

EEE6212

“Semiconductor Materials” -Quantum Mechanics

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Outline

- Purpose
- Time independent Schroedinger equation
- Infinite Well – derivations
- Application – approximating a semiconductor QW
- In-plane dispersion
- Summary

Purpose

- Quantum structures are at the heart of many (most?) semiconductor electronic or photonic devices
- Will discuss how the time independent Schroedinger equation is employed to solve the modes of an infinite quantum well
- Look at the application of this approximation to a semiconductor quantum well – see xls!
- Discuss the kinetic energy of carriers in-plane

Time Independent Schroedinger

- Wave mechanics analog to Hamilton's formulation in classical mechanicsTime-independent potentials – self explanatory...
- KE and PE sum to total energy – eigenvalues of linear operators

$$\mathcal{T}\psi + \mathcal{V}\psi = E\psi$$

- Eigenfunction Ψ describes state of the system
- Operators-

- T- kinetic energy
- V- potential energy
- E- total energy
- p - linear momentum

$$\mathcal{T} = \frac{\mathcal{P}^2}{2m}$$

$$\mathcal{P} = -i\hbar\nabla = -i\hbar\left(\frac{\partial}{\partial x}\mathbf{i} + \frac{\partial}{\partial y}\mathbf{j} + \frac{\partial}{\partial z}\mathbf{k}\right)$$

TISE (2)

- Inserting more explicit form of T

$$-\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \psi + V(x, y, z)\psi = E\psi$$

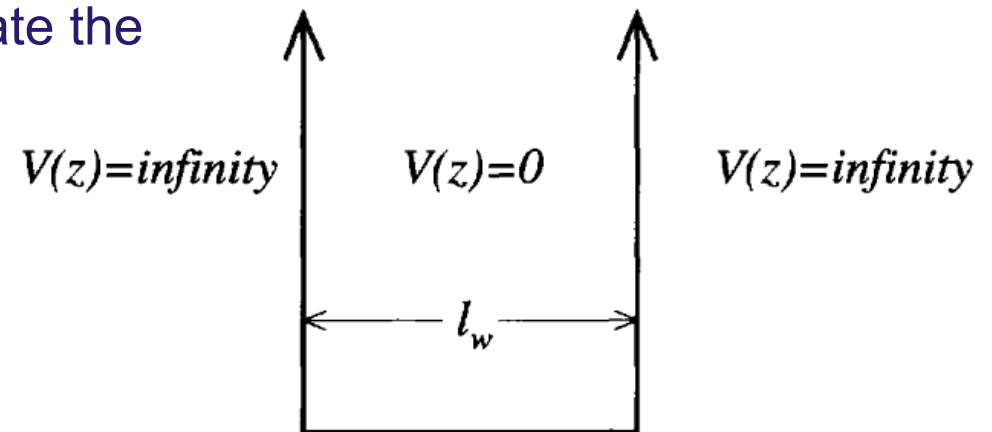
- $V(x, y, z)$ – potential energy of the system in spatial coordinates
- Restrict ourselves to 1D – 1D TISE for a particle of mass m is;-

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial z^2} \psi(z) + V(z)\psi(z) = E\psi(z)$$

Infinite Well TISE

- Outside the well, $V(z) = \infty$, So $\Psi(z) = 0$
– wavefunction does not penetrate the barriers

- Inside the well, set datum level
 $V(z) = 0$



- TISE simplifies to....

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial z^2} \psi(z) = E\psi(z)$$

Infinite Well TISE (2)

- Need to remember maths for differential equations....
- This 2nd order DE implies that the solution for Ψ is a linear combination of functions $f(z)$, which when differentiated twice, give $-f(z)$
- Let's try.... $\psi(z) = A \sin kz + B \cos kz$

$$\rightarrow \frac{\hbar^2 k^2}{2m} (A \sin kz + B \cos kz) = E (A \sin kz + B \cos kz)$$

$$\rightarrow \therefore \frac{\hbar^2 k^2}{2m} = E$$

Infinite Well TISE (3)

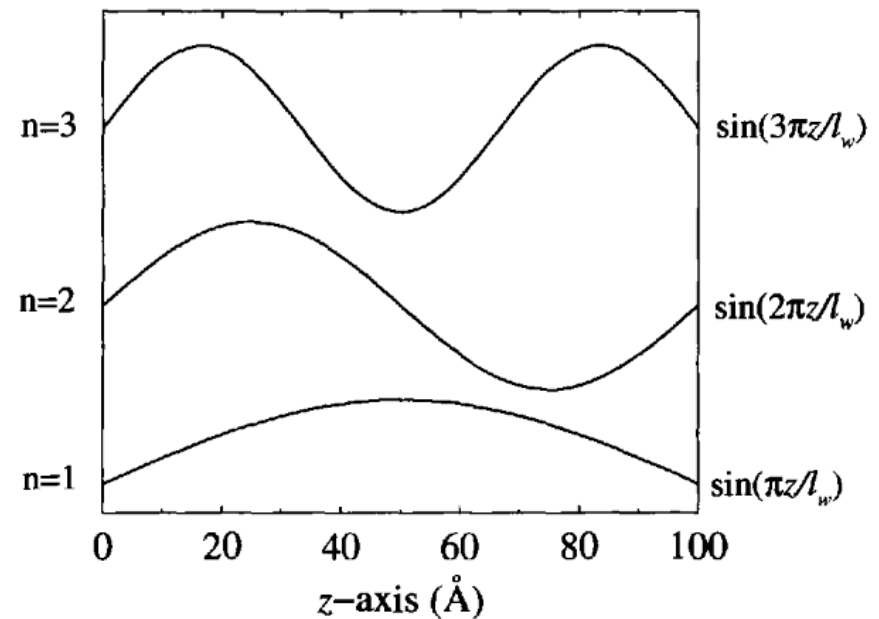
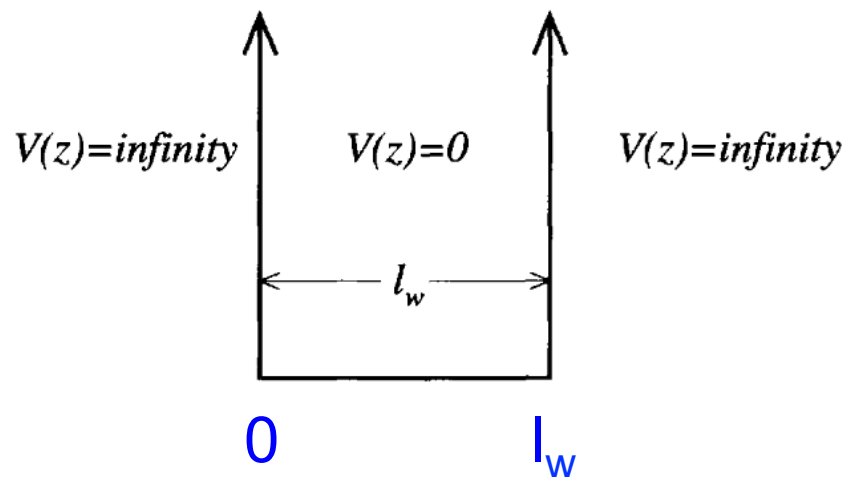
- Need to consider boundary conditions to determine the constant k

$$\mathcal{T} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial z^2} \psi(z) \quad \rightarrow \quad \mathcal{T} = -\frac{\hbar^2}{2m} \frac{\partial}{\partial z} \left(\frac{\partial}{\partial z} \psi(z) \right)$$

- $\Psi(z)$ is continuous if we have finite values for KE (if not continuous $d\Psi/dz$ contains poles)
- $\Psi(z)$ is continuous, and is zero in the barrier, so is zero at the edges of the well

Infinite Well TISE (4)

- Set origin as left hand edge of the well



$$\psi(z) = A \sin kz + B \cos kz$$

Infinite Well TISE (5)

- And $\Psi(0) = \Psi(l_w) = 0$, so $k = \frac{\pi n}{l_w}$
- Where n is an integer representing a series of solutions

- Substituting into $\frac{\hbar^2 k^2}{2m} = E$

$$E_n = \frac{\hbar^2 \pi^2 n^2}{2m l_w^2}$$

Infinite Well TISE (5)

- Last step is to determine A coefficient in $\psi(z) = A \sin kz$.

- Particle is confined to the well so $\int_0^{l_w} \psi^*(z) \psi(z) dz = 1$

- Which gives

$$A = \sqrt{(2/l_w)}$$

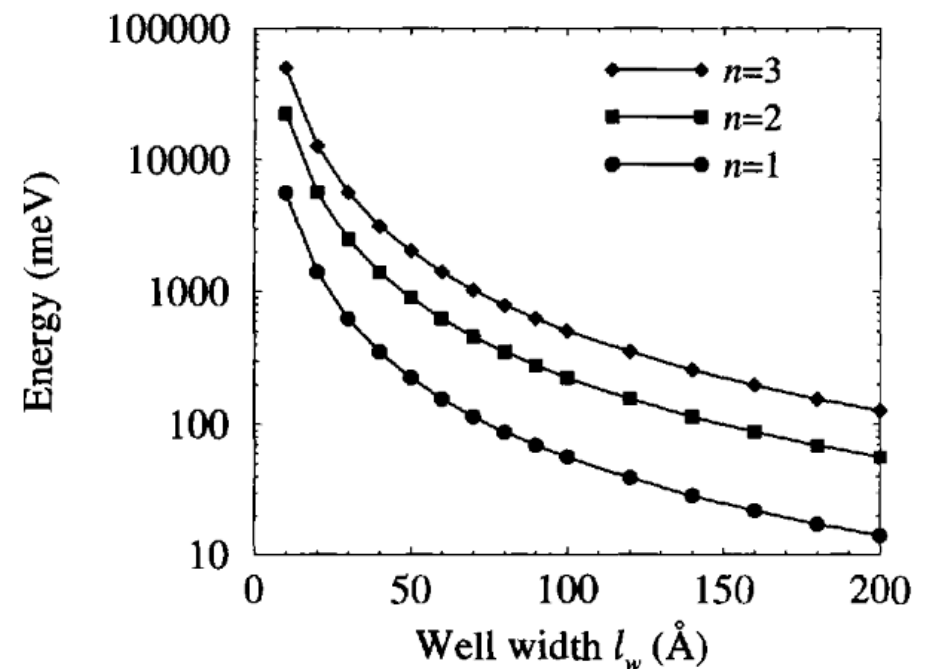
$$\psi_n(z) = \sqrt{\frac{2}{l_w}} \sin \left(\frac{\pi n z}{l_w} \right)$$

Effect of Various Parameters

- Confinement energy $E_n = \frac{\hbar^2 \pi^2 n^2}{2m l_w^2}$

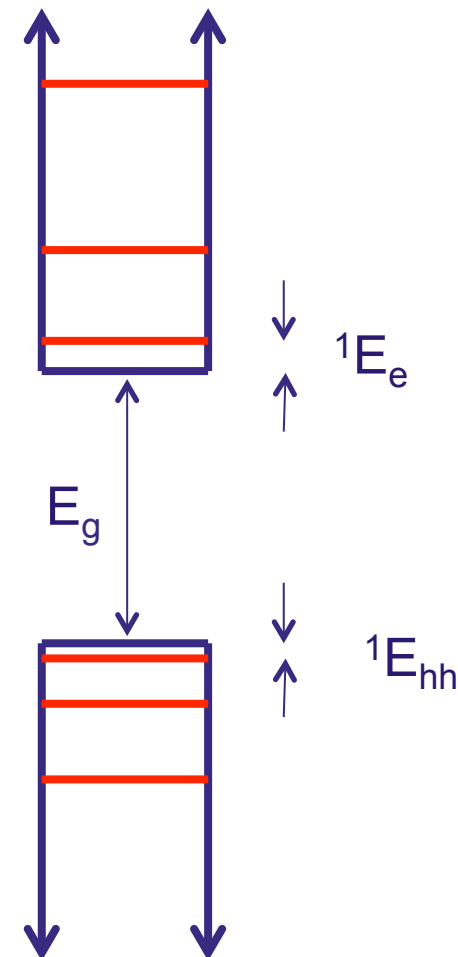
- $m \uparrow \quad E_n \downarrow$
- $l_w \uparrow \quad E_n \downarrow \downarrow$
- $n \uparrow \quad E_n \uparrow \uparrow$

Example for electrons in infinite GaAs QW



Infinite QW - Approximation

- Consider a GaAs quantum well in an infinite barrier
- What is emission/absorption energy of $e1hh1$ transition?
- $E_g = 1.42 \text{ eV}$
- $m_e = 0.06 m^*$
- $m_{hh} = 0.51 m^*$



In-plane Dispersion

- We have considered an infinite potential in z , how about the motion of the particle in x, y ? Need to consider all terms of KE operator-

$$-\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \psi + V(z)\psi = E\psi$$

- $V(z)$ can be written as the sum of independent functions
 $V = V(x) + V(y) + V(z)$, so eigenfunction of the system can be written as product

$$\psi(x, y, z) = \psi_x(x)\psi_y(y)\psi_z(z)$$

- So -

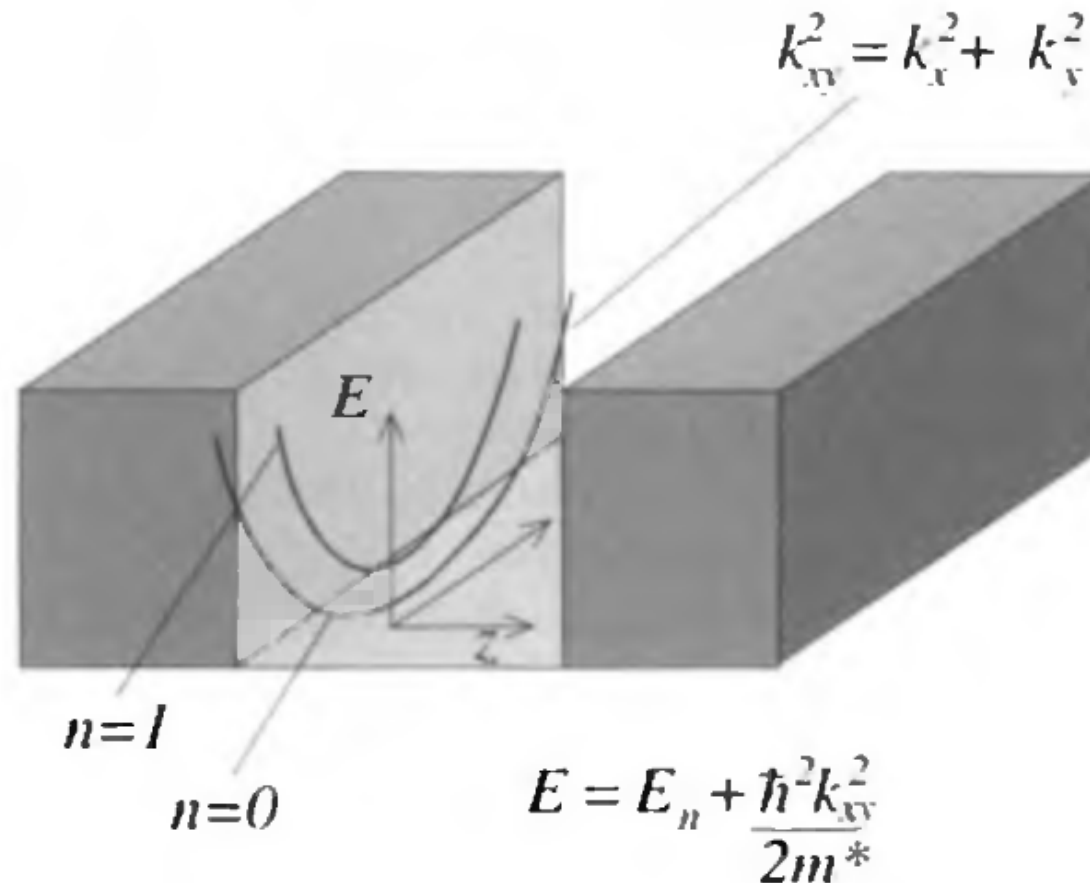
$$-\frac{\hbar^2}{2m} \left(\frac{\partial^2 \psi_x}{\partial x^2} \psi_y \psi_z + \frac{\partial^2 \psi_y}{\partial y^2} \psi_x \psi_z + \frac{\partial^2 \psi_z}{\partial z^2} \psi_x \psi_y \right) + V(z)\psi_x \psi_y \psi_z = E\psi_x \psi_y \psi_z$$

In-Plane Dispersion (2)

- Three distinct contributions to Total Energy E in each of the orthogonal axes ; $E = E_x + E_y + E_z$
- The motion is decoupled – an equation of motion for each axis

$$\begin{aligned}
 & -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_x}{\partial x^2} \psi_y \psi_z = E_x \psi_x \psi_y \psi_z \\
 & -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_y}{\partial y^2} \psi_x \psi_z = E_y \psi_x \psi_y \psi_z \\
 & -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_z}{\partial z^2} \psi_x \psi_y + V(z) \psi_x \psi_y \psi_z = E_z \psi_x \psi_y \psi_z
 \end{aligned}
 \quad \rightarrow \quad
 \begin{aligned}
 & -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_x}{\partial x^2} = E_x \psi_x \\
 & -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_y}{\partial y^2} = E_y \psi_y \\
 & -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_z}{\partial z^2} + V(z) \psi_z = E_z \psi_z
 \end{aligned}$$

In-Plane Dispersion (3)



Summary

- Introduced motivation to understand quantum structures
- Derived confinement energy and form of the wavefunction for particles in an infinite quantum well
- Applied this approximation to a semiconductor quantum well
- Noted how the transition energy/wavelength may be a strong function of well-width
- Seen how the motion (and hence KE) is decoupled in the three orthogonal axes