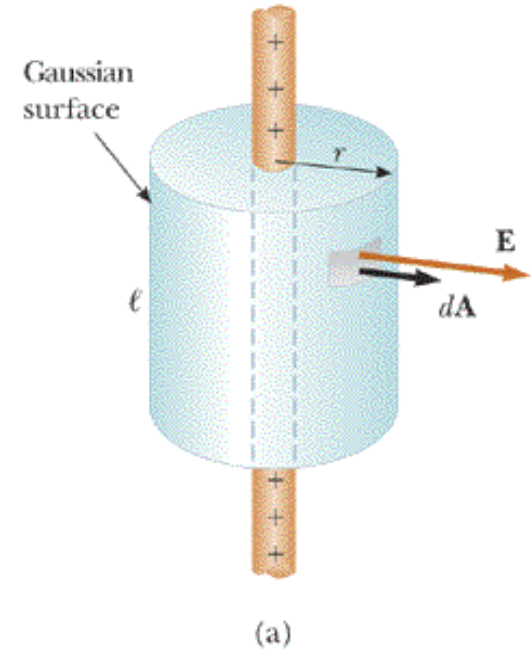
(b)



Example

A cylindrically symmetric charge distribution

Find the electric field a distance r from a line of positive charge of **infinite** length and constant charge per unit length λ .

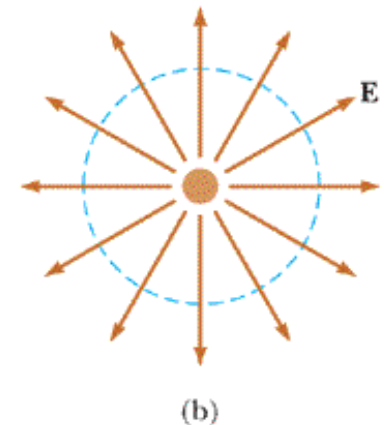
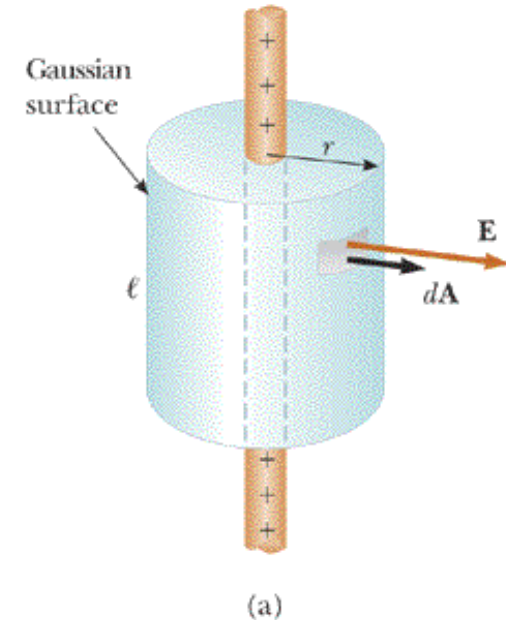


A cylindrically symmetric charge distribution

Solution: Select a **cylindrical** Gaussian surface of radius **r** and length **l** that is coaxial with the line charge.

$$\begin{aligned}
 \oint_S \vec{E} \cdot d\vec{A} &= \iint_{\text{side surface}} \vec{E} \cdot d\vec{A} + \iint_{\text{top and bottom}} \vec{E} \cdot d\vec{A} \\
 &= \iint_{\text{side surface}} \vec{E} \cdot d\vec{A} = E(2\pi rl) \\
 &= \frac{q_{in}}{\epsilon_0} = \frac{\lambda l}{\epsilon_0}
 \end{aligned}$$

$$E = \frac{1}{2\pi\epsilon_0} \frac{\lambda}{r}$$

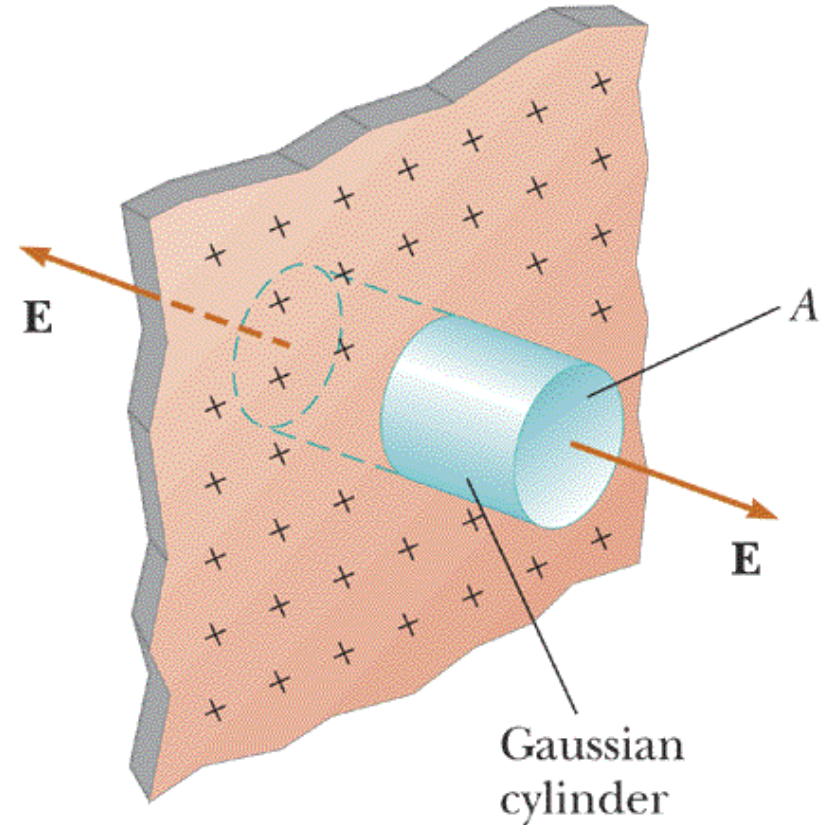


Example



A Nonconducting plane sheet of charge

Find the electric field due to a non-conducting, infinite plane with uniform surface charge density σ .

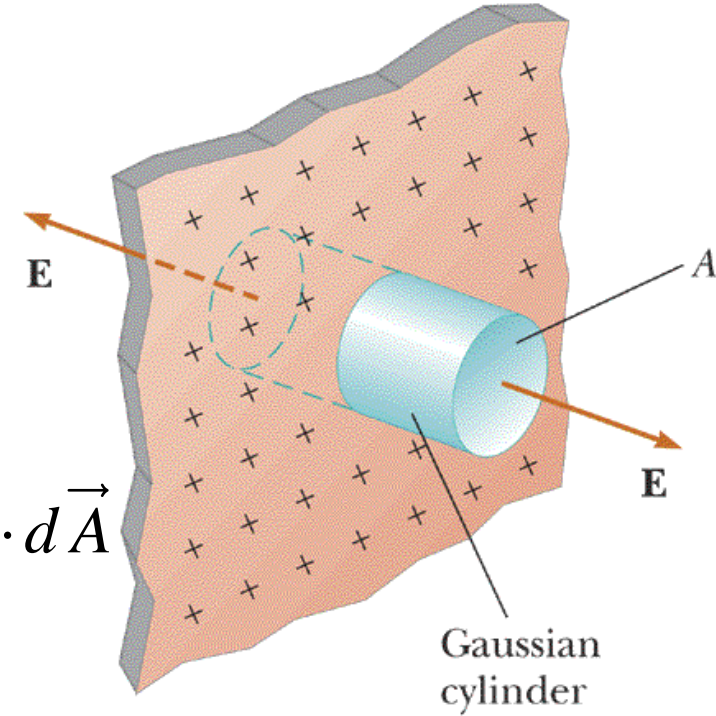


Example



A Nonconducting plane sheet of charge

Solution: Select the Gaussian surface to be a cylinder whose axis is perpendicular to the plane and whose ends each have an area A and are equidistant from the plane.



$$\begin{aligned}\oiint_S \vec{E} \cdot d\vec{A} &= \iint_{\text{side surface}} \vec{E} \cdot d\vec{A} + \iint_{\text{two ends of the cylinder}} \vec{E} \cdot d\vec{A} \\ &= \iint_{\text{two ends of the cylinder}} \vec{E} \cdot d\vec{A} = E(2A) \\ &= \frac{q_{in}}{\epsilon_0} = \frac{\sigma A}{\epsilon_0}\end{aligned}$$

$$E = \frac{\sigma}{2\epsilon_0}$$

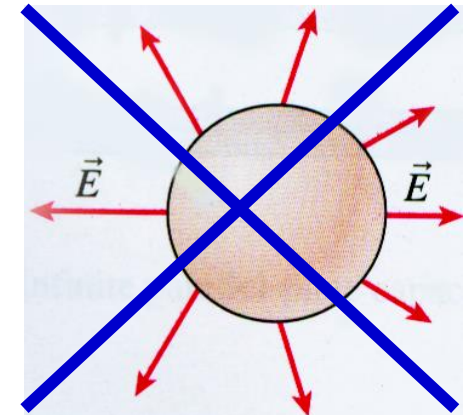
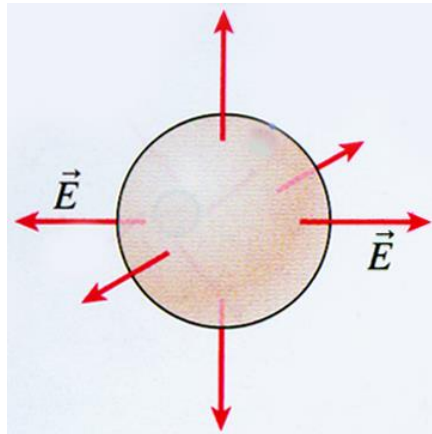
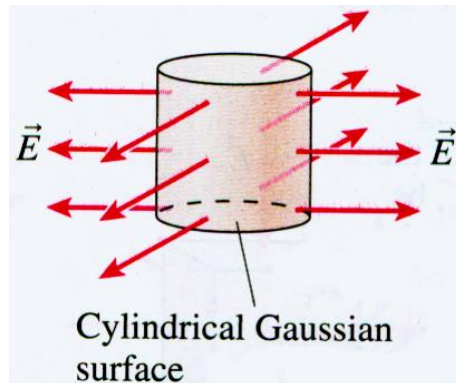
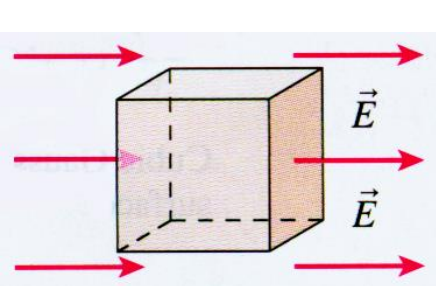
How to Choose the Gaussian Surfaces?

$$\Phi_E = \oiint_S \vec{E} \cdot d\vec{A} = \frac{q_{inside}}{\epsilon_0}$$

■ How to choose the gaussian surfaces?

- The choice of appropriate gaussian surface that allows E to be **removed** from the **integral** in Gauss's law is the key problem.

- With the appropriate gaussian surface, the dot product $\vec{E} \cdot d\vec{A}$ should be zero or equal to $E dA$, with the magnitude of E being constant.



Problems



P499, 500
Ch20 Prob. 26, 33, 34