

EEE345 : Summary of important equations and relationships

Maxwell equations

- i) Coulomb's Law : $\text{div } \underline{D} = \rho_{\text{free}}$
- ii) no magnetic monopoles : $\text{div } \underline{B} = 0$
- iii) Faraday's Law : $\text{rot } \underline{E} = -\frac{\partial \underline{B}}{\partial t}$
- iv) Ampere-Maxwell Law : $\text{rot } \underline{H} = \underline{j} + \frac{\partial \underline{D}}{\partial t}$

flux equations

- i) $\underline{D} = \epsilon_0 \epsilon_r \underline{E}$
- ii) $\underline{B} = \mu_0 \mu_r \underline{H}$
- iii) $\underline{j} = \sigma \underline{E}$

mathematical theorems:

a) Gauss's theorem: $\iiint_V \text{div } \underline{A} dV = \oint_S \underline{A} \cdot d\underline{S}$

b) Stoke's theorem: $\iint_S \text{rot } \underline{A} \cdot d\underline{S} = \oint_C \underline{A} \cdot d\underline{r}$

operator algebra: $\nabla^2 = \text{div grad}$, $\text{rot rot} = \text{grad div} - \text{div grad}$

device electrostatics:

- (i) Poisson's equation: $\nabla^2 V = -\frac{\rho}{\epsilon_0}$ (if $\rho=0$, then Laplace: $\nabla^2 V = 0$)
- (ii) $\underline{E} = -\text{grad } V$

general field equations:

a) $\underline{E} = -\text{grad } V - \frac{\partial \underline{A}}{\partial t}$

b) $\underline{B} = \text{rot } \underline{A}$

wave equation: $\ddot{\underline{f}} = v^2 \nabla^2 \underline{f}$ with velocity v , $\underline{f} = \underline{A}$ or \underline{B} or \underline{E}

propagation speed: $v = \frac{\omega}{k}$ (in vacuum: $= \frac{2\pi f}{2\pi/\lambda} = \lambda f = c$), $c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$

transmission lines

(i) voltage drop along lines: $-\frac{\partial V}{\partial x} = R \cdot I + L \cdot \frac{\partial I}{\partial t}$

(ii) leakage current between lines: $-\frac{\partial I}{\partial x} = G \cdot V + C \cdot \frac{\partial V}{\partial t}$

lossless line: $k_0 = \omega \sqrt{L^* C^*}$

lossy line: $\tilde{k} = \sqrt{\tilde{Z}^* \cdot \tilde{Y}^*} = \sqrt{R^* + j\omega L^*} \sqrt{G^* + j\omega C^*} = \omega \sqrt{L^* C^*} \left[1 - \frac{j}{2\omega} \left(\frac{G^*}{C^*} + \frac{R^*}{L^*} \right) \right]$

characteristic impedance: $Z_0 = \frac{\tilde{Z}^*}{j\tilde{k}} = \sqrt{\frac{L^*}{C^*}}$

voltage reflection coefficient: $\Gamma = \frac{Z_L - Z_0}{Z_L + Z_0}$

i) matched impedance: $Z_L = Z_0 \Rightarrow \Gamma = 0$

ii) open line: $Z_L = \infty \Rightarrow \Gamma = +1$

iii) short-circuited line: $Z_L = 0 \Rightarrow \Gamma = -1$

voltage standing wave ratio: $VSRW = \frac{1 + |\Gamma|}{1 - |\Gamma|}$

apparent impedance at source:

$Z_A = Z_0 \frac{\tilde{Z}_L + j Z_0 \tan \beta l}{Z_0 + j \tilde{Z}_L \tan \beta l}$

where $V(x,t)$ is given by $\frac{V_L}{2}$
 $|V(x)| = V_0^+ [1 + |\Gamma|^2 + 2|\Gamma| \cos(2kx + \psi)]$
 where ψ is angle in $\Gamma = |\Gamma| e^{j\psi}$
 at position $x=d$

coax cables: $C^* = \frac{C}{l} = \frac{2\pi \epsilon_0 \epsilon_r}{\ln \frac{R}{r}}$, $L^* = \frac{L}{l} = \frac{\mu_0 \mu_r}{2\pi} \ln \frac{R}{r}$, $v = \frac{c}{\sqrt{\epsilon_r \mu_r}}$

parallel plates: $C^* = \epsilon_0 \epsilon_r \frac{W}{d}$, $L^* = \mu_0 \mu_r \frac{d}{W}$, v as above, $Z_0 = \sqrt{\frac{L^*}{C^*}}$

absorption in materials

$\frac{E}{H} = \frac{E_0}{H_0} = \frac{W}{d} Z_0 = \sqrt{\frac{\mu_0 \mu_r}{\epsilon_0 \epsilon_r}}$ ($= 377 \Omega$ in vacuo)

energy density of \underline{E} field: $\tilde{w}_{elec} = \frac{1}{2} \epsilon_0 \epsilon_r E^2 = \frac{1}{2} \underline{E} \underline{D}$

" " " \underline{B} " $\tilde{w}_{mag} = \frac{1}{2} \mu_0 \mu_r H^2 = \frac{1}{2} \underline{H} \underline{B}$

" " " electromagnetic field: $w = \frac{1}{2} (\underline{E} \underline{D} + \underline{H} \underline{B})$

Poynting vector: $\underline{S} = \underline{E} \times \underline{H}$ where $|\underline{S}| = |\underline{E}| \cdot |\underline{H}| = \text{power density}$

specific conductance: $G^* = \frac{G}{L} = \frac{\sigma W}{d}$

absorption length: $\beta = \frac{\sigma}{2} \sqrt{\frac{\mu_0 \mu_r}{\epsilon_0 \epsilon_r}}$ so that $|\underline{S}| \propto \exp(-2\beta x)$

→ optics

weak absorption: $\sigma \ll \omega \epsilon_0 \epsilon_r \Rightarrow k = k_0 (1 - j\beta)$

strong " : $\sigma \gg \omega \epsilon_0 \epsilon_r \Rightarrow k = (1 - j) \sqrt{\frac{\omega \sigma \mu_0 \mu_r}{2}}$

skin depth : $\delta = \sqrt{\frac{2}{\omega \sigma \mu_0 \mu_r}}$

complex permittivity: $\epsilon_r = \epsilon_r' + j\epsilon_r''$ with $\epsilon_r'' = -\frac{\sigma}{\omega \epsilon_0}$

refractive index: $n = n' + j\kappa$ with $n' = \sqrt{\epsilon_r'}$

→ $\epsilon_r' = n'^2 - \kappa^2$

$\epsilon_r'' = 2n'\kappa$

reflection: $R = \left(\frac{|\underline{E}_r|}{|\underline{E}_i|} \right)^2$, transmission: $T = \left(\frac{|\underline{E}_t|}{|\underline{E}_i|} \right)^2$

absorption: $A = 1 - R$, $\underline{E} = c \underline{B} \times \underline{e}_r$, $\underline{B} = \frac{1}{c} \underline{e}_r \times \underline{E}$ orthonormal

all depend on polarisation

$R_{\perp} = \frac{(n_1 \cos \theta_1 - n_2 \cos \theta_2)^2 + (\kappa_1 \cos \theta_1 - \kappa_2 \cos \theta_2)^2}{(n_1 \cos \theta_1 + n_2 \cos \theta_2)^2 + (\kappa_1 \cos \theta_1 + \kappa_2 \cos \theta_2)^2}$

monotonic in θ_1

$R_{\parallel} = \frac{(n_2 \cos \theta_1 - n_1 \cos \theta_2)^2 + (\kappa_2 \cos \theta_1 - \kappa_1 \cos \theta_2)^2}{(n_2 \cos \theta_1 + n_1 \cos \theta_2)^2 + (\kappa_2 \cos \theta_1 + \kappa_1 \cos \theta_2)^2}$

has minimum at

Brewster angle $\theta_1 = \arctan \frac{n_2}{n_1}$

for unpolarised light: $R = \langle R \rangle = \frac{1}{2} (R_{\perp} + R_{\parallel})$

for vertical incidence ($\theta_1 = 0$):

$R = R_{\perp} = R_{\parallel} = \frac{(n_1 - n_2)^2 + (\kappa_1 - \kappa_2)^2}{(n_1 + n_2)^2 + (\kappa_1 + \kappa_2)^2}$

$T = T_{\perp} = T_{\parallel} = \frac{4(n_1^2 + \kappa_1^2)}{(n_1 + n_2)^2 + (\kappa_1 + \kappa_2)^2}$

$\kappa = \frac{1}{4\delta}$

→ dielectrics:

$\underline{D} = \epsilon \epsilon_r \underline{E} = \epsilon_0 \underline{E} + \underline{P}$ where \underline{P} is Polarisation

macroscopic: $\underline{P} = \epsilon_0 (\epsilon_r - 1) \underline{E}$

microscopic: $\underline{P} = n \underline{p}$ where $\underline{p} = q \underline{d}$ is dipole $\underline{p} \propto \frac{1}{r^3}$ and n density

electrostatic: $V = \frac{1}{4\pi\epsilon_0} \frac{\underline{p} \cdot \underline{r}}{r^3} \Rightarrow \underline{E} = \frac{1}{4\pi\epsilon_0} \left(\frac{3(\underline{p} \cdot \underline{r})\underline{r}}{r^5} - \frac{\underline{p}}{r^3} \right)$

electrodynamics: $V = \frac{1}{4\pi\epsilon_0} \left(\frac{\underline{p} \cdot \underline{r}}{cr^2} + \frac{\underline{p} \cdot \underline{r}}{r^3} \right)$

dipole far-field: $\underline{S} = \frac{\mu_0}{16\pi^2 c} \frac{\dot{\underline{p}}^2 \sin^2 \theta}{r^4} \underline{e}_r \propto \sin^2 \theta$ is no longer radial