

# TTK4225 - Systems Theory, Autumn 2020

Damiano Varagnolo



# Introduction to Laplace transforms

Laplace who?



**Figure:** Pierre-Simon Laplace, Beaumont-en-Auge, 23 March 1749 - Paris, 5 March 1827

# Roadmap

- why?
- what?
- how, at least intuitively?
- an essential example

## Why?

Remember: predictive control of the type

$$\mathbf{u}^{\star} = \arg \min_{\mathbf{u} \in \mathcal{U}, \mathbf{f}(\mathbf{u}) \in \mathcal{F}} \text{Cost}(\mathbf{f}(\mathbf{u}), \mathbf{u})$$

requires to be able to:

- define “Cost”
- compute  $\mathbf{f}(\mathbf{u})$  rapidly
- be sure that the model does not have “nasty” properties

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Laplace, for how we see it in this course,  
helps the “compute  $\mathbf{f}(\mathbf{u})$  rapidly” part

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what will be  $y(t)$  starting from  $y_0$  and for a given  $u(t)$ ?

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alternative strategy: solve  $g * u(t)$  using Laplace transforms!

# Laplace transforms - another pillar of automatic control

in TTK4225:  $\mathcal{L}$  used to compute  $y(t)$

in TTK4230:  $\mathcal{L}$  used to determine  $u(t)$

in virtually all the other courses at ITK:  $\mathcal{L}$  used for both purposes and more

## Laplace transform - definition

Hypothesis:  $f(t)$  is a real function that is locally integrable<sup>1</sup> on  $[0, +\infty)$  or  $(-\infty, +\infty)$

---

<sup>1</sup>I.e., so that its integral is finite on every compact subset of its domain of definition

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for all the  $s \in \mathbb{C}$  for which the integral in the RHS converges (*i.e., the ROC*)

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## Laplace transform - definition (2)

Hypothesis:  $f(t)$  is a real function that is locally integrable<sup>2</sup> on  $[0, +\infty)$  or  $(-\infty, +\infty)$

bilateral Laplace transform:

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we need then  $\mathcal{L}^{-1} : \text{frequencies} \mapsto \text{time}$

# Inverse Laplace transform

Hypothesis:

- $F(s)$  complex function

Thesis: there may be a set of piecewise-continuous, exponentially-restricted real functions  $f(t)$  which have the property that

$$\mathcal{L}\{f(t)\} = F(s);$$

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Main usefulness: convolution in time transforms into multiplication in Laplace-domain, and viceversa

$$\begin{cases} H(s) = \mathcal{L}\{h(t)\} \\ U(s) = \mathcal{L}\{u(t)\} \end{cases} \implies \mathcal{L}\{h * u(t)\} = H(s)U(s) \quad (5)$$

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*Discussion:* how then may  $\mathcal{L}$  help computing  $h * u(t)$ ?

# An intuitive explanation of the usefulness of the Laplace transform in automatic control

initial value problem

$$\dot{x}(t) + x(t) = 0, \quad x(0) = 2$$



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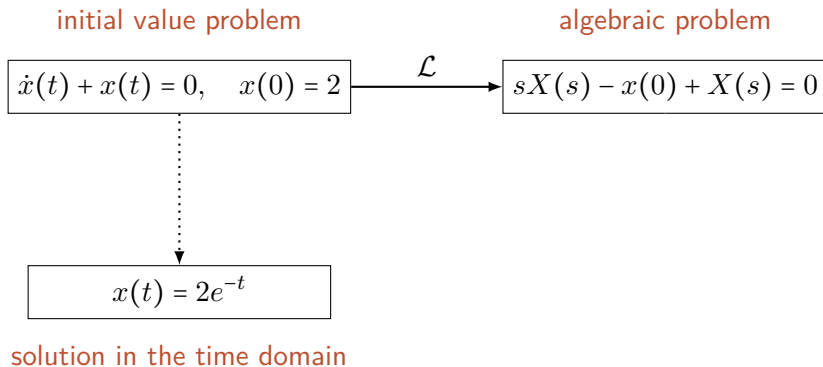
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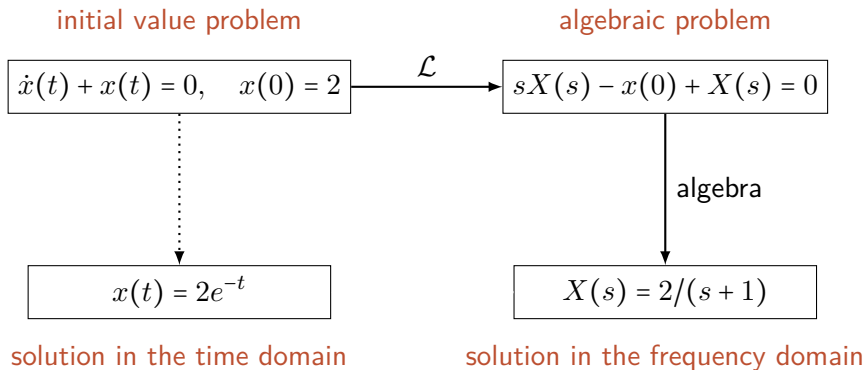
$$x(t) = 2e^{-t}$$

solution in the time domain

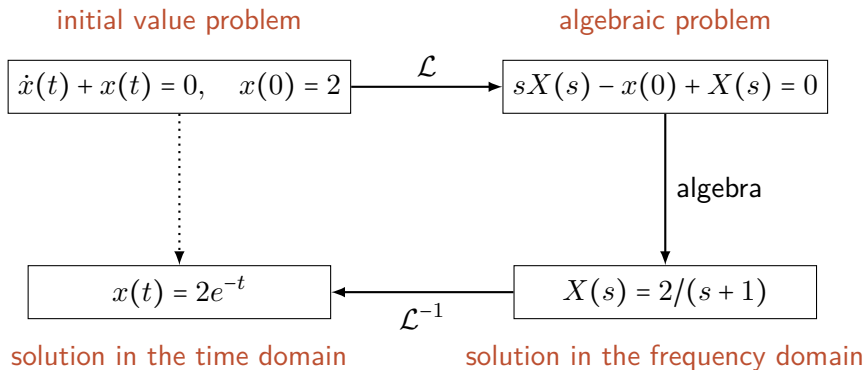
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?

next few slides:  
introduce an utterly important example,  
and with it show the “trick” mentioned before

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*Discussion:* which signal in the time domain is the most important signal you saw up to now?

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$$\dot{y} = ay + bu \quad \mapsto \quad y(t) = y_0 e^{at} + \int_0^t e^{a\tau} bu(t - \tau) d\tau$$



$$e^{at}$$

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Remember, though:  $F(s)$  is defined for all the  $s \in \mathbb{C}$  for which the integral in the RHS converges (*i.e.*, the *ROC*). Thus, *discussion*: for which  $s$  does  $F(s)$  exist?

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$$\mathcal{L}\{e^{at}\} = \frac{1}{s-a} \quad \text{for } \operatorname{Re}(s) \geq a \quad (6)$$

At the same time ...

$$\mathcal{L}^{-1} \left\{ \frac{1}{s-a} \right\} = e^{at} \quad (7)$$



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Our game, now: use our new knowledge, i.e.,  $\mathcal{L} \{ e^{at} \} = \frac{1}{s-a}$  and  $\mathcal{L}^{-1} \left\{ \frac{1}{s-a} \right\} = e^{at}$  to compute the solution to  $\dot{y} = ay$ .

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Our game, now: use our new knowledge, i.e.,  $\mathcal{L} \{ e^{at} \} = \frac{1}{s-a}$  and  $\mathcal{L}^{-1} \left\{ \frac{1}{s-a} \right\} = e^{at}$  to compute the solution to  $\dot{y} = ay$ . *Discussion:* how did we find that  $\dot{y} = ay \mapsto y(t) = y_0 e^{at}$  before?

An important property that we will prove later on

(see 4.2.1 in Balchen's book)

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$s \implies$  derivation;  $\frac{1}{s} \implies$  integration

*“mnemonics”*:  $s = \text{'derivasjon'}$

The “Laplace way” to find that  $\dot{y} = ay \mapsto y(t) = y_0 e^{at}$

(using that we know that  $\mathcal{L}\{\dot{y}\} = sY(s) - y_0$  and that  $\mathcal{L}^{-1}\left\{\frac{1}{s-a}\right\} = e^{at}$ )

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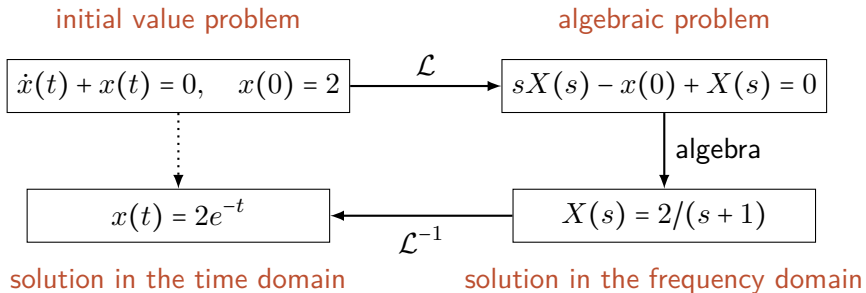
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$$\dot{y} = ay \quad \rightarrow \quad sY - y_0 = aY \quad \rightarrow \quad Y = y_0 \frac{1}{s-a}$$



?

What is a Laplace transform, actually?

# Roadmap

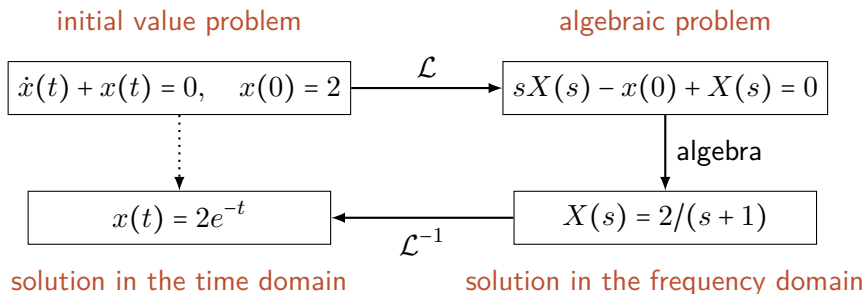
- Laplace as a generalization of Fourier
- intuitions behind Fourier transforms
- generalizing sinusoids to exponentially decaying sinusoids

In the previous unit

$$\begin{aligned}\mathcal{L}\{e^{at}\} &= \int_0^{+\infty} e^{at} e^{-st} dt \\&= \int_0^{+\infty} e^{(a-s)t} dt \\&= \left[ \frac{1}{a-s} e^{(a-s)t} \right]_0^{+\infty} \\&= \frac{1}{a-s} \left[ e^{(a-s)\infty} - e^{(a-s)0} \right] \\&= \frac{1}{s-a}\end{aligned}$$

for  $\operatorname{Re}(s) \geq a$

Main usefulness: convolution in time = multiplication in frequency,  
and this makes solving ODEs an algebraic problem



But what is a Laplace transform, eventually?

→ an opportune extension of the Fourier transform



## But what is a Fourier transform, eventually?

Definition:

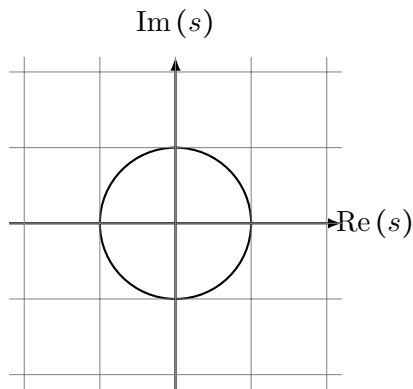
$$\mathcal{F}\{f(t)\} = \widehat{f}(\omega) := \int_{-\infty}^{+\infty} f(t)e^{-j\omega t} dt \quad (12)$$

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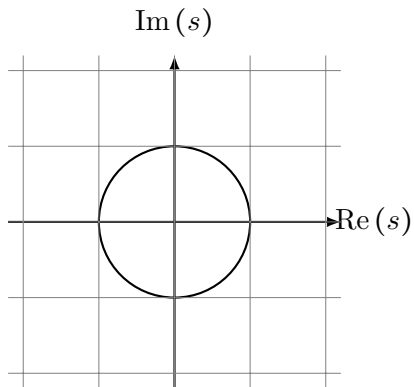
$$\mathcal{F}\{f(t)\} = \widehat{f}(\omega) := \int_{-\infty}^{+\infty} f(t)e^{-j\omega t} dt \quad (12)$$

*Discussion:* where is  $e^{-j\omega t}$ ?



Small recap: angular frequency  $\neq$  ordinary frequency

$$e^{j\omega t} = e^{j2\pi f t}$$



## Fourier transforms using Euler formula

$$\begin{aligned}\mathcal{F}\{f(t)\}(\omega) &= \widehat{f}(\omega) := \int_{-\infty}^{+\infty} f(t)e^{j\omega t} dt \\ &= \int_{-\infty}^{+\infty} f(t) \cos(\omega t) dt - j \int_{-\infty}^{+\infty} f(t) \sin(\omega t) dt\end{aligned}\tag{13}$$

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*intuition:* for each  $\omega$ ,  $\widehat{f}(\omega)$  is a couple of numbers, conveniently put together as a complex number, so that:

- $\operatorname{Re}(\widehat{f}(\omega)) = \text{area of } f(t) \cos(\omega t)$
- $\operatorname{Im}(\widehat{f}(\omega)) = \text{area of } f(t) \sin(\omega t)$

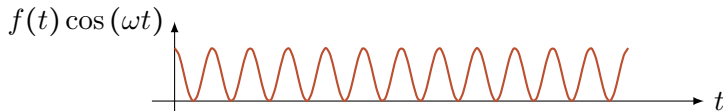
What is  $\text{Re}(\widehat{f}(\omega))$  for sinusoidal  $f$ 's?

$$\star = \int_{-\infty}^{+\infty} f(t) \cos(\omega t) dt = \int_{-\infty}^{+\infty} (\cos(\alpha t)) \cos(\omega t) dt$$

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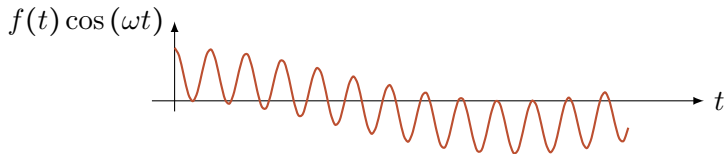
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- if  $f(t) = \cos(\omega t)$  then  $\star = +\infty$
- if  $f(t) = \cos(\alpha t)$  with  $\alpha \neq \omega$  then  $\star = 0$





What is  $\text{Re}(\widehat{f}(\omega))$  for composite sinusoidal  $f$ 's?

$$\star = \int_{-\infty}^{+\infty} f(t) \cos(\omega t) dt = \int_{-\infty}^{+\infty} \left( \sum_k \cos(\alpha_k t) \right) \cos(\omega t) dt$$

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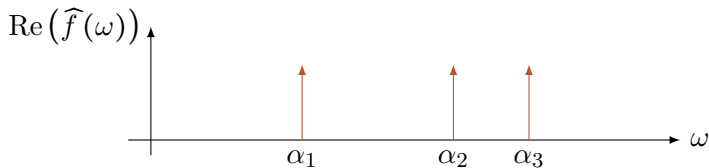
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as before: if  $\omega = \alpha_k$  then  $\star = +\infty$ , otherwise  $\star = 0$ !

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as before: if  $\omega = \alpha_k$  then  $\star = +\infty$ , otherwise  $\star = 0$ ! Thus varying  $\omega$  makes  $\text{Re}(\widehat{f}(\omega))$  kind of “comb” for the right frequency:



And what about  $\text{Im}(\widehat{f}(\omega))$ ?

$$\star = \int_{-\infty}^{+\infty} f(t) \sin(\omega t) dt$$

kind of similar, only scanning for oddity (while  $\cos$  scans for even components)  
→ see the additional material linked later for more information

And what is  $\operatorname{Re}(\widehat{f}(\omega))$  for generic non-periodic  $f$ 's?

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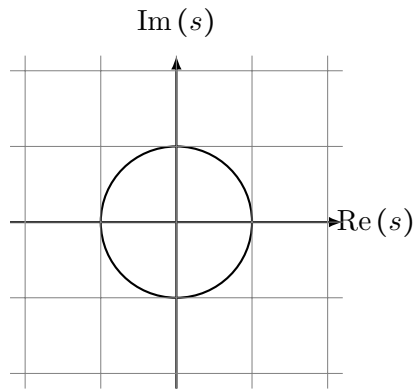
- potentially every  $\omega$  may contribute to  $\star$
- if there is no pure cos in  $f$  then each  $\omega$  contributes with a finite area

So, what is a Fourier transform, eventually?

## So, what is a Fourier transform, eventually?

= scanning the imaginary axis for sinusoidal components:

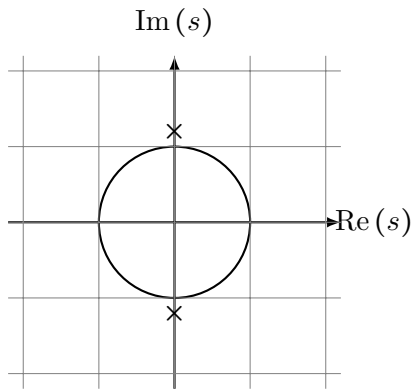
$$\begin{aligned}\widehat{f}(\omega) &= \int_{-\infty}^{+\infty} f(t) e^{-j\omega t} dt \\ &= \int_{-\infty}^{+\infty} f(t) \cos(\omega t) dt \\ &\quad + j \int_{-\infty}^{+\infty} f(t) \sin(\omega t) dt\end{aligned}$$





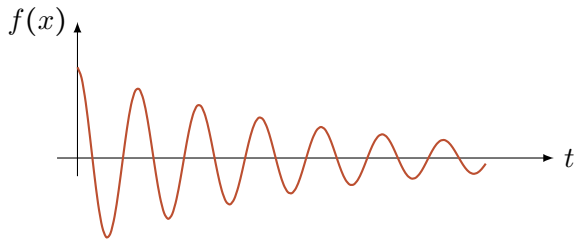
## Summarizing, graphically

$$f(t) = \cos(\omega t) = \operatorname{Re}(e^{j\omega t}) = \frac{e^{j\omega t} + e^{-j\omega t}}{2}$$



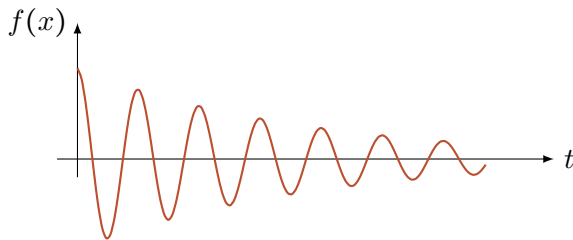
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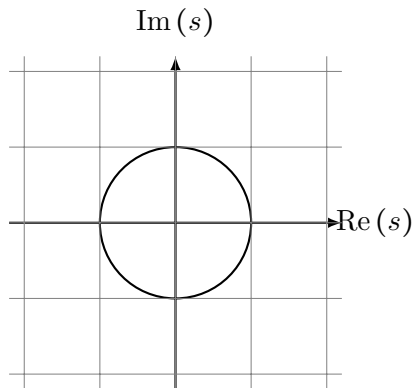
what if we consider  $\bar{f}(t) = f(t)e^{-at}$ ?

## Fourier and Laplace, side by side

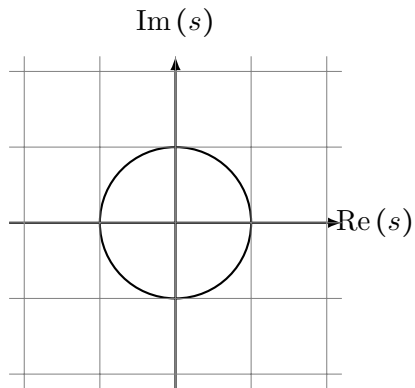
$$\mathcal{F}\{f(t)\} = \widehat{f}(\omega) := \int_{-\infty}^{+\infty} f(t)e^{-j\omega t} dt$$

$$\mathcal{L}\{f(t)\} := \int_0^{+\infty} f(t)e^{-st} dt$$

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*beware of the region of convergence!*

?

## Suggested additional material

Fourier transforms 1:

<https://www.youtube.com/watch?v=spUNpyF58BY>

Fourier transforms 2 (with some Laplace in it):

<https://www.youtube.com/watch?v=3gjJDuC AEQQ>

Laplace transforms:

<https://www.youtube.com/watch?v=n2y7n6jw5d0>



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<https://www.youtube.com/watch?v=spUNpyF58BY>

Fourier transforms 2 (with some Laplace in it):

<https://www.youtube.com/watch?v=3gjJDuC AEQQ>

Laplace transforms:

<https://www.youtube.com/watch?v=n2y7n6jw5d0>

Laplace = a pillar of automatic control

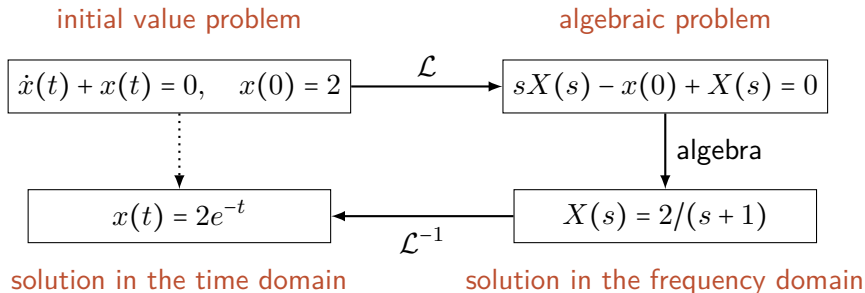
⇒ better to invest time in developing a full understanding of it

Using Laplace transforms with second order systems in free evolution

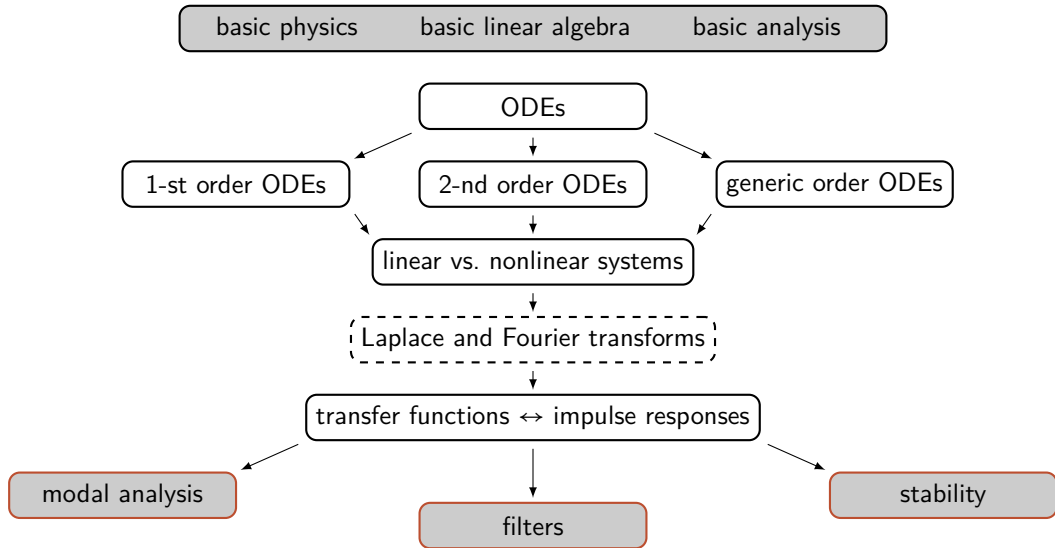
Recall: the “Laplace way” to find that  $\dot{y} = ay \mapsto y(t) = y_0 e^{at}$

$$\dot{y} = ay \quad \rightarrow \quad sY - y_0 = aY \quad \rightarrow \quad Y = y_0 \frac{1}{s - a}$$

Recall: the “Laplace way” to find that  $\dot{y} = ay \mapsto y(t) = y_0 e^{at}$



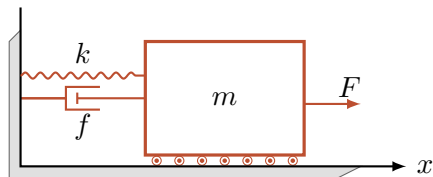
# Summary / Roadmap of TTK4225

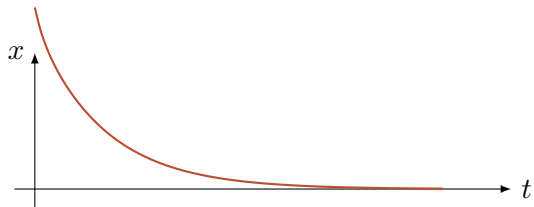


# Roadmap

- the same thing, but for second order system
- different algebraic situations  $\implies$  different damping properties
- different ways of summarizing the different damping properties

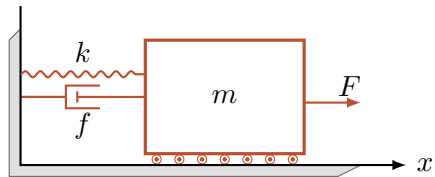
Discussion on a practical example:  
how may a second order system free-evolve?







## Dissecting the practical example “spring-mass systems”



$$m\ddot{x}(t) = -kx(t) - f\dot{x}(t) + u(t)$$

## The properties we shall use

$$\mathcal{L}\{\dot{x}\} = sX(s) - x_0 \quad (14)$$

$$\mathcal{L}\{\ddot{x}\} = s^2X(s) - sx_0 - \dot{x}_0 \quad (15)$$

$$\mathcal{L}^{-1}\left\{\frac{1}{s-a}\right\} = e^{at} \quad (16)$$

$$(17)$$

## Solution for the free evolution case

$$\ddot{x} + \frac{f}{m}\dot{x} + \frac{k}{m}x = 0$$

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we would like to use  $\mathcal{L}^{-1} \left\{ \frac{1}{s-a} \right\} = e^{at}$  to inverse-transform

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we would like to use  $\mathcal{L}^{-1} \left\{ \frac{1}{s-a} \right\} = e^{at}$  to inverse-transform  
 $\Longrightarrow$  we must try to factorize the RHS, i.e., try to write it as

$$X(s) = \frac{A}{s-\alpha} + \frac{B}{s-\beta}$$

(if possible)

## Solution for the free evolution case

Next step: factorize

$$X(s) = \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{s^2 + \frac{f}{m}s + \frac{k}{m}}$$

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$$s^2 + \frac{f}{m}s + \frac{k}{m} = 0 \quad \Longrightarrow \quad s = -\frac{f}{2m} \pm \frac{\sqrt{f^2 - 4mk}}{2m}$$



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*Discussion:* what are the different possibilities?

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*Discussion:* what are the different possibilities?

- ❶  $f^2 - 4mk > 0 \implies$  two distinct real solutions (*will imply overdamping*)
- ❷  $f^2 - 4mk = 0 \implies$  one double real solution (*will imply critical damping*)
- ❸  $f^2 - 4mk < 0 \implies$  two conjugate complex solutions (*will imply underdamping*)

$f^2 - 4mk > 0$ , i.e., overdamping in free evolution

$$X(s) = \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{s^2 + \frac{f}{m}s + \frac{k}{m}} = \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{(s - \lambda_1)(s - \lambda_2)}$$

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$$f^2 > 4mk \quad \Longrightarrow \quad \begin{cases} \lambda_1 = -\frac{f}{2m} - \frac{\sqrt{f^2 - 4mk}}{2m} \in \mathbb{R} \\ \lambda_2 = -\frac{f}{2m} + \frac{\sqrt{f^2 - 4mk}}{2m} \in \mathbb{R} \end{cases}$$

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$$\Longrightarrow \quad X(s) = \frac{A}{s - \lambda_1} + \frac{B}{s - \lambda_2}$$

*Discussion:* can we conclude immediately that

$$X(s) = \frac{A}{s - \lambda_1} + \frac{B}{s - \lambda_2} \implies x(t) = Ae^{\lambda_1 t} + Be^{\lambda_2 t}$$

or do we need to say some property about Laplace transforms that we did not say before?

We need to be sure that Laplace transforms are linear!

i.e., be sure that

$$\mathcal{L}\{\alpha x_1 + \beta x_2\} = \alpha \mathcal{L}\{x_1\} + \beta \mathcal{L}\{x_2\},$$

and that

$$\mathcal{L}^{-1}\{\alpha X_1 + \beta X_2\} = \alpha \mathcal{L}^{-1}\{X_1\} + \beta \mathcal{L}^{-1}\{X_2\}$$

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$$\mathcal{L}^{-1}\{\alpha X_1 + \beta X_2\} = \alpha \mathcal{L}^{-1}\{X_1\} + \beta \mathcal{L}^{-1}\{X_2\}$$

*Discussion:* why do the previous properties hold? (hint: recall that  $\mathcal{L}\{x(t)\} := \int_0^{+\infty} x(t)e^{-st}dt$ )



## In conclusion: overdamping in free evolution

$$f^2 > 4mk \quad \Longrightarrow \quad X(s) = \frac{A}{s - \lambda_1} + \frac{B}{s - \lambda_2} \quad \Longrightarrow \quad x(t) = Ae^{\lambda_1 t} + Be^{\lambda_2 t}$$

for opportune  $A$  and  $B$ .

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for opportune  $A$  and  $B$ . *Discussion:* what determines  $A$  and  $B$ ? Hint: recall that

$$X(s) = \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{s^2 + \frac{f}{m}s + \frac{k}{m}} = \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{(s - \lambda_1)(s - \lambda_2)}$$

important message (will be seen again later on): the denominator of the rational Laplace transform determines the behavior of the system, while the numerator determines which root is dominant and the “amplitude” of each mode

$$X(s) = \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{s^2 + \frac{f}{m}s + \frac{k}{m}} = \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{(s - \lambda_1)(s - \lambda_2)}$$

## Practical example of a overdamped system in free evolution

$$\left\{ \begin{array}{lcl} m & = & 10 \quad \text{kg} \\ f & = & 100 \quad \text{kg/s} \\ k & = & 90 \quad \text{N/s}^2 \end{array} \right. \implies \left\{ \ddot{x} + \frac{100}{10} \dot{x} + \frac{90}{10} x = 0 \right.$$

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*note that there are 2 distinct steps!!*

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ARMA-continuous-time-example.ipynb

?



$f^2 - 4mk = 0$ , i.e., critical damping in free evolution

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*Discussion:* did we learn until now in this course how to inverse-Laplace these things?

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Very interesting behavior:

$$x(t) = Ae^{\lambda t} + Bte^{\lambda t}$$

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Very interesting behavior:

$$x(t) = Ae^{\lambda t} + Bte^{\lambda t}$$

*Discussion:* what contributes to determine  $A$  and  $B$ ?

*Discussion:* recall the previous “overdamping case”, i.e.,

$$\left\{ \begin{array}{lll} X(s) & = & \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{s^2 + \frac{f}{m}s + \frac{k}{m}} \\ m & = & 10 \quad \text{kg} \\ f & = & 100 \quad \text{kg/s} \\ k & = & 90 \quad \text{N/s}^2 \\ x(0) & = & 0.16\text{m} \\ \dot{x}(0) & = & 0\text{m/s} \end{array} \right. \implies X(s) = \frac{0.16(s+10)}{s^2+10s+9}.$$

How may we change the parameters and the initial condition so to have a “critical damping” case, i.e.,

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$$\Longrightarrow \quad X(s) = \frac{A}{s - \lambda_1} + \frac{B}{s - \lambda_2}$$

with  $\lambda_1$  and  $\lambda_2$  a complex-conjugate pair

*Discussion:* did we learn until now in this course how to inverse-Laplace these things?

$$X(s) = \frac{A}{s - \lambda_1} + \frac{B}{s - \lambda_2} \text{ with } \lambda_1 \text{ and } \lambda_2 \text{ a complex-conjugate pair}$$

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How is it possible that a real system produces complex results?”

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“yes, but the corresponding exponentials come out to be complex numbers!

How is it possible that a real system produces complex results?”

Solution:

$$\mathcal{L}\{e^{at} \sin \omega t\} = \frac{\omega}{(s - a)^2 + \omega^2} \quad \mathcal{L}\{e^{at} \cos \omega t\} = \frac{s - a}{(s - a)^2 + \omega^2}$$



*Discussion:* recall the previous “overdamping case”, i.e.,

$$\left\{ \begin{array}{lll} X(s) & = & \frac{x_0 \left( \frac{f}{m} + s \right) + \dot{x}_0}{s^2 + \frac{f}{m}s + \frac{k}{m}} \\ m & = & 10 \quad \text{kg} \\ f & = & 100 \quad \text{kg/s} \\ k & = & 90 \quad \text{N/s}^2 \\ x(0) & = & 0.16\text{m} \\ \dot{x}(0) & = & 0\text{m/s} \end{array} \right. \implies X(s) = \frac{0.16(s+10)}{s^2+10s+9}.$$

How may we change the parameters and the initial condition so to have a “underdamping” case, i.e.,

$$X(s) = \frac{A\omega}{(s-a)^2 + \omega^2} + \frac{B(s-a)}{(s-a)^2 + \omega^2} ?$$

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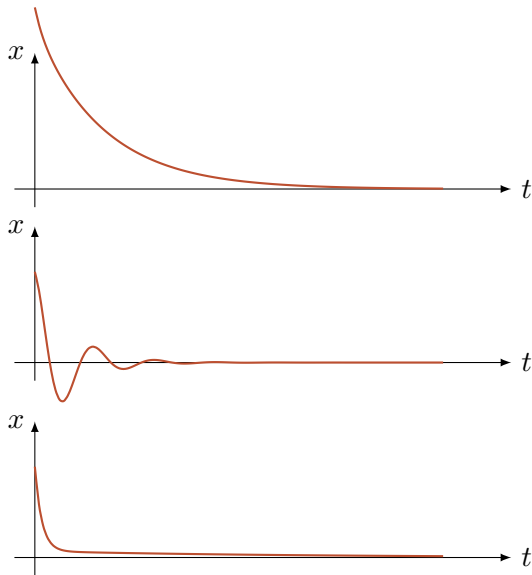
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## Recap of the possible damping modalities



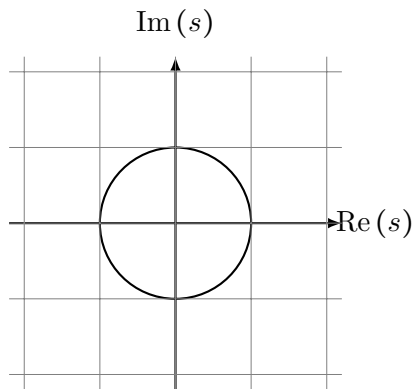
## Recap of the corresponding time-frequency pairs

overdamped:  $\mathcal{L}\{Ae^{\lambda_1 t} + Be^{\lambda_2 t}\} = \frac{A}{s - \lambda_1} + \frac{B}{s - \lambda_2}, \quad \lambda_1 \neq \lambda_2$

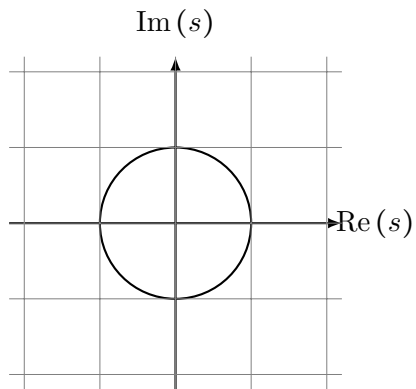
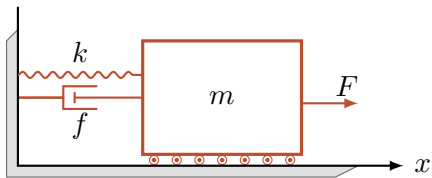
critically damped:  $\mathcal{L}\{Ae^{\lambda t} + Bte^{\lambda t}\} = \frac{A}{s - \lambda} + \frac{B}{(s - \lambda)^2}$

underdamped:  $\mathcal{L}\{e^{at}(A \sin \omega t + B \cos \omega t)\} = \frac{A\omega}{(s - a)^2 + \omega^2} + \frac{B(s - a)}{(s - a)^2 + \omega^2}$

## Recap using the “poles-zeros plot”



Do you see it?



?



The most important Laplace transforms properties for control  
people

# Roadmap

- make Laplace-transform more mathematically ground
- derive some few important examples and properties

# Laplace-properties we saw up to now:

(4.2.1 in Balchen's book)

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## Linearity

$$\mathcal{L}\{af(t) + bg(t)\} = a\mathcal{L}\{f(t)\} + b\mathcal{L}\{g(t)\} \quad (18)$$

# Laplace-properties we saw up to now:

(4.2.1 in Balchen's book)

## Derivatives

$$\mathcal{L}\{\dot{x}\} = sX(s) - x(0) \quad (19)$$

$$\mathcal{L}\{\ddot{x}\} = s^2X(s) - sx(0) - \dot{x}(0) \quad (20)$$

$$\vdots \quad (21)$$

# Laplace-properties we saw up to now:

(4.2.1 in Balchen's book)

## Integration

$$\mathcal{L}\left\{\int_0^t f(\tau)d\tau\right\} = \frac{1}{s}F(s) \quad (22)$$

## Laplace-transforms we needed up to now:

$$f(t) = \mathcal{L}^{-1}\{F(s)\} \qquad F(s) = \mathcal{L}\{f(t)\}$$

---

$$e^{at}$$

$$\frac{1}{s-a}$$

$$te^{at}$$

$$\frac{1}{(s-a)^2}$$

$$e^{at} \sin \omega t$$

$$\frac{\omega}{(s-a)^2 + \omega^2}$$

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## Important (but intuitive) theorem

If  $f(t)$  is defined and piecewise continuous for  $t > 0$ ,  
and there exist two constants  $M$  and  $k$  so that  $|f(t)| \leq Me^{kt}$  for every  $t > 0$ ,  
then  $\mathcal{L}\{f(t)\}$  is well defined for every  $s > k$

## Important (but intuitive) property: linearity

Assuming that  $f(t)$  and  $g(t)$  are so that  $\mathcal{L}\{f(t)\}$  and  $\mathcal{L}\{g(t)\}$  exists, then  $\mathcal{L}\{af(t) + bg(t)\}$  exists and is s.t.

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### Proof

Integrals are linear, thus

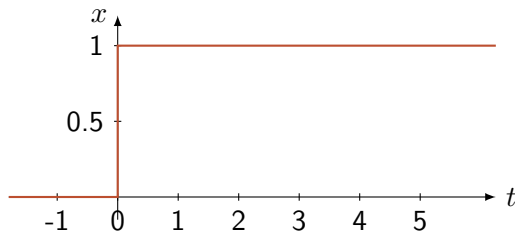
$$\begin{aligned}\mathcal{L}\{af(t) + bg(t)\} &= \int_0^{+\infty} e^{-st} [af(t) + bg(t)] dt \\ &= a \int_0^{+\infty} e^{-st} f(t) dt + b \int_0^{+\infty} e^{-st} g(t) dt \\ &= a\mathcal{L}\{f(t)\} + b\mathcal{L}\{g(t)\}\end{aligned}$$



?

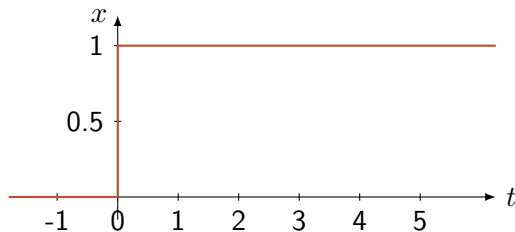
## Important example: Laplace-transforming the Heaviside-function

$$H(t) := \begin{cases} 0 & \text{for } t < 0 \\ 1 & \text{for } t > 0 \end{cases}$$



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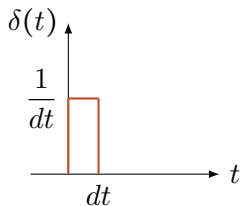
## Deriving the Heaviside function from the Dirac's impulse

$$H(t) = \int_{-\infty}^t \delta(\tau) d\tau$$

with, *very loosely*,

$$\delta(t) := \begin{cases} +\infty & \text{if } t = 0 \\ 0 & \text{otherwise,} \end{cases}$$

or, better, as the limit of concentrating an unitary point mass at 0:



## Back to the important example: Laplace-transforming the Heaviside-function

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Likely the most important example: Laplace-transforming an exponential

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$$\begin{aligned}\mathcal{L}\{t^2\} &= \int_0^{+\infty} t^2 e^{-st} dt & \begin{aligned} u &= t^2 & v &= e^{-st} \\ \dot{u} &= 2t & \dot{v} &= -s e^{-st} \end{aligned} \end{aligned} \quad (34)$$

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(this means that often one gets a ramp response and then one has to reconstruct back the impulse response)

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$$\mathcal{L}\{x^{(n)}(t)\} = s^n \mathcal{L}\{x(t)\} - s^{n-1}x(0) - s^{n-2}\dot{x}(0) - \dots - x^{(n-1)}(0) \quad (50)$$

# Translation theorems

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## Translation in the frequency domain

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## Translation in the time domain

If  $f(t)$  admits  $F(s)$ , then  $f(t - a)\mu(t - a)$  transforms into  $e^{-as}F(s)$

*Discussion:* for which  $s$  do these theorems hold?

Answer: only in the intersection of the ROCs of both “original” and “translated” transforms!

Translation theorems + Laplace-transforming polynomials = **the most useful formula that you may remember**

$$\left\{ \begin{array}{l} \mathcal{L}\{t^n\} = \frac{n!}{s^{n+1}} \\ \mathcal{L}\{e^{at}f(t)\} = F(s-a) \end{array} \right. \implies \mathcal{L}\{t^n e^{at}\} = \frac{n!}{(s-a)^{n+1}} \quad (51)$$

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Laplace-transforms of oscillatory behaviors = **the second most useful formulas that you may remember**

$$\mathcal{L}\{e^{at} \sin \omega t\} = \frac{\omega}{(s-a)^2 + \omega^2}$$

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## Laplace-transforms of oscillatory behaviors = **the second most useful formulas that you may remember**

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Limit cases:  $a = 0$ , so that

$$\mathcal{L}\{\sin \omega t\} = \frac{\omega}{s^2 + \omega^2}$$

$$\mathcal{L}\{\cos \omega t\} = \frac{s}{s^2 + \omega^2}$$

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Be careful on how to use  $s$  in this course!

Typically, in all the automatic control courses the interest is not too much on the value of  $F(s)$  for different  $s$ , but rather on the fact that  $F(s)$  enables performing algebraic operations to solve ODEs!

?

Using Laplace transforms with generic  $n$  order systems in free evolution

In the last episodes: the “Laplace way” to find that  $\dot{y} = ay \mapsto y(t) = y_0 e^{at}$ ,  
and to do something similar for second order systems in free evolution

$$\dot{y} = ay \quad \rightarrow \quad sY - y_0 = aY \quad \rightarrow \quad Y = y_0 \frac{1}{s - a}$$

# Roadmap

- the same thing, but for generic systems



## A small recap

- first order homogeneous systems:

$$\dot{x} - ax = 0 \quad \mapsto \quad x(t) = x_0 e^{at} \quad (52)$$

- first order non-homogeneous systems:

$$\dot{x} - ax = bu \quad \mapsto \quad x(t) = x(0)e^{at} + \int_0^t e^{a(t-\tau)} bu(\tau) d\tau \quad (53)$$

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*The problem:* the convolution integral gives us headaches, and we want to use Laplace to ease things up. To understand how, step-by-step path:

- 1 start from  $\dot{x} - ax = 0$
- 2 move to  $\ddot{x} - a_1\dot{x} - a_2x = 0$
- 3 move to higher order
- 4 add the “ $u$ ” part

## A more detailed roadmap for this unit:

- generalize how to do partial expansions
- solve the problem: how to use Laplace transforms with generic  $n$  order systems in free evolution

?

## Recalling some past results

$$\begin{aligned}\ddot{x} + \frac{f}{m}\dot{x} + \frac{k}{m}x &= 0 \\ \Downarrow \\ X(s) &= \frac{x_0\left(\frac{f}{m} + s\right) + \dot{x}_0}{s^2 + \frac{f}{m}s + \frac{k}{m}} \Downarrow \\ x(t) &= Ae^{\lambda_1 t} + Be^{\lambda_2 t} \text{ (and the other cases)}\end{aligned}$$

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*Discussion:* when is it that  $\ddot{x} + \frac{f}{m}\dot{x} + \frac{k}{m}x = 0$  leads to  $X(s) \neq 0$ ?

what about  $\ddot{x} + \alpha_2 \ddot{x} + \alpha_1 \dot{x} + \alpha_0 x = 0$ ?

what about  $\ddot{x} + \alpha_2 \ddot{x} + \alpha_1 \dot{x} + \alpha_0 x = 0$ ?

and what about higher orders?



# Partial fraction decomposition

## Partial fraction decomposition

case single poles: if  $\frac{N(s)}{(s-a)(s-b)(s-c)\dots}$  is s.t.  $a \neq b \neq c \neq \dots$  then there exist  $A, B, C, \dots$  s.t.

$$\frac{N(s)}{(s-a)(s-b)(s-c)\dots} = \frac{A}{s-a} + \frac{B}{s-b} + \frac{C}{s-c} + \dots \quad (54)$$

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**case repeated poles:** if some poles are repeated, then there exist  $A_1, \dots, A_{na}, B_1, \dots, B_{nb}$ , s.t.

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“But how do I compute  $A, B$ , etc.?”  $\mapsto$

`en.wikipedia.org/wiki/Partial_fraction_decomposition`

(tip: start from `en.wikipedia.org/wiki/Heaviside_cover-up_method`)

?

## Computing the free evolution for generic $n$ order LTI systems

$$\text{start from } x^{(n)} + \alpha_{n-1}x^{(n-1)} + \dots + \alpha_0x = 0 \quad (56)$$

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$$\text{for each } m \text{ apply } \mathcal{L}\{x^{(m)}(t)\} = s^m \mathcal{L}\{x(t)\} - s^{m-1}x(0) - s^{m-2}\dot{x}(0) - \dots - x^{(m-1)}(0) \quad (57)$$

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$$\text{obtain } X(s) = \frac{N(s)}{(s-a)(s-b)(s-c)\dots} \quad (58)$$



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$$\text{do a partial fraction decomposition} \quad (59)$$

# Computing the free evolution for generic $n$ order LTI systems

Two main cases:

$$x^{(n)} + \alpha_{n-1}x^{(n-1)} + \dots + \alpha_0x = 0$$

↓

$$\text{either } X(s) = \frac{A}{s-a} + \frac{B}{s-b} + \frac{C}{s-c} + \dots$$

$$\text{or } X(s) = \frac{A_1}{s-a} + \dots + \frac{A_{na}}{(s-a)^{na}} + \frac{B_1}{s-b} + \dots + \frac{B_{nb}}{(s-b)^{nb}} + \dots$$

# Computing the free evolution for generic $n$ order LTI systems

Case single poles

$$x^{(n)} + \alpha_{n-1}x^{(n-1)} + \dots + \alpha_0x = 0$$

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*Discussion:* given  $\mathcal{L}^{-1}\left\{\frac{1}{s-\alpha}\right\} = e^{\alpha t}$ , what is  $x(t)$ ?

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$$\text{remember: } \mathcal{L}\{t^n e^{at}\} = \frac{n!}{(s-a)^{n+1}}$$

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$$\text{remember: } \mathcal{L}\{t^n e^{at}\} = \frac{n!}{(s-a)^{n+1}}$$

*Discussion:* given the previous transform, what is  $x(t)$ ? Solution: a sum of terms of the type  $t^n e^{at}$

## Extremely important result

a LTI in free evolution behaves as a combination of terms  $e^{\alpha t}$ ,  $te^{\alpha t}$ ,  $t^2e^{\alpha t}$ , etc. for a set of different  $\alpha$ s and powers of  $t$ , called the *modes* of the system

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*Discussion:* assuming that we have two modes,  $e^{-0.3t}$  and  $e^{-1.6t}$ , so that

$$x(t) = \gamma_1 e^{-0.3t} + \gamma_2 e^{-1.6t}.$$

What determines  $\gamma_1$  and  $\gamma_2$ ?

comparison-of-exponentials.ipynb

?

Using Laplace transforms with second order systems in free evolution

In the last episodes: the “Laplace way” to find how LTI systems freely evolve

$$\dot{y} = ay \quad \rightarrow \quad sY - y_0 = aY \quad \rightarrow \quad Y = y_0 \frac{1}{s - a}$$

# Roadmap

- the same thing, but adding the forced response term



## A small recap

- first order homogeneous systems:

$$\dot{x} - ax = 0, x(0) = x_0 \quad \mapsto \quad x(t) = x_0 e^{at} \quad (60)$$

- first order non-homogeneous systems:

$$\dot{x} - ax = bu, x(0) = x_0 \quad \mapsto \quad x(t) = x(0)e^{at} + \int_0^t e^{a\tau} bu(t - \tau) d\tau \quad (61)$$

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*The problem, once again:* the convolution integral gives us headaches - we want to use Laplace to ease things up

## Laplace-ing everything

$$\dot{x} - ax = bu, \quad x(0) = x_0 \quad (62)$$

$$\downarrow \quad (63)$$

$$sX(s) - x(0) - aX(s) = bU(s) \quad (64)$$

$$\downarrow \quad (65)$$

$$X(s) = \frac{x(0)}{s-a} + \frac{bU(s)}{s-a} \quad (66)$$

$$\downarrow \quad (67)$$

$$x(t) = \mathcal{L}^{-1} \left\{ \frac{x(0)}{s-a} + \frac{bU(s)}{s-a} \right\} \quad (68)$$

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if  $u(t)$  has a rational Laplace transform, then

$$x(t) = \sum \mathcal{L}^{-1} \left\{ \frac{\alpha}{(s-a)^\beta} \right\}$$

## Example 1

$$\dot{x} + x = 0, \quad x(0) = 2 \quad \Longrightarrow \quad x(t > 0) = ?$$

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$$sX(s) - x(0) + X(s) = 0$$

↓

$$(s + 1)X(s) = x(0)$$

↓

$$X(s) = \frac{x(0)}{(s + 1)} = \frac{2}{s - (-1)}$$



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*Discussion:* what is the time constant in this case?

## Example 2

*Discussion:*  $\mathcal{L}\{H(t)\} = ?$

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↓

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↓

$$X(s) = \frac{2}{s+1} + \frac{1}{s(1+s)}$$

## Example 2, continued

$$X(s) = \frac{2}{s+1} + \frac{1}{s(1+s)}$$

$$\frac{1}{s(s+1)} = \frac{A}{s} + \frac{B}{s+1}$$

## Example 2, continued

$$X(s) = \frac{2}{s+1} + \frac{1}{s(1+s)}$$

$$\frac{1}{s(s+1)} = \frac{A}{s} + \frac{B}{s+1} = \frac{A(s+1) + Bs}{s(s+1)} \quad (72)$$



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$$x(t) = \mathcal{L}^{-1}\{X(s)\} = e^{-t} + \mu(t) \quad (74)$$

### Example 3

$$\dot{x} + 4x = 3e^{-4t}, \quad x(0) = 1 \quad \Longrightarrow \quad x(t > 0) = ?$$

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$$sX(s) - 1 + 4X(s) = \frac{3}{s+4} \quad (75)$$

$$\downarrow \quad (76)$$

$$X(s) = \frac{1}{s+4} + \frac{3}{(s+4)^2} \quad (77)$$

$$(78)$$

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Remember:  $\mathcal{L}\{t^n e^{at}\} = \frac{n!}{(s-a)^{n+1}}$ , thus

$$x(t) = e^{-4t} + 3te^{-4t} \quad (79)$$

?