

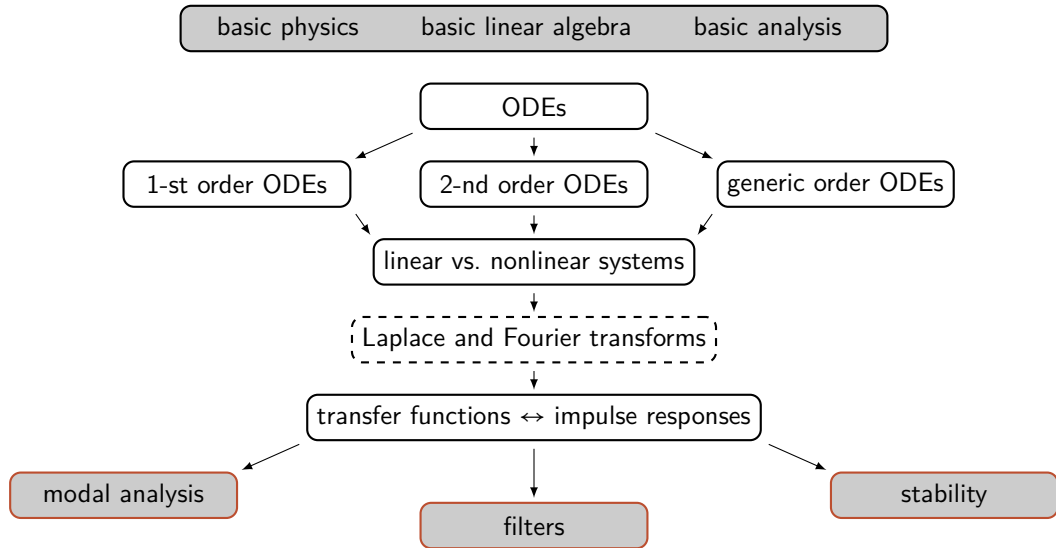
# TTK4225 - Systems Theory, Autumn 2020

Damiano Varagnolo



Some more insights on the concept of impulse response

# Summary / Roadmap of TTK4225



## Summary from the previous units

$$\dot{y} = ay + bu \quad (1)$$

**free evolution:**  $u = 0$ ,  $y_0 \neq 0$ ; implies  $y(t) = y_{\text{free}}(t) = y_0 e^{at}$

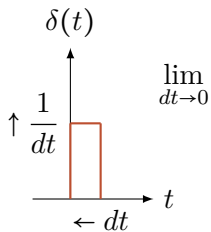
**forced response:**  $u \neq 0$ ,  $y_0 = 0$ ; implies  $y(t) = y_{\text{forced}}(t) = h * u(t)$  with  $h(t) = b e^{at}$

# Roadmap

- some more insights on the concept of “impulse response”

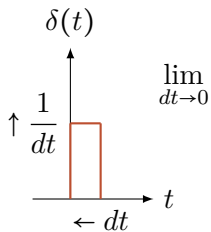
# Kronecker's delta

(i.e., informally, “pushing” an unitary mass in an infinitesimal space)



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*Discussion:*  $\int_{-\infty}^{+\infty} f(\tau) \delta(10) d\tau = ?$

## Impulse response - definition

$h(t) :=$  the system's response to an impulse  $u(t) = \delta(t)$ , since

$$y_{\delta}(t) = \int_0^t h(\tau) \delta(t - \tau) d\tau = h(t) \quad (2)$$

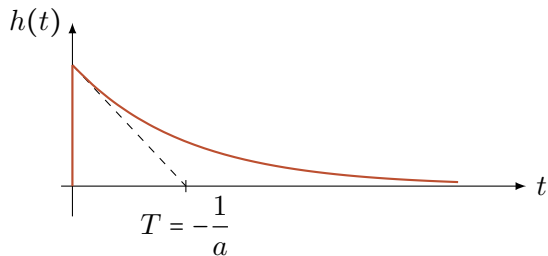


## Impulse responses of first order systems

$$\dot{y} = ay + bu \quad \mapsto \quad h(t) = be^{at}$$

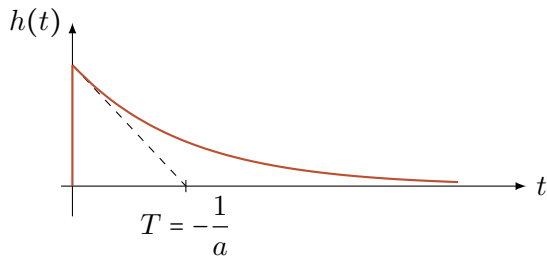
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alternative writing:  $h(t) = be^{-\frac{1}{T}t} \quad T > 0$

## What about the previous examples?

Dynamics of a cart:

$$\dot{v} = -\frac{k}{m}v(t) + \frac{k}{m}F(t) \quad (3)$$

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And for the RC-circuit?

$$\dot{v}_C = -\frac{1}{RC}v_C + \frac{1}{RC}V(t) \quad (4)$$

Introduction to first order systems - visualizing the system through  
a block scheme

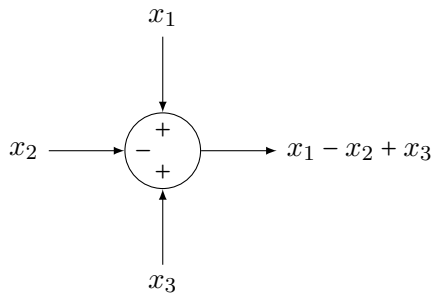
# Roadmap

- the most common block schemes
- first order systems as block schemes

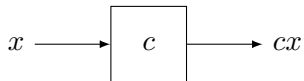
## Block diagrams - why?

- used very often in companies
- aid visualization (*until a certain complexity is reached...*)
- enable “drag & drop” way of programming
- in this course, primarily used for interpretations

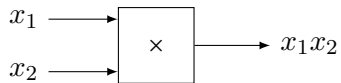
## Most common block diagrams - sum of $n$ signals



## Most common block diagrams - multiplication for a constant

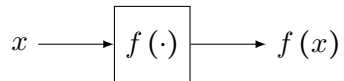


## Most common block diagrams - multiplication of two signals

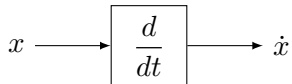




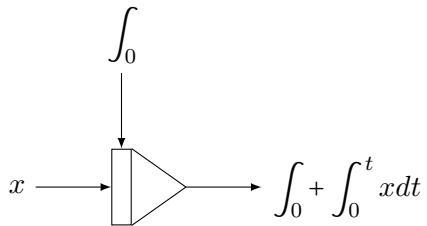
## Most common block diagrams - generic functions



## Most common block diagrams - derivatives



## Most common block diagrams - integrals

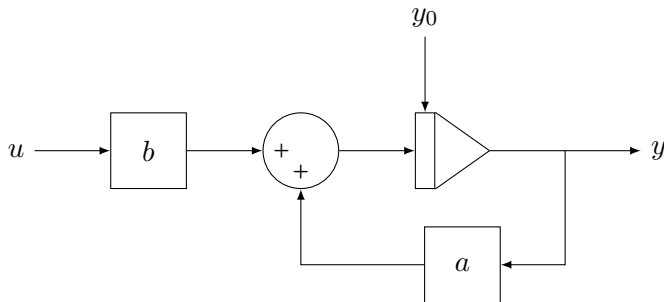


Discussion: how do we represent a first order DE with a block scheme?

$$\dot{y} = ay + bu$$

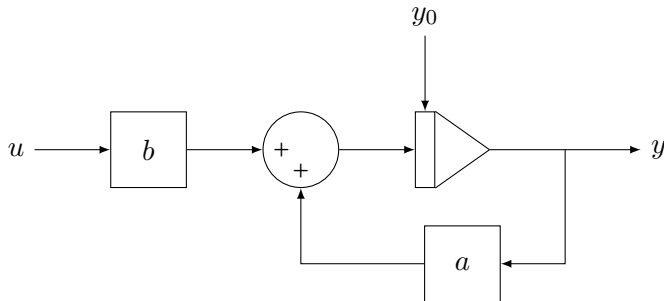
## The solution

$$\dot{y} = ay + bu$$



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*Discussion:* do you note the presence of a feedback loop?

# Introduction to Simulink

?

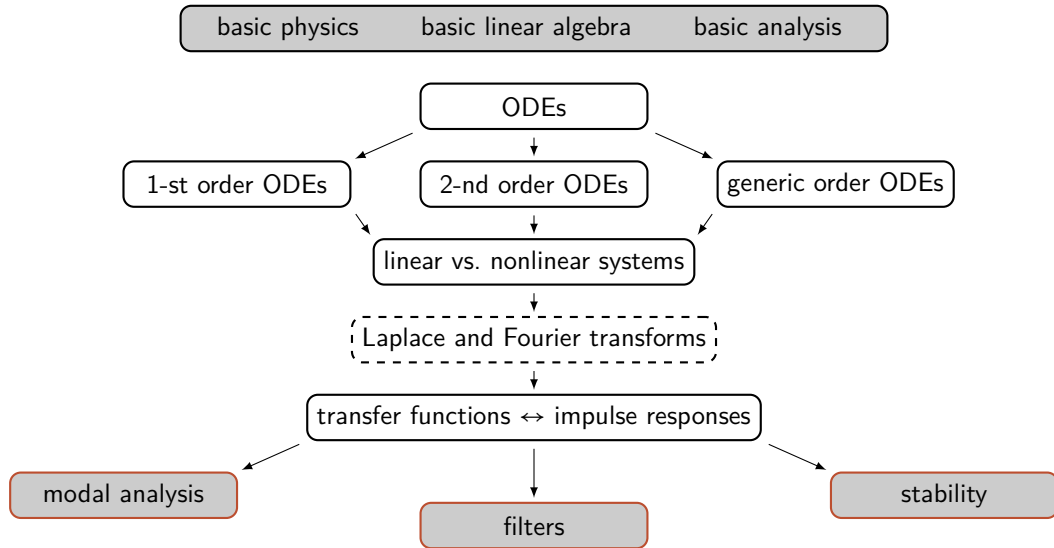


## Introduction to generic order systems

# Roadmap

- from first order to second order
- generic orders
- classification of the different types of ODEs

# Summary / Roadmap of TTK4225



## From first order to second order

*Discussion:* if we can write

$$\dot{y} = ay + bu ,$$

can we also write

$$\ddot{y} = a_1\dot{y} + a_0y + bu ?$$

## From first order to second order, third order

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And at this point also

$$\ddot{y} = a_2\ddot{y} + a_1\dot{y} + a_0y + bu ?$$

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And at this point also

$$\ddot{y} = a_2\ddot{y} + a_1\dot{y} + a_0y + bu ?$$

And why not

$$\ddot{y} = a_2\ddot{y} + a_1\dot{y} + a_0y + b_1\dot{u} + b_0u ?$$

## ARMA models

$$y^{(n)} = a_{n-1}y^{(n-1)} + \dots + a_0y + b_mu^{(m)} + \dots b_0u ?$$

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*Discussion:* and why don't we have  $a_n y^{(n)}$ , instead of only  $y^{(n)}$  on the LHS?



Discussion: can we generalize even more?

$$y^{(n)} = a_{n-1}y^{(n-1)} + \dots + a_0y + b_mu^{(m)} + \dots b_0u ?$$

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non-homogeneous:  $\dot{y} = a(t)y + b(t)u + c$

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time-varying:  $\dot{y} = a(t)y + b(t)u$

non-homogeneous:  $\dot{y} = a(t)y + b(t)u + c$

non-linear:  $\dot{y} = a\sqrt{y} + bu^{2/3}$

Discussion: can we make this linear again?

$$\dot{y} = ay + bu^{2/3}$$

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$$\dot{y} = ay + bu^{2/3}$$

*Another discussion:* can we apply the same “linearization trick” to  $\dot{y} = a\sqrt{y} + bu$ ?

main focus in this course:  
linear, time invariant, homogeneous (LTI)

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why?



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Remember: predictive control of the type

$$\mathbf{u}^* = \arg \min_{\mathbf{u} \in \mathcal{U}, \mathbf{f}(\mathbf{u}) \in \mathcal{F}} \text{Cost}(\mathbf{f}(\mathbf{u}), \mathbf{u})$$

requires to be able to:

- define “Cost”
- compute  $\mathbf{f}(\mathbf{u})$  rapidly
- be sure that the model does not have “nasty” properties

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from the previous lessons:  
1-st order LTI models  $\mapsto$  impulse responses;  
as will be clear later generic order LTI models  $\mapsto$  also impulse responses

The question we are trying to answer in this first part of the course

*“What will be the solution  $y(t)$  to a generic-order LTI ODE starting from a generic initial condition  $y(0) = y_0$  and generic input signal  $u(t)$ ?”*

*(our hope: the solution will be structurally easy,  
and fast to compute)*

?

# What is what?

(i.e., a small quiz to see if we are aligned)

	autonomous	linear	homogeneous	time constant
$\frac{\dot{x}}{x} = 1$				
$\pi = \dot{x}x$				
$2\dot{x} + \sin t = 0$				
$\sin(t\dot{x}) - x = 0$				
$t\dot{x} = x \sin t$				
$x - e^{-t} = \ddot{x} + \dot{x}$				

## Introduction to generic order systems

# Roadmap

- state space representations
- from  $n$ -th order to vectorial first order

## State space representations - Definition

mathematical model (typically but not limited to of a physical system) as a finite set of inputs, outputs and state variables related by first-order differential equations satisfying the separation principle



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## Ingredients

- finite number of inputs, outputs and state variables
- first-order differential equations
- *satisfies the separation principle*: the current value of states contains all the information necessary to forecast the future evolution of the outputs and of the states given the inputs

## State space representations - Facts

- the future output depends only on the current state and the future input

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- the future output depends on the past input only through the current state
- the state summarizes the effect of past inputs on future output  
*(sort of a “memory” of the system)*

# Example

Rechargeable flashlight:

- state = level of charge of the battery & on / off button
- output = how much light the device is producing

# State space representations - Notation

$u_1, \dots, u_m$  = inputs

$x_1, \dots, x_n$  = states

$y_1, \dots, y_p$  = outputs



## State space representations - Notation

$$\dot{x}_1 = f_1(x_1, \dots, x_n, u_1, \dots, u_m)$$

$$\vdots$$

$$\dot{x}_n = f_n(x_1, \dots, x_n, u_1, \dots, u_m)$$

$$y_1 = g_1(x_1, \dots, x_n, u_1, \dots, u_m)$$

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$$\dot{\boldsymbol{x}} = \boldsymbol{f}(\boldsymbol{x}, \boldsymbol{u})$$

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$$\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x}, \mathbf{u})$$

$$\mathbf{y} = \mathbf{g}(\mathbf{x}, \mathbf{u})$$

- $\mathbf{f}$  = state transition map
- $\mathbf{g}$  = output map

Discussion: is any  $f$  and  $g$  ok?

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$$y = g(x, u)$$

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$$\dot{x} = f(x, u)$$

$$y = g(x, u)$$

*No*: for every given  $x_0$  and  $u$  the solution must be:

- unique
- guaranteed to exist for  $t \in [0, \bar{t})$  for some  $\bar{t} > 0$

*(check “existence and uniqueness of solutions of ODEs”)*

## State space representations - Example

$$\begin{cases} \dot{x} = -x^2 \\ y = x \end{cases} \cup x_0 = -1 \implies x(t) = \frac{1}{t-1}$$

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*Discussion:* (a bit outside of the current flow of concepts on ODEs, but still relevant and very important) what happens as  $t \rightarrow 1$ ?

## State space representations - Example

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*Discussion:* (a bit outside of the current flow of concepts on ODEs, but still relevant and very important) what happens as  $t \rightarrow 1$ ?

*there is a finite escape time!*



Discussion: can every physical system be described with a state space representation?

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Noticeable exceptions:

- non-lumped parameters model (i.e., PDEs)
- continuous-time systems with delays  
*(will be discussed better later on)*

## Example: dynamics of sugar and insulin concentrations in blood

- $x_1 :=$  sugar concentration
- $x_2 :=$  insulin concentration
- $u_1 :=$  food intake
- $u_2 :=$  insulin intake
- $c :=$  sugar concentration in fasting

$$\begin{cases} \dot{x}_2 = a_{21}(x_1 - c) - a_{22}x_2 + b_2u_2 & x_1 \geq c \\ \dot{x}_2 = -a_{22}x_2 + b_2u_2 & x_1 < c \end{cases}$$

$$\begin{cases} \dot{x}_1 = -a_{11}x_1x_2 + a_{12}(c - x_1) + b_1u_1 & x_1 \geq c \\ \dot{x}_1 = -a_{11}x_1x_2 + b_1u_1 & x_1 < c \end{cases}$$

?

How can I transform a generic-order ODE into a state space representation (that requires first order DEs)?

$$\ddot{y} = \phi(\ddot{y}, \dot{y}, y, \ddot{u}, \dot{u}, u);$$

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$$\ddot{y} = \phi(\ddot{y}, \dot{y}, y, \ddot{u}, \dot{u}, u);$$

*change of variables:*

$$\mathbf{x} = \begin{bmatrix} \ddot{y} \\ \dot{y} \\ y \end{bmatrix}, \quad \mathbf{u} = \begin{bmatrix} \ddot{u} \\ \dot{u} \\ u \end{bmatrix} \quad \Rightarrow \quad \dot{\mathbf{x}} = \begin{bmatrix} \ddot{y} \\ \dot{y} \\ y \end{bmatrix} = \begin{bmatrix} \phi(\mathbf{x}, \mathbf{u}) \\ [1 \ 0 \ 0] \mathbf{x} \\ [0 \ 1 \ 0] \mathbf{x} \end{bmatrix} = \mathbf{f}(\mathbf{x}, \mathbf{u})$$

Thus, how do I transform a generic order LTI system into a first order LTI system?

$$\ddot{y} = 2\dot{y} + 3y + 4\ddot{u} + 5\dot{u} + 6u$$

## Transforming ARMA models

$$y^{(n)} = a_{n-1}y^{(n-1)} + \dots + a_0y + b_mu^{(m)} + \dots b_0u ?$$



Discussion: but then any LTI system is a first order vectorial system?

$$\dot{\mathbf{y}} = A\mathbf{y} + B\mathbf{u}$$

Discussion: but then any LTI system is a first order vectorial system?

$$\dot{\mathbf{y}} = A\mathbf{y} + B\mathbf{u}$$

YES, AND THIS IS WHY  
LINEAR ALGEBRA IS ESSENTIAL  
FOR CONTROL!

?

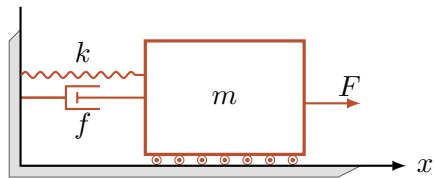
## Second order systems

# Roadmap

- interesting examples
- interesting responses
- differences between 1st order and 2nd order

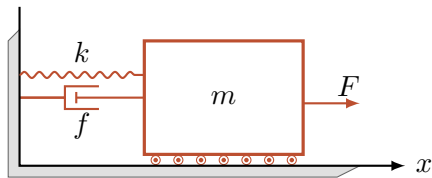
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Position of a cart fastened with a spring to a wall and subject to friction:



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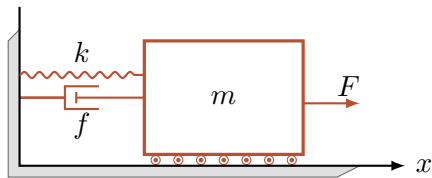


EOM:

- force from the spring:  $F_x(t) = -kx(t)$
- friction:  $F_f(t) = -f\dot{x}(t)$
- applied force:  $F(t)$
- Newton's second law:  $\sum F = m\ddot{x}(t)$

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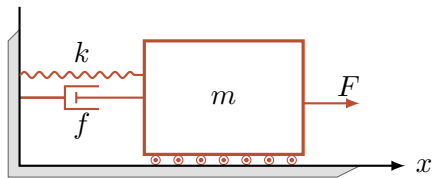
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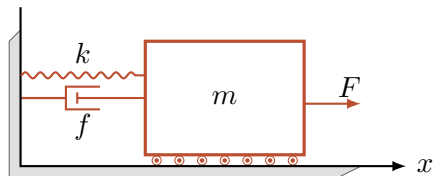
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$$m\ddot{x}(t) = F_x(t) + F_f(t) + F(t) \quad \mapsto \quad m\ddot{x}(t) = -kx(t) - f\dot{x}(t) + F(t)$$

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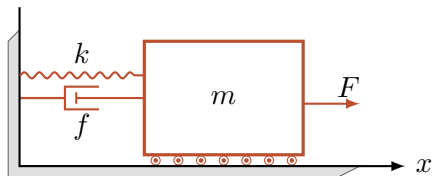
Position of a cart fastened with a spring to a wall and subject to friction:



$$\ddot{x}(t) + \frac{f}{m}\dot{x}(t) + \frac{k}{m}x(t) = \frac{1}{m}F(t) \quad (5)$$

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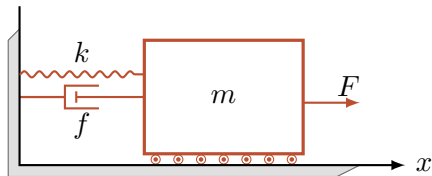
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Generalizing:

$$\ddot{x} = a_1\dot{x} + a_0x + b_0u \quad (6)$$

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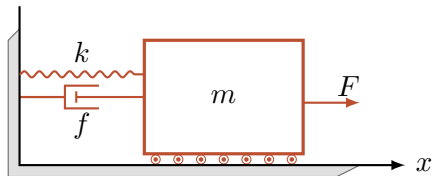
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**Discussion:** may we write this dynamics using a vectorial first order DE?

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$$\ddot{x}(t) + \frac{f}{m}\dot{x}(t) + \frac{k}{m}x(t) = \frac{1}{m}F(t) \quad (5)$$

Generalizing:

$$\ddot{x} = a_1\dot{x} + a_0x + b_0u \quad (6)$$

**Discussion:** may we write this dynamics using a vectorial first order DE?

**Discussion:** may we write this dynamics using a scalar first order DE?

## Example 1: spring-mass systems

`https://youtu.be/1ZPtFDXYQRU?t=31`

*(here there is very little friction, though)*

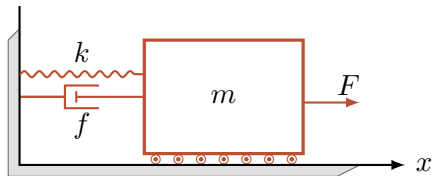
## Example 1: spring-mass systems - stability analysis from intuitive perspectives

*Discussion:*  $\dot{x} = a_0x + b_0u$  unstable if ...?

## Example 1: spring-mass systems - stability analysis from intuitive perspectives

*Discussion:*  $\dot{x} = a_0x + b_0u$  unstable if ... ? And What about, from intuitive perspectives, when considering

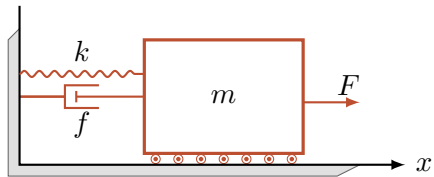
$$\ddot{x} = a_1\dot{x} + a_0x + b_0u ? \quad (7)$$



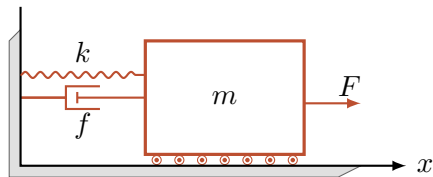


`GitHub/TTK4225/trunk/Jupyter/spring-mass-systems-introduction.ipynb`

Intuition: asymptotic stability  $\leftrightarrow$  dissipation of energy



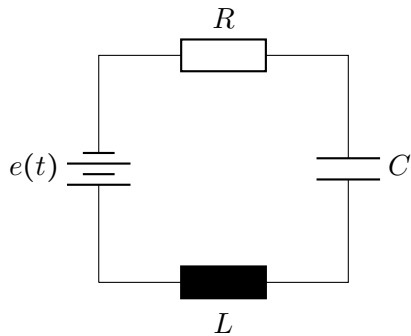
Intuition: asymptotic stability  $\leftrightarrow$  dissipation of energy



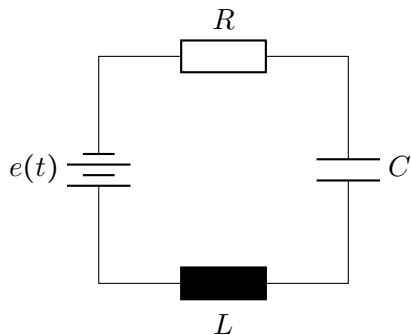
will see this MUCH better in later on courses  
when discussing about Lyapunov stability

?

## Example 2: RCL-circuit

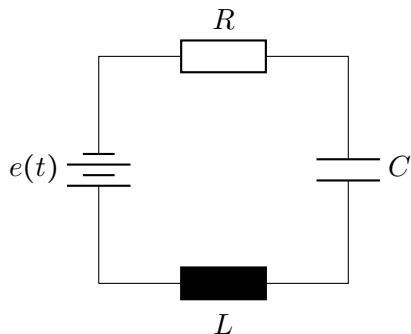


## Example 2: RCL-circuit



EOM: Kirkhoff laws

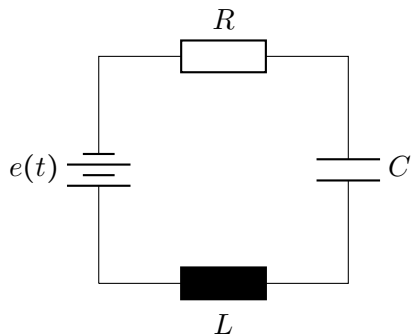
## Example 2: RCL-circuit



EOM: Kirkhoff laws

$$v_L(t) = L \frac{di(t)}{dt} \quad v_R(t) = Ri(t) \quad v_C(t) = \frac{1}{C} \int_0^t i(\tau) d\tau$$

## Example 2: RCL-circuit



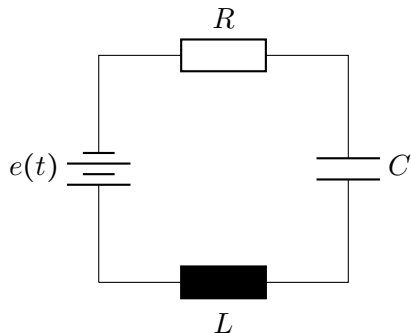
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$$e(t) = v_L(t) + v_R(t) + v_C(t)$$



## Example 2: RCL-circuit

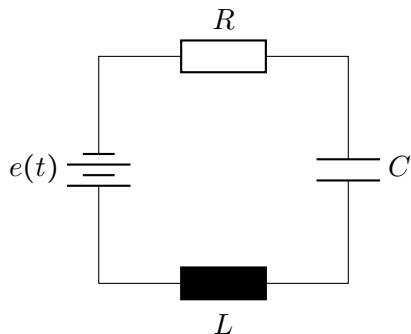


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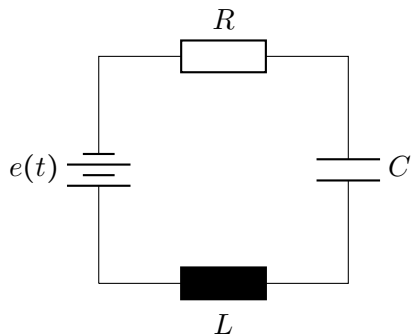
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$$e(t) = L \frac{di(t)}{dt} + Ri(t) + \frac{1}{C} \int_0^t i(\tau) d\tau$$

*Discussion:* how do we write it in terms of  $x$ ,  $\dot{x}$ , ...?

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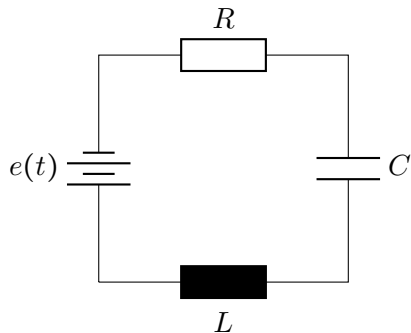


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*Discussion:* how do we write it in terms of  $x$ ,  $\dot{x}$ , ...?

$$x(t) = i(t), \quad u(t) = \dot{e}(t) \quad \implies \quad \ddot{x} = -\frac{R}{L}\dot{x} - \frac{1}{LC}x + \frac{1}{L}u \quad (9)$$

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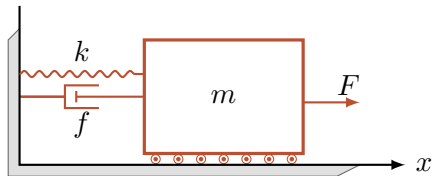


Generalizing:

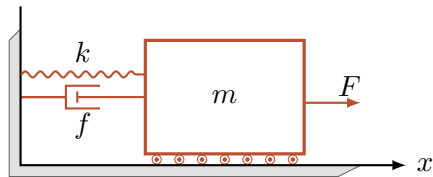
$$\ddot{x} = -\frac{R}{L}\dot{x} - \frac{1}{LC}x + \frac{1}{L}u \quad \mapsto \quad \ddot{x} = a_1\dot{x} + a_0x + b_0u \quad (10)$$

?

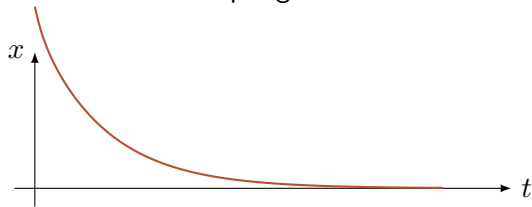
Discussion: what are the potential different ways the cart may move back to its equilibrium?



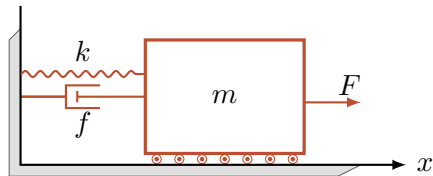
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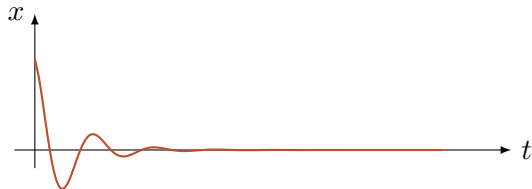
very big friction w.r.t. the spring stiffness  $\implies$  overdamping:



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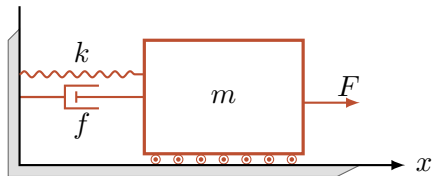


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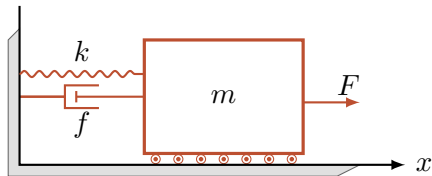
Discussion: what are the potential different ways the cart may move back to its equilibrium?



just the right amount of friction w.r.t. the spring stiffness  $\implies$  critical damping,  
i.e., *the combination that leads to the fastest decay to zero without oscillations:*



Discussion: what are the potential different ways the cart may move back to its equilibrium?

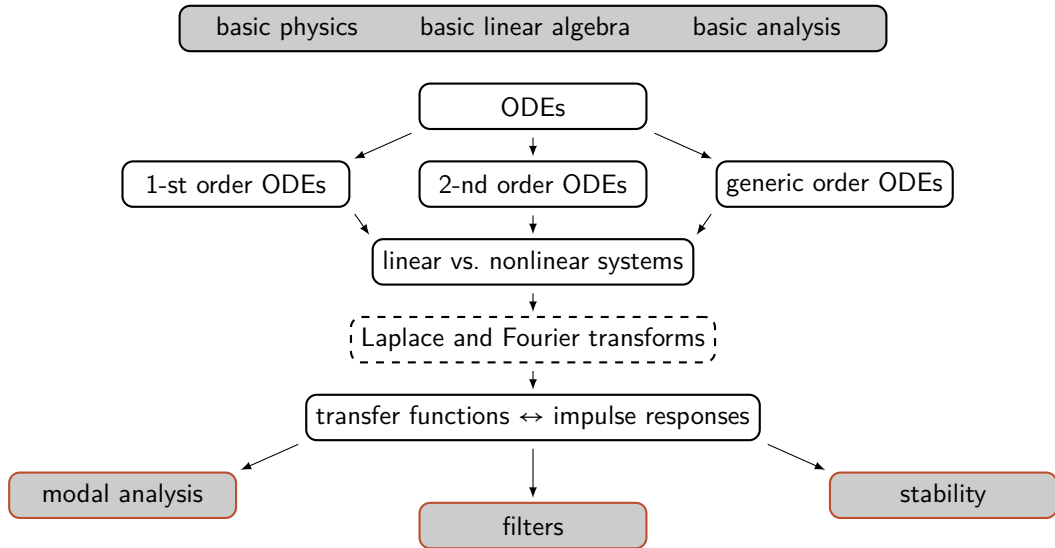


*Discussion:* can we get these behaviors with a first order system?

How do the free evolutions and forced responses look like numerically?

*we will solve this problem with the aid of Laplace transforms*

# Summary / Roadmap of TTK4225



# The essence of TTK4225

*A number like 3 is not a particular triplet of things, is about all possible triplets of things. In the same way, a model is not really about a particular system, but is rather about all the possible real-life and physical systems that behave as that model*

?

# Introduction to Laplace transforms

Laplace who?



**Figure:** Pierre-Simon Laplace, Beaumont-en-Auge, 23 March 1749 - Paris, 5 March 1827



# Roadmap

- why?
- what?
- how, at least intuitively?
- an essential example

## Why?

Remember: predictive control of the type

$$\mathbf{u}^{\star} = \arg \min_{\mathbf{u} \in \mathcal{U}, \mathbf{f}(\mathbf{u}) \in \mathcal{F}} \text{Cost}(\mathbf{f}(\mathbf{u}), \mathbf{u})$$

requires to be able to:

- define “Cost”
- compute  $\mathbf{f}(\mathbf{u})$  rapidly
- be sure that the model does not have “nasty” properties

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Laplace, for how we see it in this course,  
helps the “compute  $\mathbf{f}(\mathbf{u})$  rapidly” part

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what will be  $y(t)$  starting from  $y_0$  and for a given  $u(t)$ ?

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alternative strategy: solve  $g * u(t)$  using Laplace transforms!

# Laplace transforms - another pillar of automatic control

in TTK4225:  $\mathcal{L}$  used to compute  $y(t)$

in TTK4230:  $\mathcal{L}$  used to determine  $u(t)$

in virtually all the other courses at ITK:  $\mathcal{L}$  used for both purposes and more



## Laplace transform - definition

Hypothesis:  $f(t)$  is a real function that is locally integrable<sup>1</sup> on  $[0, +\infty)$  or  $(-\infty, +\infty)$

---

<sup>1</sup>I.e., so that its integral is finite on every compact subset of its domain of definition

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unilateral Laplace transform:

$$F(s) = \mathcal{L}\{f(t)\} := \int_0^{+\infty} e^{-st} f(t) dt \quad \text{for } t \geq 0, s = \sigma + i\omega \in \mathbb{C} \quad (12)$$

for all the  $s \in \mathbb{C}$  for which the integral in the RHS converges (*i.e., the ROC*)

---

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## Laplace transform - definition (2)

Hypothesis:  $f(t)$  is a real function that is locally integrable<sup>2</sup> on  $[0, +\infty)$  or  $(-\infty, +\infty)$

bilateral Laplace transform:

$$F(s) = \mathcal{B}\{f(t)\} := \int_{-\infty}^{+\infty} e^{-st} f(t) dt \quad \text{for } t \geq 0, s = \sigma + i\omega \in \mathbb{C} \quad (13)$$

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we need then  $\mathcal{L}^{-1} : \text{frequencies} \mapsto \text{time}$

# Inverse Laplace transform

Hypothesis:

- $F(s)$  complex function

Thesis: there may be a set of piecewise-continuous, exponentially-restricted real functions  $f(t)$  which have the property that

$$\mathcal{L}\{f(t)\} = F(s);$$

all the  $f(t)$  satisfying this equivalence differ from each other only on a subset having Lebesgue measure zero



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*Discussion:* is the existence of at least one of such  $f(t)$  guaranteed?

*Discussion:* given any  $f(t)$ , is the existence of  $F(s)$  guaranteed?

Main usefulness: convolution in time transforms into multiplication in Laplace-domain, and viceversa

$$\begin{cases} H(s) = \mathcal{L}\{h(t)\} \\ U(s) = \mathcal{L}\{u(t)\} \end{cases} \implies \mathcal{L}\{h * u(t)\} = H(s)U(s) \quad (15)$$

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*Discussion:* how then may  $\mathcal{L}$  help computing  $h * u(t)$ ?

# An intuitive explanation of the usefulness of the Laplace transform in automatic control

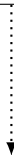
initial value problem

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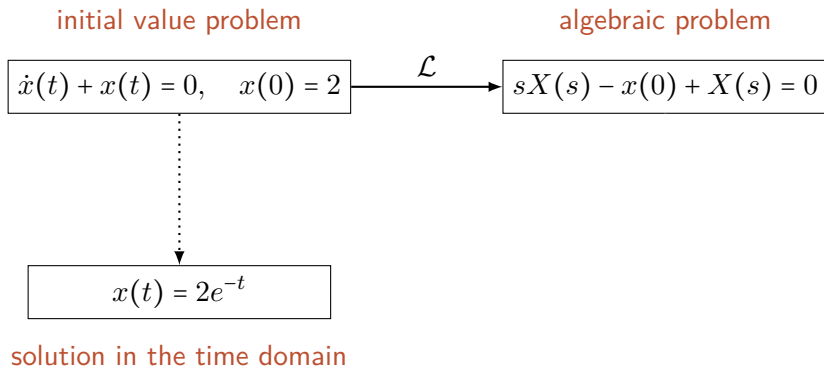
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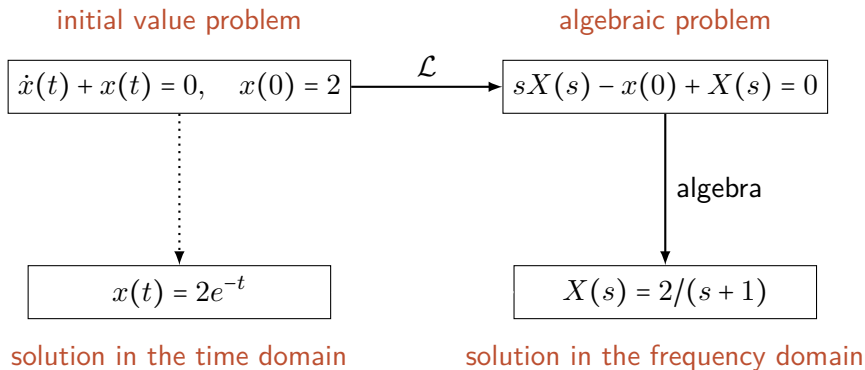
$$x(t) = 2e^{-t}$$

solution in the time domain

# An intuitive explanation of the usefulness of the Laplace transform in automatic control

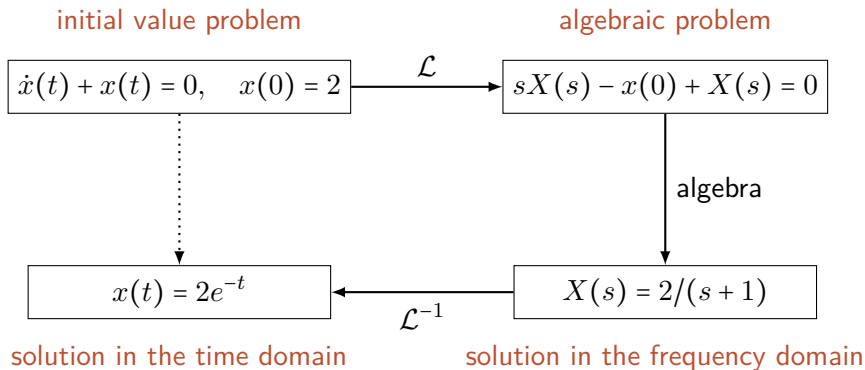


# An intuitive explanation of the usefulness of the Laplace transform in automatic control





# An intuitive explanation of the usefulness of the Laplace transform in automatic control



?

next few slides:  
introduce an utterly important example,  
and with it show the “trick” mentioned before

# The absolutely (and by far) most important Laplace transform

*Discussion:* which signal in the time domain is the most important signal you saw up to now?

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$$\dot{y} = ay + bu \quad \mapsto \quad y(t) = y_0 e^{at} + \int_0^t e^{a\tau} bu(t - \tau) d\tau$$

$$e^{at}$$

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$$\mathcal{L}\{e^{at}\}=?$$

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Remember, though:  $F(s)$  is defined for all the  $s \in \mathbb{C}$  for which the integral in the RHS converges (*i.e.*, the *ROC*). Thus, *discussion*: for which  $s$  does  $F(s)$  exist?

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$$\mathcal{L}\{e^{at}\} = \frac{1}{s-a} \quad \text{for } \operatorname{Re}(s) \geq a \quad (16)$$

At the same time ...

$$\mathcal{L}^{-1} \left\{ \frac{1}{s-a} \right\} = e^{at} \quad (17)$$

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Our game, now: use our new knowledge, i.e.,  $\mathcal{L} \{e^{at}\} = \frac{1}{s-a}$  and  $\mathcal{L}^{-1} \left\{ \frac{1}{s-a} \right\} = e^{at}$  to compute the solution to  $\dot{y} = ay$ .

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An important property that we will prove later on

(see 4.2.1 in Balchen's book)

$$\mathcal{L}\{\dot{x}\} = sX(s) - x_0 \quad (18)$$



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$$\mathcal{L}\{\dot{x}\} = sX(s) - x_0 \quad (18)$$

$$\mathcal{L}\{\ddot{x}\} = s^2X(s) - sx_0 - \dot{x}_0 \quad (19)$$

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$$\mathcal{L}\{\ddot{\ddot{x}}\} = s^3X(s) - s^2x_0 - s\dot{x}_0 - \ddot{x}_0 \quad (20)$$

$$\vdots \quad (21)$$

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$$\vdots \quad (21)$$

$s \implies \text{derivation}; \frac{1}{s} \implies \text{integration}$

*“mnemonics”*:  $s = \text{'derivasjon'}$

The “Laplace way” to find that  $\dot{y} = ay \mapsto y(t) = y_0 e^{at}$

(using that we know that  $\mathcal{L}\{\dot{y}\} = sY(s) - y_0$  and that  $\mathcal{L}^{-1}\left\{\frac{1}{s-a}\right\} = e^{at}$ )

$$\dot{y} = ay$$

The “Laplace way” to find that  $\dot{y} = ay \mapsto y(t) = y_0 e^{at}$

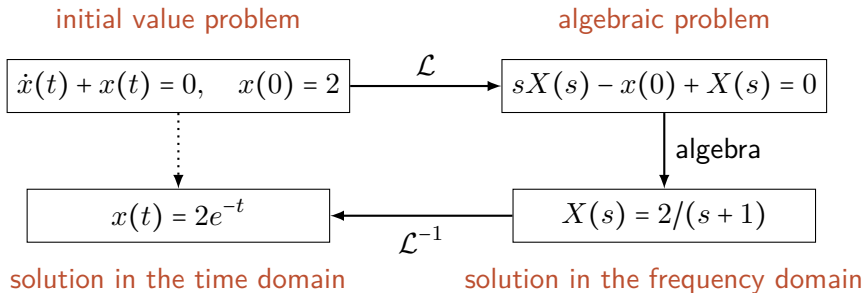
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$$\dot{y} = ay \quad \rightarrow \quad sY - y_0 = aY$$

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(using that we know that  $\mathcal{L}\{\dot{y}\} = sY(s) - y_0$  and that  $\mathcal{L}^{-1}\left\{\frac{1}{s-a}\right\} = e^{at}$ )

$$\dot{y} = ay \quad \rightarrow \quad sY - y_0 = aY \quad \rightarrow \quad Y = y_0 \frac{1}{s-a}$$



?



What is a Laplace transform, actually?

# Roadmap

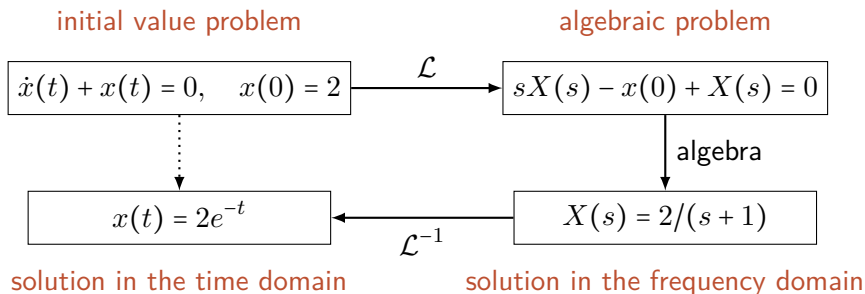
- Laplace as a generalization of Fourier
- intuitions behind Fourier transforms
- generalizing sinusoids to exponentially decaying sinusoids

In the previous unit

$$\begin{aligned}\mathcal{L}\{e^{at}\} &= \int_0^{+\infty} e^{at} e^{-st} dt \\ &= \int_0^{+\infty} e^{(a-s)t} dt \\ &= \left[ \frac{1}{a-s} e^{(a-s)t} \right]_0^{+\infty} \\ &= \frac{1}{a-s} \left[ e^{(a-s)\infty} - e^{(a-s)0} \right] \\ &= \frac{1}{s-a}\end{aligned}$$

for  $\operatorname{Re}(s) \geq a$

Main usefulness: convolution in time = multiplication in frequency,  
and this makes solving ODEs an algebraic problem



But what is a Laplace transform, eventually?

→ an opportune extension of the Fourier transform

## But what is a Fourier transform, eventually?

Definition:

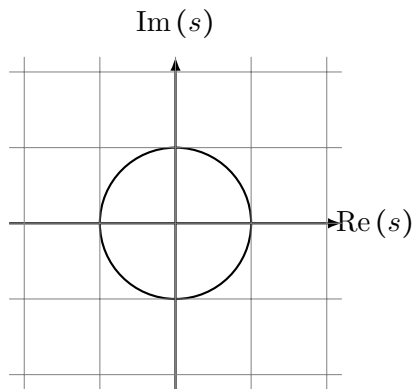
$$\mathcal{F}\{f(t)\} = \widehat{f}(\omega) := \int_{-\infty}^{+\infty} f(t)e^{-j\omega t} dt \quad (22)$$

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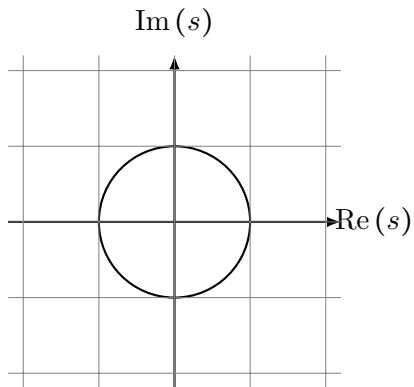
$$\mathcal{F}\{f(t)\} = \widehat{f}(\omega) := \int_{-\infty}^{+\infty} f(t)e^{-j\omega t} dt \quad (22)$$

*Discussion:* where is  $e^{-j\omega t}$ ?



Small recap: angular frequency  $\neq$  ordinary frequency

$$e^{j\omega t} = e^{j2\pi f t}$$





## Fourier transforms using Euler formula

$$\begin{aligned}\mathcal{F}\{f(t)\}(\omega) &= \widehat{f}(\omega) := \int_{-\infty}^{+\infty} f(t)e^{j\omega t} dt \\ &= \int_{-\infty}^{+\infty} f(t) \cos(\omega t) dt - j \int_{-\infty}^{+\infty} f(t) \sin(\omega t) dt\end{aligned}\tag{23}$$

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*intuition:* for each  $\omega$ ,  $\widehat{f}(\omega)$  is a couple of numbers, conveniently put together as a complex number, so that:

- $\operatorname{Re}(\widehat{f}(\omega)) = \text{area of } f(t) \cos(\omega t)$
- $\operatorname{Im}(\widehat{f}(\omega)) = \text{area of } f(t) \sin(\omega t)$

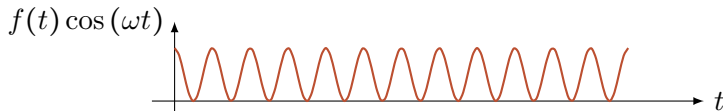
What is  $\text{Re}(\widehat{f}(\omega))$  for sinusoidal  $f$ 's?

$$\star = \int_{-\infty}^{+\infty} f(t) \cos(\omega t) dt = \int_{-\infty}^{+\infty} (\cos(\alpha t)) \cos(\omega t) dt$$

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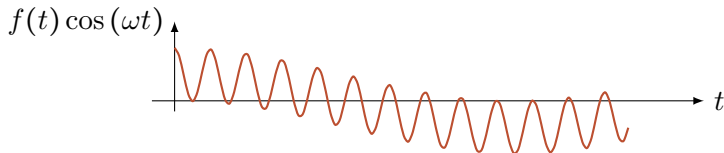
- if  $f(t) = \cos(\omega t)$  then  $\star = +\infty$



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- if  $f(t) = \cos(\omega t)$  then  $\star = +\infty$
- if  $f(t) = \cos(\alpha t)$  with  $\alpha \neq \omega$  then  $\star = 0$



What is  $\text{Re}(\widehat{f}(\omega))$  for composite sinusoidal  $f$ 's?

$$\star = \int_{-\infty}^{+\infty} f(t) \cos(\omega t) dt = \int_{-\infty}^{+\infty} \left( \sum_k \cos(\alpha_k t) \right) \cos(\omega t) dt$$

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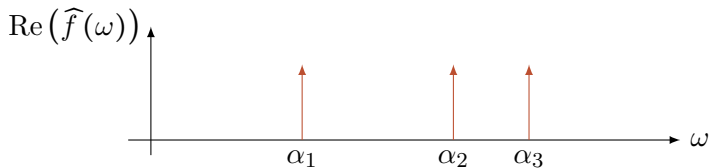
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as before: if  $\omega = \alpha_k$  then  $\star = +\infty$ , otherwise  $\star = 0$ !

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as before: if  $\omega = \alpha_k$  then  $\star = +\infty$ , otherwise  $\star = 0$ ! Thus varying  $\omega$  makes  $\text{Re}(\widehat{f}(\omega))$  kind of “comb” for the right frequency:





And what about  $\text{Im}(\widehat{f}(\omega))$ ?

$$\star = \int_{-\infty}^{+\infty} f(t) \sin(\omega t) dt$$

kind of similar, only scanning for oddity (while `cos` scans for even components)  
→ see the additional material linked later for more information

And what is  $\operatorname{Re}(\widehat{f}(\omega))$  for generic non-periodic  $f$ 's?

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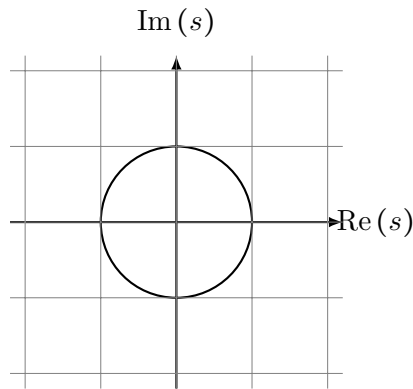
- potentially every  $\omega$  may contribute to  $\star$
- if there is no pure cos in  $f$  then each  $\omega$  contributes with a finite area

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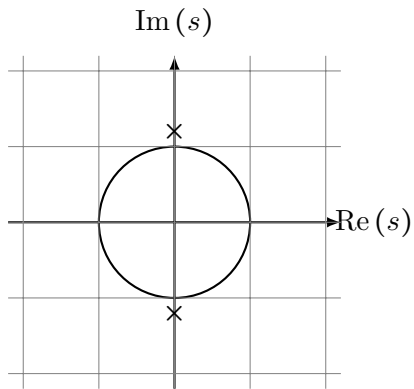
= scanning the imaginary axis for sinusoidal components:

$$\begin{aligned}\widehat{f}(\omega) &= \int_{-\infty}^{+\infty} f(t) e^{-j\omega t} dt \\ &= \int_{-\infty}^{+\infty} f(t) \cos(\omega t) dt \\ &\quad + j \int_{-\infty}^{+\infty} f(t) \sin(\omega t) dt\end{aligned}$$



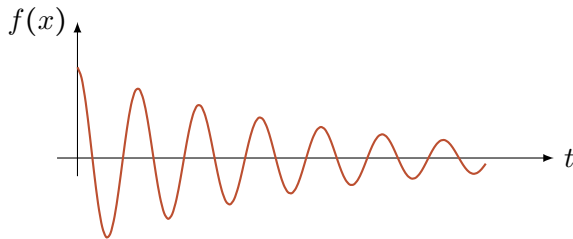
## Summarizing, graphically

$$f(t) = \cos(\omega t) = \operatorname{Re}(e^{j\omega t}) = \frac{e^{j\omega t} + e^{-j\omega t}}{2}$$



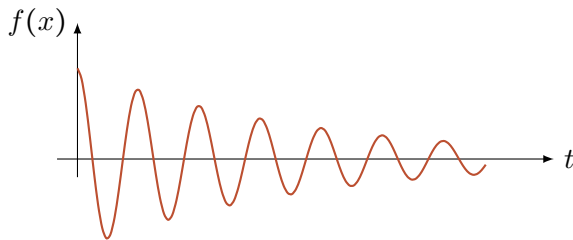
But what about exponentially decaying sinusoids?

$$f(t) = e^{at} \cos(\omega t)$$



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what if we consider  $\bar{f}(t) = f(t)e^{-at}$ ?

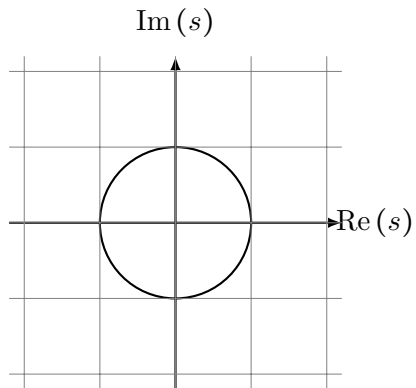


## Fourier and Laplace, side by side

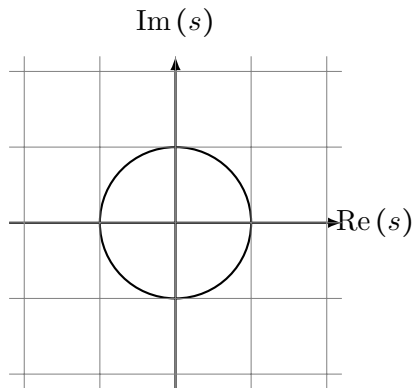
$$\mathcal{F}\{f(t)\} = \widehat{f}(\omega) := \int_{-\infty}^{+\infty} f(t)e^{-j\omega t} dt$$

$$\mathcal{L}\{f(t)\} := \int_0^{+\infty} f(t)e^{-st} dt$$

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So, what is a Laplace transform, eventually?



*beware of the region of convergence!*

?

## Suggested additional material

Fourier transforms 1:

<https://www.youtube.com/watch?v=spUNpyF58BY>

Fourier transforms 2 (with some Laplace in it):

<https://www.youtube.com/watch?v=3gjJDuC AEQQ>

Laplace transforms:

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Laplace = a pillar of automatic control

⇒ better to invest time in developing a full understanding of it

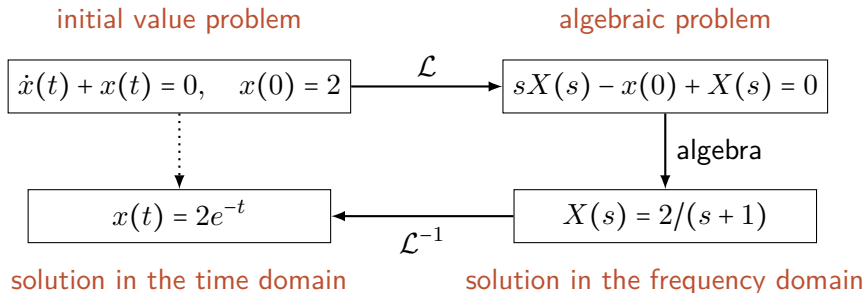
Using Laplace transforms with second order systems in free evolution

Recall: the “Laplace way” to find that  $\dot{y} = ay \mapsto y(t) = y_0 e^{at}$

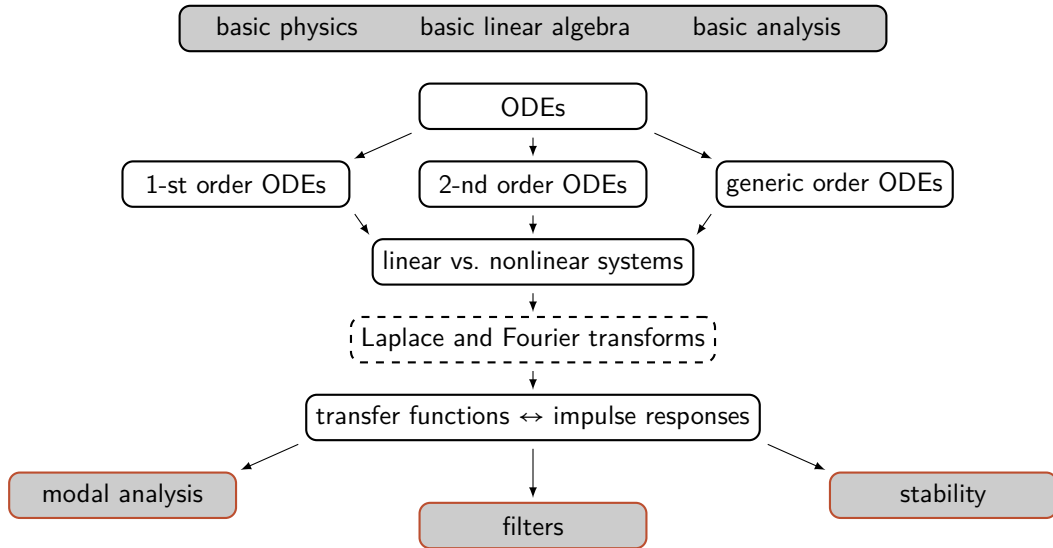
$$\dot{y} = ay \quad \rightarrow \quad sY - y_0 = aY \quad \rightarrow \quad Y = y_0 \frac{1}{s - a}$$



Recall: the “Laplace way” to find that  $\dot{y} = ay \mapsto y(t) = y_0 e^{at}$



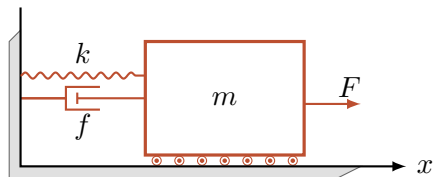
# Summary / Roadmap of TTK4225

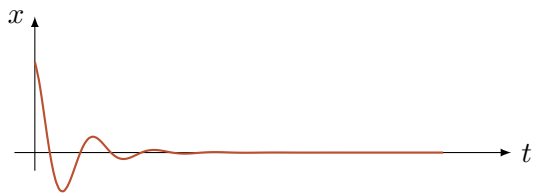
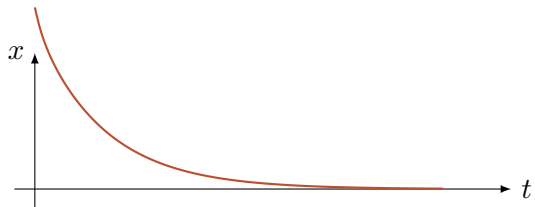


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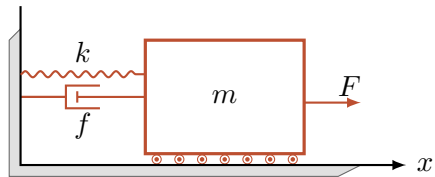
- the same thing, but for second order system
- different algebraic situations  $\implies$  different damping properties
- different ways of summarizing the different damping properties

Discussion on a practical example:  
how may a second order system free-evolve?





## Dissecting the practical example “spring-mass systems”



$$m\ddot{x}(t) = -kx(t) - f\dot{x}(t) + u(t)$$

## The properties we shall use

$$\mathcal{L}\{\dot{x}\} = sX(s) - x_0 \quad (24)$$

$$\mathcal{L}\{\ddot{x}\} = s^2X(s) - sx_0 - \dot{x}_0 \quad (25)$$

$$\mathcal{L}^{-1}\left\{\frac{1}{s-a}\right\} = e^{at} \quad (26)$$

$$(27)$$

## Solution for the free evolution case

$$\ddot{x} + \frac{f}{m}\dot{x} + \frac{k}{m}x = 0$$



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we would like to use  $\mathcal{L}^{-1} \left\{ \frac{1}{s-a} \right\} = e^{at}$  to inverse-transform  
 $\Longrightarrow$  we must try to factorize the RHS, i.e., try to write it as

$$X(s) = \frac{A}{s-\alpha} + \frac{B}{s-\beta}$$

(if possible)

## Solution for the free evolution case

Next step: factorize

$$X(s) = \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{s^2 + \frac{f}{m}s + \frac{k}{m}}$$

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*Discussion:* what are the different possibilities?

- ❶  $f^2 - 4mk > 0 \implies$  two distinct real solutions (*will imply overdamping*)
- ❷  $f^2 - 4mk = 0 \implies$  one double real solution (*will imply critical damping*)
- ❸  $f^2 - 4mk < 0 \implies$  two conjugate complex solutions (*will imply underdamping*)

$f^2 - 4mk > 0$ , i.e., overdamping in free evolution

$$X(s) = \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{s^2 + \frac{f}{m}s + \frac{k}{m}} = \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{(s - \lambda_1)(s - \lambda_2)}$$



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$$f^2 > 4mk \quad \Longrightarrow \quad \begin{cases} \lambda_1 = -\frac{f}{2m} - \frac{\sqrt{f^2 - 4mk}}{2m} \in \mathbb{R} \\ \lambda_2 = -\frac{f}{2m} + \frac{\sqrt{f^2 - 4mk}}{2m} \in \mathbb{R} \end{cases}$$

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$$\Longrightarrow \quad X(s) = \frac{A}{s - \lambda_1} + \frac{B}{s - \lambda_2}$$

*Discussion:* can we conclude immediately that

$$X(s) = \frac{A}{s - \lambda_1} + \frac{B}{s - \lambda_2} \implies x(t) = Ae^{\lambda_1 t} + Be^{\lambda_2 t}$$

or do we need to say some property about Laplace transforms that we did not say before?

We need to be sure that Laplace transforms are linear!

i.e., be sure that

$$\mathcal{L}\{\alpha x_1 + \beta x_2\} = \alpha \mathcal{L}\{x_1\} + \beta \mathcal{L}\{x_2\},$$

and that

$$\mathcal{L}^{-1}\{\alpha X_1 + \beta X_2\} = \alpha \mathcal{L}^{-1}\{X_1\} + \beta \mathcal{L}^{-1}\{X_2\}$$

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*Discussion:* why do the previous properties hold? (hint: recall that  $\mathcal{L}\{x(t)\} := \int_0^{+\infty} x(t)e^{-st}dt$ )

## In conclusion: overdamping in free evolution

$$f^2 > 4mk \quad \Longrightarrow \quad X(s) = \frac{A}{s - \lambda_1} + \frac{B}{s - \lambda_2} \quad \Longrightarrow \quad x(t) = Ae^{\lambda_1 t} + Be^{\lambda_2 t}$$

for opportune  $A$  and  $B$ .

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for opportune  $A$  and  $B$ . *Discussion:* what determines  $A$  and  $B$ ? Hint: recall that

$$X(s) = \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{s^2 + \frac{f}{m}s + \frac{k}{m}} = \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{(s - \lambda_1)(s - \lambda_2)}$$

important message (will be seen again later on): the denominator of the rational Laplace transform determines the behavior of the system, while the numerator determines which root is dominant and the “amplitude” of each mode

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## Practical example of a overdamped system in free evolution

$$\left\{ \begin{array}{lcl} m & = & 10 \quad \text{kg} \\ f & = & 100 \quad \text{kg/s} \\ d & = & 90 \quad \text{N/s}^2 \end{array} \right. \implies \left\{ \ddot{x} + \frac{100}{10} \dot{x} + \frac{90}{10} x = 0 \right.$$

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*note that there are 2 distinct steps!!*

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ARMA-continuous-time-example.ipynb

?

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*Discussion:* did we learn until now in this course how to inverse-Laplace these things?

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Very interesting behavior:

$$x(t) = Ae^{\lambda t} + Bte^{\lambda t}$$

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Very interesting behavior:

$$x(t) = Ae^{\lambda t} + Bte^{\lambda t}$$

*Discussion:* what contributes to determine  $A$  and  $B$ ?

*Discussion:* recall the previous “overdamping case”, i.e.,

$$\left\{ \begin{array}{lll} X(s) & = & \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{s^2 + \frac{f}{m}s + \frac{k}{m}} \\ m & = & 10 \quad \text{kg} \\ f & = & 100 \quad \text{kg/s} \\ d & = & 90 \quad \text{N/s}^2 \\ x(0) & = & 0.16\text{m} \\ \dot{x}(0) & = & 0\text{m/s} \end{array} \right. \implies X(s) = \frac{0.16(s+10)}{s^2+10s+9}.$$

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ARMA-continuous-time-example.ipynb

?

$f^2 - 4mk < 0$ , i.e., underdamping in free evolution

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$$\Longrightarrow \quad X(s) = \frac{A}{s - \lambda_1} + \frac{B}{s - \lambda_2}$$

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How is it possible that a real system produces complex results?”

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Solution:

$$\mathcal{L}\{e^{at} \sin \omega t\} = \frac{\omega}{(s - a)^2 + \omega^2}$$

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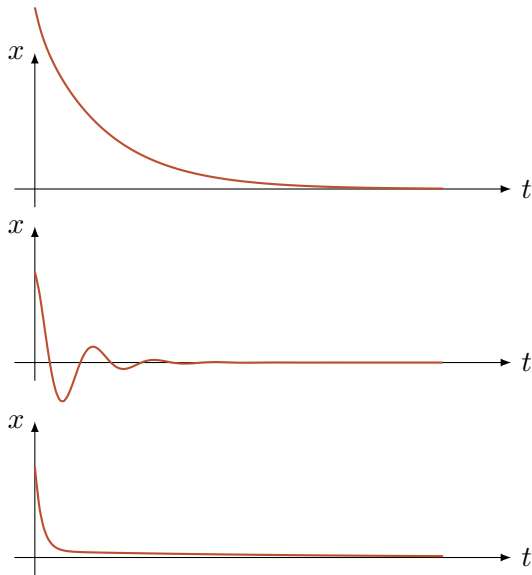
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?



## Recap of the possible damping modalities



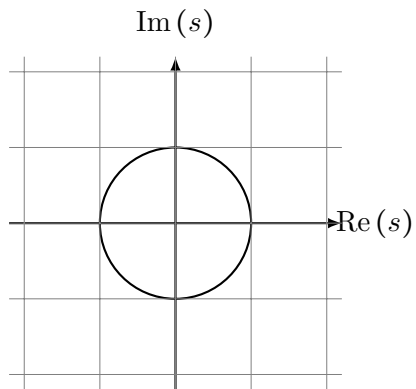
## Recap of the corresponding time-frequency pairs

overdamped:  $\mathcal{L}\{Ae^{\lambda_1 t} + Bte^{\lambda_2 t}\} = \frac{A}{s - \lambda_1} + \frac{B}{s - \lambda_2}, \quad \lambda_1 \neq \lambda_2$

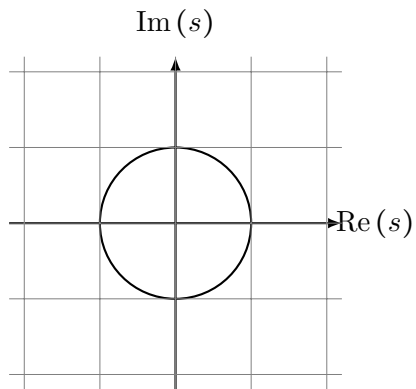
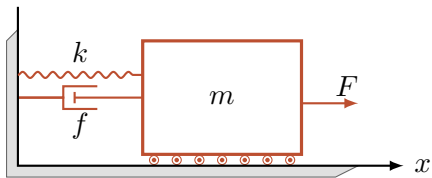
critically damped:  $\mathcal{L}\{Ae^{\lambda t} + Bte^{\lambda t}\} = \frac{A}{s - \lambda} + \frac{B}{(s - \lambda)^2}$

underdamped:  $\mathcal{L}\{e^{at} \sin \omega t\} = \frac{\omega}{(s - a)^2 + \omega^2}$

## Recap using the “poles-zeros plot”



Do you see it?



?

Using Laplace transforms with generic  $n$  order systems in free evolution

In the last episodes: the “Laplace way” to find that  $\dot{y} = ay \mapsto y(t) = y_0 e^{at}$ ,  
and to do something similar for second order systems in free evolution

$$\dot{y} = ay \quad \rightarrow \quad sY - y_0 = aY \quad \rightarrow \quad Y = y_0 \frac{1}{s - a}$$

# Roadmap

- the same thing, but for generic systems



## A small recap

- first order homogeneous systems:

$$\dot{x} - ax = 0 \quad \mapsto \quad x(t) = x_0 e^{at} \quad (28)$$

- first order non-homogeneous systems:

$$\dot{x} - ax = bu \quad \mapsto \quad x(t) = x(0)e^{at} + \int_0^t e^{a(t-\tau)} bu(\tau) d\tau \quad (29)$$

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*The problem:* the convolution integral gives us headaches, and we want to use Laplace to ease things up. To understand how, step-by-step path:

- 1 start from  $\dot{x} - ax = 0$
- 2 move to  $\ddot{x} - a_1\dot{x} - a_2x = 0$
- 3 move to higher order
- 4 add the “ $u$ ” part

# Laplace-properties we saw up to now:

(4.2.1 in Balchen's book)

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## Linearity

$$\mathcal{L}\{af(t) + bg(t)\} = a\mathcal{L}\{f(t)\} + b\mathcal{L}\{g(t)\} \quad (30)$$

# Laplace-properties we saw up to now:

(4.2.1 in Balchen's book)

## Derivatives

$$\mathcal{L}\{\dot{x}\} = sX(s) - x(0) \quad (31)$$

$$\mathcal{L}\{\ddot{x}\} = s^2X(s) - sx(0) - \dot{x}(0) \quad (32)$$

$$\vdots \quad (33)$$

## Laplace-properties we saw up to now:

(4.2.1 in Balchen's book)

### Integration

$$\mathcal{L}\left\{\int_0^t f(\tau)d\tau\right\} = \frac{1}{s}F(s) \quad (34)$$

## Laplace-transforms we needed up to now:

$$f(t) = \mathcal{L}^{-1} \{F(s)\} \qquad F(s) = \mathcal{L} \{f(t)\}$$

---

$$e^{at}$$

$$\frac{1}{s-a}$$

$$e^{at} \sin \omega t$$

$$\frac{\omega}{(s-a)^2 + \omega^2}$$

$$te^{at}$$

$$\frac{1}{(s-a)^2}$$

## A more detailed roadmap for this unit:

- make Laplace-transform more mathematically ground
- derive some few very important transforms
- generalize how to do partial expansions
- solve the problem: how to use Laplace transforms with generic  $n$  order systems in free evolution



?

## Laplace transforms, take 2

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## Important (but intuitive) theorem

If  $f(t)$  is defined and piecewise continuous for  $t > 0$ ,  
and there exist two constants  $M$  and  $k$  so that  $|f(t)| \leq Me^{kt}$  for every  $t > 0$ ,  
then  $\mathcal{L}\{f(t)\}$  is well defined for every  $s > k$



## Important (but intuitive) property: linearity

Assuming that  $f(t)$  and  $g(t)$  are so that  $\mathcal{L}\{f(t)\}$  and  $\mathcal{L}\{g(t)\}$  exists, then  $\mathcal{L}\{af(t) + bg(t)\}$  exists and is s.t.

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### Proof

Integrals are linear, thus

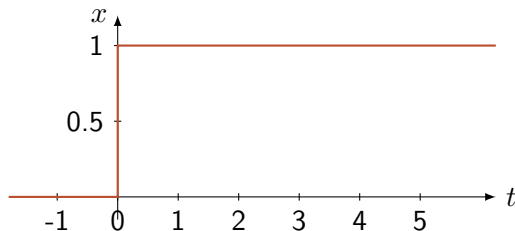
$$\begin{aligned}\mathcal{L}\{af(t) + bg(t)\} &= \int_0^{+\infty} e^{-st} [af(t) + bg(t)] dt \\ &= a \int_0^{+\infty} e^{-st} f(t) dt + b \int_0^{+\infty} e^{-st} g(t) dt \\ &= a\mathcal{L}\{f(t)\} + b\mathcal{L}\{g(t)\}\end{aligned}$$

?

# Important example: Laplace-transforming the Heaviside-function

(right now on the “deriving some few very important transforms” of the unit)

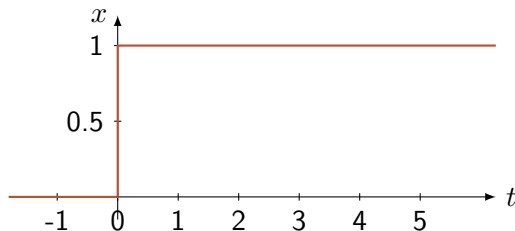
$$H(t) := \begin{cases} 0 & \text{for } t < 0 \\ 1 & \text{for } t > 0 \end{cases}$$



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*Discussion:* is  $H(0)$  defined?

# Deriving the Heaviside function from the Dirac's impulse

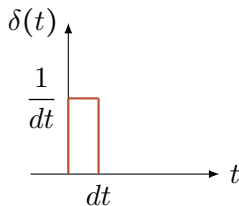
(right now on the “deriving some few very important transforms” of the unit)

$$H(t) = \int_{-\infty}^t \delta(\tau) d\tau$$

with, *very loosely*,

$$\delta(t) := \begin{cases} +\infty & \text{if } t = 0 \\ 0 & \text{otherwise,} \end{cases}$$

or, better, as the limit of concentrating an unitary point mass at 0:



## Back to the important example: Laplace-transforming the Heaviside-function

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$$\int_0^{+\infty} u \dot{v} dt = [uv]_0^{+\infty} - \int_0^{+\infty} \dot{u} v dt \quad (41)$$

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$$\mathcal{L}\{t\} = \int_0^{+\infty} t e^{-st} dt = \star \quad \begin{array}{ll} u = t & v = e^{-st} \\ \dot{u} = 1 & \dot{v} = -s e^{-st} \end{array} \quad (42)$$

$$\star = \frac{1}{-s} \int_0^{+\infty} t (-s e^{-st}) dt \quad (43)$$

# On with the important examples: Laplace-transforming polynomials

(right now on the “deriving some few very important transforms” of the unit)

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$$\begin{aligned}\mathcal{L}\{t^2\} &= \int_0^{+\infty} t^2 e^{-st} dt & \begin{aligned} u &= t^2 & v &= e^{-st} \\ \dot{u} &= 2t & \dot{v} &= -s e^{-st} \end{aligned} \end{aligned} \quad (46)$$

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$$(50)$$

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$$\mathcal{L}\{t^n\} = \frac{(n-1)!}{s^n} \quad (51)$$

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*Discussion:* what is the ROC?  $\operatorname{Re}(s) > 0$

## On with the important examples: Laplace-transforms of derivatives

(right now on the “deriving some few very important transforms” of the unit)

$$\mathcal{L}\{\dot{x}\} = \int_0^{+\infty} e^{-st} \dot{x} dt \quad (52)$$



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$$(61)$$



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$$\mathcal{L}\left\{x^{(n)}(t)\right\} = s^n \mathcal{L}\{x(t)\} - s^{n-1}x(0) - s^{n-2}\dot{x}(0) - \dots - x^{(n-1)}(0) \quad (62)$$

?

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The potential points of view of this course:

simply a complex number, i.e.,  $s = \sigma + j\omega$ , that gives rise to a *complex frequency* through the complex exponential  $e^{-st} = e^{-\sigma t} e^{-j\omega t}$

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Be careful on how to use  $s$  in this course!

Typically, in all the automatic control courses the interest is not too much on the value of  $F(s)$  for different  $s$ , but rather on the fact that  $F(s)$  enables performing algebraic operations to solve ODEs!



Indeed...

$$\begin{aligned}\ddot{x} + \frac{f}{m}\dot{x} + \frac{k}{m}x &= 0 \\ \Downarrow \\ X(s) &= \frac{x_0\left(\frac{f}{m} + s\right) + \dot{x}_0}{s^2 + \frac{f}{m}s + \frac{k}{m}} \Downarrow \\ x(t) &= Ae^{\lambda_1 t} + Be^{\lambda_2 t} \text{ (and the other cases)}\end{aligned}$$

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*Discussion:* why was it that  $\ddot{x} + \frac{f}{m}\dot{x} + \frac{k}{m}x = 0$  leads to  $X(s) \neq 0$ ?

what about  $\ddot{x} + \alpha_2 \ddot{x} + \alpha_1 \dot{x} + \alpha_0 x = 0$ ?

what about  $\ddot{x} + \alpha_2 \ddot{x} + \alpha_1 \dot{x} + \alpha_0 x = 0$ ?

and what about higher orders?

# Partial fraction decomposition

## Partial fraction decomposition

case single poles: if  $\frac{N(s)}{(s-a)(s-b)(s-c)\dots}$  is s.t.  $a \neq b \neq c \neq \dots$  then there exist  $A, B, C, \dots$  s.t.

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“But how do I compute  $A, B$ , etc.?”  $\mapsto$

`en.wikipedia.org/wiki/Partial_fraction_decomposition`

(tip: start from `en.wikipedia.org/wiki/Heaviside_cover-up_method`)



?

## Computing the free evolution for generic $n$ order LTI systems

$$\text{start from } x^{(n)} + \alpha_{n-1}x^{(n-1)} + \dots + \alpha_0x = 0 \quad (65)$$

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$$\text{do a partial fraction decomposition} \quad (68)$$

# Computing the free evolution for generic $n$ order LTI systems

Two main cases:

$$x^{(n)} + \alpha_{n-1}x^{(n-1)} + \dots + \alpha_0x = 0$$

↓

$$\text{either } X(s) = \frac{A}{s-a} + \frac{B}{s-b} + \frac{C}{s-c} + \dots$$

$$\text{or } X(s) = \frac{A_1}{s-a} + \dots + \frac{A_{na}}{(s-a)^{na}} + \frac{B_1}{s-b} + \dots + \frac{B_{nb}}{(s-b)^{nb}} + \dots$$

# Computing the free evolution for generic $n$ order LTI systems

Case single poles

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*Discussion:* given  $\mathcal{L}^{-1}\left\{\frac{1}{s-\alpha}\right\} = e^{\alpha t}$ , what is  $x(t)$ ?



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$$\text{generalization: } \mathcal{L}\left\{\frac{n!}{(s-\alpha)^{n+1}}\right\} = t^n e^{\alpha t}$$

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*Discussion:* given the previous inverse transform, what is  $x(t)$ ?

## Extremely important result

a LTI in free evolution behaves as a combination of terms  $e^{\alpha t}$ ,  $te^{\alpha t}$ ,  $t^2e^{\alpha t}$ , etc. for a set of different  $\alpha$ s and powers of  $t$ , called the *modes* of the system

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*Discussion:* assuming that we have two modes,  $e^{-0.3t}$  and  $e^{-1.6t}$ , so that

$$x(t) = \gamma_1 e^{-0.3t} + \gamma_2 e^{-1.6t}.$$

What determines  $\gamma_1$  and  $\gamma_2$ ?

?