

Indefinite Canonical Systems. Theory and Examples

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joint work with Michael Kaltenböck,
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CONTENTS

- Review of the positive definite theory
- Some examples of canonical systems
- Indefinite analogue of canonical systems
- A short account on the literature

REVIEW OF THE POSITIVE DEFINITE THEORY

Hamiltonians and canonical systems

A *Hamiltonian* is a function

- $H : [\sigma_0, \sigma_1) \rightarrow \mathbb{R}^{2 \times 2}$ defined a.e., measurable;
- $H(t) \geq 0$, $H \in L^1_{loc}((\sigma_0, \sigma_1))$;
- $\int_{\sigma_0}^{\sigma_0 + \epsilon} \text{tr } H(t) dt < \infty$ (initial value problem);
- H does not vanish on any set of positive measure.

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The *canonical system* with Hamiltonian H is the differential equation

$$y'(x) = z \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} H(x)y(x), \quad x \in [\sigma_0, \sigma_1).$$

limit circle case vs. limit point case

A Hamiltonian H is said to be in the

- *limit circle case* if $\int_{\sigma_1-\epsilon}^{\sigma_1} \text{tr } H(t) dt < +\infty$;
- *limit point case* if $\int_{\sigma_1-\epsilon}^{\sigma_1} \text{tr } H(t) dt = +\infty$.

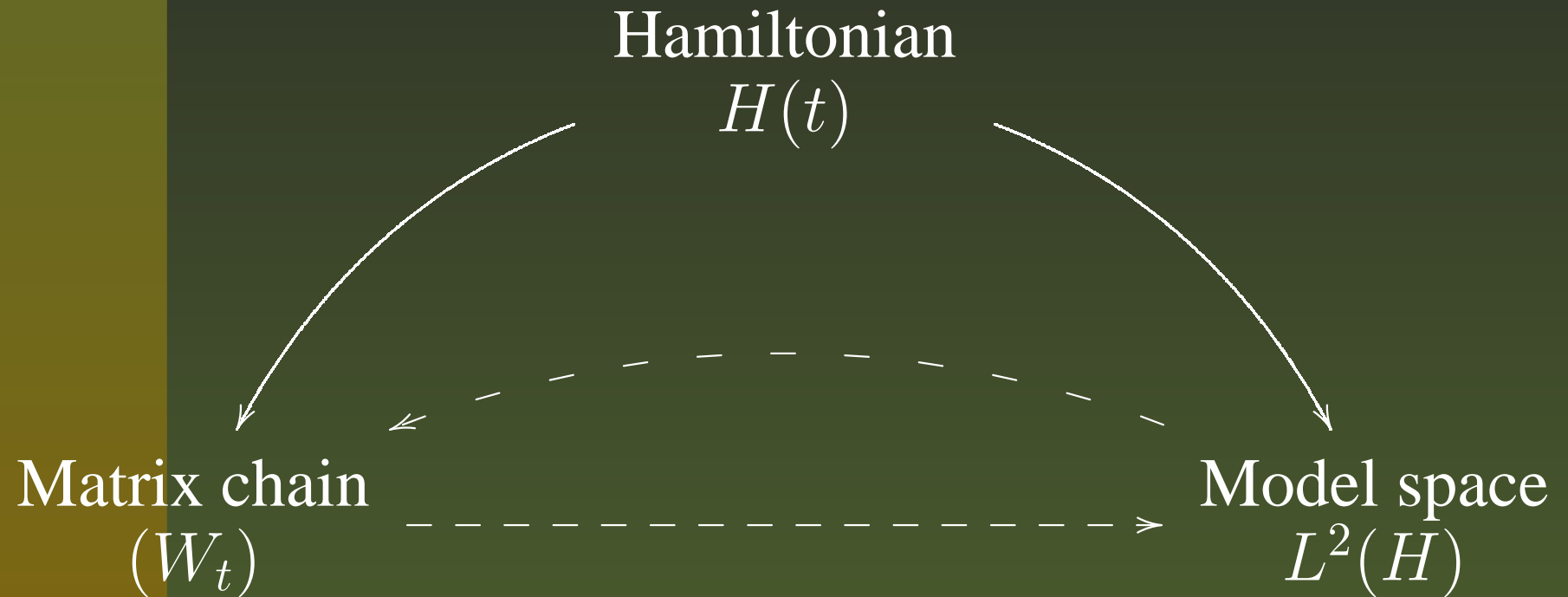
Summary

Hamiltonian
 $H(t)$

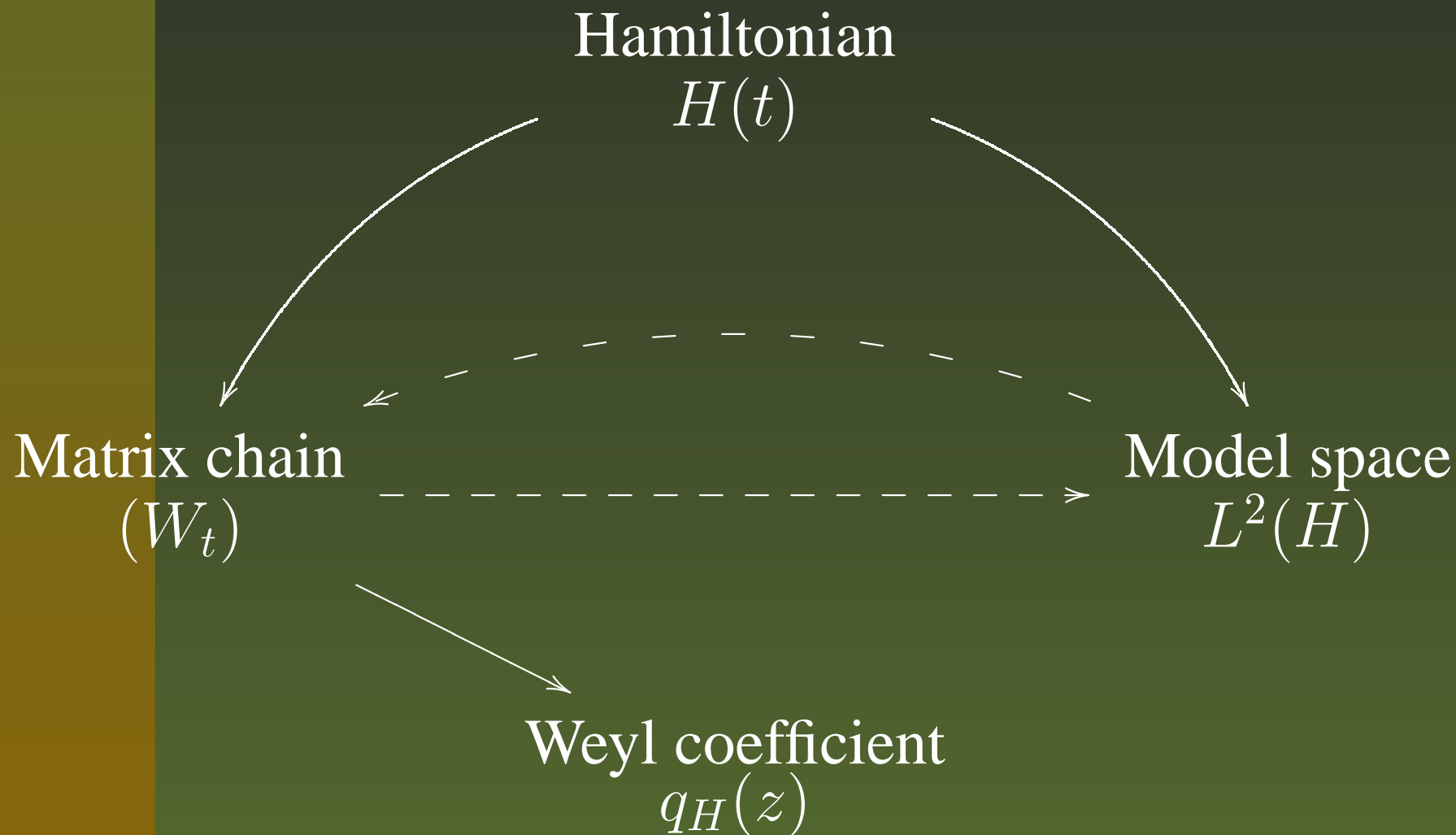
Summary



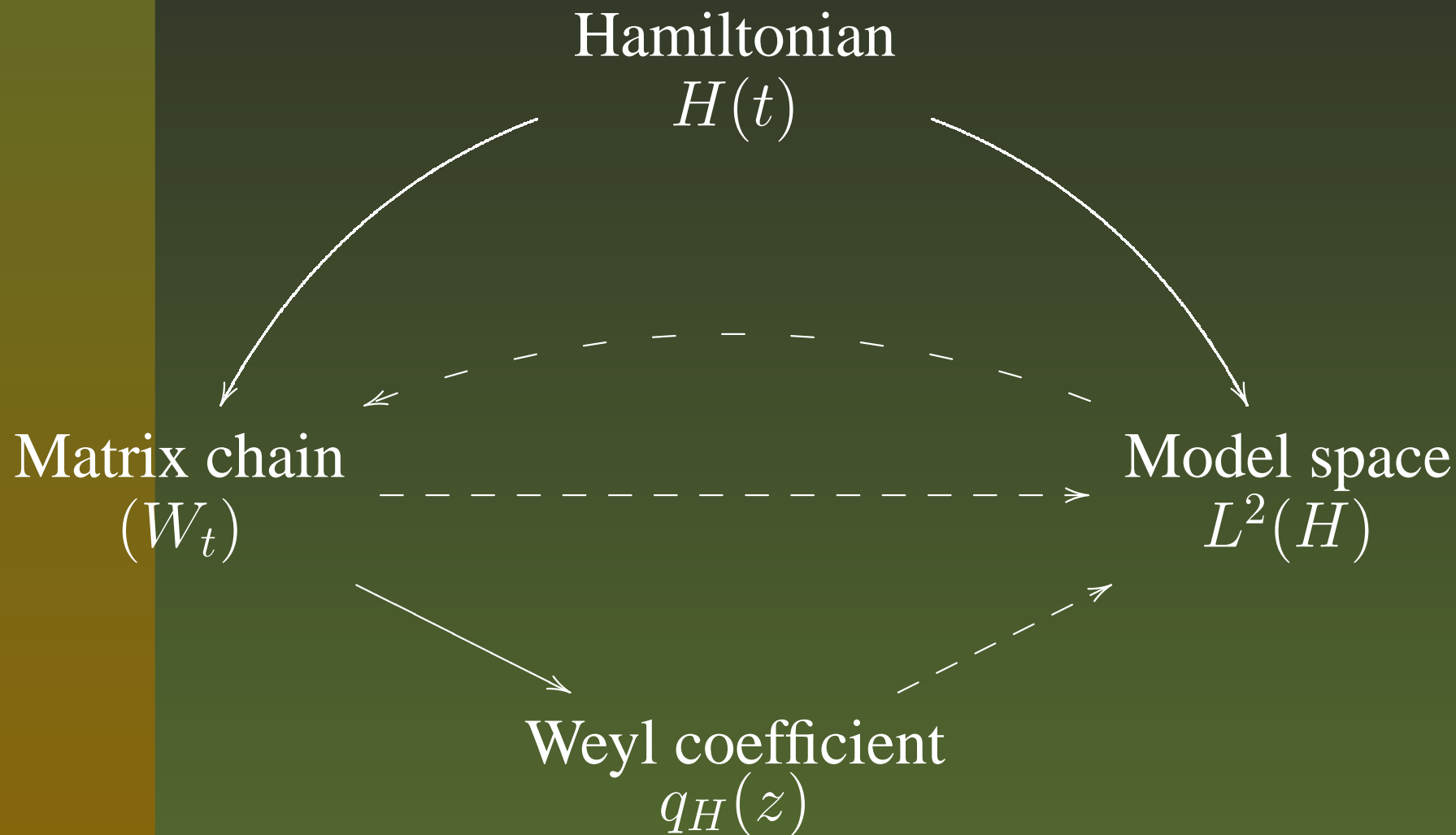
Summary



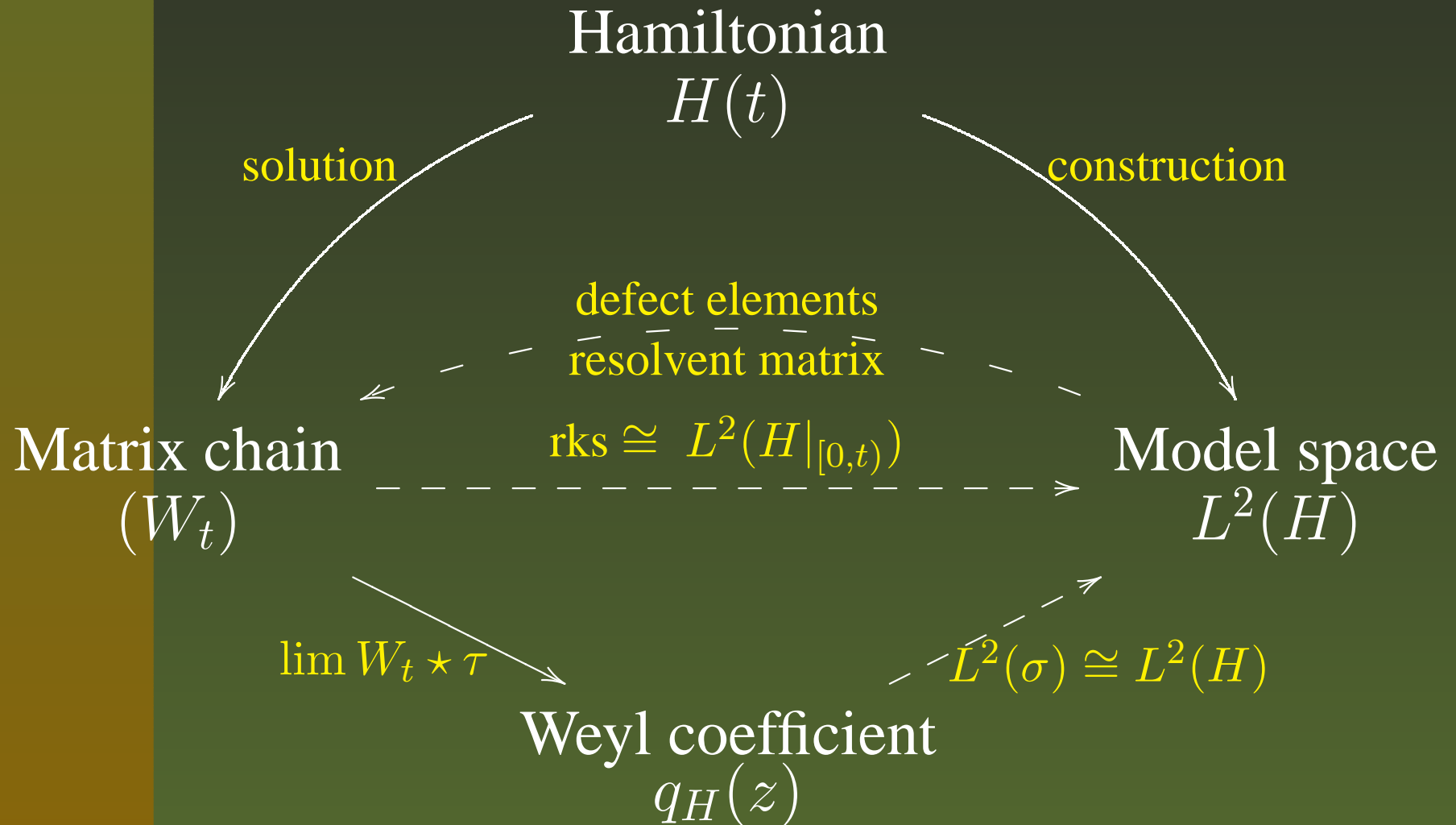
Summary (limit point case)



Summary (limit point case)



Summary (limit point case)



The Inverse Spectral Theorem

The assignment

$$H(t) \mapsto q_H(z)$$

yields a bijection between the set of all Hamiltonians (up to reparameterization) and the Nevanlinna class \mathcal{N}_0 .

SOME EXAMPLES OF CANONICAL SYSTEMS

Positive definite functions

Let $a \in (0, \infty)$. A function $f : (-2a, 2a) \rightarrow \mathbb{C}$ is called positive definite, if $f(-t) = \overline{f(t)}$ and if the kernel

$$K_f(s, t) = f(t - s), \quad s, t \in (-a, a),$$

is positive definite. The set of all continuous positive definite functions on the interval $(-2a, 2a)$ is denoted by $\mathcal{P}_{0,a}$.

Positive definite functions

Continuation problem: Let $f \in \mathcal{P}_{0,a}$. Do there exist continuations $\tilde{f} \in \mathcal{P}_{0,\infty}$?

Positive definite functions

Continuation problem: Let $f \in \mathcal{P}_{0,a}$. Do there exist continuations $\tilde{f} \in \mathcal{P}_{0,\infty}$?

Solution: There exists either a unique continuation or infinitely many continuations. In the second case the set of all continuations is parameterized by

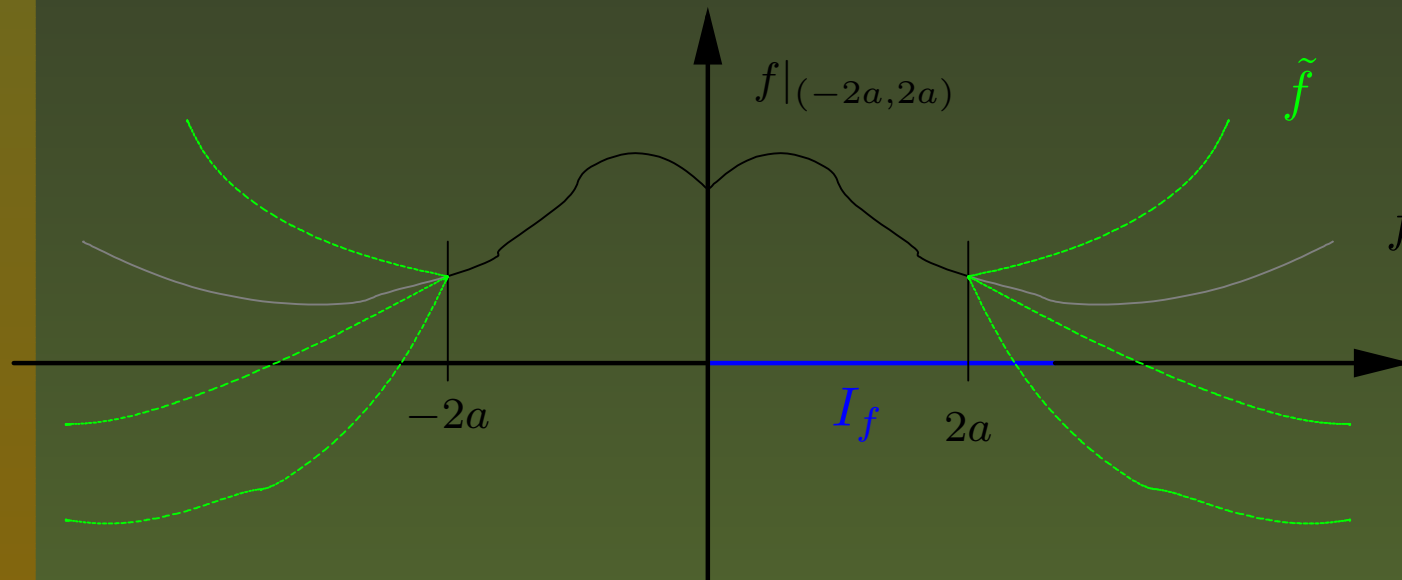
$$i \int_0^\infty e^{izt} \tilde{f}(t) dt = W_f(z) \star \tau(z)$$

where W_f is a certain entire 2×2 -matrix function and the parameter τ runs through the Nevanlinna class \mathcal{N}_0 .

Pos.def. functions & can. systems

Let $f \in \mathcal{P}_{0,\infty}$. Assume that the set

$I_f := \{a > 0 : f|_{(-2a,2a)} \text{ has infinitely many continuations}\}$
is nonempty.



Pos.def. functions & can. systems

Then the family

$$W_t(z) := \begin{cases} \begin{pmatrix} 1 & 0 \\ -(f(0)^{-1} + t)z & 1 \end{pmatrix} & , t \in [-f(0)^{-1}, 0) \\ W_{f|_{(-2t, 2t)}}(z) & , t \in I_f \end{cases}$$

is the matrix chain of a certain Hamiltonian H_f .

The Bessel equation

The Bessel equation is the eigenvalue problem with singular endpoint 0

$$-u''(x) + \frac{\nu^2 - \frac{1}{4}}{x^2}u(x) = \lambda u(x), \quad x > 0$$

Here ν is a parameter $\nu > \frac{1}{2}$ and λ is the eigenvalue parameter.

Bessel equation & can. systems

Rewriting this equation as a first-order-system, making a substitution in the independent variable, and setting $\alpha := 2\nu - 1$, $\lambda = z^2$, yields an equation of the form of a canonical system with

$$H_\alpha(x) = \begin{pmatrix} x^\alpha & 0 \\ 0 & x^{-\alpha} \end{pmatrix}$$

In order that H_α is integrable at 0, we need that $\alpha < 1$, i.e. $\nu < 1$. In this case the matrix chain $(W_{\alpha,t})_{t \in [0, \infty)}$ and the Weyl coefficient q_{H_α} can be computed explicitly:

Bessel equation & can. systems

Let $\alpha \in (0, 1)$. Then

$$W_{\alpha,t}(z) =$$
$$= \begin{pmatrix} 2^{\nu_1} \Gamma(\nu) z^{-\nu_1} t^{-\nu_1} J_{\nu_1}(tz) & 2^{\nu_1} \Gamma(\nu) z^{-\nu_1} t^{\nu} J_{\nu}(tz) \\ -2^{-\nu} \Gamma(-\nu_1) z^{\nu} t^{-\nu_1} J_{-\nu_1}(tz) & 2^{-\nu} \Gamma(-\nu_1) z^{\nu} t^{\nu} J_{-\nu}(tz) \end{pmatrix}$$

with $\nu_1 := \frac{\alpha-1}{2} = \nu - 1$, and

$$q_{H_{\alpha}}(z) = c_{\alpha} z^{-\alpha}$$

with $c_{\alpha} := \frac{2^{\alpha}}{\pi} \Gamma(\nu)^2 \sin \nu e^{i\nu\pi}$.

INDEFINITE ANALOGUE OF CANONICAL SYSTEMS

Indefinite can.systems / Motivation

CONCEPTS	POSITIVE DEFINITE	
	Hamiltonian	
	Matrix chain	
	Boundary triplet	
	Nevanlinna class \mathcal{N}_0	

Indefinite can.systems / Motivation

	POSITIVE DEFINITE	
CONCEPTS	Hamiltonian	
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I.S.T.	$\{\text{Hamiltonians}\} \overset{\text{WC}}{\longleftrightarrow} \mathcal{N}_0$	

Indefinite can.systems / Motivation

POSITIVE DEFINITE		
CONCEPTS	Hamiltonian	
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	Nevanlinna class \mathcal{N}_0	
I.S.T.	$\{\text{Hamiltonians}\} \longleftrightarrow \mathcal{N}_0$	
EXAMPLES	Positive def.functions	
	Bessel equation $\alpha < 1$	

Indefinite can.systems / Motivation

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EXAMPLES	Positive def.functions Bessel equation $\alpha < 1$ Moment problems Strings	

Indefinite can.systems / Motivation

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Indefinite can.systems / Motivation

	POSITIVE DEFINITE	INDEFINITE
CONCEPTS	Hamiltonian Matrix chain Boundary triplet Nevanlinna class \mathcal{N}_0	generalized Nev.class \mathcal{N}_κ
I.S.T.	$\{\text{Hamiltonians}\} \longleftrightarrow \mathcal{N}_0$	
EXAMPLES	Positive def.functions Bessel equation $\alpha < 1$ Moment problems Strings	

Indefinite can.systems / Motivation

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CONCEPTS	Hamiltonian Matrix chain Boundary triplet Nevanlinna class \mathcal{N}_0	B.triplet in Pontryagin space generalized Nev.class \mathcal{N}_κ
I.S.T.	$\{\text{Hamiltonians}\} \longleftrightarrow \mathcal{N}_0$	
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I.S.T.	$\{\text{Hamiltonians}\} \longleftrightarrow \mathcal{N}_0$	
EXAMPLES	Positive def.functions Bessel equation $\alpha < 1$ Moment problems Strings	Hermitian indef.functions

Indefinite can.systems / Motivation

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I.S.T.	$\{\text{Hamiltonians}\} \longleftrightarrow \mathcal{N}_0$	
EXAMPLES	Positive def.functions Bessel equation $\alpha < 1$ Moment problems Strings	Hermitian indef.functions Bessel equation $\alpha \in \mathbb{R}^+$

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Indefinite can.systems / Motivation

	POSITIVE DEFINITE	INDEFINITE
CONCEPTS	Hamiltonian Matrix chain Boundary triplet Nevanlinna class \mathcal{N}_0	general Hamiltonian maximal chain of matrices B.triplet in Pontryagin space generalized Nev.class \mathcal{N}_κ
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EXAMPLES	Positive def.functions Bessel equation $\alpha < 1$ Moment problems Strings	Hermitian indef.functions ✓ Bessel equation $\alpha \in \mathbb{R}^+$ ✓ indefinite moment problems ✓ generalized strings ✓

General Hamiltonians

A *general Hamiltonian* consists of the data

$$\sigma_0, \dots, \sigma_{n+1} \in \mathbb{R} \cup \{\pm\infty\}, \quad \sigma_0 < \sigma_1 < \dots < \sigma_{n+1},$$

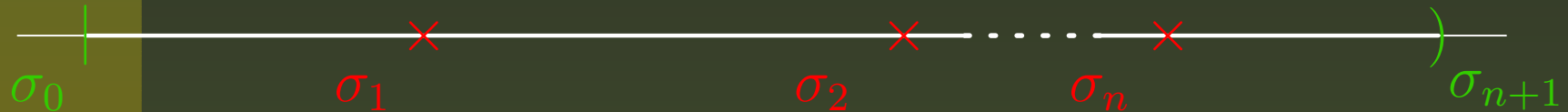
$$H_i : (\sigma_i, \sigma_{i+1}) \rightarrow \mathbb{R}^{2 \times 2}, \quad i = 0, \dots, n,$$

$$E \subseteq \bigcup_{i=0}^n (\sigma_i, \sigma_{i+1}) \cup \{\sigma_0, \sigma_{n+1}\} \text{ finite}$$

$$d_{i,0}, \dots, d_{i,2\Delta_i-1} \in \mathbb{R}, \quad \ddot{o}_i \in \mathbb{N}_0, b_{i,1}, \dots, b_{i,\ddot{o}_i+1} \in \mathbb{R}$$

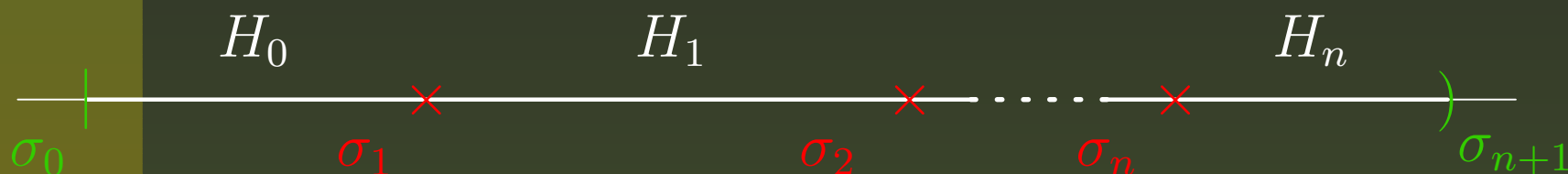
subject to certain conditions.

General Hamiltonians



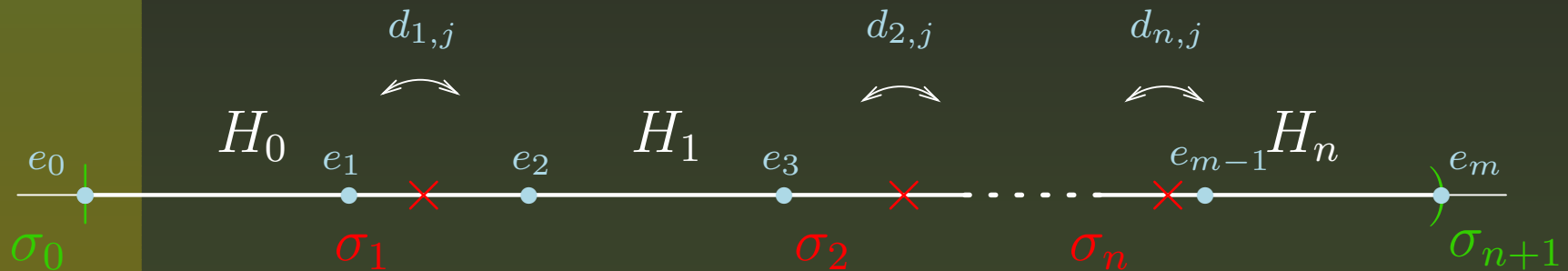
- σ_0 = starting point
- σ_{n+1} = endpoint
- $\sigma_1, \dots, \sigma_n$ = singularities

General Hamiltonians



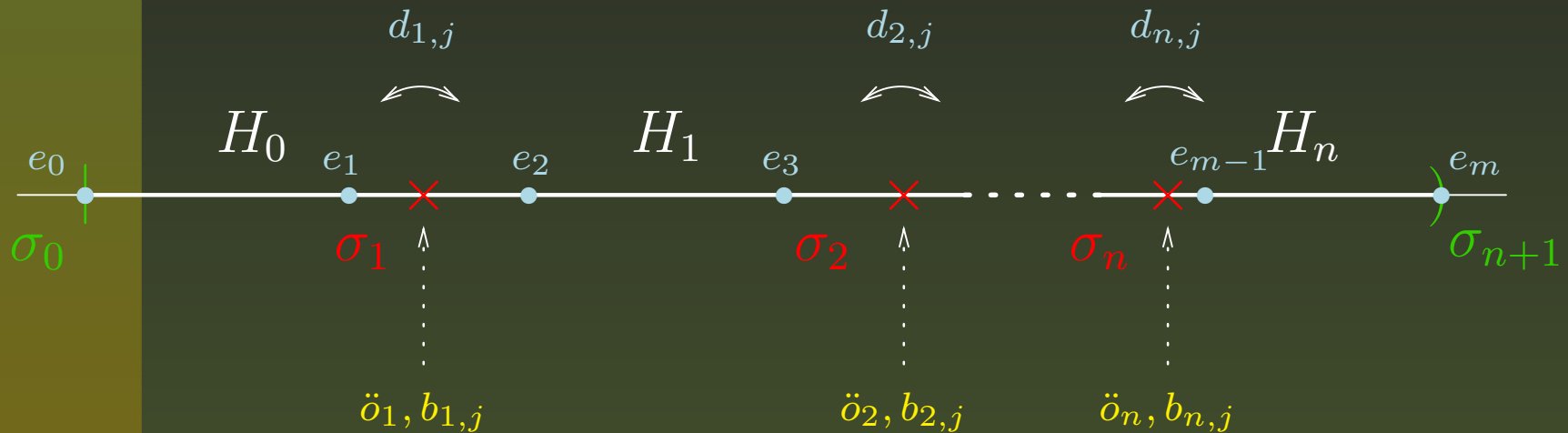
- $H_0, \dots, H_n =$ Hamiltonians, not integrable at $\sigma_1, \dots, \sigma_n$ ($\sigma_1, \dots, \sigma_n$ singularities)
- H_0 integrable at 0 (initial value problem)
- H_n integrable/not at σ_{n+1} (limit circle/point case)
- growth of H_i towards singularity is restricted

General Hamiltonians



- $d_{i,j}$ = interface conditions at a singularity
- E quantitative measurement of ‘local at a singularity’

General Hamiltonians



- $\ddot{o}_i, b_{i,j}$ = contribution concentrated in the singularity

Maximal chains of matrices

Axiomatization of ‘fundamental solution’

- $W_0 = I$ and $W_t \in \bigcup_{\kappa \geq 0} \mathcal{M}_\kappa$ for
 $t \in [\sigma_0, \sigma_1) \cup (\sigma_1, \sigma_2) \cup \dots \cup (\sigma_n, \sigma_{n+1})$
- $W_s^{-1}W_t \in \bigcup_{\kappa \geq 0} \mathcal{M}_\kappa$ and
$$\text{ind}_- W_t = \text{ind}_- W_s + \text{ind}_- W_s^{-1}W_t$$
- If $W \in \bigcup_{\kappa \geq 0} \mathcal{M}_\kappa$, $W^{-1}W_t \in \bigcup_{\kappa \geq 0} \mathcal{M}_\kappa$ and
$$\text{ind}_- W_t = \text{ind}_- W + \text{ind}_- W^{-1}W_t$$

then $W = W_t$ for some t
- some technical conditions

Theory of indefinite can.systems

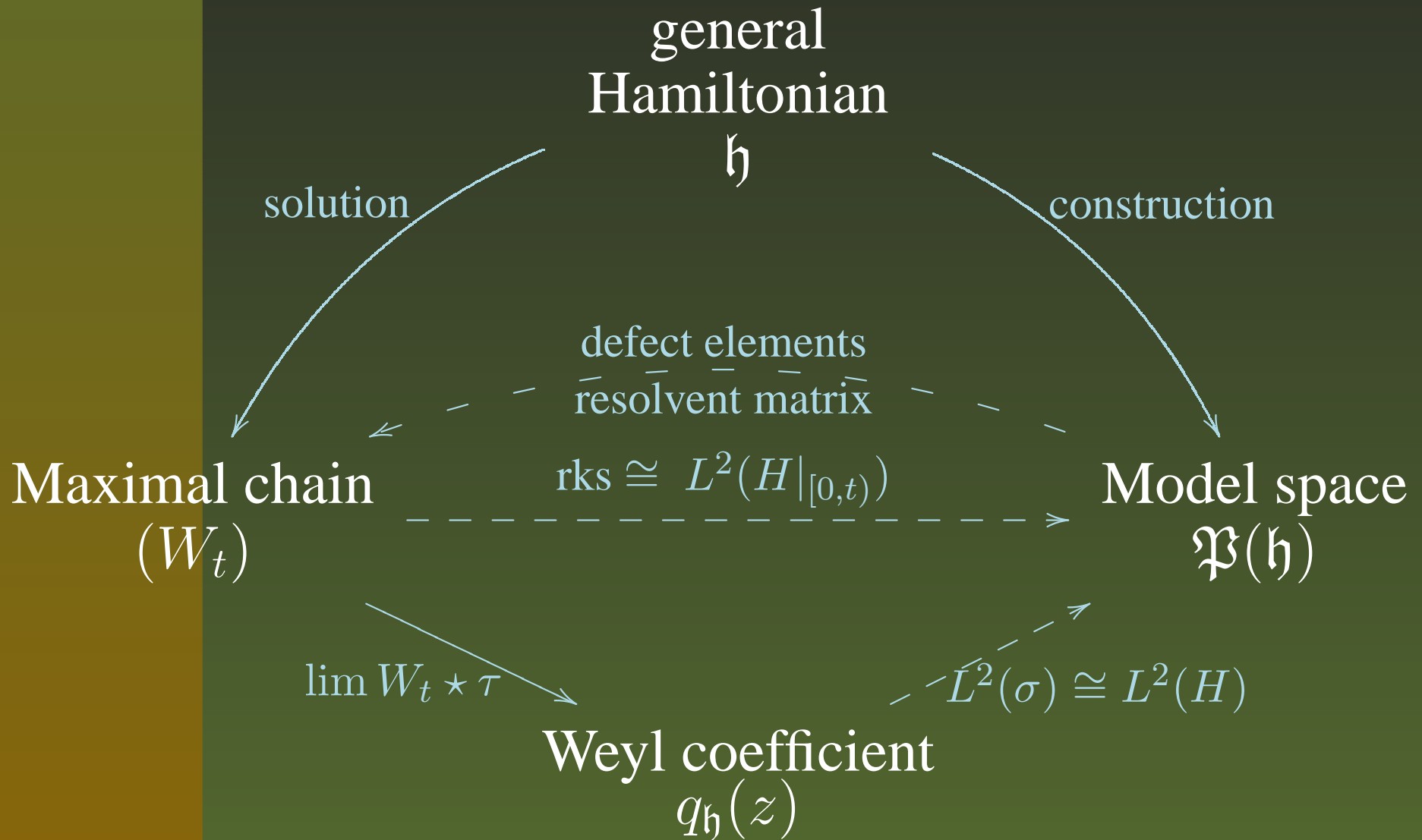
general
Hamiltonian
 \mathfrak{h}

Maximal chain
 (W_t)

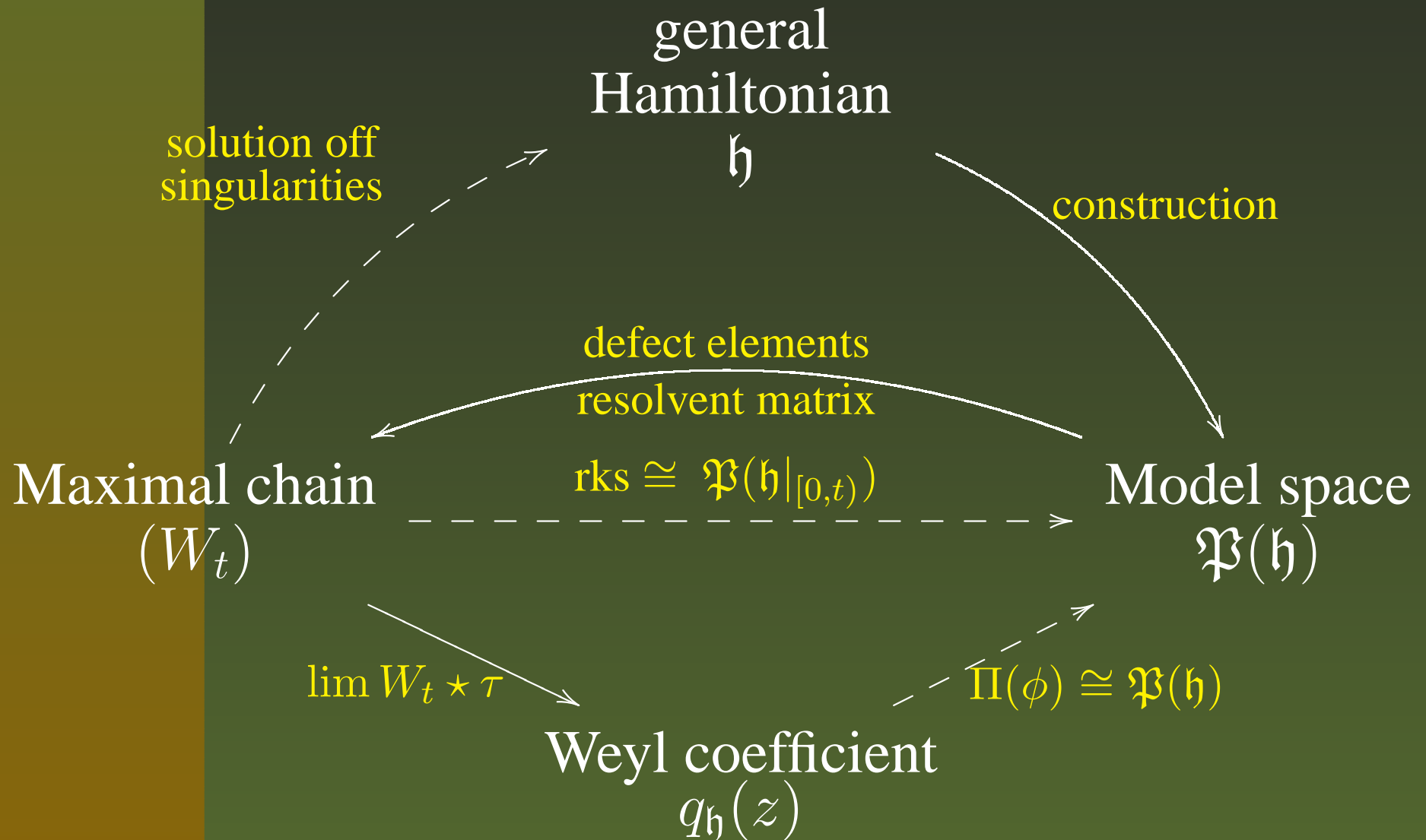
Model space
 $\mathfrak{P}(\mathfrak{h})$

Weyl coefficient
 $q_{\mathfrak{h}}(z)$

Theory of indefinite can.systems



Theory of indefinite can.systems



The Inverse Spectral Theorem

The assignment

$$\mathfrak{h} \mapsto q_{\mathfrak{h}}(z)$$

yields a bijection between the set of all general Hamiltonians (up to reparameterization) and $\bigcup_{\kappa \geq 0} \mathcal{N}_{\kappa}$.

Fitting our examples

In our examples we had obtained generalized Nevanlinna function which seemed to be the ‘Weyl coefficient of the underlying indefinite canonical system’, namely:

- $q_f(z) = \frac{i}{z^2} - \frac{1}{z} \in \mathcal{N}_1$ from the hermitian indefinite function $f(t) = 1 - |t|$.
- $q_\alpha(z) = c_\alpha z^{-\alpha} \in \mathcal{N}_{\kappa(\alpha)}$ where $\kappa(\alpha) := \left[\frac{\alpha+1}{2} \right]$, from the Bessel equation with parameter $\alpha \in \mathbb{R}^+ \setminus \{1, 3, 5, \dots\}$.

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- $q_\alpha(z) = c_\alpha z^{-\alpha} \in \mathcal{N}_{\kappa(\alpha)}$ where $\kappa(\alpha) := \left[\frac{\alpha+1}{2} \right]$, from the Bessel equation with parameter $\alpha \in \mathbb{R}^+ \setminus \{1, 3, 5, \dots\}$.

In order to fit these examples, we have to find the general Hamiltonian whose Weyl coefficient is q_f or q_α , and see how it is related to the ‘Hamiltonians’ H_f and H_α .

A SHORT ACCOUNT ON THE LITERATURE

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THE END

The matrix chain (W_t)

Let H be a Hamiltonian defined on $[\sigma_0, \sigma_1)$. Then W_t , $t \in [\sigma_0, \sigma_1)$, denotes the unique solution of the initial value problem

$$\frac{d}{dt}W_t(z) \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} = zW_t(z)H(t), \quad x \in [\sigma_0, \sigma_1),$$

$$W_0(z) = I.$$

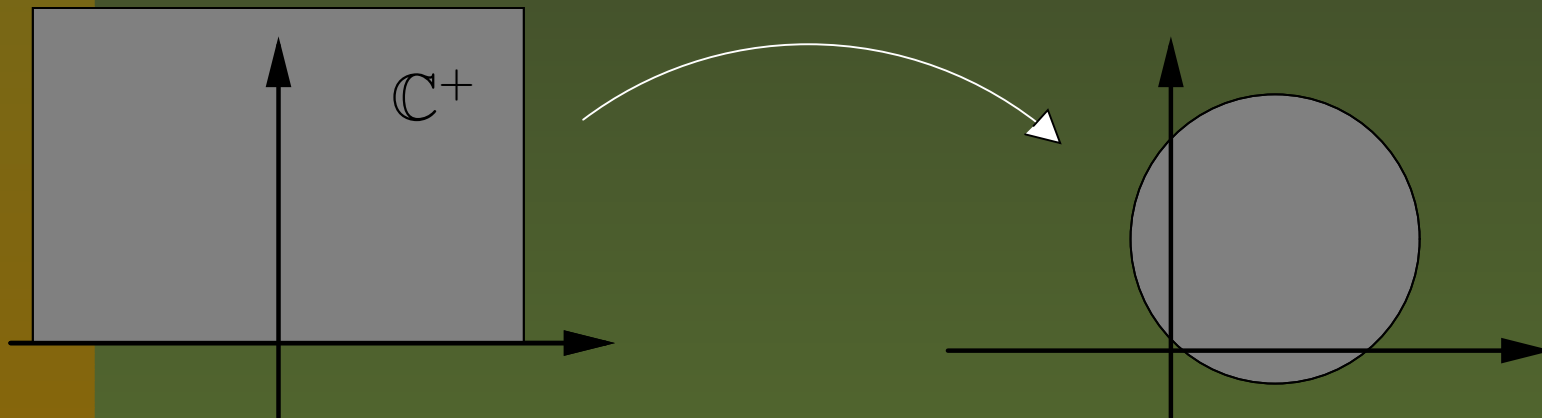


The Weyl coefficient $q_H(z)$

For $W = (w_{ij})_{i,j=1}^2 \in \mathbb{C}^{2 \times 2}$ and $\tau \in \mathbb{C}$ denote

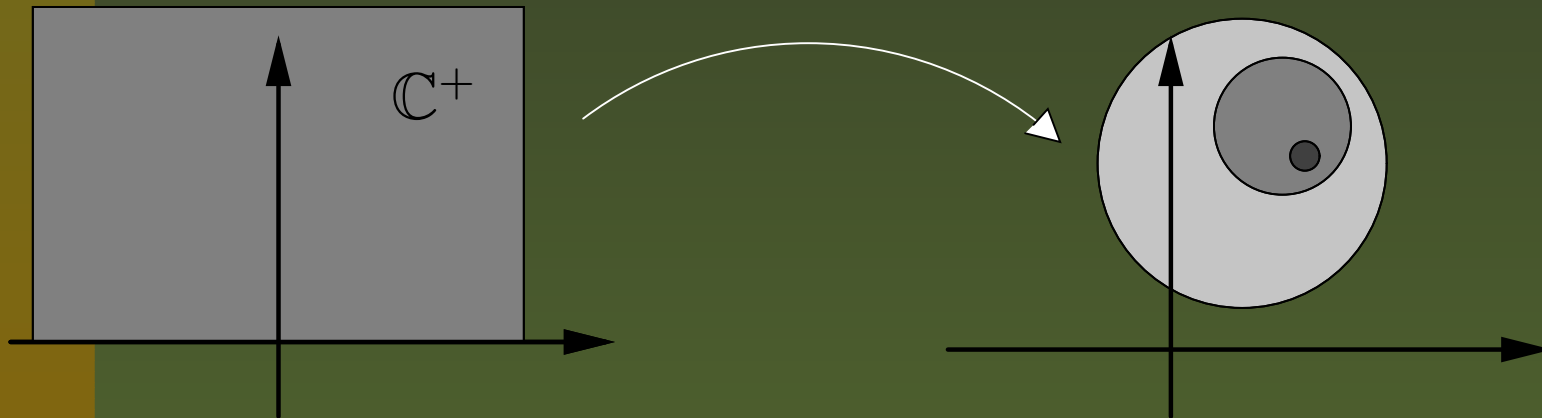
$$W \star \tau := \frac{w_{11}\tau + w_{12}}{w_{21}\tau + w_{22}}$$

The assignment $\tau \mapsto W \star \tau$ maps the upper half plane to some (general) disk:



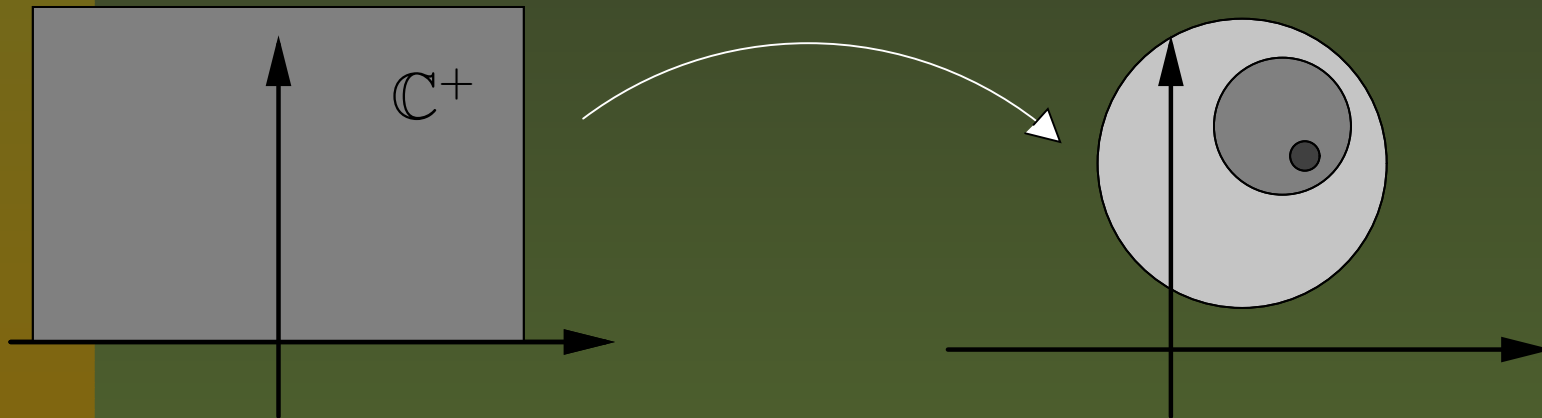
The Weyl coefficient $q_H(z)$

Let $(W_t)_{t \in [\sigma_0, \sigma_1)}$ be the matrix chain associated with the Hamiltonian H . The assignments $\tau \mapsto W_t \star \tau$ map \mathbb{C}^+ to a nested sequence of disks contained in \mathbb{C}^+ . The disk $W_t \star \mathbb{C}^+$ is contained in the upper half plane and its radius is $[\int_{\sigma_0}^t \text{tr } H(x) dx]^{-1}$.



The Weyl coefficient $q_H(z)$

Let $(W_t)_{t \in [\sigma_0, \sigma_1)}$ be the matrix chain associated with the Hamiltonian H . The assignments $\tau \mapsto W_t \star \tau$ map \mathbb{C}^+ to a nested sequence of disks contained in \mathbb{C}^+ . The disk $W_t \star \mathbb{C}^+$ is contained in the upper half plane and its radius is $[\int_{\sigma_0}^t \text{tr } H(x) dx]^{-1}$.



Thus the limit $q_H(z) := \lim_{t \nearrow \sigma_1} W_t(z) \star \tau$ exists, does not depend on $\tau \in \mathbb{C}^+$, and belongs to \mathcal{N}_0 .



The model space $L^2(H)$

Supressing some technicalities which arise from ‘indivisible intervals’, we have

$$L^2(H) := \left\{ f : (\sigma_0, \sigma_1) \rightarrow \mathbb{C}^2 : \int_{\sigma_0}^{\sigma_1} f(t)^T H(t) f(t) dt < \infty \right\}$$

$$T_{max}(H) := \left\{ (f; g) \in L^2(H)^2 : f \text{ absolutely continuous,} \right. \\ \left. f(t) = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} H(t) g(t), \text{ a.e.} \right\}$$

$$\Gamma(H)(f; g) := f(\sigma_0), (f; g) \in T_{max}(H)$$



W_t from defect elements

Let $y_1(z, x) = (y_1(z, x)_2, y_1(z, x)_2)^T$ and $y_2(z, x) = (y_2(z, x)_2, y_2(z, x)_2)^T$ be the elements of $\ker(T_{max}(H|_{(\sigma_0, t)} - z))$, such that $y_1(z, \sigma_0) = (1, 0)^T$ and $y_2(z, \sigma_0) = (0, 1)^T$. Then

$$W_t(z) = \begin{pmatrix} y_1(z, t)_1 & y_1(z, t)_2 \\ y_2(z, t)_1 & y_2(z, t)_2 \end{pmatrix}$$



W_t as resolvent matrix

Consider

$$S_1 := \left\{ (x; y) \in T_{\max}(H|_{(\sigma_0, t)}) : \right. \\ \left. \pi_{l,1}\Gamma(H|_{(\sigma_0, t)})(x; y) = 0, \pi_r\Gamma(H|_{(\sigma_0, t)})(x; y) = 0 \right\}$$

$$u : (x; y) \mapsto \pi_{l,2}\Gamma(H|_{(\sigma_0, t)})(x; y), \quad (x; y) \in T_{\max}(H|_{(\sigma_0, t)})$$

Then S_1 is symmetric with defect 1 and $u|_{S_1^*}$ is continuous. The matrix function W_t is a u -resolvent matrix of S_1 .



The reproducing kernel space of W_t

The kernel

$$K_{W_t}(w, z) := \frac{W_t(z) \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} W_t(w)^* - \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}}{z - \overline{w}}$$

is positive definite, thus generates a reproducing kernel Hilbert space $\mathfrak{K}(W_t)$. The elements of $\mathfrak{K}(W_t)$ are entire 2-vector-functions.

The operator $\mathcal{S}(W_t)$ of multiplication by z is a symmetry with defect 2. The map $\Gamma(W_t) : f \mapsto f(0)$ is a boundary map for $\mathcal{S}(W_t)$.

The reproducing kernel space of W_t

The boundary triplet

$\langle L^2(H|_{(\sigma_0,t)}), T_{\min}(H|_{(\sigma_0,t)}), \Gamma(H|_{(\sigma_0,t)}) \rangle$ is isomorphic to $\langle \mathfrak{K}(W_t), \mathcal{S}(W_t), \Gamma(W_t) \rangle$. The isomorphism of $L^2(H|_{(\sigma_0,t)})$ to $\mathfrak{K}(W_t)$ is given by

$$f(x) \mapsto \int_{\sigma_0}^t W_x(z) H(x) f(x) dx .$$



The Fourier transform

The Weyl coefficient admits an integral representation of the form

$$q_H(z) = a + bz + \int_{\mathbb{R}} \left(\frac{1}{t - z} - \frac{t}{1 + t^2} \right) d\sigma(t) .$$

Assuming $b = 0$, the map

$$f(x) \mapsto \int_{\sigma_0}^{\sigma_1} (0, 1) W_x(z) H(x) f(x) dx$$

is an isomorphism of $L^2(H)$ onto $L^2(\sigma)$.



The class \mathcal{N}_0

Denote by \mathcal{N}_0 the set of all functions τ , which are analytic in $\mathbb{C} \setminus \mathbb{R}$, satisfy $\tau(\bar{z}) = \overline{\tau(z)}$, and are such that the Nevanlinna kernel

$$Q_\tau(w, z) := \frac{\tau(z) - \overline{\tau(w)}}{z - \bar{w}}$$

is nonnegative definite. This means that each of the quadratic forms

$$q_\tau(\xi_1, \dots, \xi_m) := \sum_{i,j=1}^m Q_f(z_j, z_i) \xi_i \overline{\xi_j}$$

is nonnegative definite.

The class \mathcal{N}_0

A more classical approach to this class of functions is given by the following result:

A function τ which is analytic in $\mathbb{C} \setminus \mathbb{R}$ and satisfies $\tau(\bar{z}) = \overline{\tau(z)}$ belongs to the class \mathcal{N}_0 , if and only if it maps the open upper half plane into the closed upper half plane.



Hermitian indefinite functions

Let $a \in (0, \infty)$. A function $f : (-2a, 2a) \rightarrow \mathbb{C}$ is called hermitian indefinite, if $f(-t) = \overline{f(t)}$ and if the kernel

$$K_f(s, t) = f(t - s), \quad s, t \in (-a, a),$$

has a finite number of negative squares. The set of all continuous hermitian indefinite functions with κ negative squares on the interval $(-2a, 2a)$ is denoted by $\mathcal{P}_{\kappa, a}$.

Hermitian indefinite functions

Continuation problem: Let $f \in \mathcal{P}_{\kappa_0, a}$. Do there exist continuations $\tilde{f} \in \mathcal{P}_{\kappa, \infty}$?

Clearly, for the existence of a continuation $\tilde{f} \in \mathcal{P}_{\kappa, \infty}$, it is necessary that $\kappa \geq \kappa_0$.

Hermitian indefinite functions

Solution: There exists a number $\Delta(f) \in \mathbb{N} \cup \{0, \infty\}$:

- If $\Delta(f) = 0$, then f has infinitely many continuations in each of the classes $\mathcal{P}_{\kappa, \infty}$, $\kappa \geq \kappa_0$.
- If $0 < \Delta(f) < \infty$, then f has a unique continuation in $\mathcal{P}_{\kappa_0, \infty}$, no continuations in $\mathcal{P}_{\kappa, \infty}$ with $\kappa_0 < \kappa < \kappa_0 + \Delta(f)$, and infinitely many continuations in each of the classes $\mathcal{P}_{\kappa, \infty}$, $\kappa \geq \kappa_0 + \Delta(f)$.
- If $\Delta(f) = \infty$, then f has a unique continuation in $\mathcal{P}_{\kappa_0, \infty}$, and no continuations in any of the classes $\mathcal{P}_{\kappa, \infty}$, $\kappa > \kappa_0$.

Hermitian indefinite functions

Assume that $\Delta(f) < \infty$. Then there exists an entire 2×2 -matrix function W_f such that the formula

$$i \int_0^\infty e^{izt} \tilde{f}(t) dt = W_f(z) \star \tau(z)$$

parameterizes the continuations of f in $\bigcup_{\kappa \geq \kappa_0} \mathcal{P}_{\kappa, \infty}$.
Thereby continuations $\tilde{f} \in \mathcal{P}_{\kappa, \infty}$ correspond to parameters τ in the class $\mathcal{K}_{\kappa - \kappa_0}^{\Delta(f)}$. If $\Delta(f) > 0$, the unique solution in $\mathcal{P}_{\kappa_0, \infty}$ is given by the parameter $\tau = \infty$.

Her. indef. functions & can. systems

An example: The function $f(t) := 1 - |t|$ belongs to $\mathcal{P}_{1,\infty}$. Again consider the restrictions $f|_{(-2t,2t)}$. Then

$$\Delta(f|_{(-2t,2t)}) = \begin{cases} 0 & , 0 < t < 1 \text{ or } t > 1 \\ 1 & , t = 1 \end{cases}$$

$f|_{(-2t,2t)} \in \mathcal{P}_{0,t}$, $0 < t < 1$, $f|_{(-2t,2t)} \in \mathcal{P}_{1,t}$, $t > 1$, and

$$\begin{aligned} W_{f|_{(-2t,2t)}}(z) &= \\ &= \begin{pmatrix} \frac{\sin tz - z \cos tz}{(t-1)z} & \left(\frac{1}{z^2} - (t-1)\right) \sin tz - \frac{t \cos tz}{z} \\ \frac{z \cos tz}{t-1} & (t-1)z \sin tz + \cos tz \end{pmatrix} \end{aligned}$$

Her. indef. functions & can. systems

The family

$$W_t(z) := \begin{cases} \begin{pmatrix} 1 & 0 \\ -(1+t)z & 1 \end{pmatrix} & , t \in [-1, 0] \\ W_{f|_{(-2t, 2t)}}(z) & , t \in (0, 1) \cup (1, \infty) \end{cases}$$

satisfies a differential equation of the form of a canonical system with

$$H_f(t) = \begin{cases} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} & , t \in (-1, 0) \\ \begin{pmatrix} (t-1)^2 & 0 \\ 0 & \frac{1}{(t-1)^2} \end{pmatrix} & , t \in (0, 1) \cup (1, \infty) \end{cases}$$

Her. indef. functions & can. systems

The function H_f is locally integrable on $[-1, 1) \cup (1, \infty)$, but NOT at the point 1. Moreover, $\int_T^\infty \text{tr } H_f(x) dx = +\infty$ for $T > 1$, i.e. the ‘limit point case’ prevails at infinity.

If we formally carry out the construction of the Weyl coefficient, we obtain

$$q_{H_f}(z) = \frac{i}{z^2} - \frac{1}{z}, \quad z \in \mathbb{C}^+$$

This function belongs to \mathcal{N}_1 .



Bessel equation ($\alpha \geq 1$)

The function $H_\alpha(t)$ and the matrices $W_{\alpha,t}(z)$ are well-defined for $\alpha \in \mathbb{R}^+ \setminus \{1, 3, 5, \dots\}$. Moreover, $W_{\alpha,t}$ satisfies a differential equation of the form of a canonical system with H_α . The function H_α is locally integrable on $(0, \infty)$, but NOT at the point 0. Moreover, $\int_T^\infty \text{tr } H_\alpha(x) dx = +\infty$ for $T > 0$, i.e. the ‘limit point case’ prevails at infinity.

If we formally carry out the construction of the Weyl coefficient, we obtain

$$q_{H_\alpha}(z) = c_\alpha z^{-\alpha}, \quad z \in \mathbb{C}^+$$

This function belongs to $\mathcal{N}_{\kappa(\alpha)}$ with $\kappa(\alpha) := \left\lceil \frac{\alpha+1}{2} \right\rceil$.



The class \mathcal{K}_ν^Δ

For $\nu, \Delta \in \mathbb{N}_0$, denote by \mathcal{K}_ν^Δ the set of all functions τ , which are meromorphic in $\mathbb{C} \setminus \mathbb{R}$, satisfy $\tau(\bar{z}) = \overline{\tau(z)}$, and are such that the maximal number of negative squares of quadratic forms

$$\begin{aligned} q_\tau(\xi_1, \dots, \xi_m; \eta_0, \dots, \eta_{\Delta-1}) &:= \\ &= \sum_{i,j=1}^m \frac{\tau(z_i) - \overline{\tau(z_j)}}{z_i - \bar{z}_j} \xi_i \bar{\xi}_j + \sum_{k=0}^{\Delta-1} \sum_{i=1}^m \operatorname{Re}(z_i^k \xi_i \bar{\eta}_k) \end{aligned}$$

is ν . Note that $\mathcal{K}_\nu^\Delta = \mathcal{N}_\nu$.



The class \mathcal{N}_κ

For $\kappa \in \mathbb{N}_0$, denote by \mathcal{N}_κ the set of all functions τ , which are meromorphic in $\mathbb{C} \setminus \mathbb{R}$, satisfy $\tau(\bar{z}) = \overline{\tau(z)}$, and are such that the Nevanlinna kernel

$$Q_f(w, z) := \frac{\tau(z) - \overline{\tau(w)}}{z - \bar{w}}$$

has κ negative squares. This means that the maximal number of negative squares of quadratic forms

$$q_\tau(\xi_1, \dots, \xi_m) = \sum_{i,j=1}^m Q_f(z_j, z_i) \xi_i \overline{\xi_j}$$

is equal to κ .



The Weyl coefficient $q_{\mathfrak{h}}(z)$

The limit

$$q_{\mathfrak{h}}(z) := \lim_{t \nearrow \sigma_n} W_t(z) \star \tau$$

exists as a meromorphic function locally uniformly on $\mathbb{C} \setminus \mathbb{R}$ and does not depend on $\tau \in \mathbb{C}^+$.



The maximal chain (W_t)

The matrix function W_t , $t \in \bigcup_{i=1}^n (\sigma_{i-1}, \sigma_i)$, is for every $i = 1, \dots, n$ a solution of the differential equation

$$\frac{d}{dt} W_t(z) \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} = z W_t(z) H_i(t), \quad x \in (\sigma_{i-1}, \sigma_i)$$

On the interval $[\sigma_0, \sigma_1)$ it is uniquely determined by its initial value $W_{\sigma_0} = I$.



The reproducing kernel space of W_t

The kernel

$$K_{W_t}(w, z) := \frac{W_t(z) \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} W_t(w)^* - \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}}{z - \bar{w}}$$

has a finite number of negative squares, thus generates a reproducing kernel Pontryagin space $\mathfrak{K}(W_t)$. The elements of $\mathfrak{K}(W_t)$ are entire 2-vector-functions.

The operator $\mathcal{S}(W_t)$ of multiplication by z is a symmetry with defect 2. The map $\Gamma(W_t) : f \mapsto f(0)$ is a boundary map for $\mathcal{S}(W_t)$.

The reproducing kernel space of W_t

There exists an isomorphism Φ_t of the boundary triplets $\langle \mathcal{K}(W_t), \mathcal{S}(W_t), \Gamma(W_t) \rangle$ and $\langle \mathfrak{P}^2(\mathfrak{h}|_{(\sigma_0, t)}), \mathcal{S}(\mathfrak{h}|_{(\sigma_0, t)}), \Gamma(\mathfrak{h}|_{(\sigma_0, t)}) \rangle$

If $J := [s_-, s_+] \subseteq (\sigma_{i-1}, \sigma_i)$, then the map

$$\lambda_J : f(x) \mapsto \int_{s_-}^{s_+} W_x(z) H(x) f(x) dx$$

is an isomorphism of $L^2(H_i|_{[s_-, s_+]})$ onto $\mathcal{K}(W_{s_+})[-]\mathcal{K}(W_{s_-})$. We have

$$\begin{array}{ccc} L^2(H_i|_J) & \xrightarrow{\iota_J} & \mathfrak{P}(\mathfrak{h}|_{(\sigma_0, s_+)}) \\ \lambda_J \downarrow & \nearrow \Phi_{s_+} & \\ \mathcal{K}(W_{s_+})[-]\mathcal{K}(W_{s_-}) & & \end{array}$$



The Fourier transform

The Weyl coefficient admits a representation of the form

$$q_{\mathfrak{h}}(z) = \phi\left(\frac{1}{t - z}\right)$$

with some distribution on $\overline{\mathbb{R}}$. This distribution generates a Pontryagin space $\Pi(\phi)$. There exists an isomorphism of $\mathfrak{P}(\mathfrak{h})$ onto $\Pi(\phi)$.



W_t as resolvent matrix

For $t \in I$ consider

$$S_1 := \left\{ (x; y) \in T(\mathfrak{h}|_{(\sigma_0, t)}) : \right. \\ \left. \pi_{l,1} \Gamma(\mathfrak{h}|_{(\sigma_0, t)})(x; y) = 0, \pi_r \Gamma(\mathfrak{h}|_{(\sigma_0, t)})(x; y) = 0 \right\}$$

$$u : (x; y) \mapsto \pi_{l,2} \Gamma(\mathfrak{h}|_{(\sigma_0, t)})(x; y), \quad (x; y) \in T(\mathfrak{h}|_{(\sigma_0, t)})$$

Then S_1 is symmetric with defect 1 and $u|_{S_1^*}$ is continuous. The matrix function W_t is a u -resolvent matrix of S_1 .



W_t from defect elements

Let $\phi_z, \psi_z \in \ker(T(\mathfrak{h}|_{(\sigma_0,t)}))$ be such that

$$\pi_l \Gamma(\mathfrak{h}|_{(\sigma_0,t)})(\phi_z; z\phi_z) = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \pi_l \Gamma(\mathfrak{h}|_{(\sigma_0,t)})(\psi_z; z\psi_z) = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Then

$$W_t(z) = \begin{pmatrix} \pi_r \Gamma(\mathfrak{h}|_{(\sigma_0,t)})(\phi_z; z\phi_z)^T \\ \pi_r \Gamma(\mathfrak{h}|_{(\sigma_0,t)})(\psi_z; z\psi_z)^T \end{pmatrix}$$



The model space $\mathfrak{P}(\mathfrak{h})$

Given a general Hamiltonian \mathfrak{h} we construct an operator model, which is a Pontryagin space boundary triplet

$$\langle \mathfrak{P}(\mathfrak{h}), T(\mathfrak{h}), \Gamma(\mathfrak{h}) \rangle$$

The actual construction is quite involved and too complicated to be elaborated here.

The model space $\mathfrak{P}(\mathfrak{h})$

If $J = [s_-, s_+] \subseteq (\sigma_i, \sigma_{i+1})$, there exists an isometric and homeomorphic embedding

$$\iota_J : L^2(H_i|_J) \rightarrow \mathfrak{P}(\mathfrak{h})$$

If $J \subseteq J'$, then

$$\begin{array}{ccc} L^2(H_i|_J) & \xrightarrow{\iota_J} & \mathfrak{P}(\mathfrak{h}) \\ \subseteq \downarrow & \nearrow \iota_{J'} & \\ L^2(H_i|_{J'}) & & \end{array}$$



Hamiltonian for q_f

The general Hamiltonian made up of the data

$$\sigma_0 = -1, \sigma_1 = 1, \sigma_2 = +\infty, \quad E = \{-1, 0, 2, +\infty\}$$

$$H_0(t) = \begin{cases} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} & , t \in (-1, 0) \\ \begin{pmatrix} (t-1)^2 & 0 \\ 0 & (t-1)^{-2} \end{pmatrix} & , t \in (0, 1) \end{cases}$$

$$H_1(t) = \begin{pmatrix} (t-1)^2 & 0 \\ 0 & (t-1)^{-2} \end{pmatrix}$$

$$\ddot{o}_1 = 1, b_{1,1} = 2, b_{1,2} = 0, \quad d_0 = -2, d_1 = 0$$

has Weyl coefficient q_f .



Hamiltonian for q_α

The general Hamiltonian made up of the data

$$\sigma_0 = -1, \sigma_1 = 0, \sigma_2 = +\infty, \quad E = \{-1, 1, +\infty\}$$

$$H_0(t) = \frac{1}{t^2} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}, \quad H_1(t) = \begin{pmatrix} t^\alpha & 0 \\ 0 & t^{-\alpha} \end{pmatrix}$$

$$\ddot{o}_1 = 0, \quad d_0 = \frac{1}{\alpha - 1}, d_1 = 0$$

has Weyl coefficient q_α .



The class \mathcal{M}_κ

$W \in \mathcal{M}_\kappa$ if

- W is entire 2×2 -matrix function
- $W(0) = I$
- The kernel

$$K_W(w, z) := \frac{W(z) \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} W(w)^* - \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}}{z - \overline{w}}$$

has κ negative squares

