Two-Body Problem

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1 Introduction

The two body problem is an isolated system of two particles which interact through a central potential. For two particles with mass m_1 and m_2 , the equations of motion is

$$m_1\ddot{r_1} = F_{21}; \ m_2\ddot{r_2} = F_{12}.$$
 (1)

In Newtonian gravity, the potentials are therefore

$$U_{12} = u_{21} = \frac{Gm_1m_2}{|r_1 - r_2|^2}. (2)$$

Since the 2 particles interact via a central potential, the total momentum is conserved. We consider the center of mass,

$$R = \frac{m_1 r_1 + m_2 r_2}{M} \tag{3}$$

where $M = m_1 + m_2$.

The time derivative is therefore

$$V_{CM} = \frac{m_1 v_1 + m_2 v_2}{M}. (4)$$

Since our system is isolated, the net force and total acceleration is zero, thus V_{CM} is a constant. The motion of the center of mass is then

$$R(t) = V_{CM}t. (5)$$

We define the vector of distance between the particles:

$$r = r_1 - r_2. (6)$$

Thus,

$$r_1 = R + \frac{m_2 r}{m_1 + m_2}; \ r_2 = R - \frac{m_1 r}{m_1 + m_2}.$$
 (7)

Since we already know R, our problem reduces to finding r.

If we multiply the masses to ??, we get

$$m_1 m_2 \ddot{r_1} = m_2 F_{21}; \ m_1 m_2 \ddot{r_2} = m_1 F_{12}.$$
 (8)

Subtract the second from the first equation and simplify, we get

$$\frac{m_1 m_2}{m_2 + m_1} \ddot{r} = F_{21} \tag{9}$$

The force derived from the central potential only depends on distance between the 2 particles.

$$F_{21} = -\frac{\partial}{\partial r_1} U_{12}(|r_1 - r_2|) = \nabla U_{12}$$
 (10)

From above, we can summarize them as

$$m_*\ddot{r} = -\frac{\partial}{\partial r}U(|r|) = F(r)$$
 (11)

where

$$m_* = \frac{m_1 m_2}{m_1 + m_2} \tag{12}$$

is the reduced mass.

2 Conservation of Angular Momentum

In this section, we will rewrite the system above in terms of the angular momentum.

$$L = m_* r \times V; \ V = \dot{r}. \tag{13}$$

Since F is in the r direction, the torque has to be 0, and the angular momentum must be constant.

$$\frac{d}{dt}L = 0\tag{14}$$

Next, we use the scalar triple product identity,

$$a \cdot (b \times c) = b \cdot (c \times a) = c \cdot (a \times b) \tag{15}$$

and apply it to the angular momentum, we get

$$r \cdot L = m_* r \cdot (r \times V) = m_* V \cdot (r \times r) = 0. \tag{16}$$

Hence r is always perpendicular to angular momentum L.

Because the angular momentum is constant, there must be a fixed vector of which r is perpendicular to in space. Since the position vector is always perpendicular to a certain direction in space, we can reduce the probelm into a 2D plane in polar coordinates where

$$r_x = r\cos\theta; \ r_y = r\sin\theta.$$
 (17)

We choose the x--y plane for the particles to be on in accordance to conventions, and the angular momentum is oriented to the z direction. The velocity in θ direction is always positive, therefore the particles always rotates around the center of our coordinate system in the same direction.

$$\dot{\theta} = \frac{|L|}{m_* r^2} \tag{18}$$

3 Conservation of Energy

To begin with, we know that Energy is the sum of kinetic and potential energies.

$$E = \frac{1}{2}m_*v^2 + U(r) \tag{19}$$

In polar coordinates, the velocity is

$$v^2 = \dot{x}^2 + \dot{y}^2 = \dot{r}^2 + r^2 \dot{\theta}^2. \tag{20}$$

The energy conservation becomes

$$E = \frac{1}{2}m_*\dot{r}^2 + \frac{1}{2}m_*r^2\dot{\theta}^2 + U(r).$$
 (21)

4 Solutions for Two-Body Problem

Since $t\dot{he}ta = \frac{L}{mr^2}$, we get the acceleration in θ direction

$$\ddot{\theta} = -\frac{\dot{r}L}{mr^3} = -\frac{2\dot{r}\dot{\theta}}{r}.\tag{22}$$

Acceleration in r direction is

$$\ddot{r} = \frac{F}{m} = r\dot{\theta}^2 - \frac{Gm}{r^2} \tag{23}$$

because GM/r^2 and $r = \frac{x}{\cos \theta}$.

Our problem then becomes a system of two second-order ordinary differential equations stated above. To numerically solve it using RK-4,

we have to use 4 first-order differential equations such as

$$V_r = \dot{r} \tag{24}$$

$$V_{\theta} = \dot{\theta} \tag{25}$$

$$\dot{V}_r = rV_\theta^2 - \frac{Gm}{r^2} \tag{26}$$

$$\dot{V}_{\theta} = -\frac{V_r L}{mr^3} \tag{27}$$