

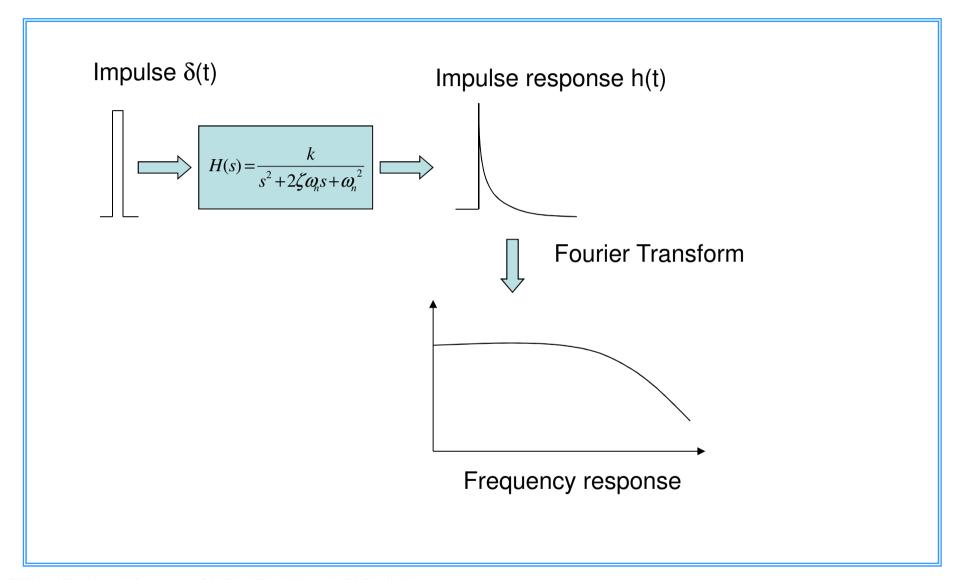
Lecture content



Relationship between Fourier Series and Fourier Transform

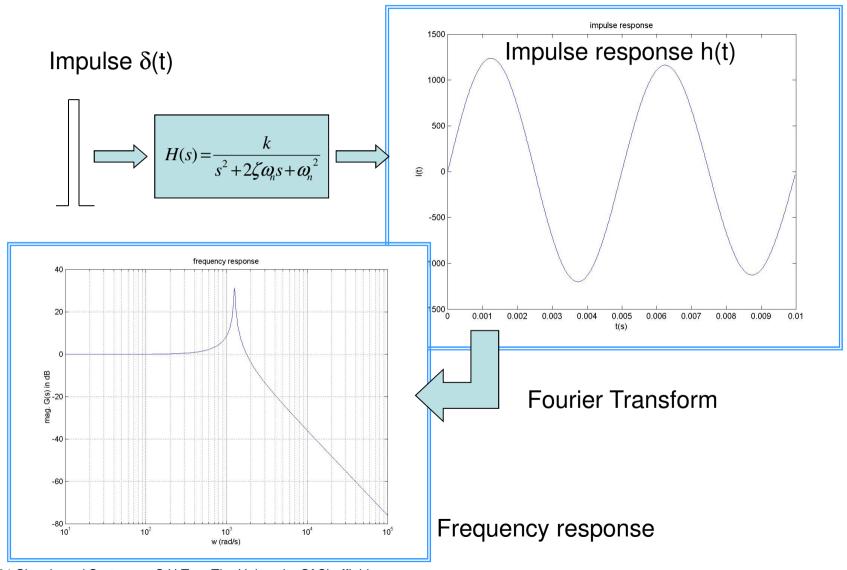


Why Fourier Transform





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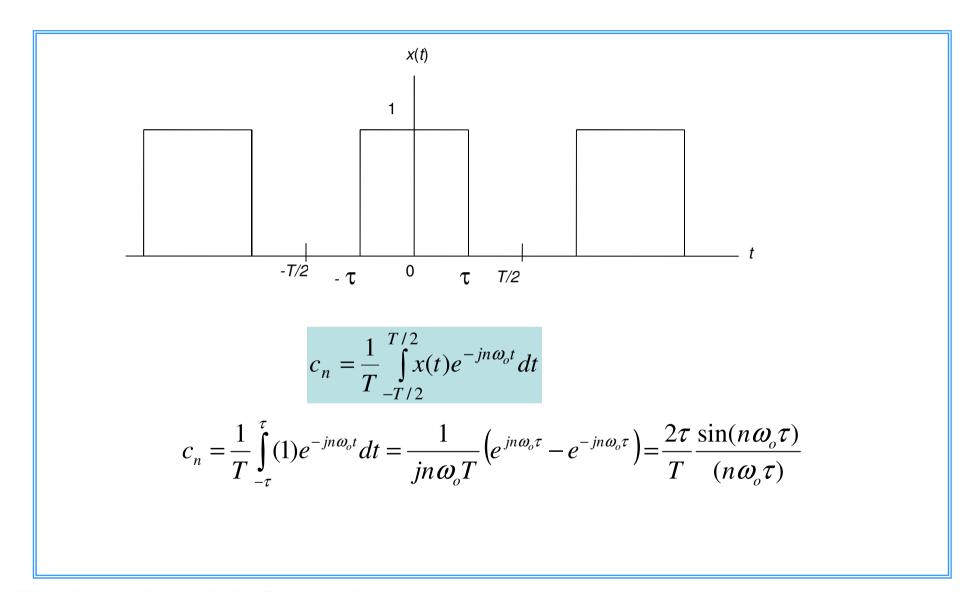
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The Fourier Series representation is applicable to periodic signals with infinite duration but many practical signals are non-periodic (or aperiodic) and have finite duration. We shall modify the Fourier Series so that it is applicable to aperiodic signals as well. The signal x(t) in figure 1 can be expressed as

$$x(t) = \sum_{n = -\infty}^{\infty} c_n e^{jn\omega_o t}$$

$$c_n = \frac{1}{T} \int_{-T/2}^{T/2} x(t)e^{-jn\omega_o t} dt$$





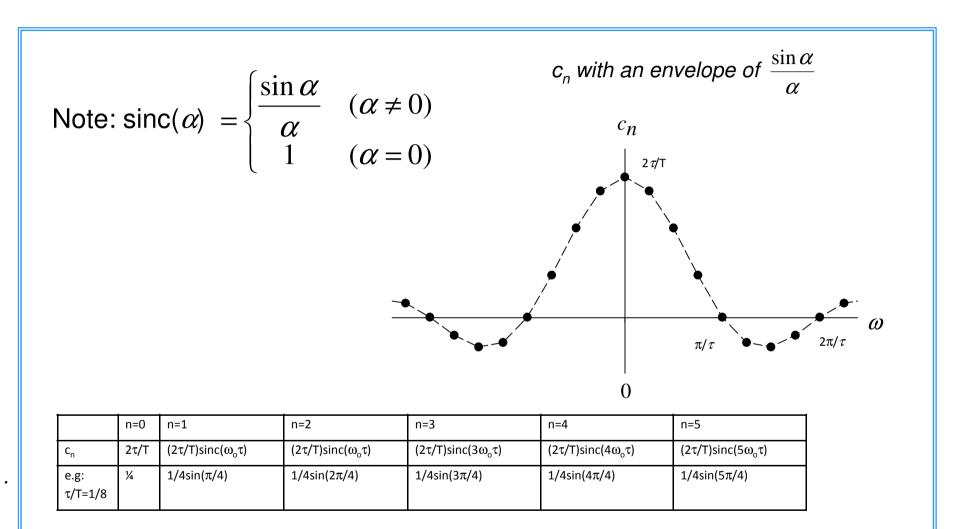
 c_n has magnitude with an envelope of $\frac{\sin \alpha}{\alpha}$ where $\alpha = n\omega_o \tau$ and a peak magnitude of $2\pi T$, as shown in figure 2.

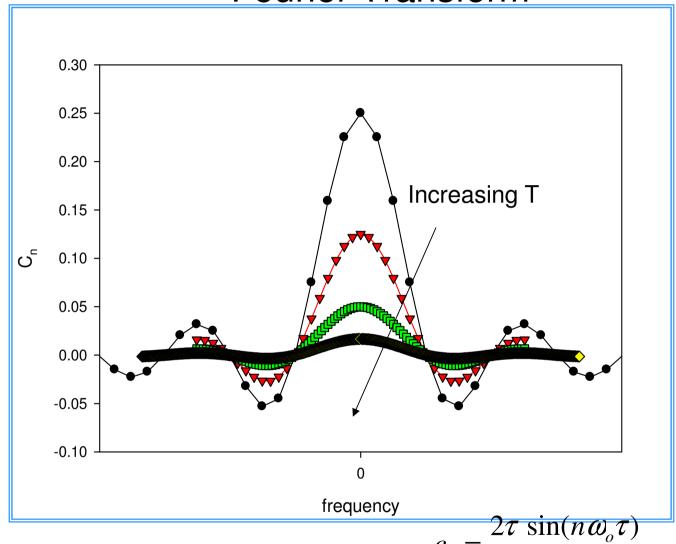
Note that the function $\frac{\sin\alpha}{\alpha}$ is sampled every ω_o rad/s (i.e the frequency of the harmonics). The function $\frac{\sin\alpha}{\alpha}$ is a sinc function and it has a peak magnitude of 1 at $\alpha=0$.

(Use l'Hopital's rule:
$$\lim_{\alpha \to 0} \frac{\sin \alpha}{\alpha} = \lim_{\alpha \to 0} \frac{\cos \alpha}{1} = 1$$
).

Nulls of occur when $\sin \alpha = 0$, that is when $\alpha = m\pi$ where m is integer to denote the nulls. Hence, the nulls are at $\omega = m\pi/\tau$ and we have the 1st null at π/τ , the 2nd null at $2\pi/\tau$ and so on.







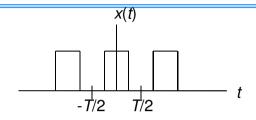
$$c_n = \frac{2\tau}{T} \frac{\sin(n\omega_o \tau)}{(n\omega_o \tau)}$$

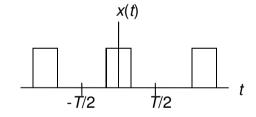


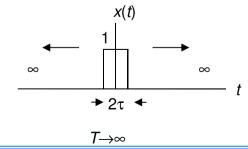
Now, consider the envelope function

$$Tc_n = \frac{2\tau \sin \omega \tau}{\omega \tau} \bigg|_{\omega = n\omega_o}$$

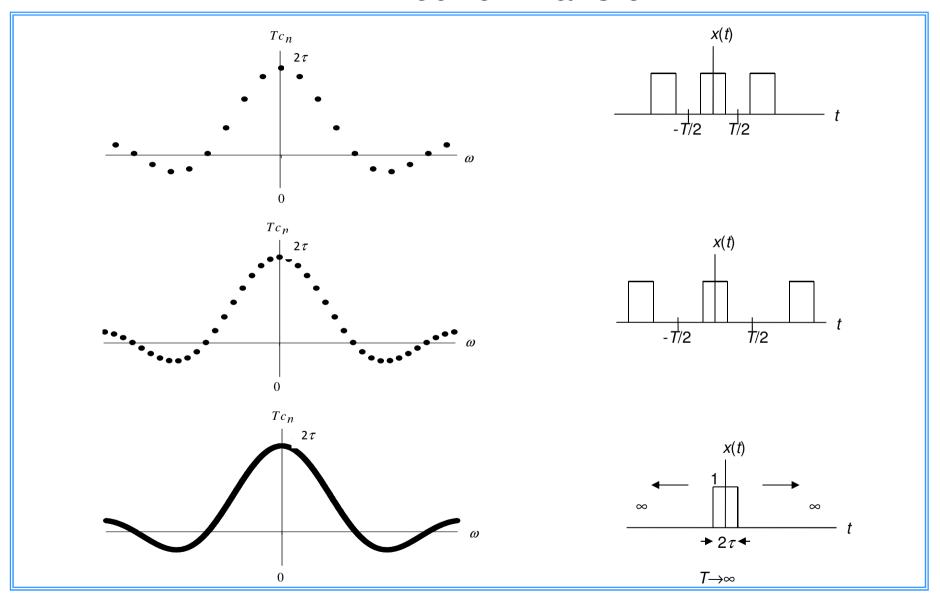
By changing T we can investigate the changes in the magnitude spectrum.











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We shall now develop the Fourier Transform of a rectangular pulse x(t). Let

$$c_n = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-jn\omega_o t} dt$$

$$X(\omega) = Tc_n = \int_{-T/2}^{T/2} x(t)e^{-j\omega t}dt$$

We know that

$$x(t) = \frac{1}{T} \sum_{n=-\infty}^{\infty} T c_n e^{jn\omega_o t} = \frac{1}{2\pi} \sum_{n=-\infty}^{\infty} X(\omega) e^{jn\omega_o t} \omega_o$$

As $T \to \infty$, $\omega_o \to 0$ so that ω becomes a continuum and ω_o can be written as $d\omega$. The summation becomes an integration and hence we have

$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\omega) e^{j\omega t} d\omega \quad \text{and} \quad X(\omega) = \int_{-\infty}^{\infty} x(t) e^{-j\omega t} dt$$

Inverse FT of $X(\omega)$

FT of x(t)



If the symmetry of the signal x(t) is known we can simplify the Fourier Transform integral to

$$X(\omega) = 2\int_{0}^{\infty} x(t)\cos \omega t dt$$

if x(t) has an even symmetry and

$$X(\omega) = -j2\int_{0}^{\infty} x(t)\sin \omega t dt$$

if x(t) has an odd symmetry.



1. Obtain the Fourier Transform of the rectangular window function in figure 4.

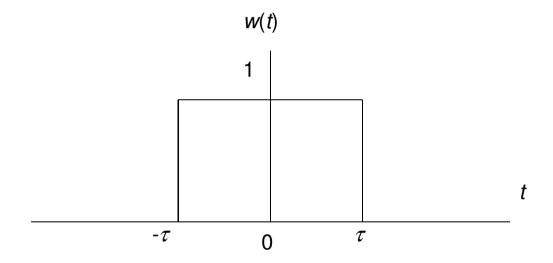
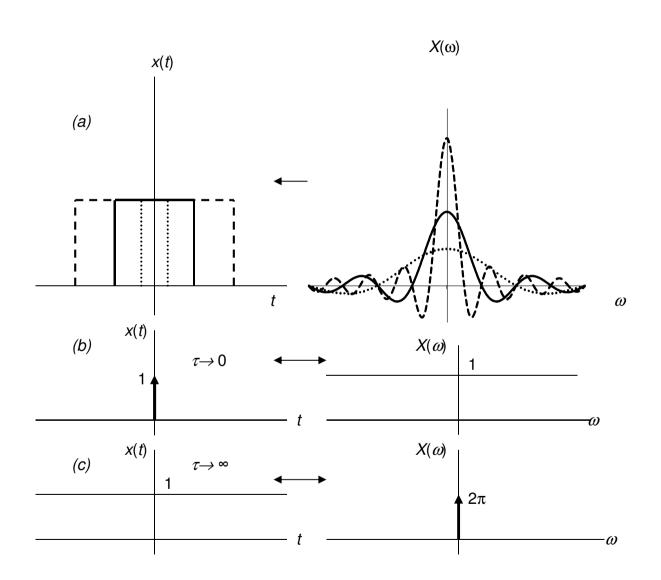


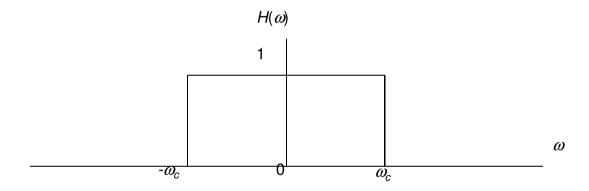
Figure 4: A rectangular window function with a duration of 2τ.





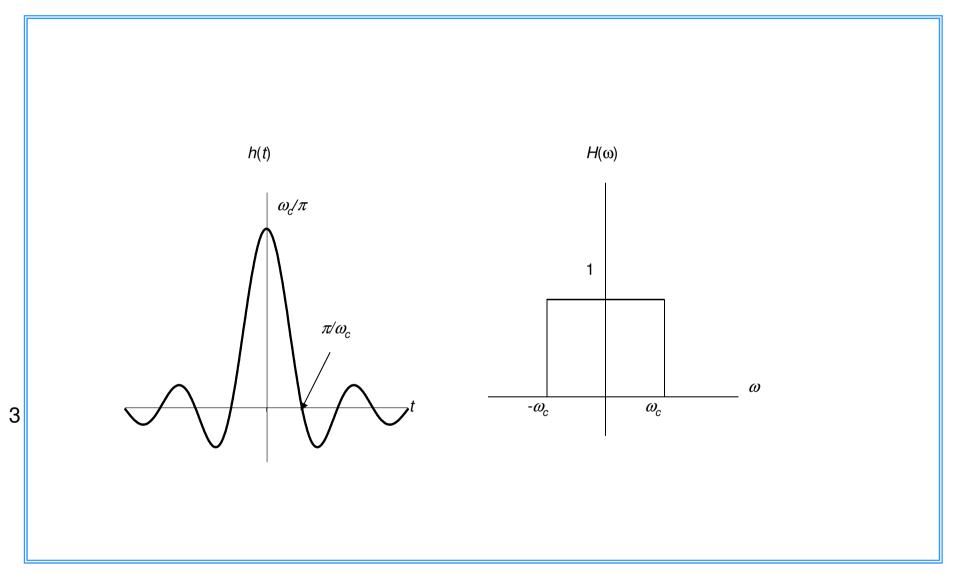


2. Compute the time function that has the magnitude spectrum (the positive half of the spectrum, $0 \le \omega \le \omega_c$, is an ideal low pass filter) shown in figure 6.



A rectangular spectrum defined by $H(\omega) = 1$ for $|\omega| \le \omega_c$ and zero otherwise.







3. Verify the Fourier Transform pair $x(t) = e^{-at}u(t) \leftrightarrow$, a > 0.