

Autonomous Determination of an Orbit around an irregular-shaped Celestial Body

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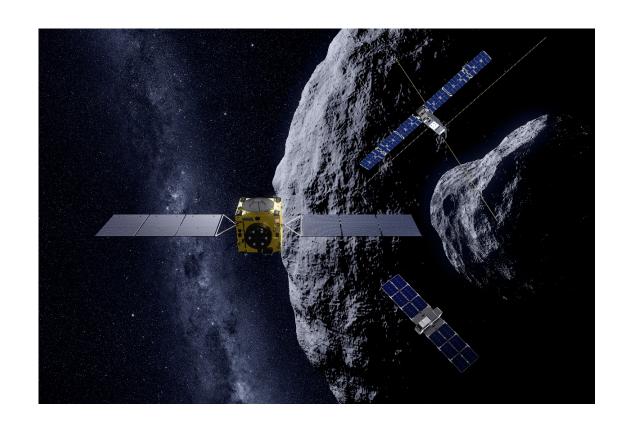






Introduction

- There is a growing interest in deep space missions for the exploration of asteroids and comets.
- NEAR-Shoemaker [Eros, 2001],
 Rosetta [Churyumov-Gerasimenko, 2014],
 Hayabusa [Itokawa, 2005],
 Hayabusa-2 [Ryugu, 2018],
 OSIRIS-Rex [Bennu, 2018-21],
 demonstrated close proximity operations near
 asteroids.
- Trend toward more autonomous functionalities for guidance, navigation and control (GNC)



ESA's Hera mission seen with its CubeSats in orbit around its target asteroid. Launch planned for October 2024.

Credit: ESA/Science Office







Small Celestial Bodies

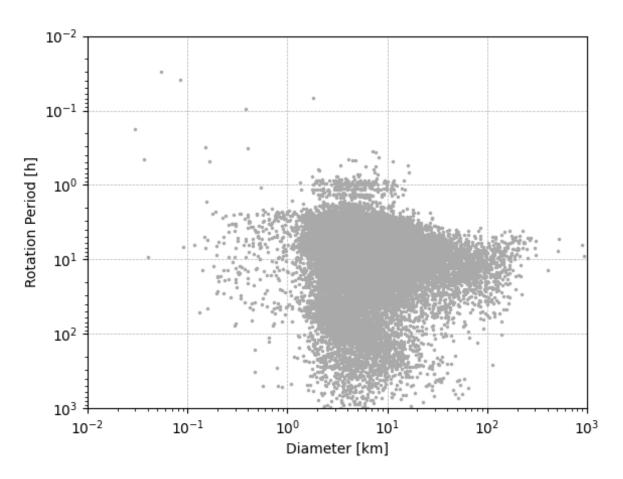


Fig 2: Distribution of small bodies (rotation period vs. diameter)

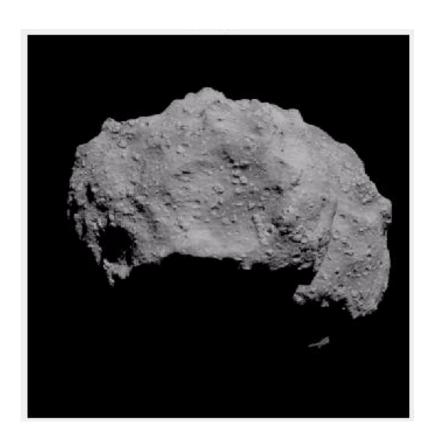


Fig 3: Animation of 67P/Churyumov-Gerasimenko in Blender



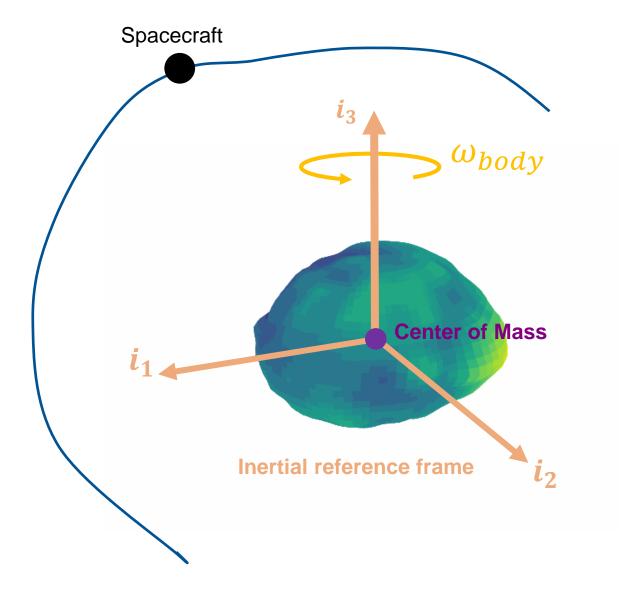




Problem Setting

Given a spacecraft in a bounded orbit around a <u>rotating</u> small celestial body.

- 1. What is the detailed shape of the non-Keplerian orbit?
- 2. What is the shape of the celestial body?
- 3. Where is its center of mass?
- 4. What is its rotation axis?
 - → reference frame









Traditional Deep-Space Missions

- They require communication with an Earth ground-segment for navigation
- Appart from some autonomous functionality, all exisiting spacecraft operate by uploaded command sequences
- During cruise and approach phase, navigation & control heavily relies on radiometric tracking from Earth

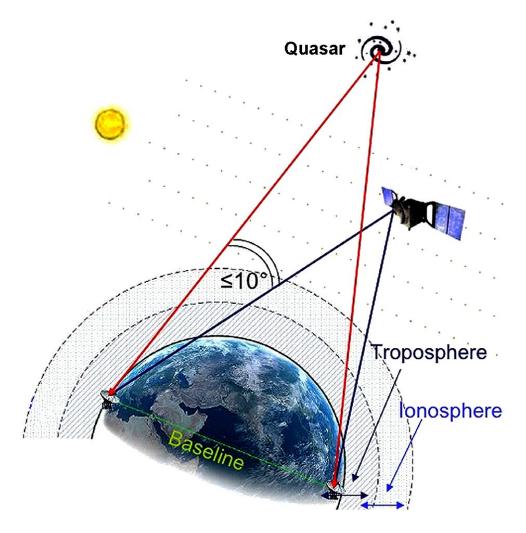


Fig 4: Concept of ground-based radiometric tracking (Delta-DOR)

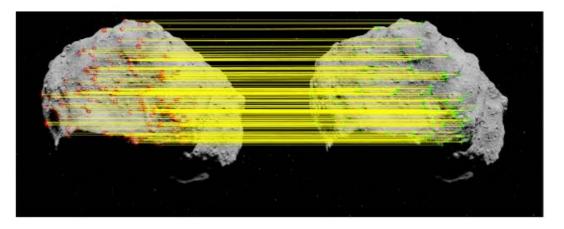






Optical Navigation

- Extract observables from images using image-processing techniques:
 - ✓ Line-of-Sight vector
 - ✓ Apparent diameter
- Observables are usually fed into a state estimation filter using a dynamical model
- For closer distance, the body becomes resolved: Feature-based methods
- 3D-shape of the body (polyhedron)
- rotation axis



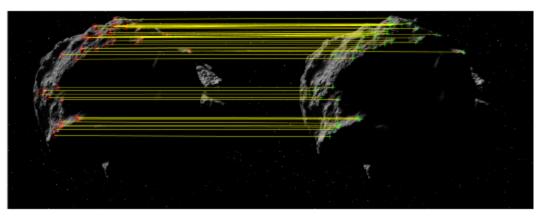


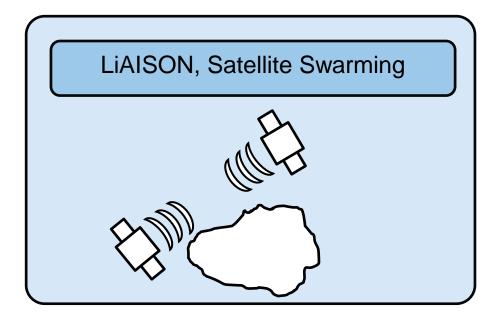
Fig 5: Examples of feature-matching with different lighting conditions





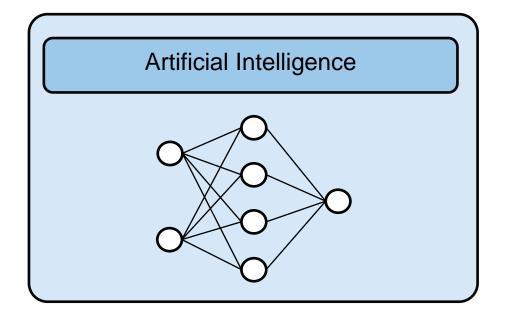


Further Studies on Automous Navigation



- Orbit determination by inter-satellite measurements
- Multiple satellites are utilized

LiAISON – Linked Autonomous Inter Satellite Orbit Navigation

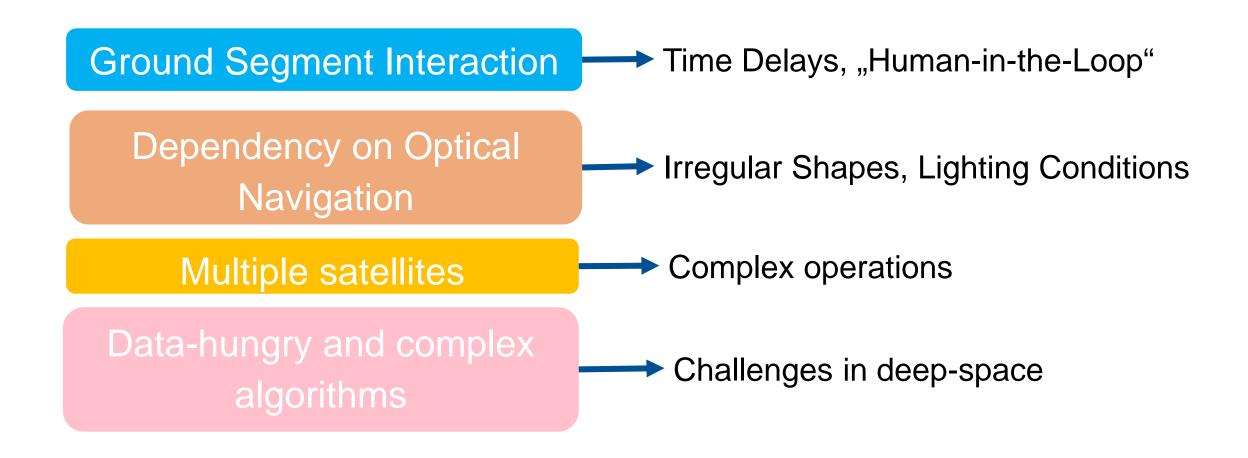


- Navigation techniques benefit from artificial intelligence
- Can be used in an end-to-end manner or for a specific low-level task





Limitations and Opportunities



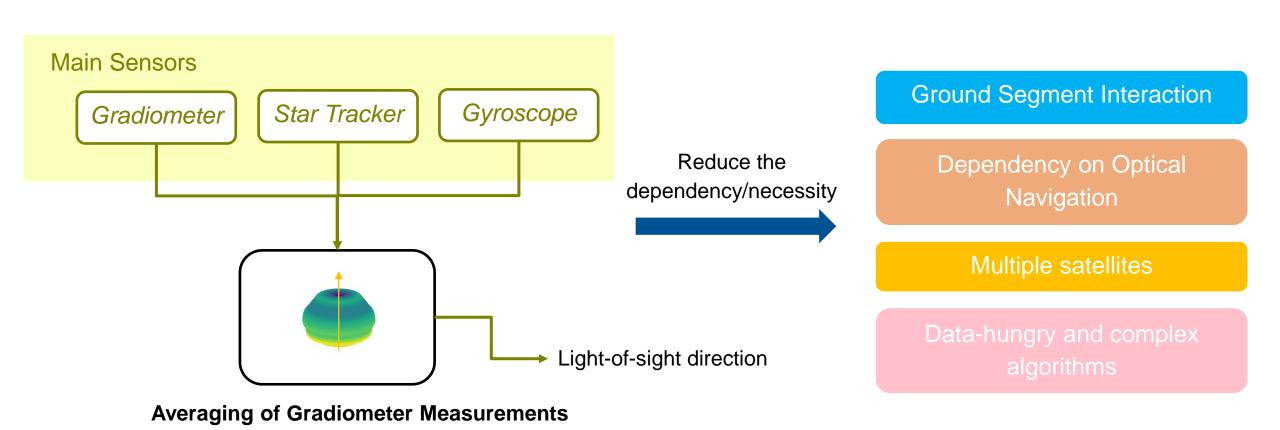






Proposed new Technology

Use Gradiometer to support autonomous navigation near irregular-shaped bodies





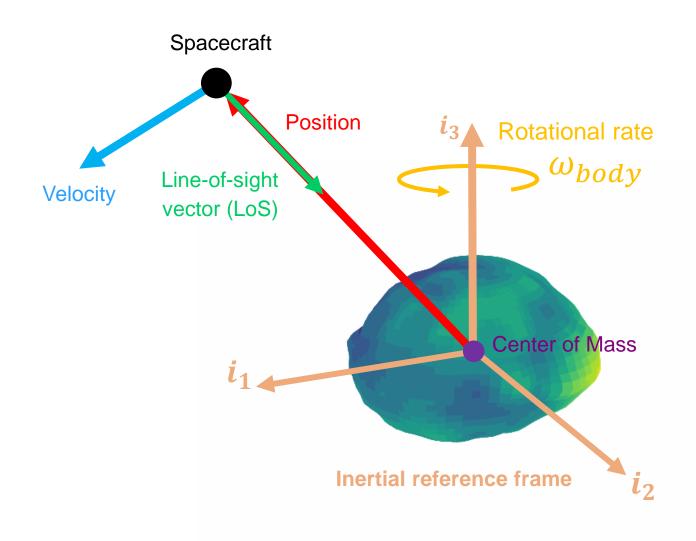




Solution Approach

The spacecraft state (state vector) is defined by 6 degrees of freedom: position and velocity in a given inertial reference frame

 We focus on determining the line-of-sight vector to the center of mass.









Approach and Assumptions

Goal: Extract Line-of-Sight vector from measurements

 $\langle 1 \rangle$: Average using body's rotational rate

$$P_{body} \ll P_{orbit}$$







$$\epsilon = rac{r}{R}$$
 r Orbital radius R Mean radius of body

Trade-off: Best distance to body?

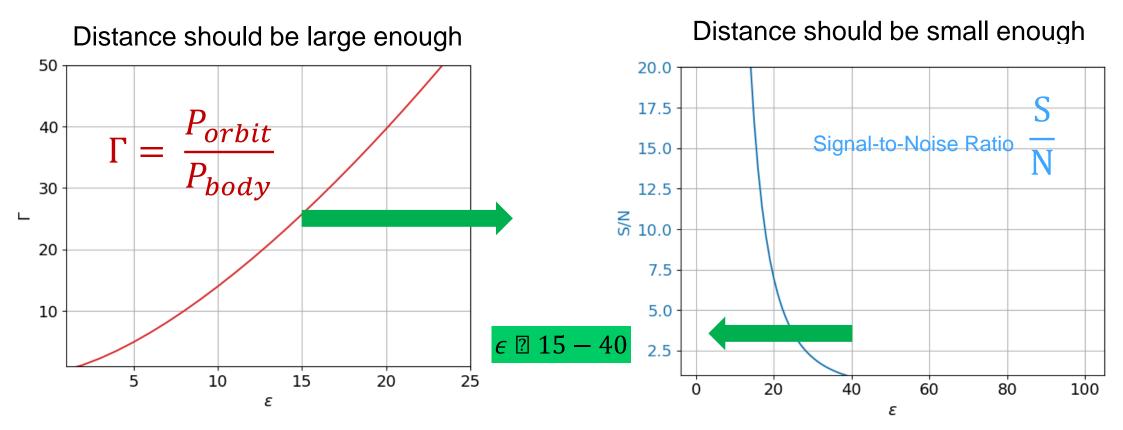


Fig 6: Example how the period ratio depends on the distance

Fig 7: Example how the signal-to-noise ratio depends on the distance

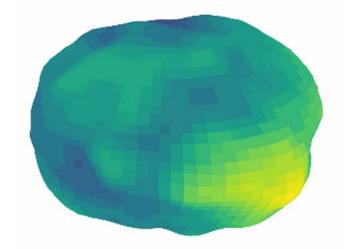




Gravitational Field Representation

Different representations:

- Polyedron Model
- Mascon Model
- **Spherical Harmonics**



Gravitational Potential

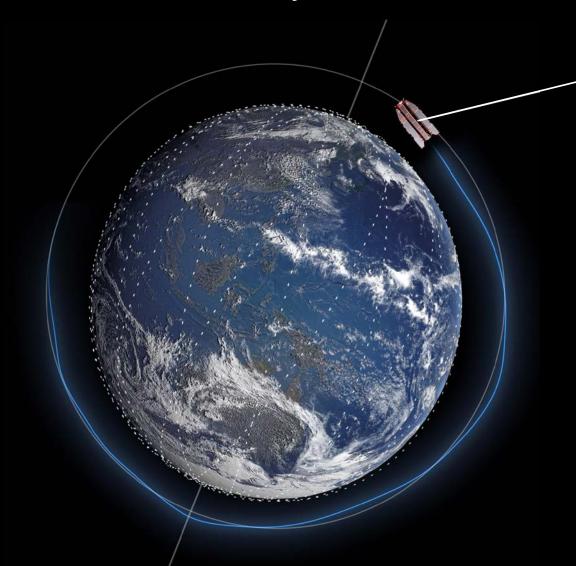
$$U(r,\lambda,\phi) = \frac{\mu}{r} \sum_{n=0}^{\infty} \sum_{m=0}^{n} \frac{R^n}{r^n} (C_{nm} \cos m\lambda + S_{nm} \sin m\lambda) P_{nm}(\sin \phi)$$

$$T = \nabla \nabla U = \begin{bmatrix} \frac{\partial^2 U}{\partial^2 x} & \frac{\partial^2 U}{\partial x \partial y} & \frac{\partial^2 U}{\partial x \partial z} \\ \frac{\partial^2 U}{\partial y \partial x} & \frac{\partial^2 U}{\partial^2 x} & \frac{\partial^2 U}{\partial y \partial z} \\ \frac{\partial^2 U}{\partial z} & \frac{\partial^2 U}{\partial z} & \frac{\partial^2 U}{\partial z} \end{bmatrix} = \begin{bmatrix} U_{xx} & U_{xy} & U_{xz} \\ U_{yx} & U_{yy} & U_{yz} \\ U_{zx} & U_{zy} & U_{zz} \end{bmatrix}$$
 Gravitational Gradient Tensor

 r, λ, ϕ Radius, geocentric longitude, latitude Gravity parameter Reference radius Associated Legendre polynomials $P_{nm}(\sin \phi)$ C_{nm} , S_{nm} Gravitational coefficients Degree, order

Satellite Gradiometry – GOCE







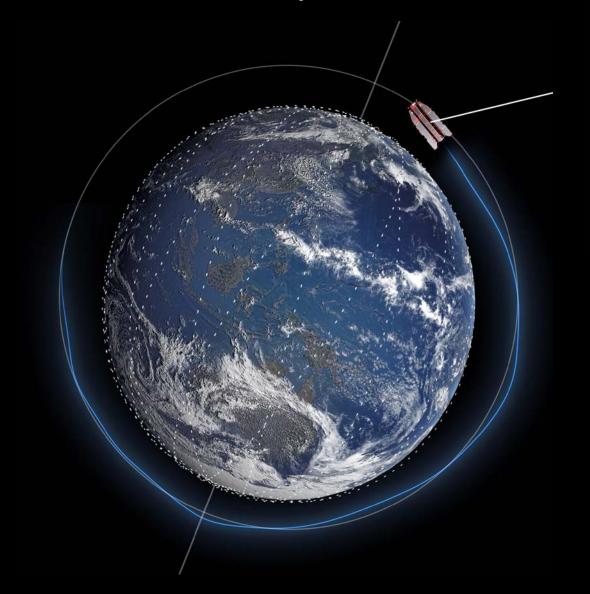
3 pairs of proof masses. A capacitive measuring principle measures the differential acceleration in one spatial direction = 3D tidal force at one orbital point.

→ Medium and short wavelength part of the gravitational potential

GPS-Measurement

→ Long wavelength portion

Satellite Gradiometry – GOCE







Satellite Gradiometry

3 pairs of proof masses. A capacitive measuring principle measures the differential acceleration in one spatial direction = 3D tidal force at one orbital point.

→ Medium and short wavelength part of the gravitational potential

GPS-Measurement

→ Long wavelength portion





Measurement Principle

Measurements from six accelerometers:

$$a_i = -Tr + \dot{\omega} \times r + \omega \times (\omega \times r)$$
Euler Centrifugal acceleration acceleration

- \boldsymbol{a}_i Measurered acceleration vector
- T Gravity gradient tensor
- r Distance to center of mass
- ω Angular velocity

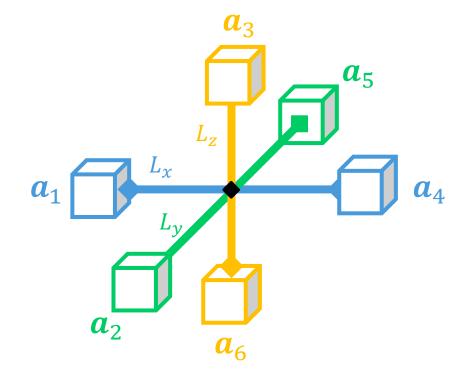


Fig 8: Concept of a Gradiometer







Common-Mode and Differential-Mode Accelerations

Can be extracted from gradiometer data

$$a_i = -Tr + \dot{\omega} \times r + \omega \times (\dot{\omega} \times r) \longrightarrow T$$

Common-mode accelerations:

$$a_{c,1,4,i} = \frac{1}{2} (a_{1,i} + a_{4,i}) = 0$$

$$a_{c,2,5,i} = \frac{1}{2} (a_{2,i} + a_{5,i}) = 0$$

$$a_{c,3,6,i} = \frac{1}{2} (a_{3,i} + a_{6,i}) = 0$$

Vanish for non-gravitational accelerations (e.g. drag and solar radiation pressure)

Differential-mode accelerations:

$$a_{c,1,4,i} = \frac{1}{2} (a_{1,i} - a_{4,i}) = f(\mathbf{T}, \boldsymbol{\omega}, \dot{\boldsymbol{\omega}})$$

Can be extracted from gyroscope/startracker data

$$a_{c,2,5,i} = \frac{1}{2} (a_{2,i} - a_{5,i}) = f(\mathbf{T}, \boldsymbol{\omega}, \dot{\boldsymbol{\omega}})$$

$$a_{c,3,6,i} = \frac{1}{2} (a_{3,i} - a_{6,i}) = f(\mathbf{T}, \boldsymbol{\omega}, \dot{\boldsymbol{\omega}})$$

T can be exstracted using angular rates/accelerations

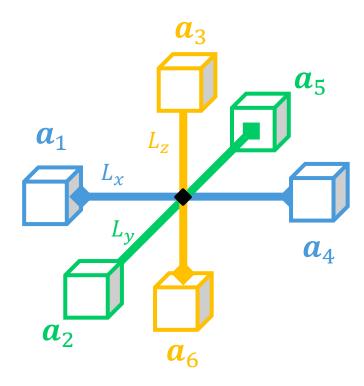


Fig 8: Concept of a Gradiometer



Eigendecomposition

$$\mathbf{T}^{i} = \begin{bmatrix} U_{xx} & U_{xy} & U_{xz} \\ U_{yx} & U_{yy} & U_{yz} \\ U_{zx} & U_{zy} & U_{zz} \end{bmatrix} = \mathbf{V} \mathbf{\Lambda} \mathbf{V}^{T} = \begin{bmatrix} \mathbf{v}_{1} & \mathbf{v}_{2} & \mathbf{v}_{3} \end{bmatrix}^{T}$$
$$\begin{bmatrix} \mathbf{v}_{1} & \mathbf{v}_{2} & \mathbf{v}_{3} \end{bmatrix} \begin{bmatrix} \lambda_{1} & 0 & 0 \\ 0 & \lambda_{2} & 0 \\ 0 & 0 & \lambda_{3} \end{bmatrix} [\mathbf{v}_{1} & \mathbf{v}_{2} & \mathbf{v}_{3}]^{T}$$

 T^i Gravity Gradient Tensor resolved in \mathcal{F}_i

 $\lambda_1, \lambda_2, \lambda_3$ Eigenvalues

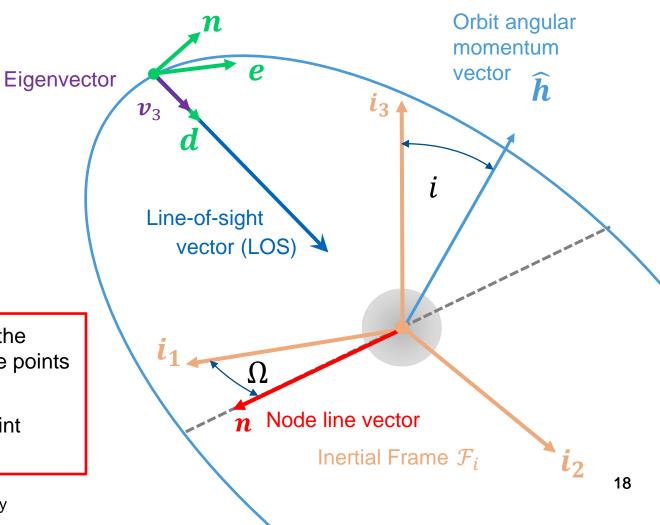
 $oldsymbol{v}_1 \quad oldsymbol{v}_2 \quad oldsymbol{v}_3 \quad ext{ Eigenvectors resolved in } \mathcal{F}_i$

For point mass: $\mathbf{T}^n = \frac{\mu}{r^3} \begin{bmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 2 \end{bmatrix}$ Eigenvector with the largest eigenvalue points to body's CM

For any irregular- Eigenvector does not necessarily point to body's center of mass

North-East-Down (NED) Frame \mathcal{F}_n

- **n** North direction
- e East direction
- **d** Down direction



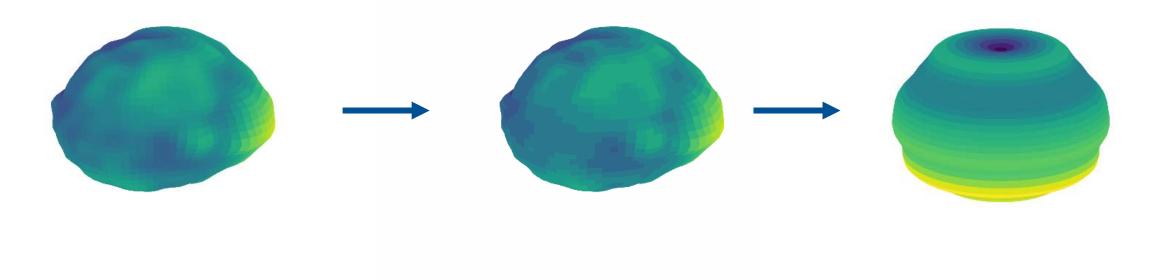






(1): Average using Body's Rotational Rate

- Assume: $P_{body} \ll P_{orbit}$
- Collect gradiometer measurements: $T_j, T_{j+1}, ...$
- Take the average over period P_{body} : $\langle T \rangle$



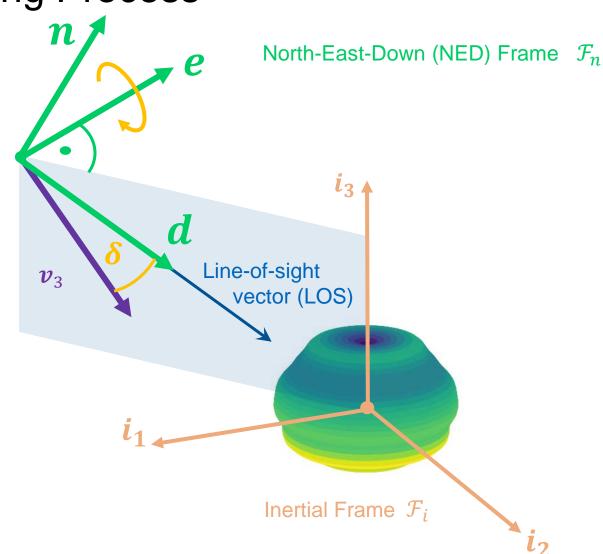


Eigendecomposition after Averaging Process

$$\mathbf{T}^{n} = \begin{bmatrix} U_{xx} & 0 & U_{xz} \\ 0 & U_{yy} & 0 \\ U_{zx} & 0 & U_{zz} \end{bmatrix} \\
= \begin{bmatrix} \cos \delta & 0 & -\sin \delta \\ 0 & 1 & 0 \\ \sin \delta & 0 & \cos \delta \end{bmatrix} \begin{bmatrix} \lambda_{1} & 0 & 0 \\ 0 & \lambda_{2} & 0 \\ 0 & 0 & \lambda_{3} \end{bmatrix} \begin{bmatrix} \cos \delta & 0 & \sin \delta \\ 0 & 1 & 0 \\ -\sin \delta & 0 & \cos \delta \end{bmatrix}$$

 $m{T}^i$ Gravity Gradient Tensor resolved in \mathcal{F}_n $\lambda_1,\lambda_2,\lambda_3$ Eigenvalues δ Eigenvectors resolved in \mathcal{F}_n

The Eigenvector v_3 lies in the n, d-plane and is rotated by the angle δ





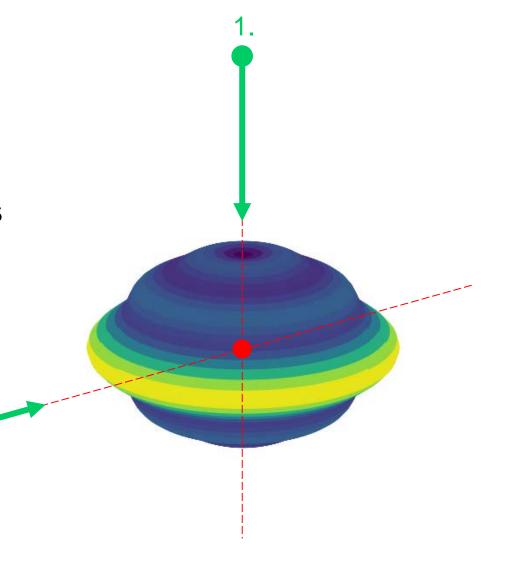


When is $\delta = 0$?

Condition $\delta = 0$ holds for:

- 1. Position on rotation axis (above the poles)
- 2. In equatorial plane for even zonal harmonics

3. Far away from body









Approach and Assumptions

Goal: Extract Line-of-Sight vector from measurements

 $\langle 1 \rangle$: Average using body's rotational rate

(2): Average using orbital rate

$$P_{body} \ll P_{orbit}$$

Forced circular polar motion of spacecraft

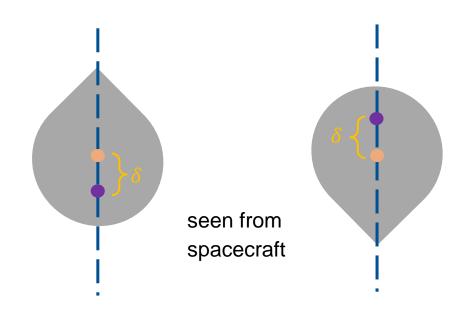


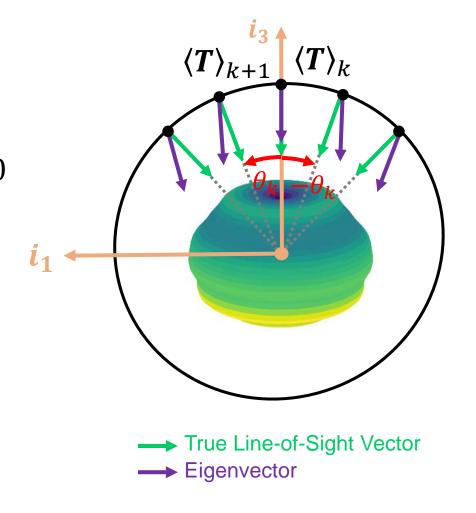




(2): Average by spacecraft polar orbit

- Assume polar orbital motion: $-\delta(\theta) = \delta(-\theta)$
- Collect measurements: $\langle T \rangle_k$, $\langle T \rangle_{k+1}$, ...
- Take the average over period: $P_{orbit} \rightarrow \langle \delta \rangle = 0$



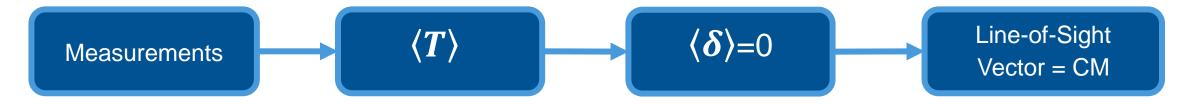




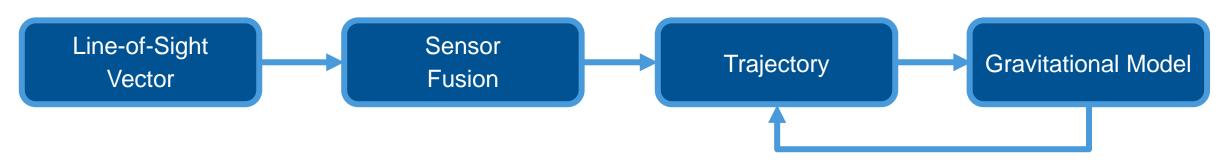




Summary of Approach



Future potentials to the complete full trajectory and gravitational model:

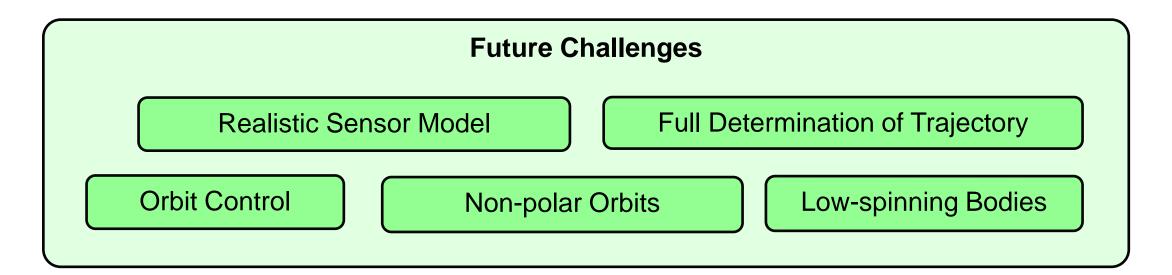




Conclusion

Major Findings

- Gradiometry allows to determine the line-of-sight vector → CM for an irregular-shaped body without knowing the body's shape
- Determination is accomplished by averaging the measurements
- Approach addresses shortcomings of optical navigation







Thank you for your attention!

