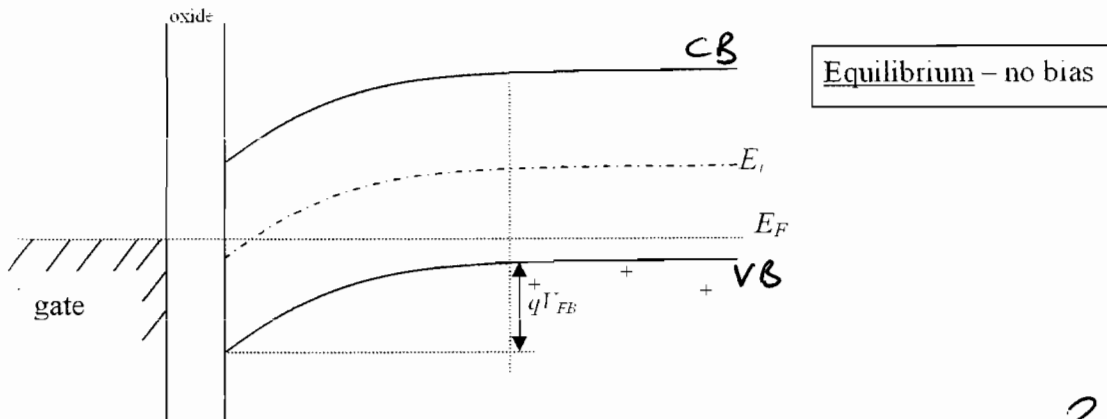


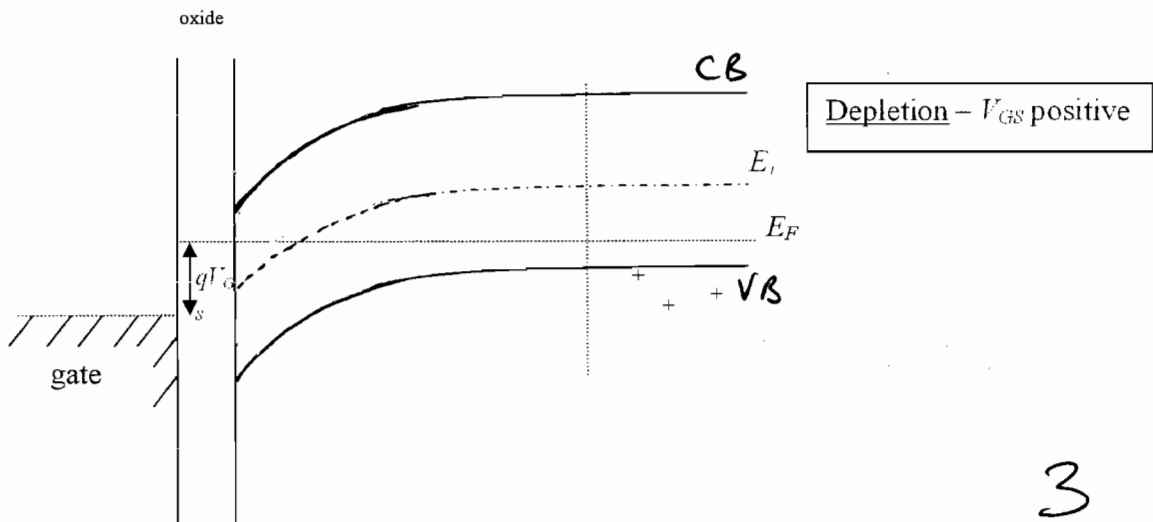
EE416/6040 08/09  
High Speed Electronic Devices  
Solutions

①

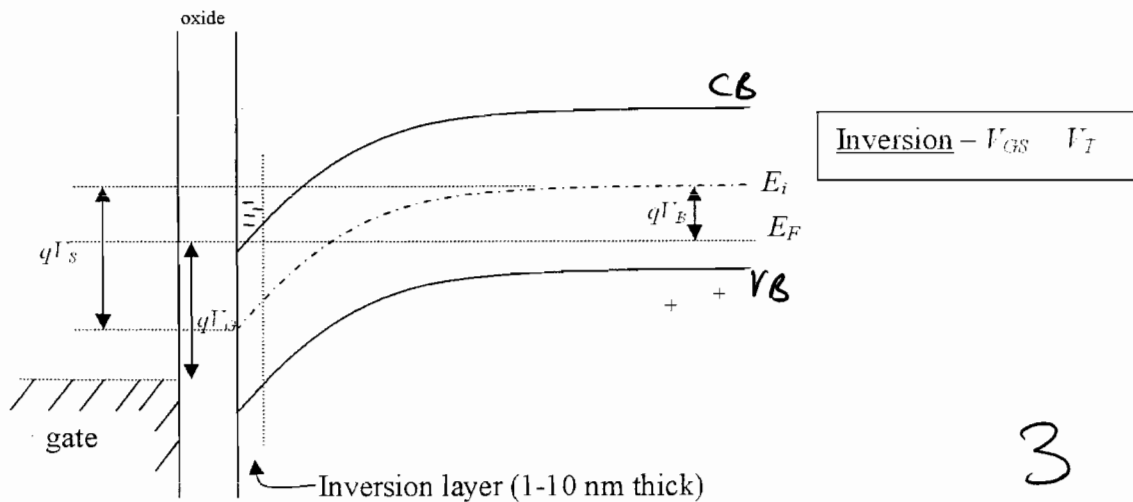
Q1 a)



3



3



3

Q1 (cont.)

(b) Main variation in  $V_T$  is through

$\frac{N_A^{1/2}}{C_{ox}}$  term.  $C_{ox}$  increases as the oxide thickness reduces due to scaling.

$N_A$  must increase to reduce depletion regions in line with scaling. Overall  $\frac{N_A^{1/2}}{C_{ox}}$  reduces which reduces  $V_T$ . Both

$V_{FB}$  and  $V_B$  are only weakly dependent on  $N_A$  and therefore do not influence  $V_T$ . 4

$V_T$  needs to be reduced during scaling to minimise power dissipation during switching. 2

As the third term in  $V_T$  expression is reduced due to scaling the other terms start to dominate and hence  $V_T$  reduces very little. 2

(c)

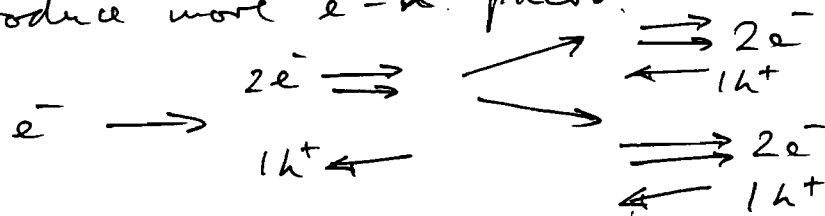
Change in threshold voltage due to oxide charge

$$\Delta V_T = \frac{Q_{ss}}{C_{ox}} = \frac{5 \times 10^{-7}}{\epsilon/d_{ox}}$$

$$= \frac{5 \times 10^{-7} \times 5 \times 10^{-8}}{3.45 \times 10^{-13}}$$

$$= \underline{\underline{0.072 V}}$$

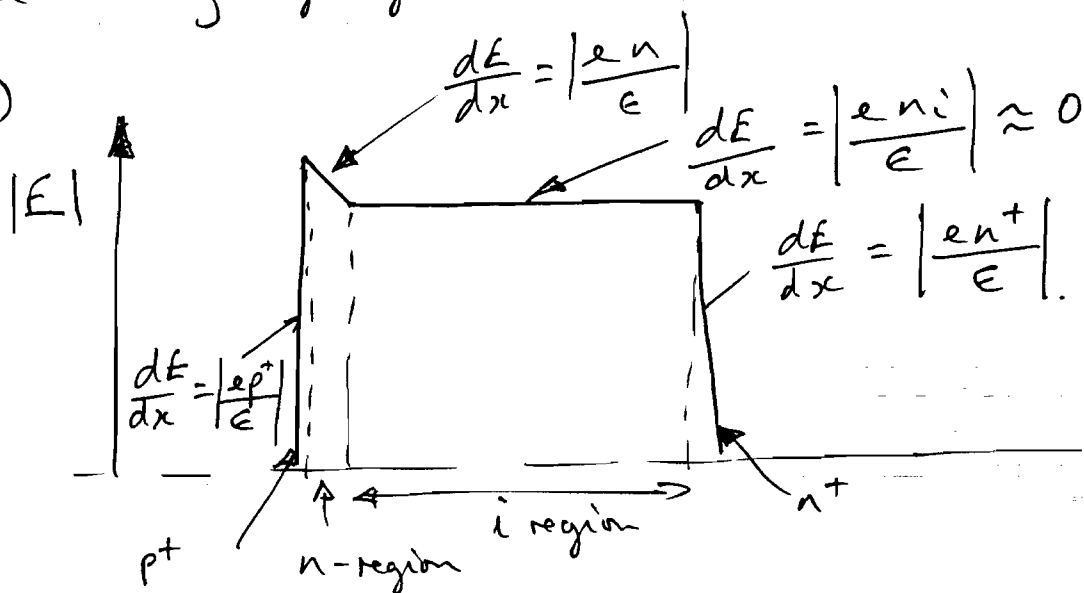
Q2 (a) When electrons gain sufficient energy from an electric field they can create e-h pairs by transferring their kinetic energy to break a crystal bond. Secondary electrons and holes so created are accelerated in turn to produce more e-h pairs.



3

This process is important for small devices since very high fields are encountered.

(b)



4

Assumptions from question  $\alpha = \beta$  and  $= \text{const.}$

$$\therefore 1 - \frac{1}{M_n} = \int_0^W \alpha dx = \alpha W$$

$$\therefore \alpha W = 1 \text{ for breakdown (i.e. } M_n \rightarrow \infty)$$

$$\text{i.e. } 3.8 \times 10^7 \exp\left(\frac{-3 \times 10^6}{E_n}\right) \times 0.05 \times 10^{-6} = 1$$

$$\Rightarrow \frac{3 \times 10^6}{E_n} = \ln[3.8 \times 10^7 \times 0.05 \times 10^{-6}]$$

$$\Rightarrow E_n = \frac{3 \times 10^6}{\ln[3.8 \times 10^7 \times 0.05 \times 10^{-6}]} = 4.67 \times 10^6 \text{ V/m}$$

2

Q2 (cont.)

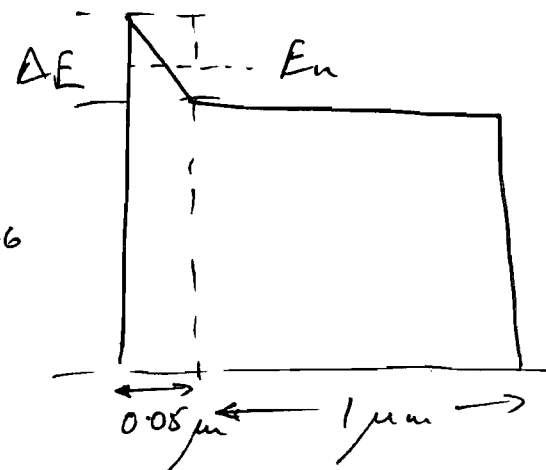
Now  $\Delta E = \frac{e n}{\epsilon} dx$

(from Poisson's Equation)

$$= \frac{1.6 \times 10^{-19} \times 3 \times 10^{21} \times 0.05 \times 10^{-6}}{1.17 \times 10^{-10}}$$

$$= \underline{2.05 \times 10^6 \text{ V/m}}$$

2



Breakdown voltage = area under  $E(x)$  diagram

$$= E_n \times 0.05 \times 10^{-6} + \left( E_n - \frac{\Delta E}{2} \right) \times 1 \times 10^{-6}$$

$$= 0.2 + 3.65 = \underline{\underline{3.85 \text{ V}}}$$

4

frequency - take electron velocity  $1 \times 10^5 \text{ m s}^{-1}$

$$\therefore \underline{\text{frequency}} = \frac{1}{\tau} = \frac{1 \times 10^5}{1 \times 10^{-6}} = 10^{11} \text{ Hz}$$

$$= \underline{\underline{100 \text{ GHz}}}$$

2

Q3

(a)  $r_{bb'}$  = base access resistance (lateral resistance between base contact and centre of the emitter)



$$r_{\pi} = \frac{\partial I_E}{\partial V_{BE}} = \frac{kT}{qI_E} = \text{input dynamic resistance of EB junction}$$

$$r_o = \text{output conductance} \left( \frac{\partial I_C}{\partial V_{CE}} \right)^{-1} - \text{due to Early Effect}$$

In an HBT structure:

$r_{bb'}$  is greatly reduced since very high doping in the base is possible without reducing the gain.

$r_o$  is increased since depletion of the base (Early effect) greatly reduced due to high doping there.

$r_{\pi}$  is unaffected.

Benefits

$f_{max}$  is increased due to reduced  $r_{bb'}$

Increased  $r_o$  results in a more linear amplification

2

Q3 (cont.)

(b)  $\tau_{BE}$  - delay in charging  $C_{BE}$  capacitance through intrinsic diode (emitter-base) resistance.

$$\tau_{BE} = \frac{dQ_{BE}}{dI_E} = C_{BE} \frac{\partial V_{BE}}{\partial I_E} = \frac{kT}{qI_E} \cdot C_{BE}$$

$\tau_B$  - base transit delay - time for electrons to diffuse through the base.

equivalent capacitance  $C_B = \frac{\partial Q_B}{\partial V_{BE}}$  where  $Q_B$  is the excess base charge.

$\tau_{BC}$  - delay in charging  $C_{BC}$  through the intrinsic emitter-base diode resistance plus emitter and collector resistance since current must also flow through these resistors

$$\text{i.e. } \tau_{BC} = \left( \frac{kT}{qI_E} + r_e + r_c \right) C_{BC}$$

$\tau_C$  - collector transit time - same argument as  $\tau_B$  i.e. equivalent capacitance  $= \frac{\partial Q_C}{\partial V_{BE}}$   $Q_C$  - charge in transit in collector

Since if this charge is driven by  $V_{BE}$ , the total delay can be represented by  $C_{b'e}$  which includes the sum of all the capacitance delays above.

4 (a) transconductance  $g_m = \frac{\partial I_D}{\partial V_{GS}}$  (definition)

$$= \frac{\cancel{Z} \mu C_{ox}}{L} (V_{GS} - V_T) \quad 2$$

$f_T$  gain =  $\frac{\text{output current}}{\text{input current}} = \frac{g_m V_{GS}}{V_{GS} W C_{ox} \cancel{Z} \times L}$

= 1 at  $f = f_T$

$$\Rightarrow f_T = \frac{g_m}{2\pi C_{ox} \cancel{Z} L} = \frac{\cancel{Z} \mu C_{ox}}{2\pi C_{ox} \cancel{Z} L^2} (V_{GS} - V_T)$$

$$= \frac{\mu}{2\pi L^2} (V_{GS} - V_T) \quad 2$$

These expressions assume that the velocity of the electrons will continue to rise as  $v = \mu E$  but in fact the velocity saturates at high fields when gate lengths are short, hence  $g_m$  and  $f_T$  will not rise as fast as predicted above as  $L$  shrinks. 2

Reductions in  $I_D$  and  $g_m$  due to velocity saturation can be compensated by increasing  $C_{ox}$ . Also the threshold voltage can be reduced as  $C_{ox}$  increases since the former will increase as the channel doping increases during the shrinking process.

(b) Transit-time  $f_T = \frac{1}{2\pi \tau} \quad \tau = \frac{v}{L}$

Just at the point of pinch-off (saturation)

$$V_{DS} = V_{GS} - V_T = 5 - 2 = 3 \text{ V}$$

5  $\mu\text{m}$  gate Electric field  $E = \frac{3}{5 \times 10^{-6}}$

$$= 6 \times 10^5 \text{ Vm}^{-1} = 6 \text{ kV/cm}$$

4 (cont.) velocity  $v = 1 \times 10^5 \left( \frac{6}{7+6} \right) \text{ m s}^{-1} = 0.46 \times 10^5 \text{ m s}^{-1}$

$$\therefore f_T (5 \mu\text{m}) = \frac{v}{2\pi L} = \frac{0.46 \times 10^5}{2\pi \times 5 \times 10^{-6}} = \underline{1.46 \text{ GHz}} \quad 2$$

0.1  $\mu\text{m}$  gate

$$E = \frac{3}{1 \times 10^{-7}} = 300 \text{ kV/cm}$$

$$v = 1.5 \times \frac{300}{307} = 0.977 \times 10^5 \text{ m s}^{-1}$$

$$\therefore f_T (0.1 \mu\text{m}) = \frac{0.977 \times 10^5}{2\pi \times 0.1 \times 10^{-6}} = \underline{155 \text{ GHz}} \quad 2$$

$f_T$  values increase faster than  $\propto \frac{1}{L}$  due to velocity increasing towards saturation.

C. Full marks to this section given for reasoned arguments showing knowledge of the issues (2 only required).

May include:

2 marks each point

- Cost of setting up process equipment and facilities
- Intolerable S-D leakage currents for  $L < 9 \text{ nm}$
- Inability to develop "high- $k$ " dielectrics of sufficient quality and reliability.
- Control of small dimensions to required tolerances becomes impossible.
- Non-available high conductivity interconnect metal and "low- $k$ " dielectrics.
- Accurate high frequency device and circuit models not available.
- Front-end processing models not available.