

$$Q1 \quad (i) \quad q = f_0 / \Delta f = \frac{1.59 \text{ kHz}}{199 \text{ Hz}} = \underline{\underline{8.0}}$$

$$(ii) \quad f_0 = \frac{1}{2\pi\sqrt{LC}} = \frac{1}{2\pi\sqrt{L \cdot 100 \text{ nF}}} = 1.59 \text{ kHz}$$

$$\text{or } L = \frac{1}{10^{-7} \times (2\pi \cdot 1.59 \text{ kHz})^2}$$

$$= \underline{\underline{100 \text{ mH}}}$$

$$(iii) \quad q = \frac{1}{R_T} \sqrt{\frac{L}{C}} = \frac{1}{R_T} \sqrt{\frac{100 \text{ mH}}{100 \text{ nF}}} = \frac{10^3}{R_T} = 8$$

$$\therefore R_T = R + R_L = 100 + R_L = 125 \Omega$$

$$\therefore \underline{\underline{R_L = 25 \Omega}}$$

(iv) Voltage measured at V_R would be

$$\frac{V_s R}{R + R_L} = \frac{V_s \cdot 100}{100 + 25} = \frac{4V_s}{5} = \underline{\underline{0.8 V_s}}$$

Voltage measured at V_C would be.

$$|V_C| \text{ at res} = q V_s = \underline{\underline{8 V_s}}$$

(v). |Voltage across ideal bit of L | = $8 V_s$.

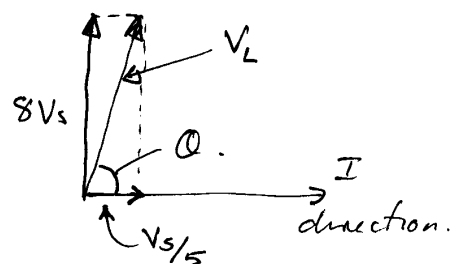
$$\text{Voltage across } R_L = V_s \frac{R_L}{R + R_L} = \frac{V_s}{5}$$

$$|V_L| = \sqrt{(8V_s)^2 + (V_s/5)^2}$$

$$\approx \underline{\underline{8 V_s}}$$

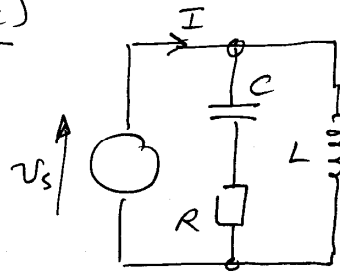
$$\angle V_L = \tan^{-1} \frac{8V_s}{V_s/5} = \tan^{-1} 40$$

$$= \underline{\underline{+88.6^\circ}}$$



Q2 $Z = \frac{V_s}{I} = \frac{j\omega L(R + \frac{1}{j\omega C})}{j\omega L + R + \frac{1}{j\omega C}}$

$$= \frac{j\omega L(1 + j\omega CR)}{1 + j\omega CR - \omega^2 LC}$$



Resonance occurs when Z is real - i.e. j terms vanish

$$\frac{(-\omega^2 LCR + j\omega L)((1 - \omega^2 LC) - j\omega CR)}{\text{real.}}$$

extracting j terms

$$j\omega(CR^2LC\omega^2 + L(1 - \omega^2 LC)) = 0$$

$$\therefore C^2R^2\omega^2L + L - \omega^2L^2C = 0$$

$$\omega^2(C^2R^2L - L^2C) = -L$$

$$\text{or } \omega^2(L^2C - C^2R^2L) = L$$

$$\omega^2 = \frac{L}{L^2C - C^2R^2L} = \frac{1}{LC - C^2R^2}$$

$$\text{or } \omega = \frac{1}{\sqrt{LC - C^2R^2}} \quad \text{or } \underline{\underline{f = \frac{1}{2\pi} \cdot \frac{1}{\sqrt{LC - C^2R^2}}}}$$

Q3. The impedance of the network is

$$Z = \frac{(R_C + \frac{1}{j\omega C})(R_L + j\omega L)}{R_C + \frac{1}{j\omega C} + R_L + j\omega L}$$

$$= \frac{R_C R_L + \frac{R_L}{j\omega C} + R_C j\omega L + \frac{L}{C}}{R_C + R_L + j(\omega L - \frac{1}{\omega C})}$$

$$Q3 \text{ cont.} \quad = \frac{(R_C R_L + \frac{L}{C}) + j(\omega L R_C - \frac{R_L}{\omega C})}{(R_C + R_L) + j(\omega L - \frac{1}{\omega C})}$$

We can make the imaginary parts of the numerator + denominator the same, except for a factor, if we make $R_C = R_L = R$. Z then becomes

$$Z = \frac{(R^2 + \frac{L}{C}) + R j(\omega L - \frac{1}{\omega C})}{2R + j(\omega L - \frac{1}{\omega C})}$$

(This is the argument needed by the real experts....
... the experts start from here -----)

We also need to take R out of the real part of the numerator and if we do this we get.

$$Z = R \cdot \frac{(R + \frac{L}{CR}) + j(\omega L - \frac{1}{\omega C})}{2R + j(\omega L - \frac{1}{\omega C})}$$

and the complex numerator and denominator can be made to cancel if the real parts are the same (the imaginary parts are already the same), i.e.

$$\text{if } R + \frac{L}{CR} = 2R$$

$$\text{or } \frac{L}{CR} = R \quad \text{or } \underline{\underline{R = \sqrt{\frac{L}{C}}}}$$

This will give

$$\begin{aligned} Z &= R \cdot \frac{(R + \frac{R^2}{R}) + j(\omega L - \frac{1}{\omega C})}{2R + j(\omega L - \frac{1}{\omega C})} = R \\ &= \text{purely real} = \text{resonant.} \end{aligned}$$

Q4 (i) at $t=0^-$, $I=0$ since $\frac{dV}{dt}$ must be zero.

$t=0^+$, $I = \frac{10 - V_C}{2k\Omega} = \frac{10 - 0}{2k\Omega} = \underline{5mA}$ since at $t=0^+$ there has not been any time to allow charge to build up in C.

$t \Rightarrow \infty$ $I=0$ since, once again, $\frac{dV}{dt}$ must be zero.

(ii) at $t=0^-$, $V_C=0$ since $V_S=0$ and there is no voltage drop across R (since $I=0$).

$t=0^+$, $V_C=0$ since there has been insufficient time for the charge in C to change.

$t \Rightarrow \infty$, $V_C \Rightarrow 10V$ since all transient effects will have settled down, $\frac{dV_C}{dt} = 0$ so $I_C = 0$ so no voltage drop across R.

Q5 (i) I_L at $t=0^- \dots = \frac{3V}{2k\Omega} = \underline{1.5mA}$. Since at $t=0^-$ the circuit is at a steady state - ie all transient effects have died away, $I_C=0$ and $V_L=0$ and so all V_S appears across $2k\Omega$.

at $t=0^+$ $I_L = 1.5mA$. Since $I = \frac{1}{L} \int V dt$ I must be continuous over an infinitesimally small time interval unless V can be infinitely big for an infinitely small time.

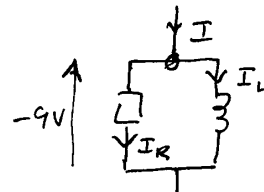
at $t \Rightarrow \infty$, $I_L = -\frac{6V}{2k\Omega} = -3mA$. Again, all the transient effects have settled at $t \Rightarrow \infty$, L looks like a short circuit and all of V_S appears across the $2k\Omega$.

(ii) at $t=0^-$, $I = I_L = 1.5mA$ since $V_L=0$, all of I must flow through L.

Q5 cont. at $t=0^+$ $I = 1.5\text{mA} - \frac{9\text{V}}{1\text{k}\Omega} = \underline{\underline{-7.5\text{mA}}}$.

This is the trickiest one so far. On the transient the voltage across C remains unchanged - ie LHS is 3V +ve w.r.t. RHS. But the LHS voltage changes from +3 to -6 and if the voltage across C remains unchanged the RHS must show the same change - ie from 0V to -9V. I_L at $t=0^+$ is 1.5mA and $I = I_L + I_R = 1.5\text{mA} + \frac{-9}{1\text{k}\Omega}$.

Notice that V_L can change without changing I_L over small timescales.



at $t \Rightarrow \infty$ all transients will have settled and we essentially have a d.c. problem with $V_S = -6$.
 $\therefore I_L = I$ as $t \Rightarrow \infty = \frac{-6}{2\text{k}\Omega} = \underline{\underline{-3\text{mA}}}$.

(iii) at $t=0^-$ $\underline{\underline{V_L = 0\text{V}}}$; the problem is a d.c. one.

at $t=0^+$ $\underline{\underline{V_L = -9\text{V}}}$ as described for the $t=0^+$ part of (ii).

at $t \Rightarrow \infty$ $\underline{\underline{V_L = 0\text{V}}}$ because the problem is once again a d.c. problem - all transients have settled.

Q6 (i) at $t=0^-$ the problem is a d.c. problem, $V_L = 0$
 and $I = -6 \times \frac{2\text{k}\Omega \parallel 2\text{k}\Omega}{2\text{k}\Omega + 2\text{k}\Omega \parallel 2\text{k}\Omega} \times \frac{1}{2\text{k}\Omega} = \underline{\underline{-1\text{mA}}}$.

at $t=0^+$ L maintains the level of current present at $t=0^-$, ie $I = -1\text{mA}$.

at $t \Rightarrow \infty$ the problem is once more d.c.
 $I = +12 \times \frac{2\text{k}\Omega \parallel 2\text{k}\Omega}{2\text{k}\Omega + 2\text{k}\Omega \parallel 2\text{k}\Omega} \times \frac{1}{2\text{k}\Omega} = \underline{\underline{+2\text{mA}}}$

Q6 cont... (ii) V_R at $t=0^-$ is $-6 \times \frac{2k\Omega // 2k\Omega}{2k\Omega + 2k\Omega // 2k\Omega} = \underline{\underline{-2V}}$.

at $t=0^+$ L behaves like a current source that maintains the $t=0^-$ value of I

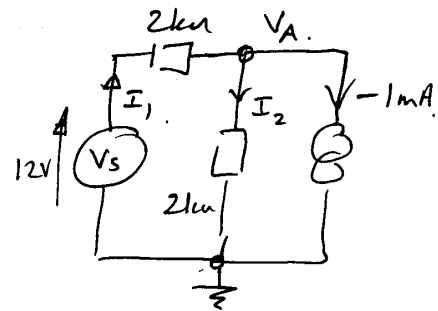
summing currents at V_A node at $t=0^+$

$$\frac{12 - V_A}{2k\Omega} = \frac{V_A}{2k\Omega} + (-1mA)$$

$$\frac{12}{2k\Omega} + 1mA = \frac{2V_A}{2k\Omega}$$

$$6mA + 1mA = \frac{V_A}{1k\Omega}$$

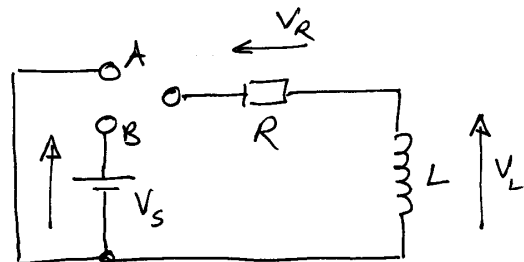
$$V_A = 7mA \times 1k\Omega = 7V = \underline{\underline{V_R}}$$



at $t \Rightarrow \infty$ the problem reverts to a dc one and

$$V_R = 12 \times \frac{2k\Omega // 2k\Omega}{2k\Omega + 2k\Omega // 2k\Omega} = \underline{\underline{4V}}$$

Q7. When switch switched to position A after having been in position B for a long time ...



$$V_R + V_L = 0$$

$$IR + L \frac{dI}{dt} = 0 \quad \text{or} \quad \frac{dI}{I} = -\frac{R}{L} dt$$

integrating both sides gives

$$\ln I = -\frac{R}{L} t + C$$

$$\text{or } I = e^{(-\frac{R}{L} t + C)} = A e^{-\frac{R}{L} t}$$

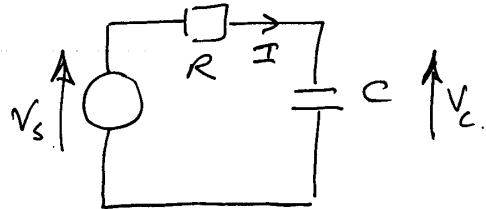
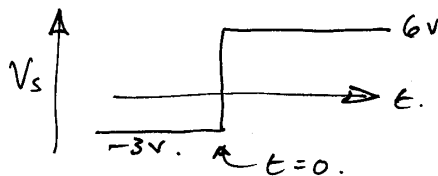
Q7 cont... when $t=0$ $I = \frac{V_s}{R}$

$$\therefore A = \frac{V_s}{R} \text{ and}$$

$$\underline{I = \frac{V_s}{R} e^{-t/(L/R)}}$$

$$\begin{aligned} V_L &= L \frac{dI}{dt} = L \frac{V_s}{R} \left(-\frac{1}{L/R} \right) e^{-t/(L/R)} \\ &= \underline{\underline{-V_s e^{-t/(L/R)}}} \end{aligned}$$

Q8



$V_s = V_R + V_c = IR + \frac{1}{C} \int I dt + V_{c(0)}$
 to get this equation in differential form, differentiate both sides...

$$0 = R \frac{dI}{dt} + \frac{I}{C} + 0$$

$$\text{or } \frac{dI}{I} = -\frac{dt}{RC}$$

integrating both sides

$$\ln I = -t/RC + C$$

$$\begin{aligned} \therefore I &= e^{(-t/RC + C)} \\ &= A e^{-t/RC} \end{aligned}$$

This is the initial condition for V_c to which the integrated current adds.

$$\text{when } t=0^+, I = \frac{6 - V_{c(0)}}{R} = \frac{6 - (-3)}{R} = \frac{9}{R} = A.$$

$$\text{so } \underline{\underline{I = \frac{9}{R} e^{-t/RC}}}$$

Q8 cont... The question actually asks you to solve the problem by developing an equation in V_c

$$V_c(t) + RC \frac{dV_c(t)}{dt} = V_s$$

\uparrow \uparrow
 V_c $V_R = IR$
 \uparrow
 $C \frac{dV_c}{dt}$

$$\text{or } -\frac{dt}{RC} = \frac{dV_c(t)}{(V_c(t) - V_s)} \quad \text{or} \quad C - \frac{t}{RC} = \ln(V_c(t) - V_s)$$

$$\text{or } e^{-t/RC} \cdot e^C = V_c(t) - V_s = A e^{-t/RC}$$

$$\text{When } t=0^+ \quad V_c(0^+) = -3 \quad \text{and} \quad V_s = 6$$

$$\therefore A = -9.$$

$$V_c(t) - 6 = -9 e^{-t/RC}$$

$$\text{or } V_c(t) = 6 - 9 e^{-t/RC} = \underline{\underline{9(1 - e^{-t/RC}) - 3}}$$

$I_c(t)$ can be derived from this result...

$$I_c(t) = C \frac{dV_c(t)}{dt} = C(-1)\left(-\frac{1}{RC}\right) 9 e^{-t/RC}$$

$$= \frac{9}{R} e^{-t/RC}$$

which agrees with the direct solution of $I_c(t)$.

Q9 When switch is moved to A,

$$I_1 = -I_L$$

and

$$V_{R_1} = V_L + V_{R_2}$$

$$-I_L(t) R_1 = L \frac{dI_L(t)}{dt} + I_L(t) R_2.$$

$$\text{or } -I_L(t) [R_1 + R_2] = L \frac{dI_L(t)}{dt}$$

$$\text{or } -\frac{R_1 + R_2}{L} dt = \frac{dI_L(t)}{I_L(t)}$$

integrating both sides gives.....

$$-\frac{R_1 + R_2}{L} t + C = \ln I_L(t).$$

$$\text{or } A e^{-\frac{R_1 + R_2}{L} t} = I_L(t).$$

$$\text{When } t=0, I_L(t) = V_s / R_2 = \frac{10}{R_2}.$$

$$\therefore \underline{\underline{\frac{10}{R_2} e^{-\frac{R_1 + R_2}{L} t} = I_L(t)}}.$$

$$V_{R_1} = -I_L(t) \cdot R_1 = -\frac{10 R_1}{R_2} e^{-\frac{R_1 + R_2}{L} t}$$

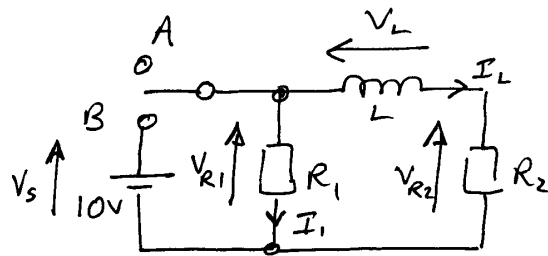
peak value when $t=0$ because this is the biggest value of $e^{-\frac{R_1 + R_2}{L} t}$

$$\therefore \text{peak } V_{R_1} = -10 \frac{R_1}{R_2} = \underline{\underline{-50V}}.$$

When switch goes back to position B the circuit is governed by....

$$V_s = V_L + V_{R_2}$$

$$= L \frac{dI_L(t)}{dt} + I_L(t) R_2$$



$$\frac{V_s}{R_2} = \frac{L}{R_2} \frac{dI_L(t)}{dt} + I_L(t).$$

$$-\left(I_L(t) - \frac{V_s}{R_2}\right) = \frac{L}{R_2} \frac{dI_L(t)}{dt}$$

$$-\frac{R_2}{L} dt = \frac{dI_L(t)}{I_L(t) - V_s/R_2}$$

integrating both sides

$$-\frac{R_2}{L} t + C = \ln(I_L(t) - V_s/R_2).$$

$$A e^{-\frac{R_2}{L} t} = I_L(t) - V_s/R_2.$$

$$\text{When } t = 0 \quad I_L(t) = 0 \quad \text{so } A = -V_s/R_2.$$

$$\underline{\underline{I_L(t) = \frac{V_s}{R_2} \left(1 - e^{-\frac{R_2}{L} t}\right)}}.$$

Notice that the time constant for "charging" the inductor with current, L/R_2 , is longer than that for discharging, $L/(R_1 + R_2)$.

The fact that current takes time to rise is important in electromechanical devices such as motors and solenoids because force tends to be proportional to current. There is thus a time delay between switching on a relay by driving a voltage across its coil and the relay switch contacts operating.