

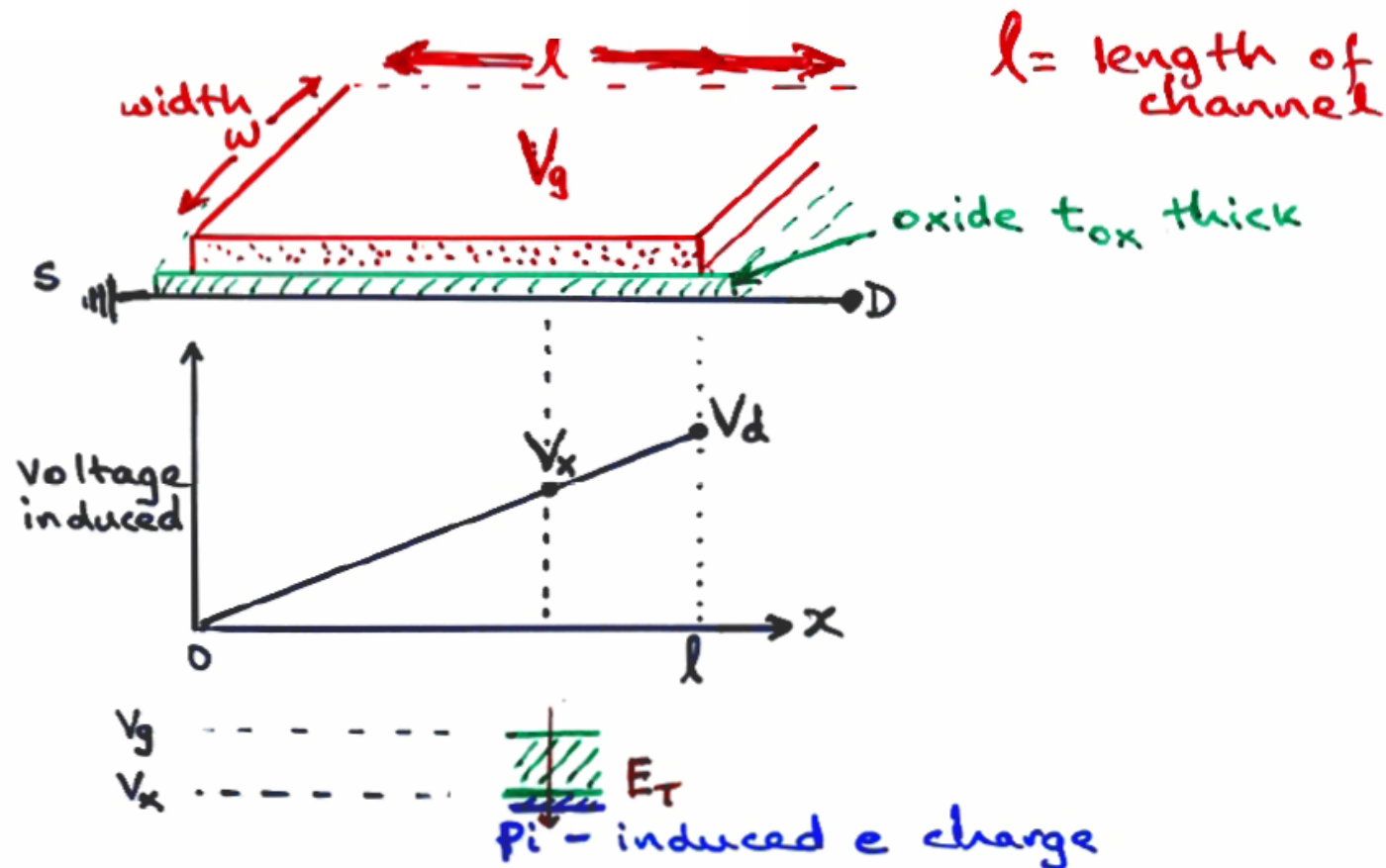


The
University
Of
Sheffield.

Metal oxide Semiconductor Transistor

I_d versus drain (V_d) characteristics

Related to the conducting inversion layer.





Metal oxide Semiconductor Transistor

Voltage across oxide at $x = V_g - V_x$

Oxide is t_{ox} thick, so transverse electric field is $E_T = (V_g - V_x) / t_{ox}$

By Gauss Law, surface charge density in channel (Q/A)

$$(\rho_i) = \epsilon_r E = \epsilon_r \epsilon_o (V_g - V_x) / t_{ox} \quad (1)$$

Not all of this is available for conduction. At low V_g only a depletion layer is formed. Requires a minimum $V = V_t$

effective voltage = $(V_g - V_x) - V_t$

$$\rho_{ieff} = \epsilon_r \epsilon_o (V_g - V_x - V_t) / t_{ox} = e \cdot \Delta n(x) \text{ (the surface electron density in the channel)} \quad (2)$$

This surface charge will behave like a capacitance (C) = Q/V $C_g = \frac{\epsilon_r \epsilon_o l w}{t_{ox}}$

$$\rho_{ieff} = e \cdot \Delta n(x) = C_g (V_g - V_x - V_t) / l w, \quad V_g - V_x > V_t \quad (3)$$

$$\rho_{ieff} = 0 \quad V_g - V_x < V_t$$

We know about charge, now what about the resistance?



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Integrate over channel length,

$$I_d \int_0^l dx = \frac{\mu_e C_g}{l} \int_0^{V_d} (V_g - V_0 - V_x) dV$$

$$I_d = \mu_e (C_g / l^2) [V_g - V_T - V_d / 2] V_d$$

V_d can be increased until $V_g - V_t = V_d$

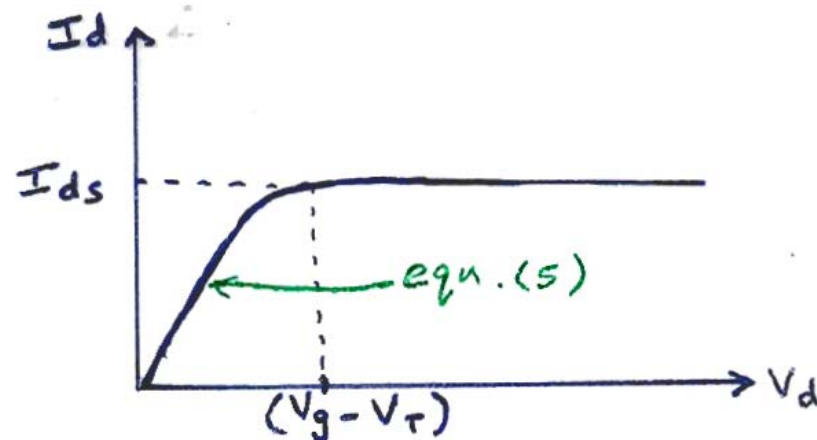
this is called the saturation condition

As V_d increases further, I_d saturates and excess voltage is dropped across the high resistance depletion layer.

Saturation current is I_{ds} .

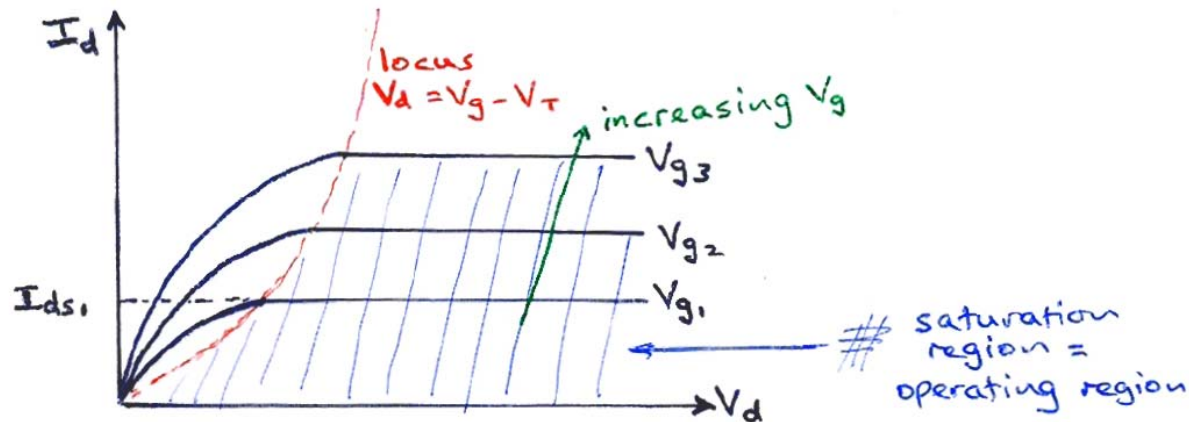
$$I_{ds} = \mu_e (C_g / l^2) [V_g - V_t]^2 / 2 \quad \text{or we can say,}$$

$$I_{ds} = \frac{C_g \mu_e V_{ds}^2}{2l^2}$$





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Another useful parameter is the Transconductance (i.e. gain)

$$g_m = \frac{dI_d}{dV_{gs}} \therefore = \frac{C_G \mu_e V_{ds}}{l^2}$$

In practice the g_m is less than this value, for a number of reasons:

We have ignored parasitic source and drain resistances (cant assume contacts are completely ohmic)

There is scattering of electrons due to the oxide, causing a reduced channel mobility for thin channel widths



Semiconductor Optoelectronics

Light emitting diodes

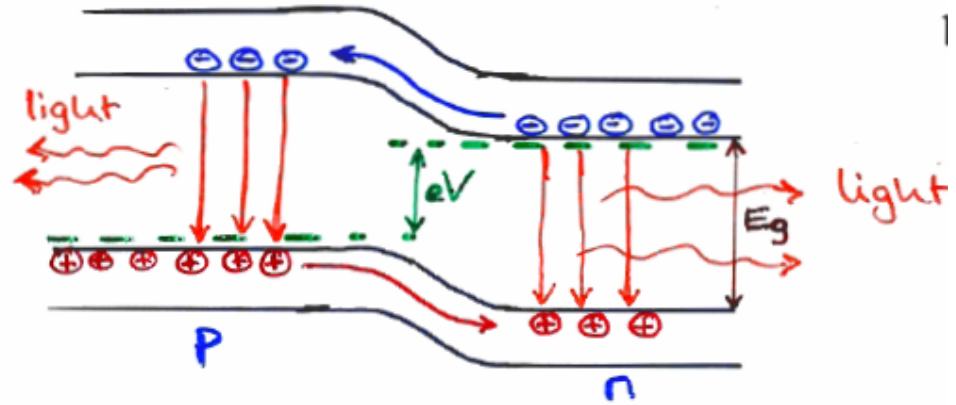
Under forward bias

Energy of Photons

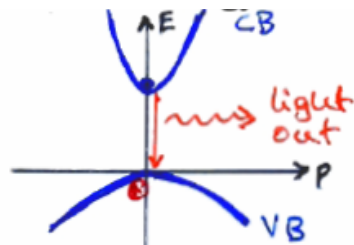
$$E_g = \frac{hc}{\lambda}$$

eg: for GaAs $E_g = 1.44\text{eV}$ – $\lambda = 860\text{nm}$

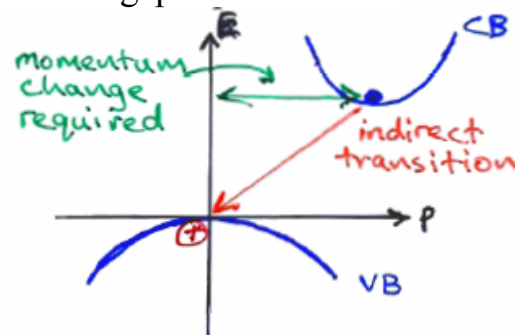
Simple rule $E_g (\text{eV}) = 1240/\lambda(\text{nm})$



Electrons and holes injected across the junction recombine with majority carriers and emit a photon of energy = bandgap



Direct band-structure
- efficient, recombination.
e.g. GaAs, InP, InGaAs



Indirect band structure
- poor recombination - need
momentum change - phonon,
interaction with lattice
e.g. Si, Ge

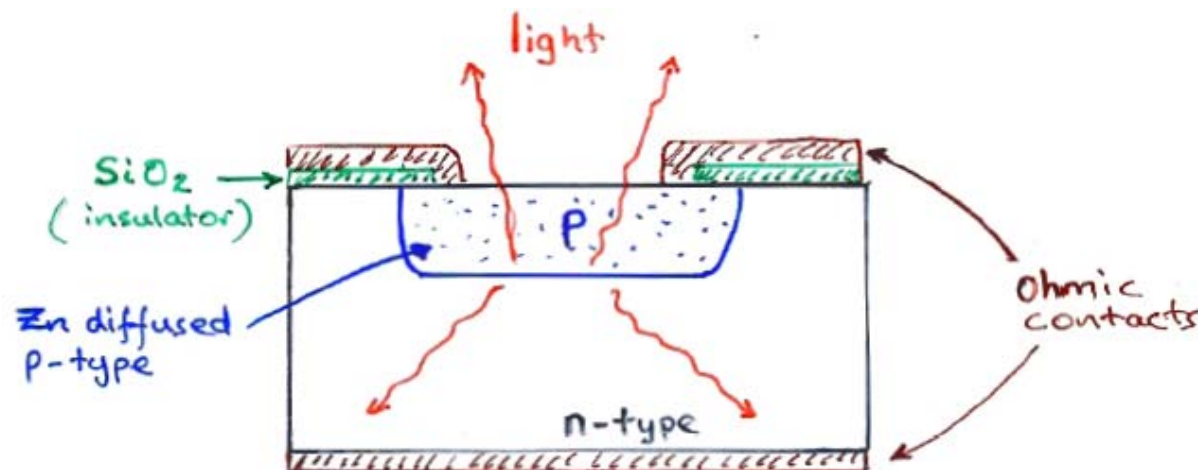


Semiconductor Optoelectronics

**Emitted wavelength
depends on the
choice of material**

InAs	3800nm	Mid IR (sensing)
InGaAs	1300-1550nm	Near IR (optical fibre telecommunications)
AlGaAs	750-800nm	Deep red (CD player)
GaP	690nm	Green (basic green LED)
GaInP	550-580nm	Orange/Yellow LED
InGaN	340-690nm	Blue/green LE, Laser (eg: Blu-ray)

Typical LED structure





Semiconductor Optoelectronics

An important property is the internal quantum efficiency = $\frac{\text{number.of .photons}}{\text{number.of .injected.e - h.pairs}}$

Some electron-hole pairs will be lost through non-radiative recombination. Better crystal quality gives

Better IQE. IQE is typically 20-50%. At > 30% LEDs become viable for lighting applications. However have other advantages such as longer lifetime, which means they are taking over in high maintenance applications (eg: traffic lights)

Power efficiency = $\frac{\phi}{IV}$ where ϕ is the optical power

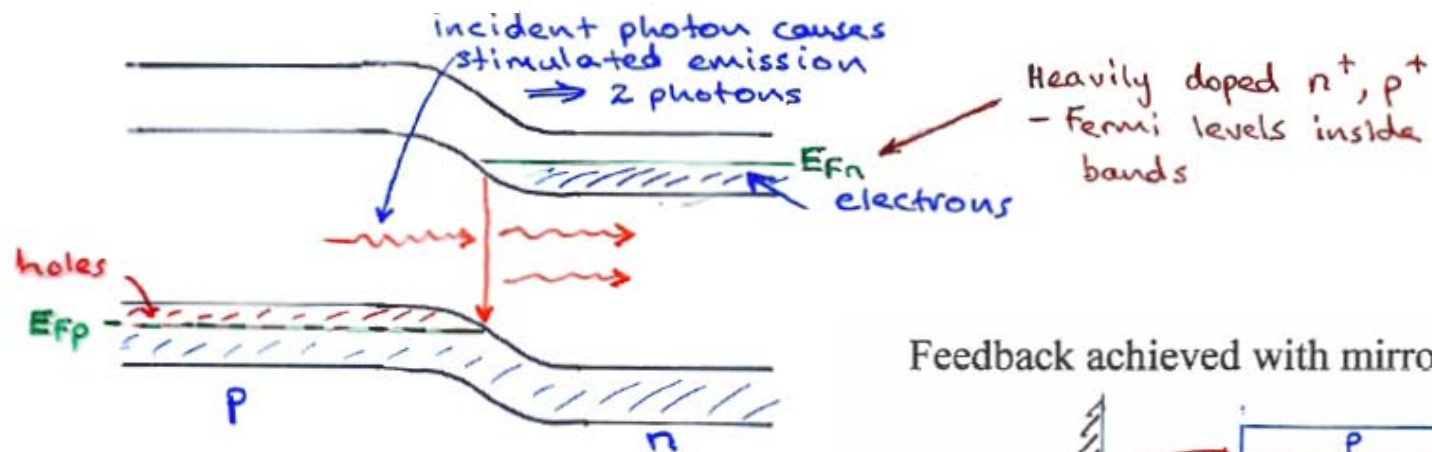
However ϕ is often quite difficult to measure. It involves the external quantum efficiency, which is a measure of how many photons escape. Many photons are lost through absorption and reflection. ExQE is often 10-20%.

Semiconductor laser

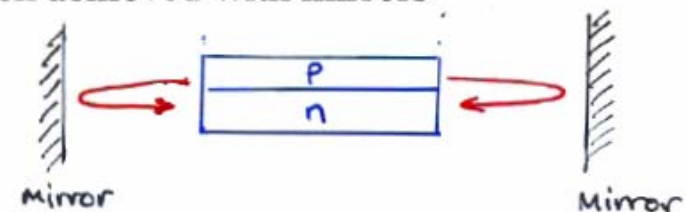
The semiconductor laser is very similar to an LED in construction. The difference here is that light is fed back using reflective mirrors (often just the facets of the semiconductor which have 30% reflectivity)



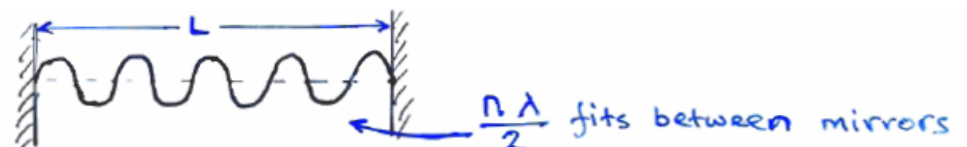
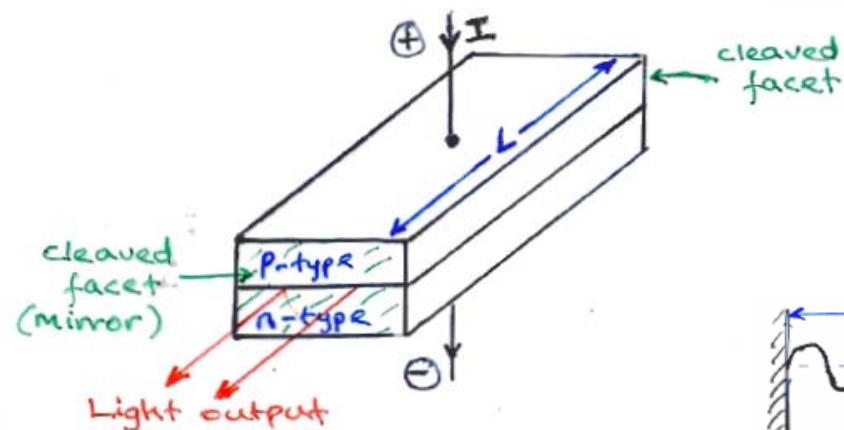
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Feedback achieved with mirrors



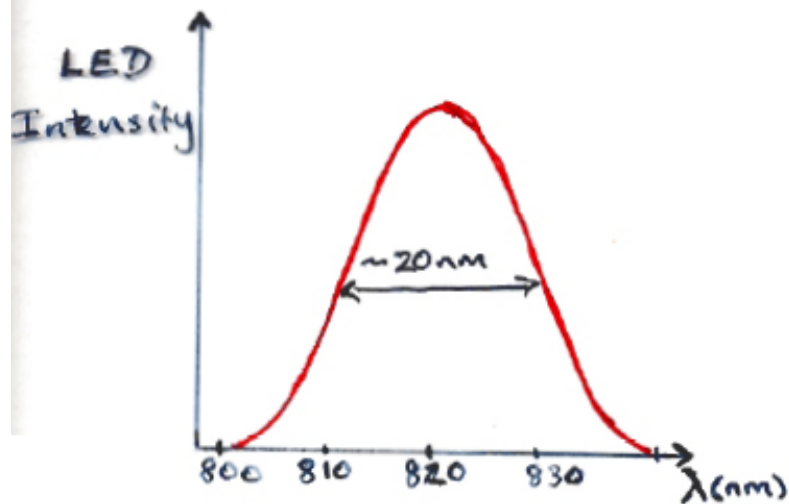
Get constructive interference when $n\lambda/2=l$
or $\lambda=2l/n$, where n is an integer





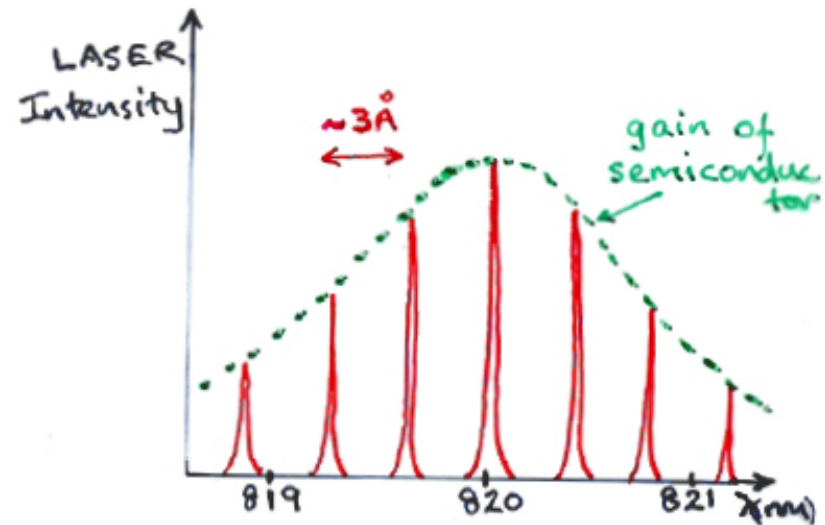
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Spectral Output



Broad linewidth centered
 \sim on material band-gap.

- Photons have a spread in energy
- Photons have no phase relationship with each other - i.e. random
- Spontaneous emission



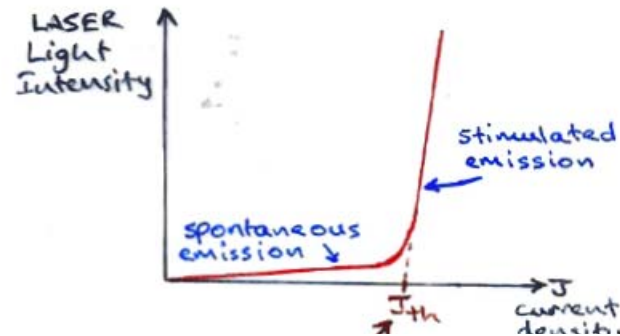
Linewidth of each mode
is very narrow compared
to LED. ~ 6 longitudinal
modes covering 2 nm.

- Photons have almost identical energies
- Photons have the same phase relationship with each other - coherent
- Stimulated emission



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L-I characteristics



At J_{th} , device gain = loss

Quantum Well Laser

Fully forward biased

Electrons and holes form quantum confined states

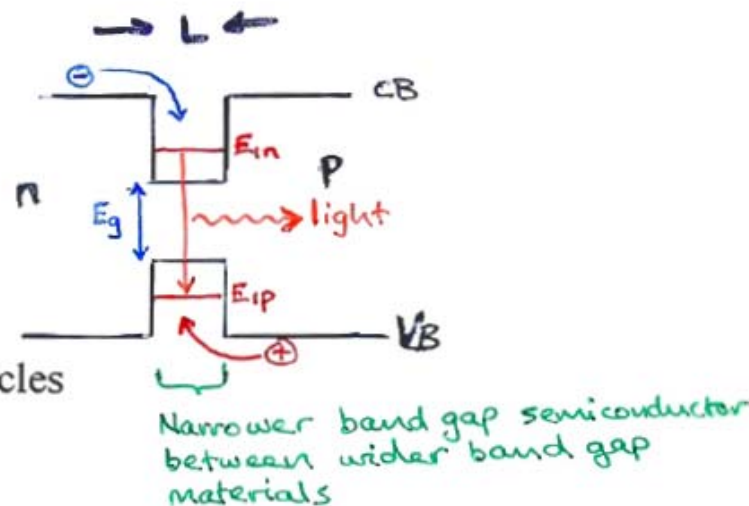
From section on Bound particles

$$E_{ln} = \frac{n^2 h^2}{8mL^2}$$

m different for e, h

Hence wavelength determined by $E_g + E_{ln} + E_{lp}$

Also quantum well lasers are much more efficient than 3-dimensional structures.



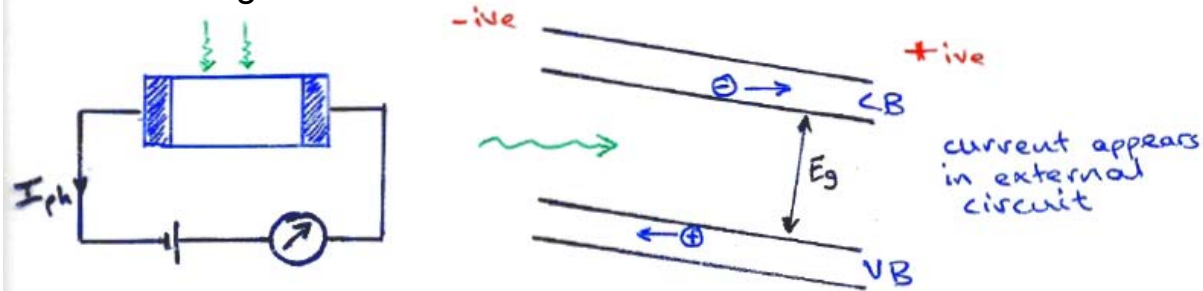


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Detectors

Two types

- (a) Photoconductive. Conductivity increases as the light increases



Needs photons of energy $> E_g$ to work. Also the conductivity of the semiconductor needs to be high in the dark ($\sigma \rightarrow 0$).

Under illumination $\sigma = e(n_e \mu_e + n_p \mu_p)$.
But 1 photon makes 1 electron and 1 hole so $\sigma = en(\mu_e + \mu_p)$.

Assume light of frequency ω (energy $h\omega$) and input optical power P (J/s).

$$\text{Average number of e-h pairs/sec} = \frac{P}{h\omega} x \eta$$

Where η is the efficiency
= no.pairs/ no..photons



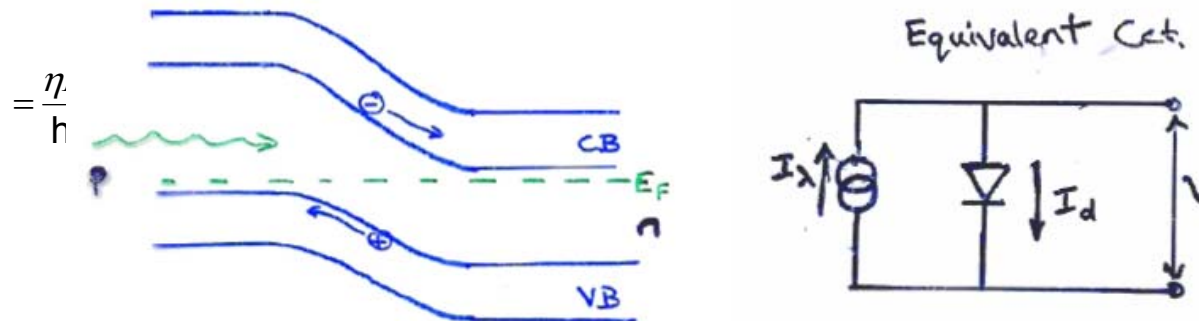
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However some of these pairs will recombine before they reach the contacts.

Hence the equilibrium number of pairs
$$= \frac{\eta P \tau}{h \omega}$$

Since $\sigma = en(\mu_e + \mu_p)$. Conductivity (s) is proportional to Input optical power (P)

(b) Photodiode



Built in field (can also be reversed biased) causes e-h pairs to separate, generating a current or voltage (open circuit).

Consider the equivalent circuit.
Remember diode equation

$$I_d = I_0 \exp\left(\frac{eV}{KT} - 1\right)$$



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For an open circuit $I_d = I_\lambda$ (input current due to arriving photons)

$$I_\lambda = \text{charge arriving per second} = \frac{\eta e P}{h \omega}$$

$$I_\lambda = \frac{\eta e P}{h \omega} = I_0 \left[\exp\left(\frac{eV}{kT}\right) - 1 \right]$$

$$V = \frac{kT}{e} \ln \left[\frac{\eta e P}{h \omega I_0} \right]$$

In photoconductive mode $I_{ext} = I_d - I_\lambda = I_0 \left[\exp\left(\frac{eV}{kT}\right) - 1 \right] - \frac{\eta e P}{h \omega}$

First term is insignificant if reverse biased $I_{ext} = \frac{\eta e P}{h \omega}$

In a **Solar cell** the photodiode is optimised for maximum power output

