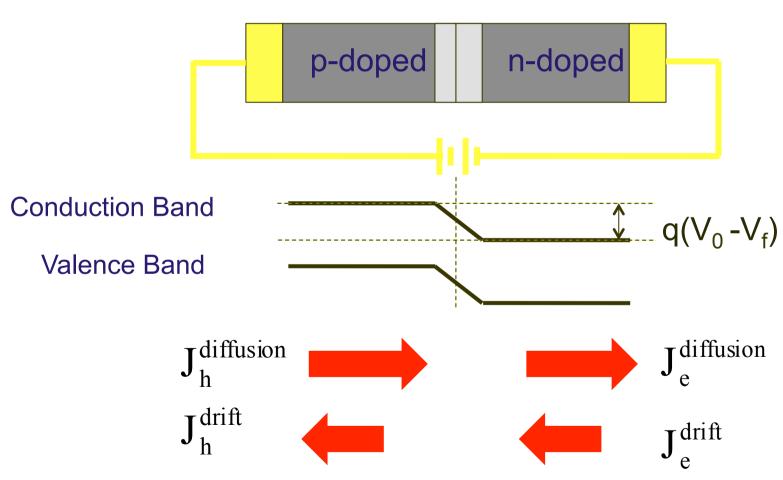


#### Lecture 12

- Forward Bias in p-n Junction Qualitative
- Quantitative derivation
- Diode equation



# Forward Bias, $V_f$





# Forward Bias, $V_f$

- Applied voltage changes the potential barrier and thus *E*-field within junction region since we have forward bias the potential barrier is reduced
- The electric field in the transition region reduces
- This reduces the transition region width (need fewer "exposed" ionized dopants to achieve this lower *E*-field
- Diffusion Current potential barrier smaller so increased diffusion current
- Drift Current essentially same as zero bias very few minority carriers to contribute to drift–so very small

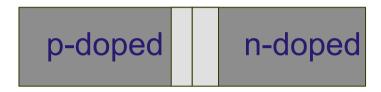


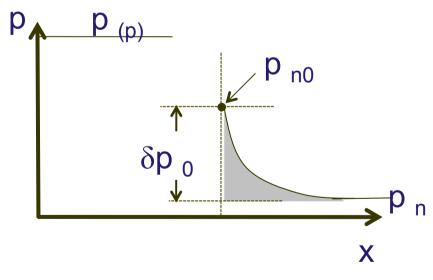
## Our Job Today.....

- Calculate diffusion current of carriers (electrons into pdoped, holes into n-type) at fixed bias
- Recall discussion on steady state injection of one (minority) carrier
- We will calculate current as a function of applied voltage
- We will calculate charge injected per unit time by electrons and holes - work on one charge carrier type and modify equations for other carrier



#### Consider Hole Diffusion





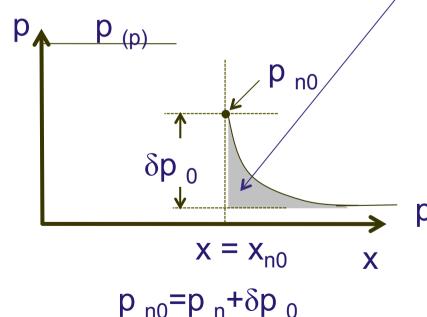
 $p_{n0} = p_n + \delta p_0$ 

- Holes diffusing over barrier become excess minority carriers where they diffuse away from junction and recombine with majority electrons
- If bias is maintained we will achieve a continuous injection of minority holes (considered before) which will sustain an exponentially decreasing excess minority hole concentration
- Similar story for electron diffusion



# Calculate $I(V_f)$





- Need to calculate area under this curve to get the charge  $Q_p$
- The diffused (excess minority) hole concentration is a function of hole concentration in p-type
- We will use minority carrier lifetime as it will describe characteristic diffusion length
- Do this for holes and electrons and add together



## Calculating $\delta p_0$

Recall Equation for built-in potential (NOTE: zero applied bias case)

$$V_0 = \frac{k_B T}{e} ln \left( \frac{p_p}{p_n} \right) \qquad \qquad p_n = p_p exp \left[ \frac{-eV_0}{k_B T} \right]$$

Modifying to include the applied forward bias,  $V_f$ , we get this in terms of  $p_{n0}$ , the minority carrier density at  $x = x_{n0}$ .

$$p_{n0} = p_p exp \left[ \frac{-e(V_0 - V_f)}{k_B T} \right]$$

Note  $p_{n0} > p_n$  due to bias forward bias  $V_f$ ,



#### Continued

$$p_{n0}=p_n+\delta p_0$$
 so  $\delta p_0=p_{n0}-p_n$ 

We have seen both of these previously, so:

Using 
$$p_n = p_p exp \left[ \frac{-eV_0}{k_B T} \right]$$

$$\partial p_0 = p_p \exp\left[\frac{-e(V_0 - V_f)}{k_B T}\right] - p_p \exp\left[\frac{-eV_0}{k_B T}\right]$$

$$\partial p_0 = p_p \exp\left[\frac{eV_f}{k_B T}\right] \times p_p \exp\left[\frac{-eV_0}{k_B T}\right] - p_p \exp\left[\frac{-eV_0}{k_B T}\right]$$



# Continued (2)

Giving

$$\partial p_0 = p_n \left| exp \left[ \frac{eV_f}{k_B T} \right] - 1 \right|$$

And with a similar treatment for electron diffusion

$$\partial n_0 = n_p \left[ exp \left[ \frac{eV_f}{k_B T} \right] - 1 \right]$$

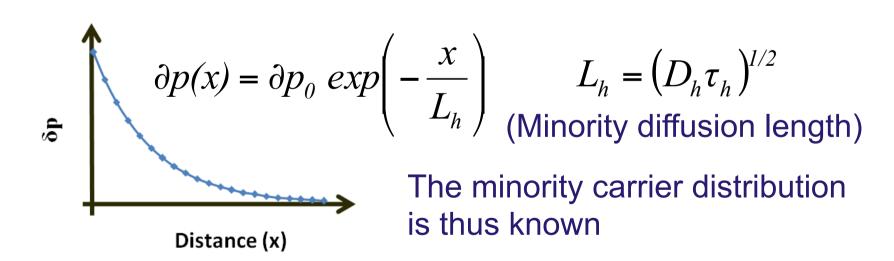
These show that the excess minority carrier concentration injected across the junction increases exponentially with forward bias  $V_f$ .



# Continue with Deriving $I(V_f)$

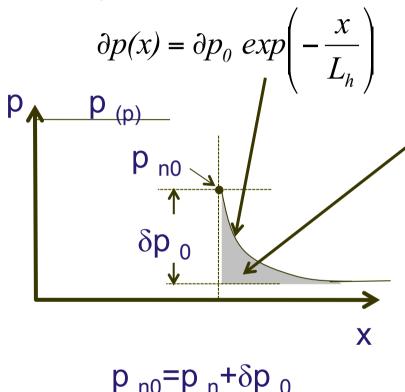
We now need to include time and distances

Remembering minority carrier diffusion length – we can relate diffusion distance to minority carrier lifetime





#### From a previous lecture



#### Injected hole charge

$$Q_{p} = eA \int_{0}^{\infty} \delta p(x) \delta x$$

$$Q_{p} = eA \delta p_{0} \int_{0}^{\infty} exp\left(-\frac{x}{L_{h}}\right) dx$$

$$Q_{p} = eA \delta p_{o} L_{h}$$

Similarly for electrons diffusing into the p-type material, charge injected,  $Q_e$ , is;

$$Q_e = eA \delta n_o L_e$$



### Definite Integral Step

$$Q_p = eA\delta p_0 \int_0^\infty exp\left(-\frac{x}{L_h}\right) dx$$

$$Q_{p} = qA\delta p_{0} \left[ \left[ -L_{h} exp \left( -\frac{x}{L_{h}} \right) \right]_{x=\infty} - \left[ -L_{h} exp \left( -\frac{x}{L_{h}} \right) \right]_{x=0} \right)$$

$$\sum_{x} \underline{Lim}_{-\infty} \exp(x) = 0$$

$$\lim_{x \to 0} \exp(x) = 1$$

$$Q_p = eA\delta p_0 L_h$$



## Putting It All Together

Diode current

$$I = I_e + I_h = \frac{Q_e}{\tau_e} + \frac{Q_h}{\tau_h} = \frac{eAL_e\delta n_0}{\tau_e} + \frac{eAL_h\delta p_0}{\tau_h}$$

Assumes excess charge replenished every  $\tau_{e,h}$  seconds

using

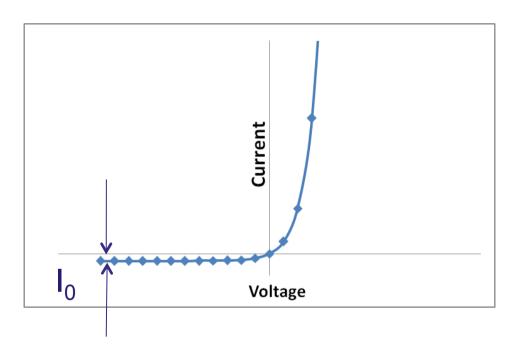
$$\partial p_0 = p_n \left[ exp \left[ \frac{eV_f}{k_B T} \right] - 1 \right]$$
 and  $\partial n_0 = n_p \left[ exp \left[ \frac{eV_f}{k_B T} \right] - 1 \right]$ 

$$I = \left[\frac{eAL_{e}n_{p}}{\tau_{e}} + \frac{eAL_{h}p_{n}}{\tau_{h}}\right] \left[exp\left(\frac{eV_{f}}{k_{B}T}\right) - 1\right]$$



Can be + or - to reflect forward and reverse bias

### **Diode Equation**



In terms of current density, J

$$I = I_0 \left[ exp\left(\frac{eV}{k_B T}\right) - 1 \right]$$

- Exponential increase in I with  $V = V_f$
- For reverse bias  $V \rightarrow$
- $-V >> k_B T$  (large and negative)

$$I \rightarrow I_0$$

• *I*<sub>0</sub> is the (reverse) saturation current

$$J = J_0 \left[ exp \left( \frac{eV}{k_B T} \right) - 1 \right]$$



#### Saturation or Reverse Current

$$I_0 = I_{e0} + I_{h0} = eA \left[ \frac{L_e n_p}{\tau_e} + \frac{L_h p_n}{\tau_h} \right]$$

Can be rewritten as 
$$I_0 = eAn_i^2 \left| \frac{D_e}{L_e N_A} + \frac{D_h}{L_h N_D} \right|$$

using 
$$L_h = (I$$

using 
$$L_h = (D_h \tau_h)^{1/2}$$
  $n_i^2 = p_p n_p = N_A n_p$   $n_i^2 = n_n p_n = N_D p_n$ 



## Summary

- The p-n junction under zero bias has a built-in potential preventing carrier diffusion
- Under forward bias the built-in potential is reduced, allowing carriers to diffuse more readily
- A continuous diffusion process across the junction is set up under a constant forward bias
- The recombination of diffusing excess minority carriers with majority carriers results in a current flow
- Forward current varies exponentially with applied bias the diode equation
- Under reverse bias the (reverse or saturation) current is small and constant