

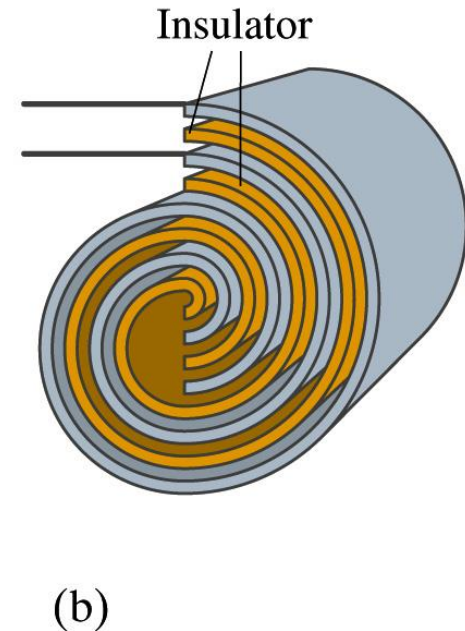
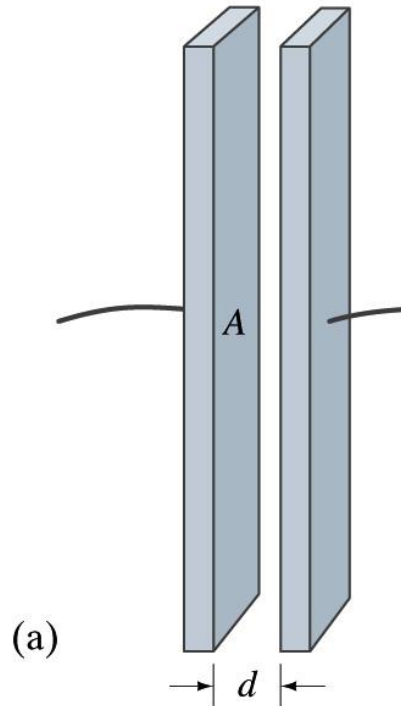
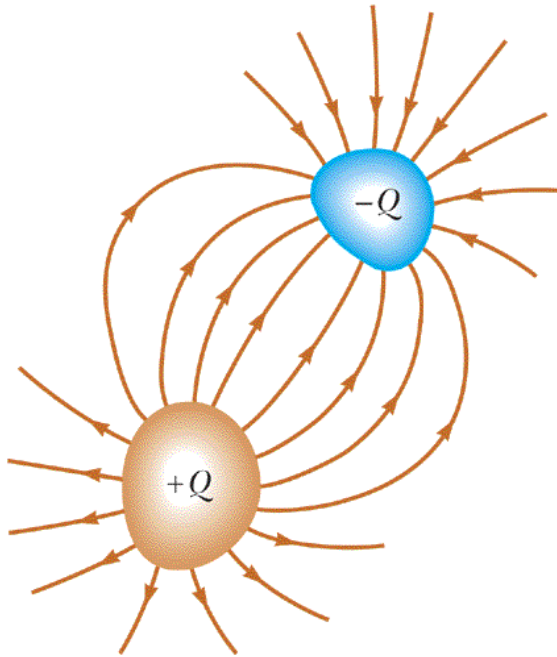
# Chapter 22 Capacitance, Dielectrics, Electric Energy Storage



## § 22-1 Capacitance (P525)

### ■ Capacitors

- Any two **conductors** separated by an insulator (or a vacuum) form a **capacitor**, which can store amount of **charge**.



# Capacitance of a capacitor



- ➔ The capacitance  $C$  of a capacitor

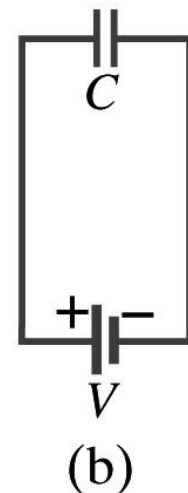
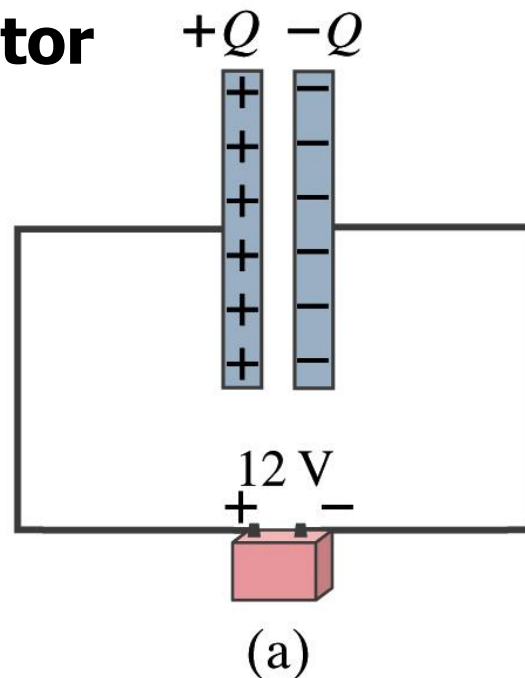
$$Q = C(\Delta V)$$

$$C \equiv \frac{Q}{\Delta V}$$

Farad

$\mu\text{F}(10^{-6}\text{F})$

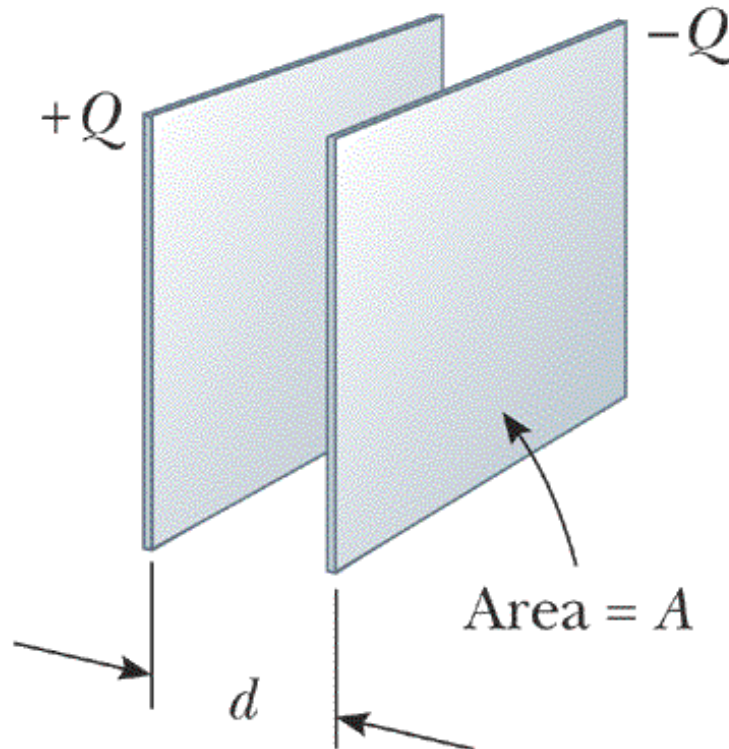
$\text{pF}(10^{-12}\text{F})$



- ➔ The capacitance of a capacitor **depends on** the **geometric arrangement** of the conductors, and is **independent of** the charge  $Q$  or the potential difference  $\Delta V$ . Because the potential difference is proportional to the charge, the ratio  $Q/\Delta V$  is constant for a given capacitor.

### Problem-Solving Strategy :

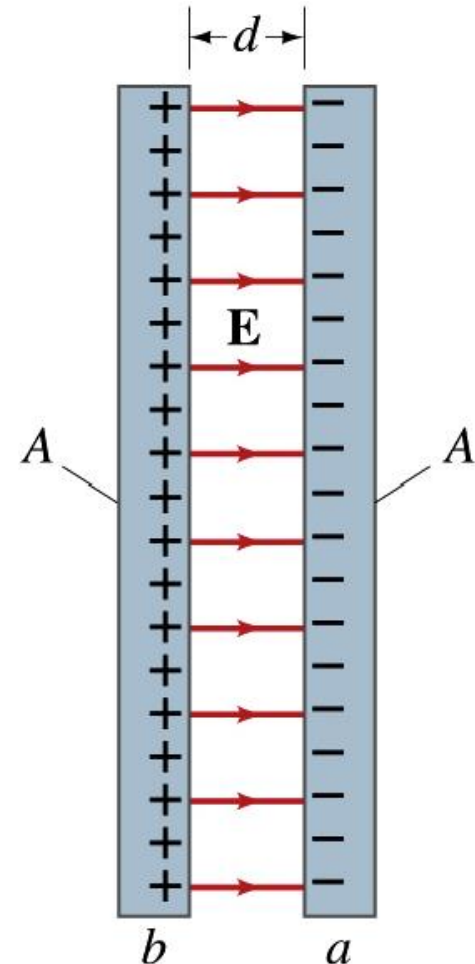
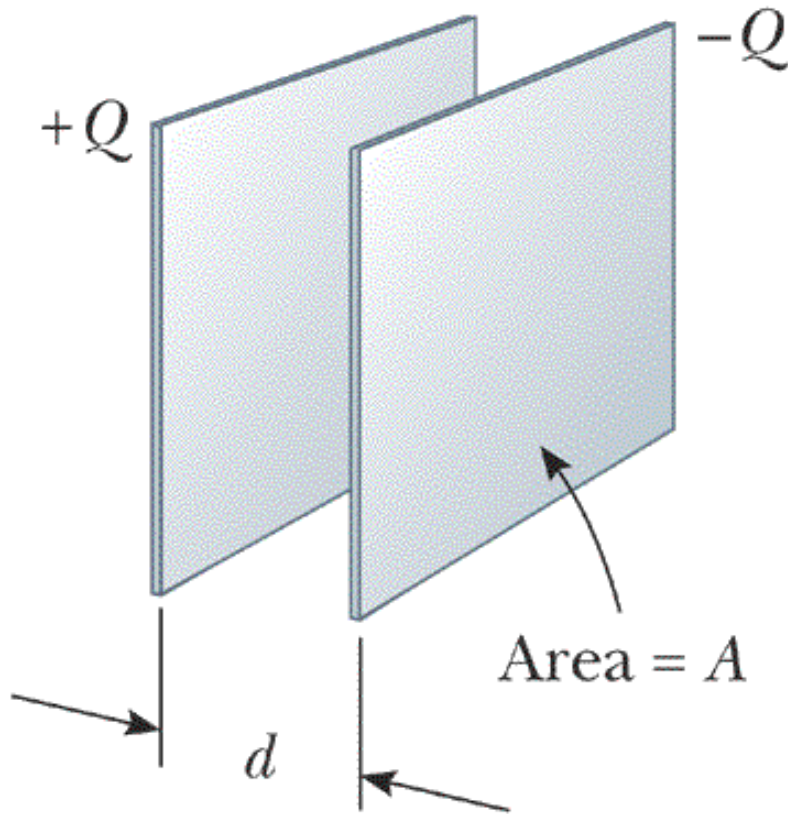
- A convenient charge of magnitude  $Q$  is assumed.
- The potential difference  $\Delta V$  is calculated.
- Use  $C=Q/\Delta V$  to evaluate the capacitance.



## The parallel-plate capacitor (P527)



A parallel-plate capacitor consists of two parallel plates of equal area  $A$ , separated by a distance  $d$ . Find the **capacitance**.



## The parallel-plate capacitor



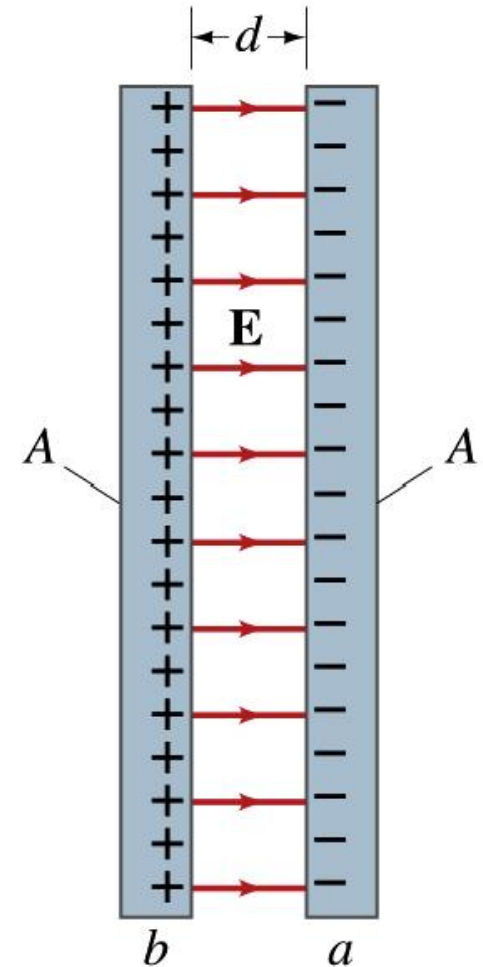
**Solution: Assume the two plates have opposite charges  $+Q$  and  $-Q$ . An uniform electric field is:**

$$E = \frac{\sigma}{\epsilon_0} = \frac{Q}{\epsilon_0 A}$$

**The potential difference:**

$$\Delta V = \int_+^- \vec{E} \cdot d\vec{l} = Ed = \frac{Qd}{\epsilon_0 A}$$

$$C = \frac{Q}{\Delta V} = \frac{\epsilon_0 A}{d}$$

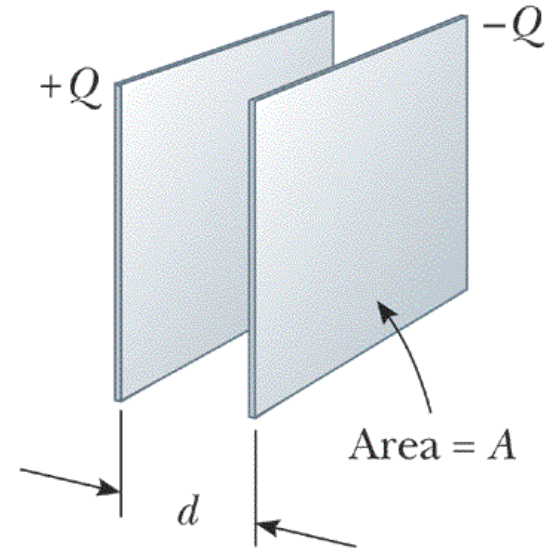


# The parallel-plate capacitor



$$C = \epsilon_0 \frac{A}{d}$$

- ➔ The capacitance of a parallel-plate capacitor is proportional to the **area** of its plates and inversely proportional to the plate **separation**, which are the geometrical factors.
- ➔ The capacitance does **not** depend on the potential difference or the **charge** carried by the plates.
- ➔ The **capacitance** has form of  $\epsilon_0$  times a quantity with the dimension of **length** ( $A/d$ ), which is essential form for all the capacitors.

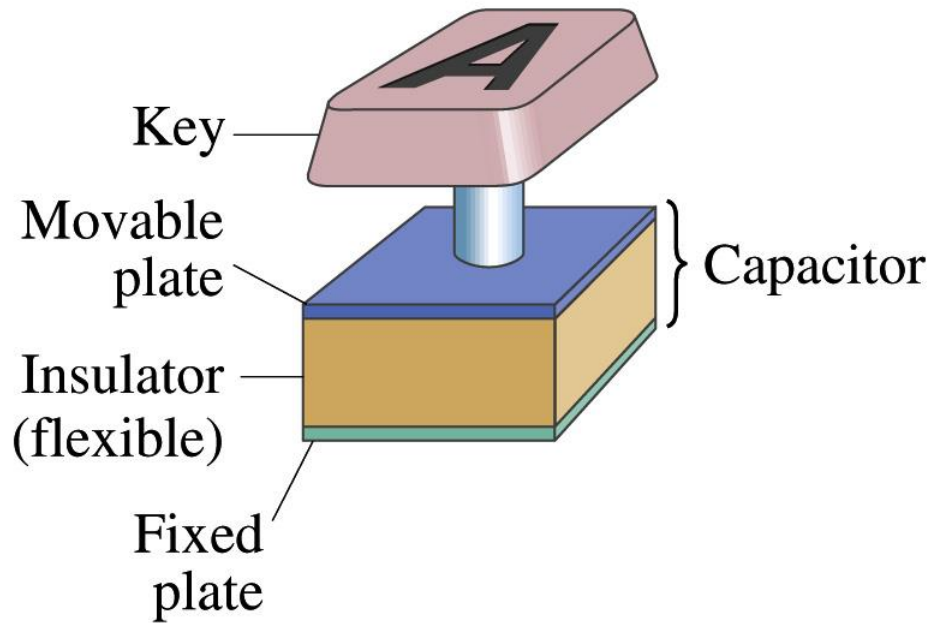


$$\epsilon_0 = 8.85 \times 10^{-12} \text{ F/m} = 8.85 \text{ pF/m}$$

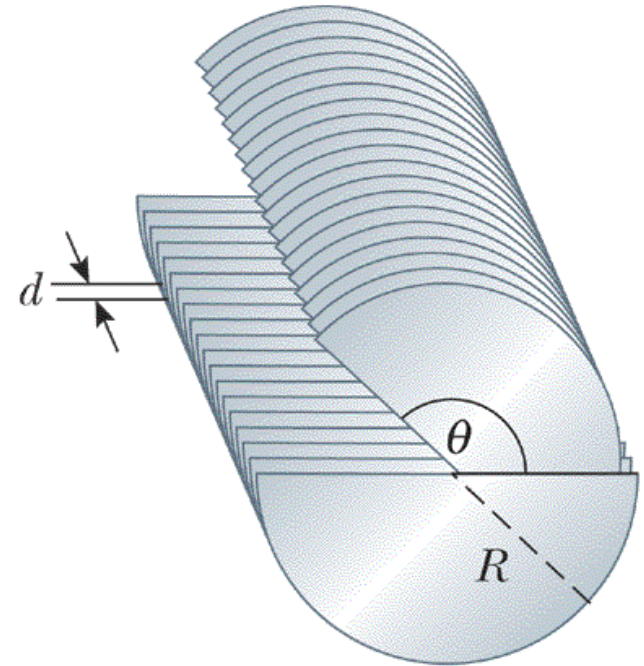
# The parallel-plate capacitor



$$C = \frac{\epsilon_0 A}{d}$$



**Key on a computer keyboard**



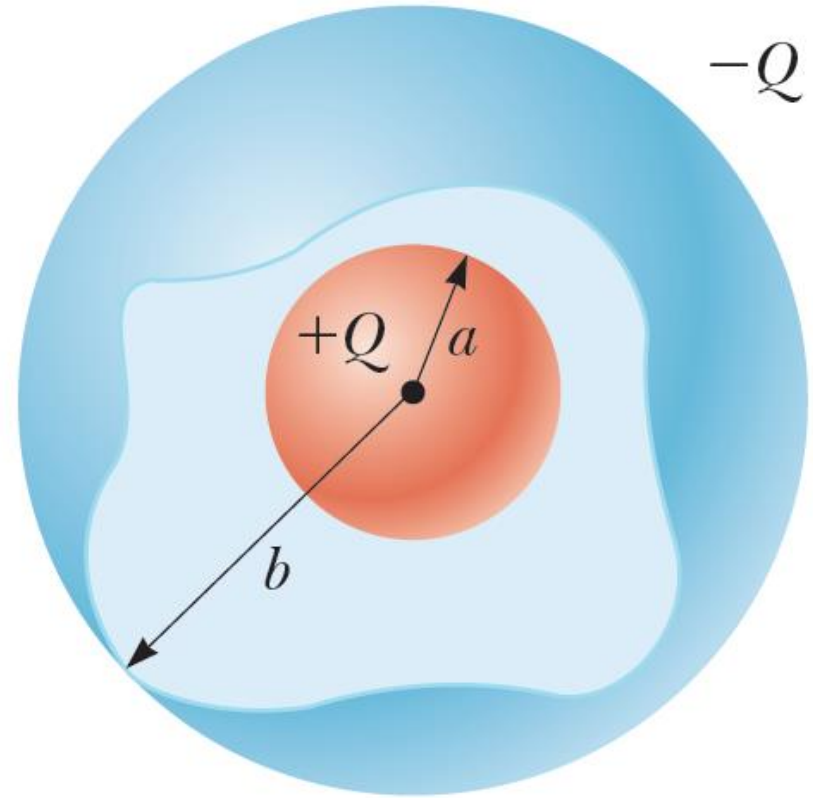
**A radio tuner**



## The Spherical Capacitor (P529 Ex. 22-3)



A spherical capacitor in which the inner conductor is a solid sphere of radius  $a$ , and outer conductor is a hollow spherical shell of inner radius  $b$ . Find the capacitance.





## The Spherical Capacitor



**Solution:** Assume the inner and outer sphere have opposite charges  $+Q$  and  $-Q$ . In the region  $a < r < b$ , we can use **Gauss' law** to determine:

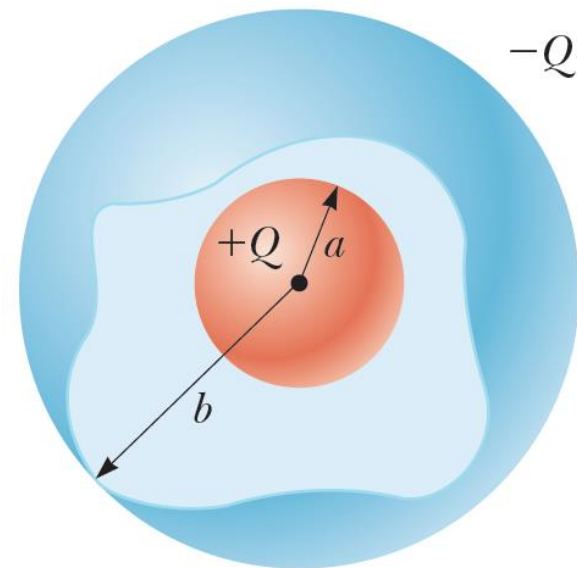
$$E = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}, \quad (a < r < b)$$

**The potential difference:**

$$\Delta V = \int_+^- \vec{E} \cdot d\vec{s} = \frac{Q}{4\pi\epsilon_0} \int_{r_a}^{r_b} \frac{dr}{r^2} = \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{a} - \frac{1}{b} \right)$$

$$= \frac{Q}{4\pi\epsilon_0} \frac{b-a}{ab},$$

$$C = \frac{Q}{\Delta V} = 4\pi\epsilon_0 \frac{ab}{b-a}$$



# The Spherical Capacitor

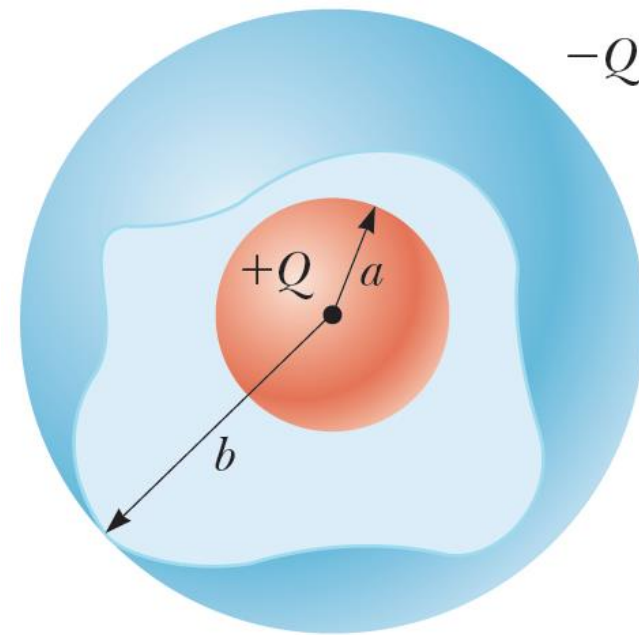


$$C = 4\pi\epsilon_0 \frac{ab}{b-a}$$

➡  $\epsilon_0$  times a quantity with dimension of **length**.

➡ When  $b \rightarrow \infty$ ,  $C = 4\pi\epsilon_0 a$  (**isolated** conducting sphere )

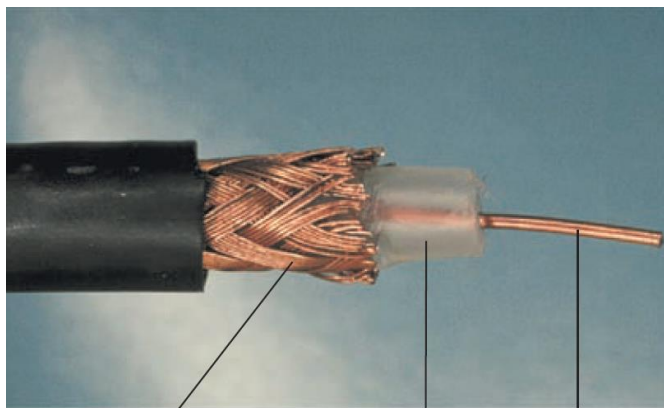
➡ When  $b-a \ll a$ ,  $ab \approx a^2$ ,  $d=b-a$ ,  $A=4\pi a^2$ ,  $C = \epsilon_0 A/d$   
(**parallel-plate** capacitor)



## The Cylindrical Capacitor (P528 Ex.22-2)

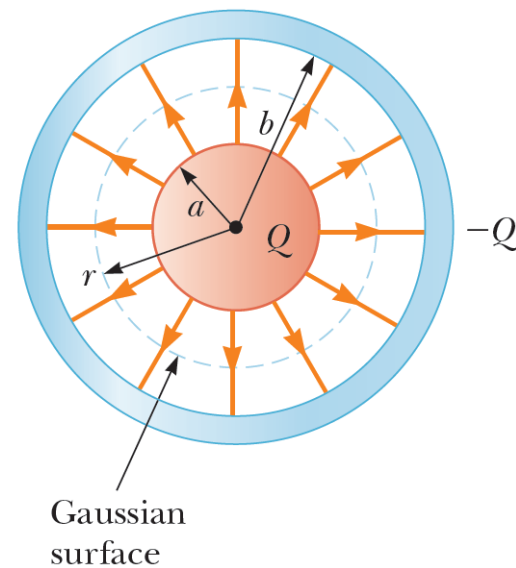
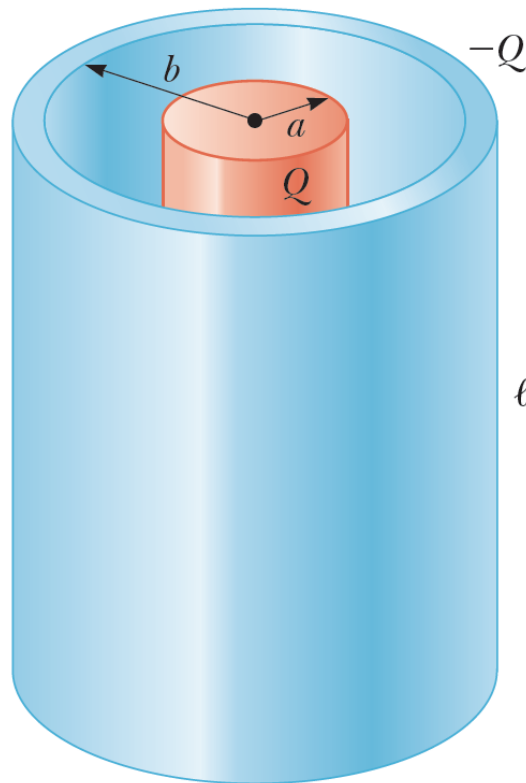


A cylindrical capacitor consists of a cylindrical conductor of radius  $a$  coaxial with a larger cylindrical shell of radius  $b$ . Find the **capacitance** of this device if its length is  $l$ .



Hollow conducting cylinder    Insulator    Central wire

**coaxial cable**



# The Cylindrical Capacitor



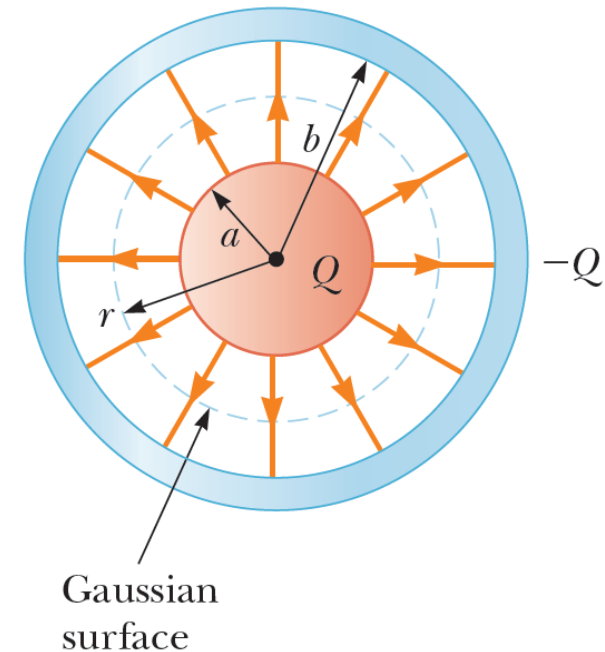
**Solution:** Assume the inner and outer conductors have opposite charges  $+Q$  and  $-Q$ . In the region  $a < r < b$ , we can use **Gauss' law** to determine:

$$\oiint_S \vec{E} \cdot d\vec{A} = E(2\pi r l) = \frac{\lambda l}{\epsilon_0}, \quad E = \frac{\lambda}{2\pi\epsilon_0 r}$$

**The potential difference:**

$$\Delta V = \int_+^- \vec{E} \cdot d\vec{s} = \int_a^b \frac{\lambda}{2\pi\epsilon_0} \frac{dr}{r} = \frac{\lambda}{2\pi\epsilon_0} \ln\left(\frac{b}{a}\right)$$

$$Q = \lambda l, \quad C = \frac{Q}{\Delta V} = \frac{2\pi\epsilon_0 l}{\ln(b/a)}$$



# The Cylindrical Capacitor



$$C = \frac{2\pi\epsilon_0 l}{\ln(b/a)}$$

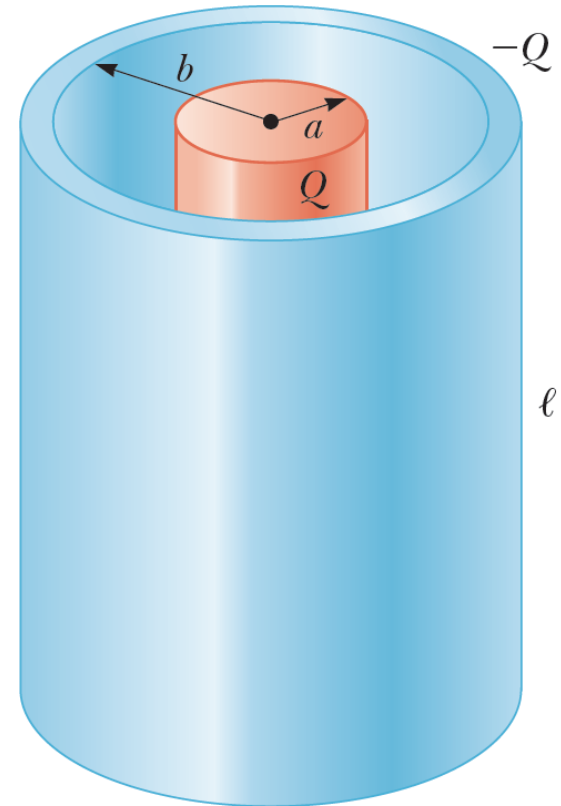
➤ The form of  $\epsilon_0$  times a quantity with dimension of **length**.

➤ When  $d=b-a \ll a$

$$\ln\left(\frac{b}{a}\right) = \ln\left(\frac{a+d}{a}\right) = \ln\left(1 + \frac{d}{a}\right) \approx \frac{d}{a}$$

$$C = \frac{2\pi\epsilon_0 l a}{d} = \epsilon_0 \frac{A}{d}, \quad A = (2\pi a)l$$

(**parallel-plate** capacitor)



## § 22-3 Combinations of Capacitor (P529)

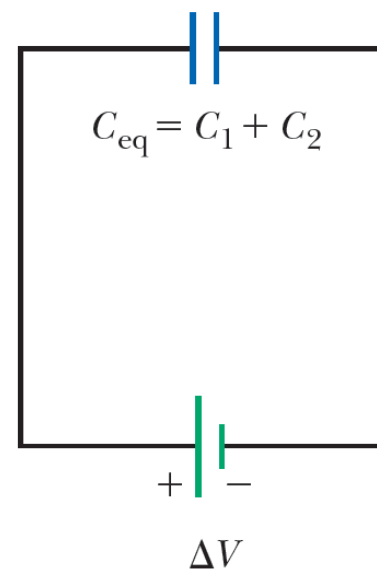
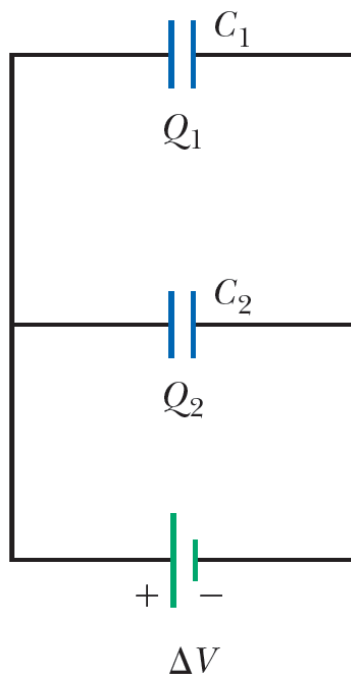
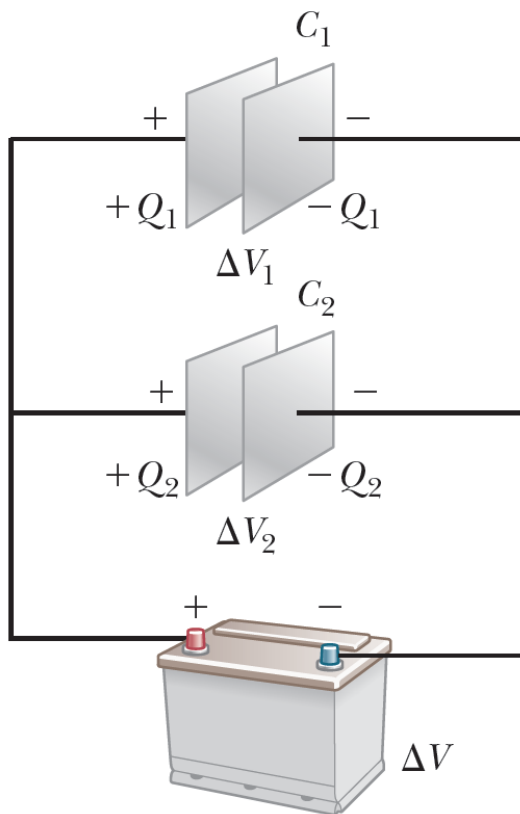


### ■ Parallel Combination

$$C_1 = \frac{Q_1}{\Delta V_1}, \quad C_2 = \frac{Q_2}{\Delta V_2}, \quad C_{eq} = \frac{Q}{\Delta V} = \frac{Q_1 + Q_2}{\Delta V} = C_1 + C_2$$

$$Q = Q_1 + Q_2, \quad \Delta V = \Delta V_1 = \Delta V_2$$

- The equivalent capacitance of a parallel combination of capacitors is the algebraic sum of the individual capacitances.



# Combinations of Capacitor



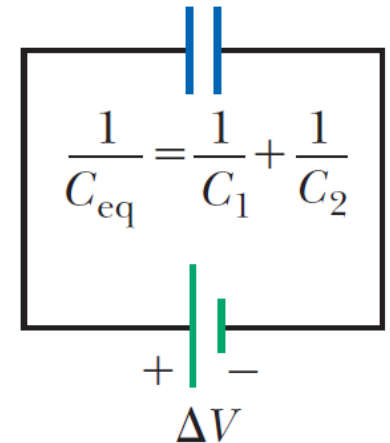
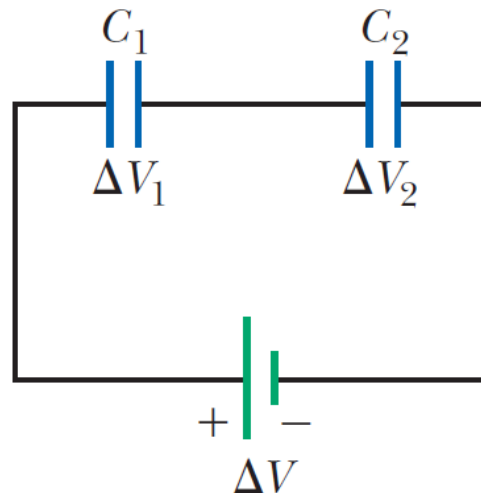
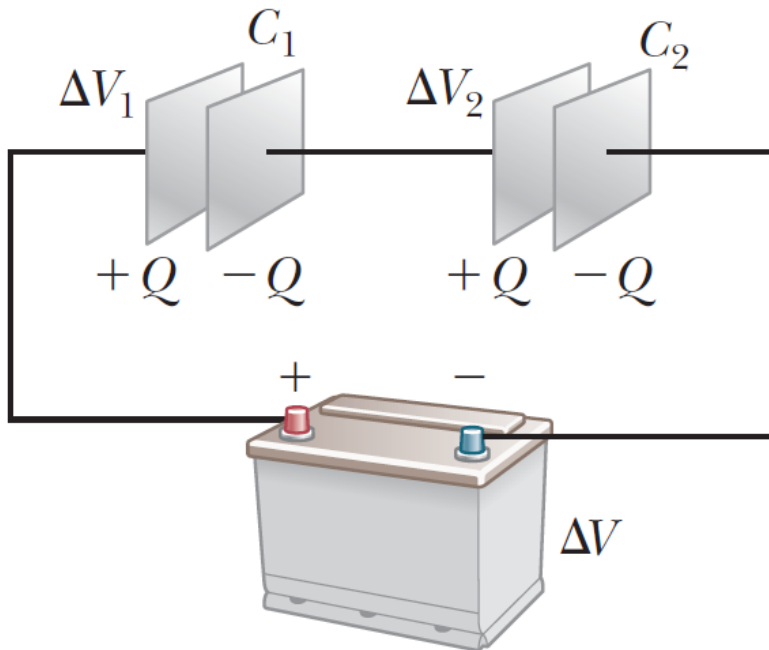
## ■ Series Combination

$$C_1 = \frac{Q_1}{\Delta V_1}, \quad C_2 = \frac{Q_2}{\Delta V_2},$$

$$\frac{1}{C_{eq}} = \frac{\Delta V}{Q} = \frac{\Delta V_1 + \Delta V_2}{Q} = \frac{1}{C_1} + \frac{1}{C_2}$$

$$Q = Q_1 = Q_2, \quad \Delta V = \Delta V_1 + \Delta V_2$$

➔ The inverse of the equivalent capacitance is the algebraic sum of the inverse of the individual capacitances.

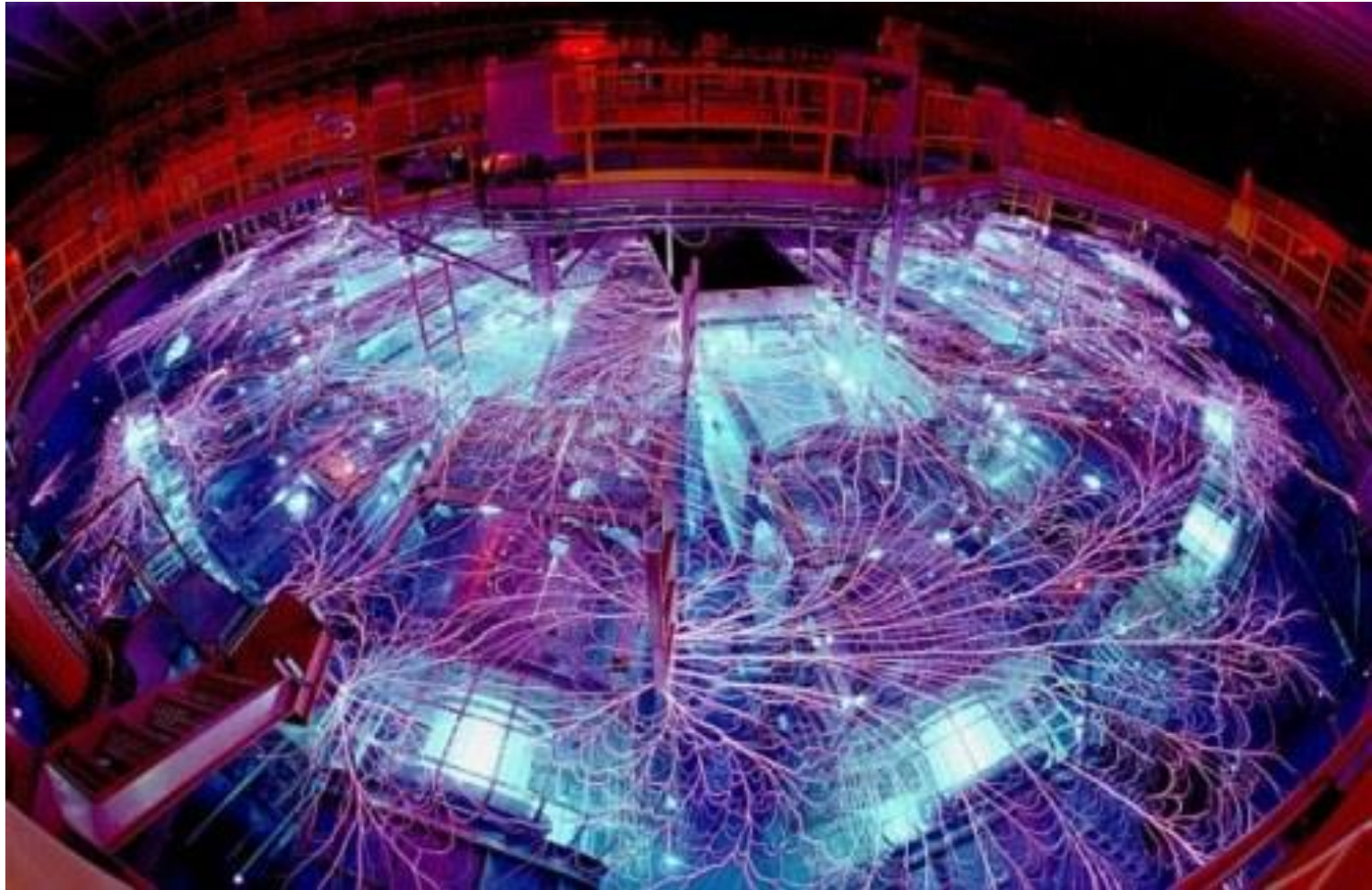




## § 22-4 Electric Energy Storage (P532)



A capacitor can store **charge**, and can also store **energy**!

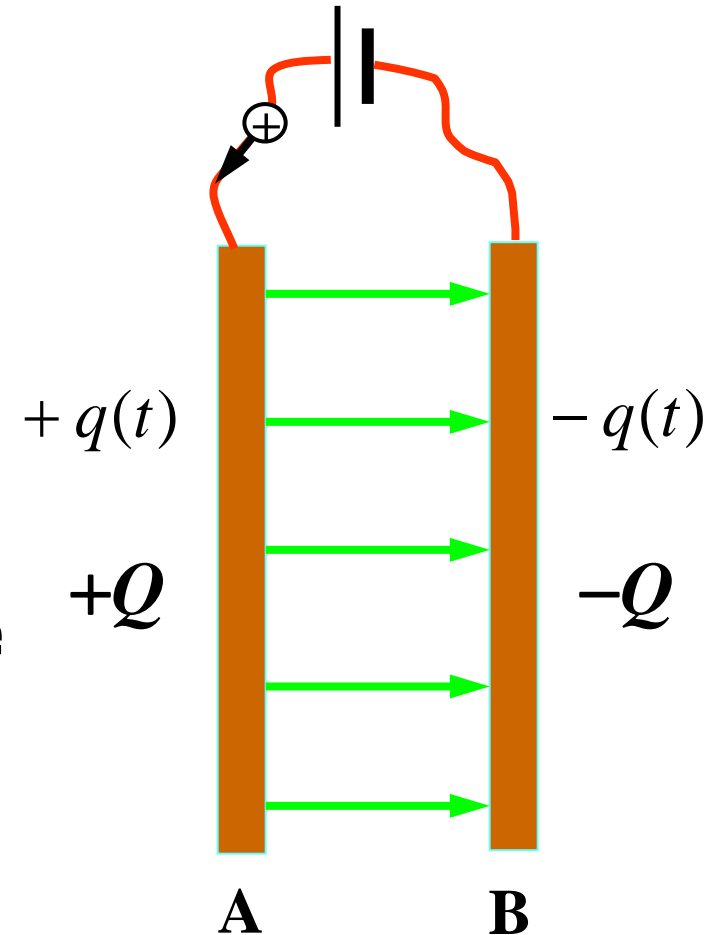


Sandia's **Z machine** is the world's most powerful and efficient laboratory radiation source.

# The potential energy of a charged capacitor



- The potential energy of a charged capacitor
  - ➡ The **energy** stored in a capacitor will be equal to the **work** done to charge it.
  - ➡ We evaluate the work of charging that an external agent continuously pulls charge  $dq$  from negative plate to positive plate until the capacitor has the opposite charge of  $\pm Q$ .

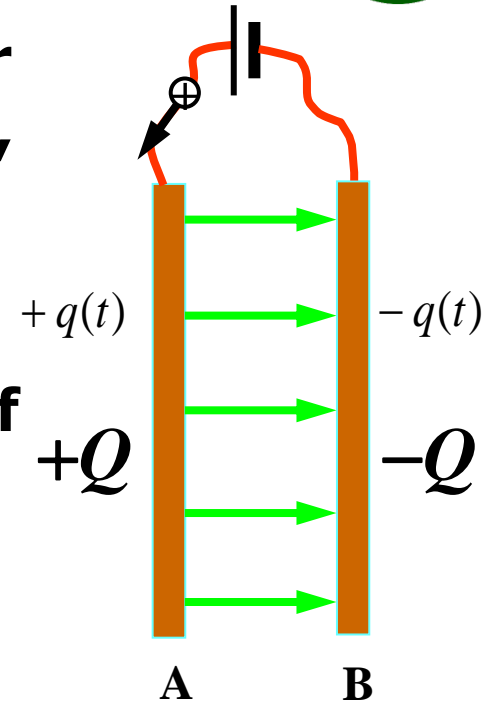


## The potential energy of a charged capacitor



- Suppose that  $q$  is the charge on the capacitor at some instant during this charging process, the potential difference across the capacitor is  $\Delta V = q/C$ . Imaging that the external agent transfers an additional increment of charge  $dq$  from the plate of charge  $-q$  to the plate of charge  $q$ , the resulting small change  $dU$  in the electric potential energy is:

$$dU = (\Delta V) dq = \frac{q}{C} dq$$



- If this process is continued until to charge the capacitor from  $q = 0$  to the final charge  $q = Q$ , the total potential energy is:

$$U = \int dU = \int_0^Q \frac{q}{C} dq = \frac{Q^2}{2C}$$



## The potential energy of a charged capacitor

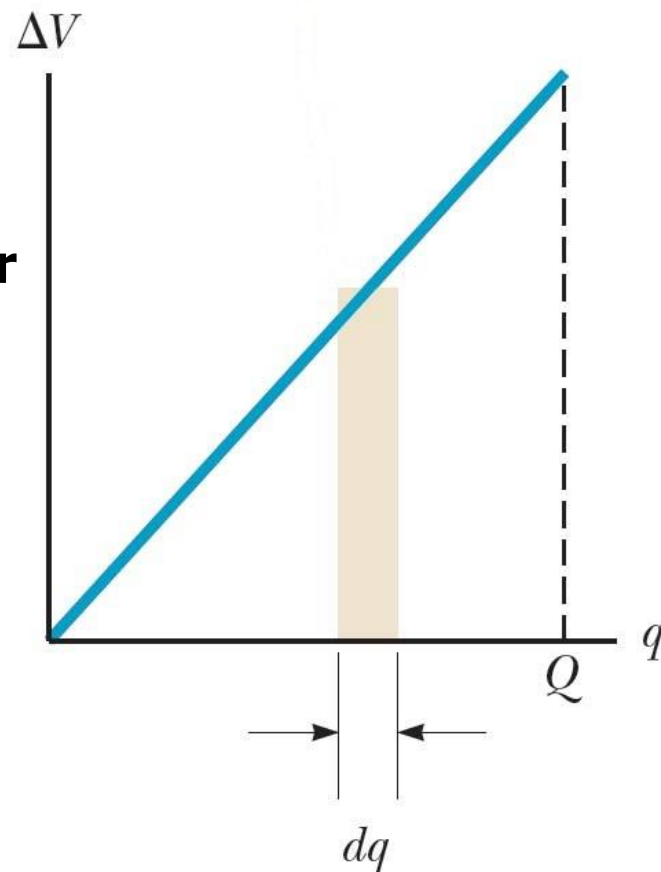


Using  $C = \frac{Q}{\Delta V}$ ,  $U = \frac{1}{2} \frac{Q^2}{C} = \frac{1}{2} C (\Delta V)^2 = \frac{1}{2} Q \Delta V$

**Graphical interpretation:**

A plot of potential difference versus charge for a capacitor is a **straight** line having slope  $1/C$ .

The total **area** under the curve is the **potential energy** stored in a charged capacitor.



## Where does the potential energy reside?

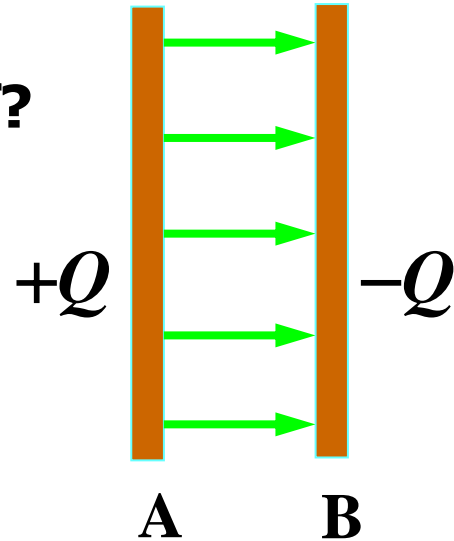


Question: Which one is the storehouse of the energy? The **charges** or the **electric field** itself?

➡ From the equation  $U = Q^2/2C$ , we conclude that the energy relates to the **charging**.

➡ Another point of view:

$$C = \epsilon_0 \frac{A}{d}, \quad U = \frac{1}{2} C (\Delta V)^2 = \left( \frac{1}{2} \epsilon_0 E^2 \right) (Ad)$$
$$\Delta V = Ed,$$



**$U$**  is proportional to the volume  **$Ad$**  between the two plates.

➡ Because the electric field is present in the space between the two plates, the energy is stored in the **electric field** that is present in this region.

## Where does the potential energy reside?



$$U = \left( \frac{1}{2} \varepsilon_0 E^2 \right) (Ad)$$

► The energy density:  $u = \frac{U}{Ad} = \frac{1}{2} \varepsilon_0 E^2$

► If an electric field  $\vec{E}$  exists at any point in empty space, we can think of that point as the site of stored energy in amount of  $\frac{1}{2} \varepsilon_0 E^2$  .

$$U = \int dU = \iiint_V u dV = \iiint_V \left( \frac{1}{2} \varepsilon_0 E^2 \right) dV$$

## Where does the potential energy reside?



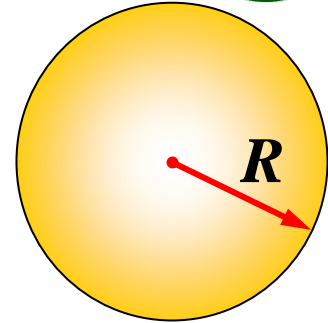
- In the case of **electrostatic** field, we can not answer which one is the storehouse of the energy.
  - ➡ Because in the case of electrostatic field, the electric **field** is always accompanied with the **charge**.
- In the case of **time-varying** electromagnetism field
  - ➡ The electromagnetic wave can exists in the vacuum, whether the charge exists or not.



## Example



How much **electrostatic energy** is stored in the electric field of an **isolated** conducting sphere of radius  $R$  and charge  $Q$ .



**Solution (I):**

The energy stored in the electric **field** is  $U = \iiint_V \left( \frac{1}{2} \epsilon_0 E^2 \right) dV$

The electric field distribution:  $E = \begin{cases} 0 & \text{if } r < R \\ \frac{Q}{4\pi\epsilon_0 r^2} & \text{if } r > R \end{cases}$

**We choose a differential spherical shell of radius  $r$  and thickness  $dr$  and integrate the energies in the shells**

$$U = \iiint \left[ \frac{1}{2} \epsilon_0 \left( \frac{Q}{4\pi\epsilon_0 r^2} \right)^2 \right] dV = \int_R^\infty \left[ \frac{1}{2} \epsilon_0 \left( \frac{Q}{4\pi\epsilon_0 r^2} \right)^2 \right] (4\pi r^2 dr) = \frac{Q^2}{8\pi\epsilon_0} \int_R^\infty \frac{dr}{r^2} = \frac{Q^2}{8\pi\epsilon_0 R}$$

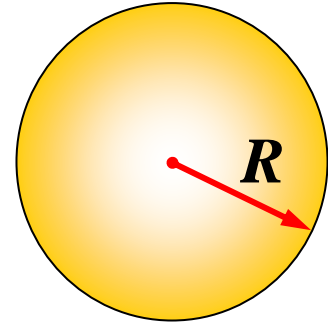
## Example



**Solution (II):**

The energy stored in the spherical **capacitor** is

$$U = \frac{Q^2}{2C}, \quad C = 4\pi\epsilon_0 R, \quad U = \frac{Q^2}{8\pi\epsilon_0 R}$$



**Solution (III):**

The **work** required to bring a differential charge  $dq$  to the sphere is

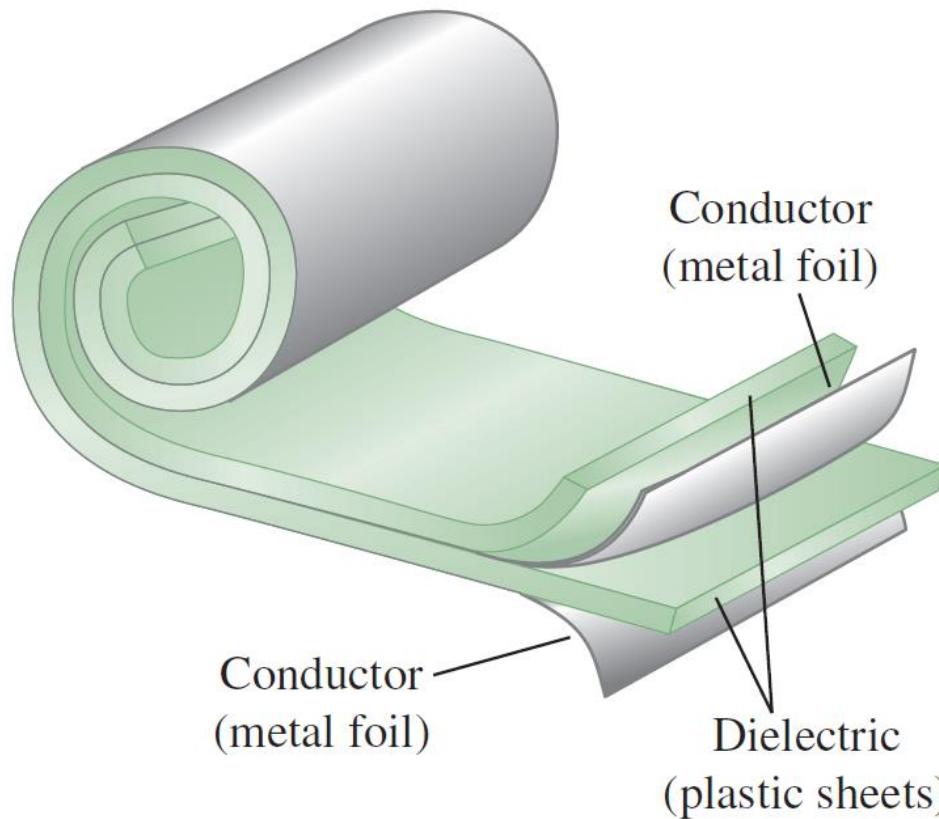
$$dU = Vdq, \quad V = \frac{1}{4\pi\epsilon_0} \frac{q}{R}$$

$$U = \int_0^Q Vdq = \frac{1}{4\pi\epsilon_0 R} \int_0^Q qdq = \frac{Q^2}{8\pi\epsilon_0 R}$$

## **Ch22 Prob. 49, 50, 85 (P543)**

- **Dielectrics** vs. **conductors**

Most capacitors have a nonconducting material, or **dielectric**, between their **conducting** plates.

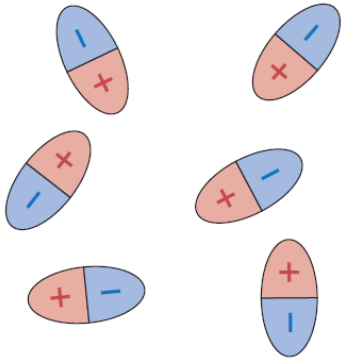


**Why?**

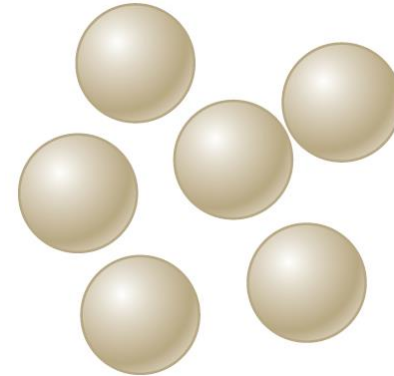
# Polar vs. nonpolar dielectric materials



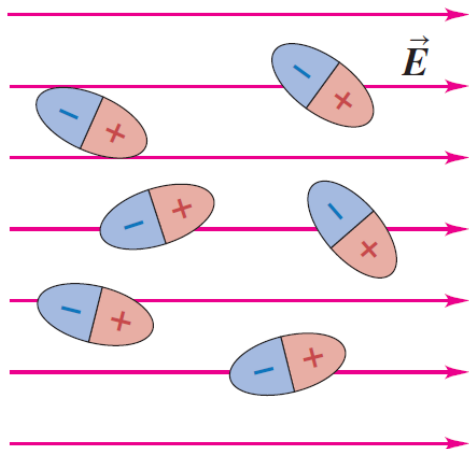
In the absence of an electric field, **Polar** molecules orient randomly.



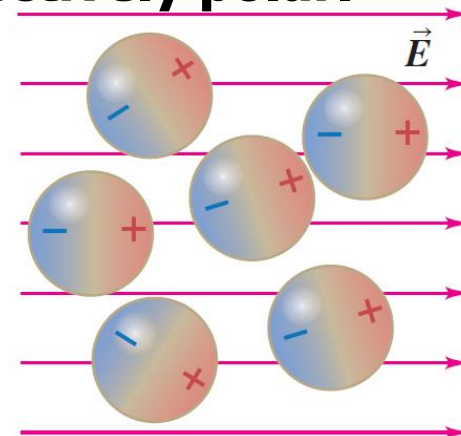
**Nonpolar** molecules are not electric dipoles.



When an electric field is applied, **polar** molecules tend to align with it.



**Nonpolar** molecules are made effectively polar.



# Polar and nonpolar dielectric materials



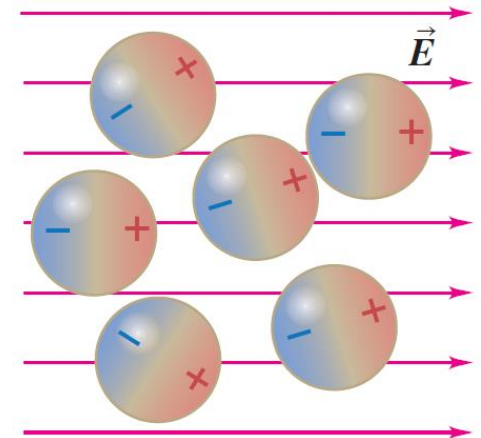
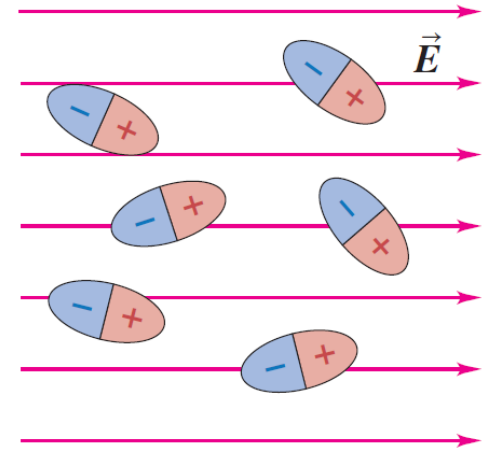
- ➔ **Polar** dielectric material — its molecule has a **permanent** electric dipole moment, such as water.

The external electric field exerts a torque on the dipole that tries to align it with the field.

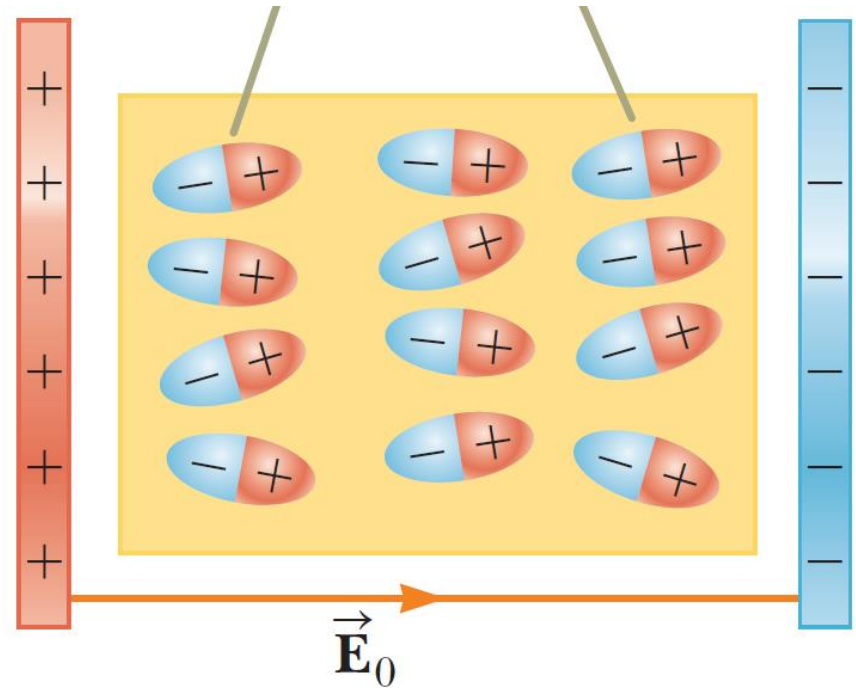
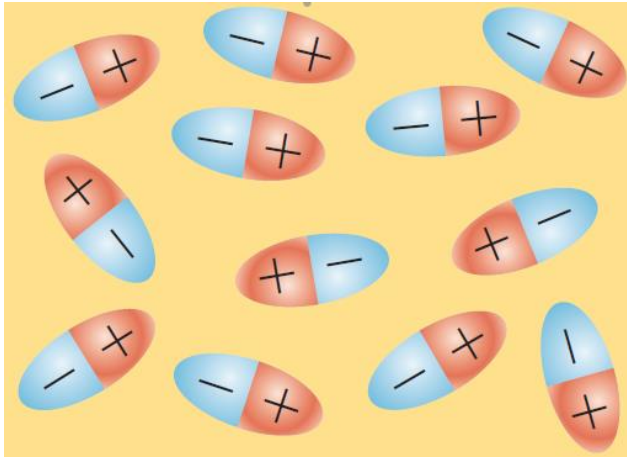
$$\vec{\tau} = \vec{p} \times \vec{E}$$

- ➔ **Nonpolar** dielectric material — its molecule has no permanent electric dipole.

The atom acquires an **induced** dipole moment when the atom is placed in an external electric field.



# Polarization of a dielectric material



Either polar or nonpolar materials are put in an **external** field.

The **induced surface charges** arise as a result of **redistribution** of positive and negative charge within the dielectric material, a phenomenon called **polarization**.



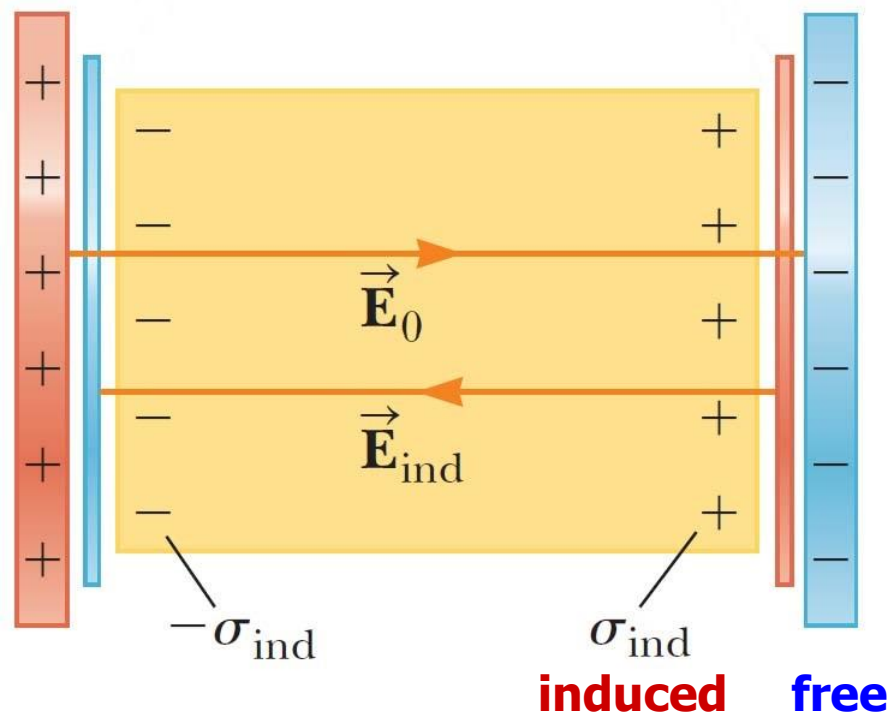
## Induced **polarization** field



- ➔ When a dielectric material is placed in an external applied field  $E_0$ , induced surface charges  $q_{\text{ind}}$  appear that tend to **weaken** the original field  $E_0$  by a polarization field  $E_{\text{ind}}$  within the material. For a **linear** material, the net field inside the material is

$$\vec{E} = \vec{E}_0 + \vec{E}_{\text{ind}}$$

- ➔ The charge  $q_0$ , the origin of  $E_0$ , that resides in the conductors is called **free charge**, and induced charge  $q_{\text{ind}}$  that resides in the surface of dielectric materials, that not free to move and bound to a molecule, is called **induced bound charge**.



# Permittivity



$$\vec{E} = \vec{E}_0 + \vec{E}_{\text{ind}},$$

$$E = \frac{E_0}{\kappa}$$

➡  $\kappa$  is called the **relativity** permittivity (dielectric constant) (相对介电常数), which is greater than 1.

**Absolute** permittivity of the dielectric:

$$\epsilon = \kappa \epsilon_0$$

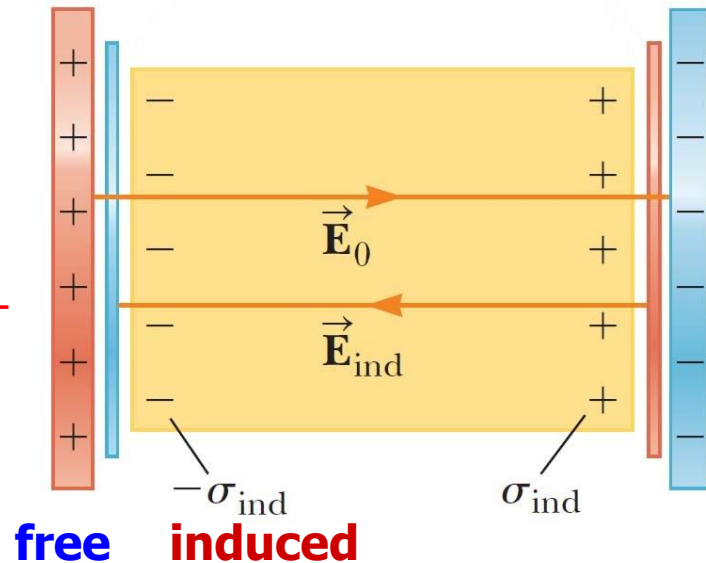
The electric field

within the dielectric:  $E = \frac{E_0}{\kappa} = \frac{\sigma_0}{\kappa \epsilon_0} = \frac{\sigma_0}{\epsilon}$

Induced charge density:

$$E = E_0 - E_{\text{ind}}$$

$$\frac{\sigma_0}{\kappa \epsilon_0} = \frac{\sigma_0}{\epsilon_0} - \frac{\sigma_{\text{ind}}}{\epsilon_0} \Rightarrow \sigma_{\text{ind}} = \left(1 - \frac{1}{\kappa}\right) \sigma_0$$



# The dielectric strength



## ➔ The dielectric strength: $E_{\text{break}}$

If we apply a large enough electric field to an insulator, we can ionize atoms or molecules of the insulator and thus create a condition for electric charge to flow, as in a conductor. The field necessary for the **breakdown** of the insulator is called the **dielectric strength**.

Material	Dielectric Constant $\kappa$	Dielectric Strength <sup>a</sup> ( $10^6$ V/m)
Air (dry)	1.000 59	3
Bakelite	4.9	24
Fused quartz	3.78	8
Mylar	3.2	7
Neoprene rubber	6.7	12
Nylon	3.4	14
Paper	3.7	16
Paraffin-impregnated paper	3.5	11
Polystyrene	2.56	24
Polyvinyl chloride	3.4	40
Porcelain	6	12
Pyrex glass	5.6	14
Silicone oil	2.5	15
Strontium titanate	233	8
Teflon	2.1	60
Vacuum	1.000 00	—
Water	80	—

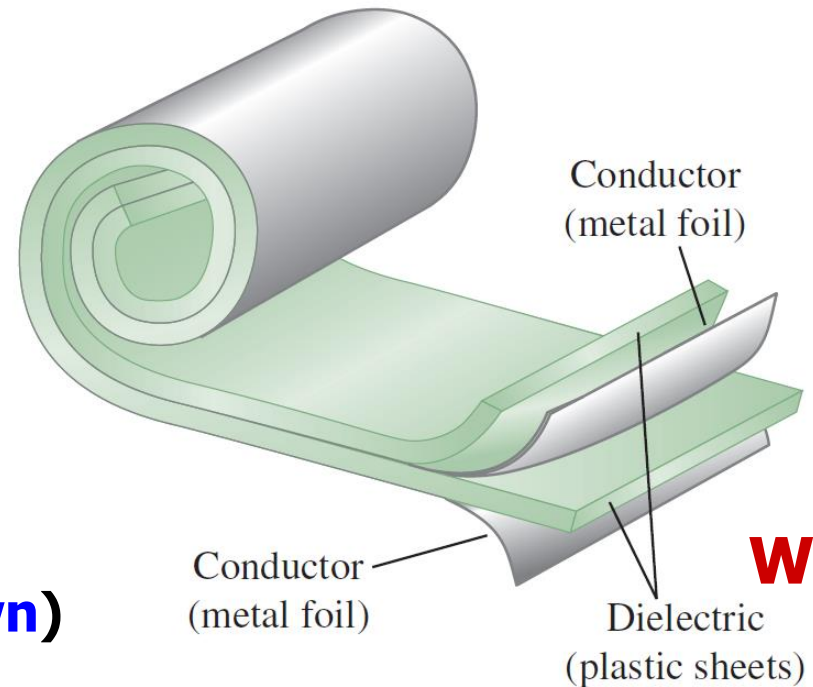
## § 7 Capacitors with Dielectrics (P534)



Placing a solid dielectric between the plates of a capacitor serves **three** functions.

First, it solves the **mechanical** problem of maintaining two large metal sheets at a very small separation without actual contact.

Second, using a dielectric increases the maximum possible potential difference between the capacitor plates. (**without dielectric breakdown**)



Third, the **capacitance** of a capacitor of given dimensions is **greater** when there is a dielectric material between the plates than when there is vacuum.

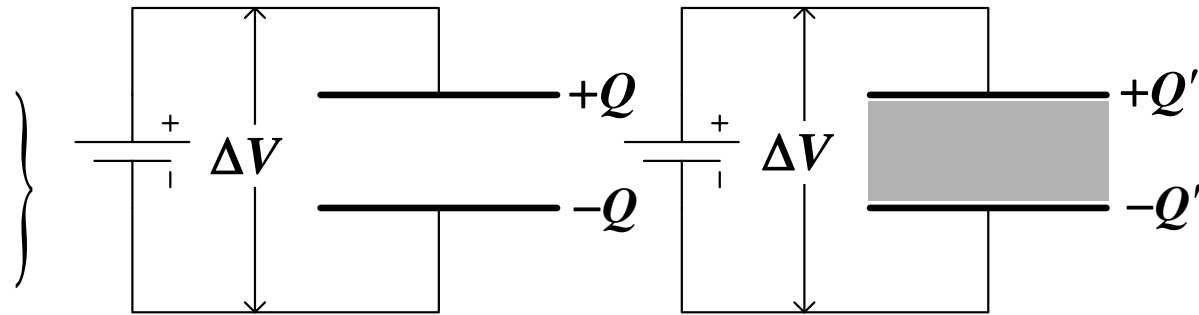
# Capacitors with Dielectrics



Two identical capacitors, filling one with a dielectric material and leaving the other with air between its plates

- When both capacitors are connected to batteries with the **same potential difference**.

$$\Delta V = \Delta V' \Rightarrow E = E'$$
$$E = \frac{Q}{\epsilon_0 A}, \quad E' = \frac{1}{\kappa} \frac{Q'}{\epsilon_0 A}$$



$$Q' = \kappa Q$$

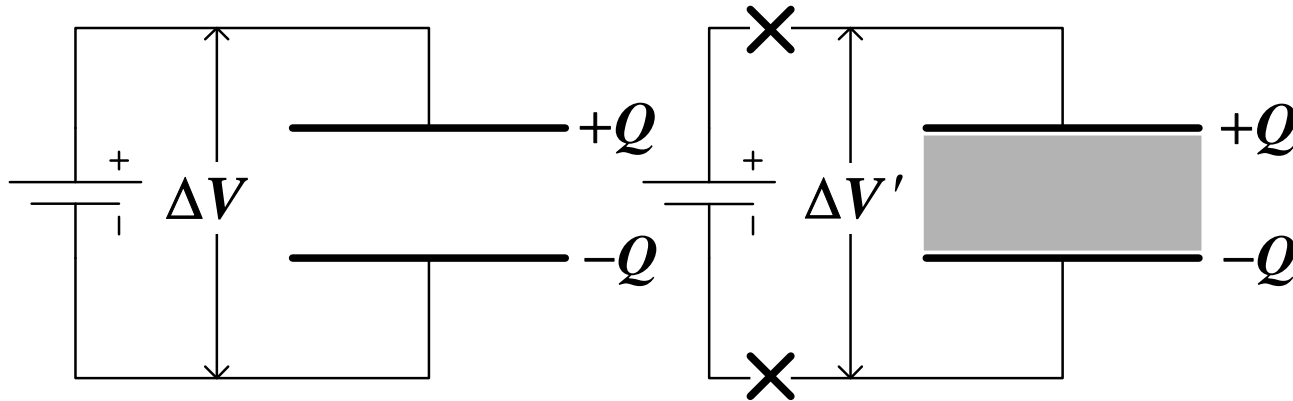
$$C' = \frac{Q'}{\Delta V'} = \frac{\kappa Q}{\Delta V} \Rightarrow C' = \kappa C \Rightarrow C' = \frac{\kappa \epsilon_0 A}{d} = \frac{\epsilon A}{d}$$

$\epsilon = \kappa \epsilon_0$  permittivity

# Capacitors with Dielectrics



- When both are disconnected the batteries with the **same charge**.



$$Q' = Q, \quad E' = \frac{E}{\kappa},$$

$$\Delta V' = E'd = \frac{Ed}{\kappa} = \frac{\Delta V}{\kappa}$$

$$C' = \frac{Q'}{\Delta V'} = \kappa \frac{Q}{\Delta V} = \kappa C = \frac{\epsilon A}{d}$$

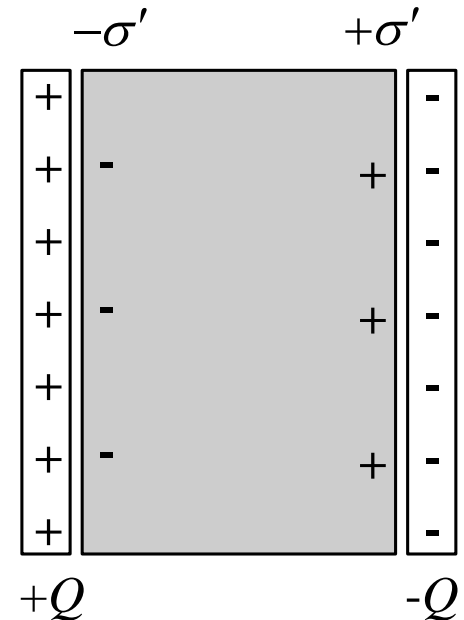
# The electric field **energy** stored in a capacitor with dielectric



## ■ The electric field energy stored in a capacitor with dielectric

$$U = \frac{Q^2}{2C} = \frac{Q^2 d}{2\kappa\epsilon_0 A} = \frac{1}{2} \kappa\epsilon_0 \left( \frac{Q}{\kappa\epsilon_0 A} \right)^2 (Ad)$$

$$E = \frac{E_0}{\kappa} = \frac{1}{\kappa} \frac{\sigma}{\epsilon_0} = \frac{Q}{\kappa\epsilon_0 A}, \quad U = \frac{1}{2} \kappa\epsilon_0 E^2 (Ad)$$



## ■ The electric field energy density in dielectric materials

$$u = \frac{1}{2} \kappa\epsilon_0 E^2 = \frac{1}{2} \epsilon E^2$$



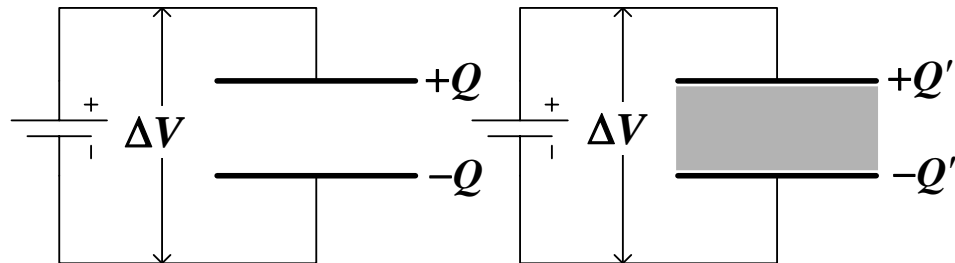
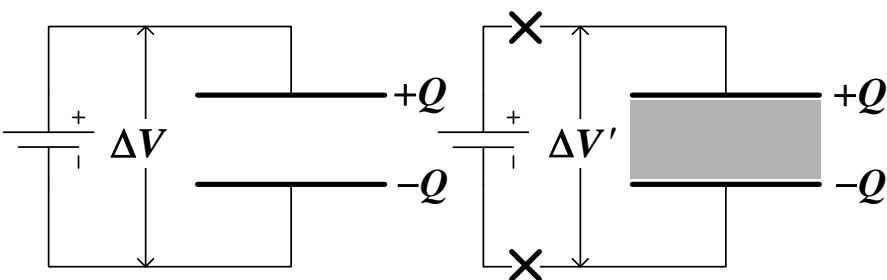
## Example



In following two cases, find the **electric field energy** stored in a parallel-plate capacitor **before** and **after** the dielectric is inserted. The capacitor without dielectric is  $C_0$ , and dielectric material has dielectric constant  $\kappa$ .

(1) At beginning, the capacitor, with empty, is connected to the battery of voltage  $\Delta V$ . The battery is then removed, and the capacitor is fill with the dielectric material.

(2) From beginning to end, the capacitor is always connected to the battery of voltage  $\Delta V$  ;

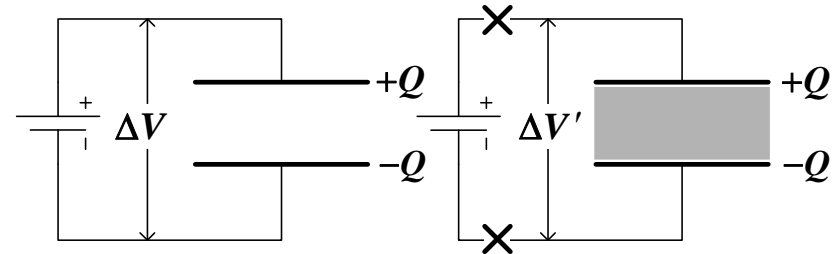


# Solution



$$U = \frac{1}{2} \frac{Q^2}{C} = \frac{1}{2} C (\Delta V)^2 = \frac{1}{2} Q \Delta V$$

(1) Before and after removing the battery, the **charges** in the capacitor are the **same** and the **capacitance increases**.



Before inserting the dielectric:  $U_{before} = \frac{1}{2} \frac{Q^2}{C_0}$

After inserting the dielectric:  $U_{after} = \frac{1}{2} \frac{Q^2}{\kappa C_0} = \frac{U_{before}}{\kappa}$

$$\Delta U = U_{after} - U_{before} = (1 - \kappa) U_{after} < 0$$

The dielectric, when inserted, **is pulled into** the device. To keep the dielectric from accelerating, an external agent must do **negative** work on the dielectric.

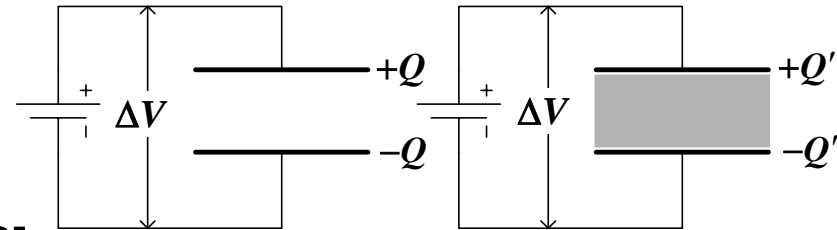
# Solution



$$U = \frac{1}{2} \frac{Q^2}{C} = \frac{1}{2} C (\Delta V)^2 = \frac{1}{2} Q \Delta V$$

**(2) Before inserting the dielectric:**

$$U_{\text{before}} = \frac{1}{2} C_0 (\Delta V)^2$$



**After inserting the dielectric:**

$$U_{\text{after}} = \frac{1}{2} \kappa C_0 (\Delta V)^2 = \kappa U_{\text{before}}$$

$$\Delta U = U_{\text{after}} - U_{\text{before}} = (\kappa - 1) U_{\text{before}} > 0$$

## **Ch22 Prob. 65, 87 (P544)**