

Interpolating conformal algebra between the instant form and the front form of relativistic dynamics

Light Cone 2021: Contributed talk

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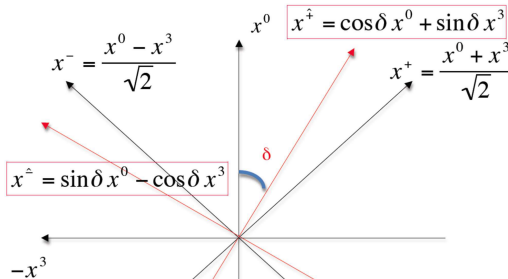
Motivation

- The instant form and the front form of relativistic dynamics introduced by Dirac in 1949 can be interpolated by introducing an interpolation angle parameter δ spanning between the instant form dynamics (IFD) at $\delta = 0$ and the front form dynamics, which is now known as the light-front dynamics (LFD) at $\delta = \frac{\pi}{4}$. **K.Hornboste, Phys. Rev. D 45, 3781 (1992).**
- In **C.-R. Ji and C. Mitchell, Phys. Rev. D 64, 085013 (2001)**, they showed the Boost K^3 is dynamical in the region where $0 \leq \delta < \frac{\pi}{4}$ but becomes kinematic in the light-front limit ($\delta = \frac{\pi}{4}$).
- We extend the Poincaré algebra interpolation between instant and light-front time quantizations to the conformal algebra.
- Among the five more generators in the conformal algebra, only one generator known as the dilatation is kinematic for the entire region of the interpolation angle ($0 \leq \delta \leq \frac{\pi}{4}$).
- We find that one more generator from the Special Conformal Transformation (SCT) becomes kinematic in the light-front limit ($\delta = \frac{\pi}{4}$), i.e. the LFD.

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Interpolation between Instant and Front Forms



K. Hornbostel, PRD45, 3781 (1992) – **RQFT**

C.Ji and S.Rey, PRD53,5815(1996) – **Chiral Anomaly**

C.Ji and C. Mitchell, PRD64,085013 (2001) – **Poincare Algebra**

C.Ji and A. Suzuki, PRD87,065015 (2013) – **Scattering Amps**

C.Ji, Z. Li and A. Suzuki, PRD91, 065020 (2015) – **EM Gauges**

Z.Li, M. An and C.Ji, PRD92, 105014 (2015) – **Spinors**

C.Ji, Z.Li, B.Ma and A.Suzuki, in prepartion – **Fermion Prop.**

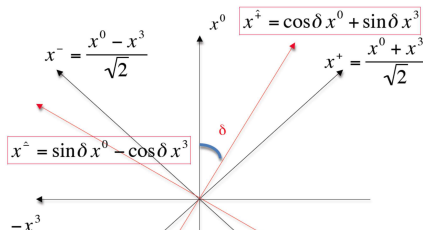
Method of Interpolation Angle

The interpolating space-time coordinates may be defined as a transformation from the ordinary space-time coordinates, $x^{\hat{\mu}} = \mathcal{R}^{\hat{\mu}}_{\nu} x^{\nu}$, i.e.

$$\begin{pmatrix} x^{\hat{+}} \\ x^{\hat{1}} \\ x^{\hat{2}} \\ x^{\hat{-}} \end{pmatrix} = \begin{pmatrix} \cos \delta & 0 & 0 & \sin \delta \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ \sin \delta & 0 & 0 & -\cos \delta \end{pmatrix} \begin{pmatrix} x^0 \\ x^1 \\ x^2 \\ x^3 \end{pmatrix}, \quad (1)$$

in which the interpolation angle is allowed to run from 0 through 45° , $0 \leq \delta \leq \frac{\pi}{4}$.

Interpolation between Instant and Front Forms



Method of Interpolation Angle

In this interpolating basis, the metric becomes

$$g^{\hat{\mu}\hat{\nu}} = g_{\hat{\mu}\hat{\nu}} = \begin{pmatrix} \mathbb{C} & 0 & 0 & \mathbb{S} \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ \mathbb{S} & 0 & 0 & -\mathbb{C} \end{pmatrix}, \quad (2)$$

where $\mathbb{S} = \sin 2\delta$ and $\mathbb{C} = \cos 2\delta$.

The Poincaré matrix

$$M^{\mu\nu} = \begin{pmatrix} 0 & K^1 & K^2 & K^3 \\ -K^1 & 0 & J^3 & -J^2 \\ -K^2 & -J^3 & 0 & J^1 \\ -K^3 & J^2 & -J^1 & 0 \end{pmatrix} \quad (3)$$

transforms as well, so that

$$M^{\hat{\mu}\hat{\nu}} = \begin{pmatrix} 0 & E^{\hat{1}} & E^{\hat{2}} & -K^3 \\ -E^{\hat{1}} & 0 & J^3 & -F^{\hat{1}} \\ -E^{\hat{2}} & -J^3 & 0 & -F^{\hat{2}} \\ K^3 & F^{\hat{1}} & F^{\hat{2}} & 0 \end{pmatrix} \quad (4)$$

where

$$\begin{aligned} E^{\hat{1}} &= J^2 \sin \delta + K^1 \cos \delta, \\ E^{\hat{2}} &= K^2 \cos \delta - J^1 \sin \delta, \\ F^{\hat{1}} &= K^1 \sin \delta - J^2 \cos \delta, \\ F^{\hat{2}} &= K^2 \sin \delta + J^1 \cos \delta. \end{aligned} \quad (5)$$

The Poincaré matrix

$$M_{\hat{\mu}\hat{\nu}} = g_{\hat{\mu}\hat{\alpha}} M^{\hat{\alpha}\hat{\beta}} g_{\hat{\beta}\hat{\nu}} = \begin{pmatrix} 0 & \mathcal{D}^{\hat{1}} & \mathcal{D}^{\hat{2}} & K^3 \\ -\mathcal{D}^{\hat{1}} & 0 & J^3 & -\mathcal{K}^{\hat{1}} \\ -\mathcal{D}^{\hat{2}} & -J^3 & 0 & -\mathcal{K}^{\hat{2}} \\ -K^3 & \mathcal{K}^{\hat{1}} & \mathcal{K}^{\hat{2}} & 0 \end{pmatrix}, \quad (6)$$

where

$$\begin{aligned} \mathcal{K}^{\hat{1}} &= -K^1 \sin \delta - J^2 \cos \delta, \\ \mathcal{K}^{\hat{2}} &= J^1 \cos \delta - K^2 \sin \delta, \\ \mathcal{D}^{\hat{1}} &= -K^1 \cos \delta + J^2 \sin \delta, \\ \mathcal{D}^{\hat{2}} &= -J^1 \sin \delta - K^2 \cos \delta. \end{aligned} \quad (7)$$

Generators of Poincaré group

$$(\text{translation}) \quad P^{\hat{\mu}} = -i\partial^{\hat{\mu}}, \quad (8)$$

$$(\text{rotation}) \quad L^{\hat{\mu}\hat{\nu}} = i(x^{\hat{\mu}}\partial^{\hat{\nu}} - x^{\hat{\nu}}\partial^{\hat{\mu}}). \quad (9)$$

In the interpolating basis, the metric becomes

$$g^{\hat{\mu}\hat{\nu}} = g_{\hat{\mu}\hat{\nu}} = \begin{pmatrix} \mathbb{C} & 0 & 0 & \mathbb{S} \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ \mathbb{S} & 0 & 0 & -\mathbb{C} \end{pmatrix}, \quad (10)$$

The Poincaré algebra (Contra-variant form) in this interpolating basis is given by

$$[P^{\hat{\mu}}, P^{\hat{\nu}}] = 0, \quad (11a)$$

$$[P^{\hat{\rho}}, L^{\hat{\mu}\hat{\nu}}] = i(g^{\hat{\rho}\hat{\mu}}P^{\hat{\nu}} - g^{\hat{\rho}\hat{\nu}}P^{\hat{\mu}}), \quad (11b)$$

$$[L^{\hat{\alpha}\hat{\beta}}, L^{\hat{\rho}\hat{\sigma}}] = -i(g^{\hat{\beta}\hat{\sigma}}L^{\hat{\alpha}\hat{\rho}} - g^{\hat{\beta}\hat{\rho}}L^{\hat{\alpha}\hat{\sigma}} + g^{\hat{\alpha}\hat{\rho}}L^{\hat{\beta}\hat{\sigma}} - g^{\hat{\alpha}\hat{\sigma}}L^{\hat{\beta}\hat{\rho}}). \quad (11c)$$

A comprehensive table of the 45 commutation relations among the co-variant components of the Poincaré' generators is presented below:

	P_+	P_i	P_z	K^3	\mathcal{D}^1	\mathcal{D}^2	J^3	\mathcal{K}^1	\mathcal{K}^2	P_-
P_+	0	0	0	$i(\mathcal{C}P_- - \mathcal{S}P_+)$	$i\mathcal{C}P_i$	$i\mathcal{C}P_z$	0	$i\mathcal{S}P_i$	$i\mathcal{S}P_z$	0
P_i	0	0	0	0	iP_+	0	$-iP_z$	iP_-	0	0
P_z	0	0	0	0	0	iP_+	iP_i	0	iP_-	0
K^3	$-i(\mathcal{C}P_- - \mathcal{S}P_+)$	0	0	0	$i\mathcal{S}\mathcal{D}^1 - i\mathcal{C}\mathcal{K}^1$	$i\mathcal{S}\mathcal{D}^2 - i\mathcal{C}\mathcal{K}^2$	0	$-i\mathcal{S}\mathcal{K}^1 - i\mathcal{C}\mathcal{D}^1$	$-i\mathcal{S}\mathcal{K}^2 - i\mathcal{C}\mathcal{D}^2$	$-i(\mathcal{S}P_- + \mathcal{C}P_+)$
\mathcal{D}^1	$-i\mathcal{C}P_i$	$-iP_+$	0	$-i\mathcal{S}\mathcal{D}^1 + i\mathcal{C}\mathcal{K}^1$	0	$-i\mathcal{C}J^3$	$-i\mathcal{D}^2$	$-i\mathcal{K}^3$	$-i\mathcal{S}J^3$	$-i\mathcal{S}P_i$
\mathcal{D}^2	$-i\mathcal{C}P_z$	0	$-iP_+$	$-i\mathcal{S}\mathcal{D}^2 + i\mathcal{C}\mathcal{K}^2$	$i\mathcal{C}J^3$	0	$i\mathcal{D}^1$	$i\mathcal{S}J^3$	$-i\mathcal{K}^3$	$-i\mathcal{S}P_z$
J^3	0	iP_z	$-iP_i$	0	$i\mathcal{D}^2$	$-i\mathcal{D}^1$	0	$i\mathcal{K}^2$	$-i\mathcal{K}^1$	0
\mathcal{K}^1	$-i\mathcal{S}P_i$	$-iP_-$	0	$i\mathcal{S}\mathcal{K}^1 + i\mathcal{C}\mathcal{D}^1$	$i\mathcal{K}^3$	$-i\mathcal{S}J^3$	$-i\mathcal{K}^2$	0	$i\mathcal{C}J^3$	$i\mathcal{C}P_i$
\mathcal{K}^2	$-i\mathcal{S}P_z$	0	$-iP_-$	$i\mathcal{S}\mathcal{K}^2 + i\mathcal{C}\mathcal{D}^2$	$i\mathcal{S}J^3$	$i\mathcal{K}^3$	$i\mathcal{K}^1$	$-i\mathcal{C}J^3$	0	$i\mathcal{C}P_z$
P_-	0	0	0	$i(\mathcal{S}P_- + \mathcal{C}P_+)$	$i\mathcal{S}P_i$	$i\mathcal{S}P_z$	0	$-i\mathcal{C}P_i$	$-i\mathcal{C}P_z$	0

Interpolation angle	Kinematic	Dynamic
$\delta = 0$	$\mathcal{K}^{\hat{1}} = -J^2, \mathcal{K}^{\hat{2}} = J^1, J^3, P^1, P^2, P^3$	$\mathcal{D}^{\hat{1}} = -K^1, \mathcal{D}^{\hat{2}} = -K^2, K^3, P^0$
$0 \leq \delta < \pi/4$	$\mathcal{K}^{\hat{1}}, \mathcal{K}^{\hat{2}}, J^3, P^1, P^2, P_-$	$\mathcal{D}^{\hat{1}}, \mathcal{D}^{\hat{2}}, K^3, P_+$
$\delta = \pi/4$	$\mathcal{K}^{\hat{1}} = -E^1, \mathcal{K}^{\hat{2}} = -E^2, J^3, K^3, P^1, P^2, P_-$	$\mathcal{D}^{\hat{1}} = -F^1, \mathcal{D}^{\hat{2}} = -F^2, P_+$

Chueng-Ryong Ji and Chad Mitchell, Phys. Rev. **D 64**, 085013 (2001).

Chueng-Ryong Ji, Ziyue Li, and Alfredo Takashi Suzuki, Phys. Rev. **D 91**, 065020 (2015).

Kinematic and dynamic generators for different interpolation angles (Phys. Rev. **D 64**, 085013 (2001); Phys. Rev. **D 91**, 065020 (2015))

Interpolation angle	Kinematic	Dynamic
$\delta = 0$	$\mathcal{K}^{\hat{1}} = -J^2, \mathcal{K}^{\hat{2}} = J^1, J^3, P^1, P^2, P^3$	$\mathcal{D}^{\hat{1}} = -K^1, \mathcal{D}^{\hat{2}} = -K^2, K^3, P^0$
$0 \leq \delta < \pi/4$	$\mathcal{K}^{\hat{1}}, \mathcal{K}^{\hat{2}}, J^3, P^1, P^2, P_{\hat{\perp}}$	$\mathcal{D}^{\hat{1}}, \mathcal{D}^{\hat{2}}, K^3, P_{\hat{\perp}}$
$\delta = \pi/4$	$\mathcal{K}^{\hat{1}} = -E^1, \mathcal{K}^{\hat{2}} = -E^2, J^3, K^3, P^1, P^2, P_{\hat{\perp}}$	$\mathcal{D}^{\hat{1}} = -F^1, \mathcal{D}^{\hat{2}} = -F^2, P_{\hat{\perp}}$

- Among the ten Poincaré generators, the six generators $(\mathcal{K}^{\hat{1}}, \mathcal{K}^{\hat{2}}, J^3, P_1, P_2, P_{\hat{\perp}})$ are always kinematic in the sense that the $x^{\hat{\perp}} = 0$ plane is intact under the transformations generated by them. The operator $K^3 = M_{\hat{\perp}\hat{\perp}}$ is dynamical in the region where $0 \leq \delta < \pi/4$ but becomes kinematic in the light-front limit ($\delta = \pi/4$).
- To understand this, note that $[P^{\hat{\perp}}, K^{\hat{3}}] = i(\mathbb{S}P^{\hat{\perp}} - \mathbb{C}P^{\hat{\perp}}) \rightarrow iP^{\hat{\perp}}$ as $\delta \rightarrow \pi/4$. Similarly we have $[x^{\hat{\perp}}, L^{\hat{\perp}\hat{\perp}}] = i(\mathbb{S}x^{\hat{\perp}} - \mathbb{C}x^{\hat{\perp}}) \rightarrow ix^{\hat{\perp}}$ as $\delta \rightarrow \pi/4$. Therefore the instant defined by $x^+ = 0$ becomes invariant under longitudinal boosts as we move to the light front.

Conformal Transformations

The Conformal transformation $x \mapsto x'$ can be defined by, (1996 CFT Book by Francesco)

$$\frac{\partial x'^{\alpha}}{\partial x^{\mu}} \frac{\partial x'^{\beta}}{\partial x^{\nu}} g'_{\alpha\beta} = F(x) g_{\mu\nu} \quad (12)$$

Consider an infinitesimal translation,

$$x'^{\mu} = x^{\mu} + \epsilon^{\mu}(x) . \quad (13)$$

The metric changes by,

$$\delta g_{\mu\nu} = \frac{\partial \epsilon_{\mu}}{\partial x^{\nu}} + \frac{\partial \epsilon_{\nu}}{\partial x^{\mu}} = \partial_{\mu} \epsilon_{\nu}(x) + \partial_{\nu} \epsilon_{\mu}(x) \quad (14)$$

Conformality then requires,

$$\boxed{\partial_{\mu} \epsilon_{\nu}(x) + \partial_{\nu} \epsilon_{\mu}(x) = F(x) \delta_{\mu\nu}} \quad \text{Conformal Killing Equation} \quad (15)$$

contraction with $\delta^{\mu\nu}$ yields

$$2 \partial^{\mu} \epsilon_{\mu} = F(x) d \quad (16)$$

$$\Rightarrow F(x) = \frac{2}{d} \partial_{\mu} \epsilon^{\mu} \quad (17)$$

Conformal Transformations

For $d \geq 3$, there are ONLY 4 classes of solutions for $\epsilon_\mu(x)$

$$(\text{Infinitesimal Translation}) \quad \epsilon^\mu(x) = a^\mu \quad (\text{constant}) \quad (18)$$

$$(\text{Infinitesimal Rotation}) \quad \epsilon^\mu(x) = L^\mu_\nu x^\nu \quad (19)$$

$$(\text{Infinitesimal Scaling}) \quad \epsilon^\mu(x) = \lambda x^\mu \quad (20)$$

$$(\text{Infinitesimal SCT}) \quad \epsilon^\mu(x) = 2(b \cdot x)x^\mu - x^2 b^\mu \quad (21)$$

The generators of conformal transformations are:

$$(\text{translation}) \quad P^\mu = -i\partial^\mu ,$$

$$(\text{dilation}) \quad D = -ix_\mu \partial^\mu ,$$

$$(\text{rotation}) \quad L^{\mu\nu} = i(x^\mu \partial^\nu - x^\nu \partial^\mu) ,$$

$$(\text{SCT}) \quad \mathcal{K}^\mu = -i(2x^\mu x_\nu \partial^\nu - x^2 \partial^\mu) .$$

Finite conformal transformations

$$(\text{translation}) \quad x'^{\mu} = x^{\mu} + a^{\mu}$$

$$(\text{dilatation}) \quad x'^{\mu} = \alpha x^{\mu}$$

$$(\text{rotation}) \quad x'^{\mu} = M^{\mu}_{\nu} x^{\nu}$$

$$(\text{SCT}) \quad x'^{\mu} = \frac{x^{\mu} - b^{\mu} x^2}{1 - 2b \cdot x + b^2 x^2}$$

Let us also note that for finite Special Conformal Transformations, we can re-write the expression as follows

$$\begin{aligned} x'^{\mu} &= \frac{x^{\mu} - b^{\mu} x^2}{1 - 2b \cdot x + b^2 x^2} = \frac{x^2}{|x - bx^2|^2} (x^{\mu} - b^{\mu} x^2), \\ \Rightarrow \frac{1}{x'^{\mu}} &= \frac{x^{\mu} - b^{\mu} x^2}{x^2}, \\ \Rightarrow \frac{x'^{\mu}}{x'^2} &= \frac{x^{\mu}}{x^2} - b^{\mu} \end{aligned}$$

SCT

From $\frac{x'^{\mu}}{x'^2} = \frac{x^{\mu}}{x^2} - b^{\mu}$, we see that the SCT can be understood as an inversion of x^{μ} , followed by a translation b^{μ} , and followed again by an inversion.

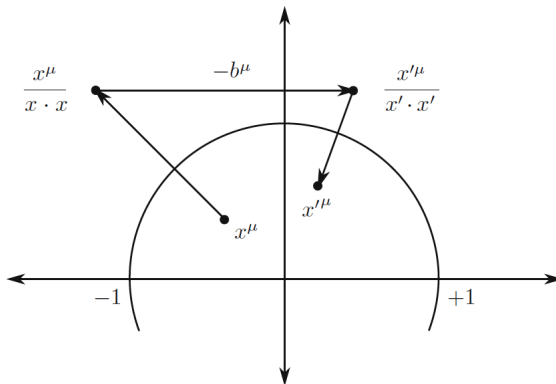


Figure: Illustration of a finite SCT

The generators of conformal transformations:

$$(\text{translation}) \quad P^{\hat{\mu}} = -i\partial^{\hat{\mu}} \quad (22)$$

$$(\text{dilation}) \quad D = -ix_{\hat{\mu}}\partial^{\hat{\mu}} \quad (23)$$

$$(\text{rotation}) \quad L^{\hat{\mu}\hat{\nu}} = i(x^{\hat{\mu}}\partial^{\hat{\nu}} - x^{\hat{\nu}}\partial^{\hat{\mu}}) \quad (24)$$

$$(\text{SCT}) \quad \mathcal{K}^{\hat{\mu}} = -i(2x^{\hat{\mu}}x_{\hat{\nu}}\partial^{\hat{\nu}} - x^2\partial^{\hat{\mu}}) \quad (25)$$

Therefore the full Conformal algebra is given by

$$[P^{\hat{\mu}}, P^{\hat{\nu}}] = 0,$$

$$[\mathcal{K}^{\hat{\mu}}, \mathcal{K}^{\hat{\nu}}] = 0,$$

$$[D, P^{\hat{\mu}}] = iP^{\hat{\mu}},$$

$$[D, \mathcal{K}^{\hat{\mu}}] = -i\mathcal{K}^{\hat{\mu}},$$

$$[P^{\hat{\rho}}, L^{\hat{\mu}\hat{\nu}}] = i(g^{\hat{\rho}\hat{\mu}}P^{\hat{\nu}} - g^{\hat{\rho}\hat{\nu}}P^{\hat{\mu}}),$$

$$[\mathcal{K}^{\hat{\rho}}, L^{\hat{\mu}\hat{\nu}}] = i(g^{\hat{\rho}\hat{\mu}}\mathcal{K}^{\hat{\nu}} - g^{\hat{\rho}\hat{\nu}}\mathcal{K}^{\hat{\mu}}),$$

$$[L^{\hat{\alpha}\hat{\beta}}, L^{\hat{\rho}\hat{\sigma}}] = -i(g^{\hat{\beta}\hat{\sigma}}L^{\hat{\alpha}\hat{\rho}} - g^{\hat{\beta}\hat{\rho}}L^{\hat{\alpha}\hat{\sigma}} + g^{\hat{\alpha}\hat{\rho}}L^{\hat{\beta}\hat{\sigma}} - g^{\hat{\alpha}\hat{\sigma}}L^{\hat{\beta}\hat{\rho}}),$$

$$[\mathcal{K}^{\hat{\mu}}, P^{\hat{\nu}}] = 2i(g^{\hat{\mu}\hat{\nu}}D - L^{\hat{\mu}\hat{\nu}}),$$

$$[D, L^{\hat{\mu}\hat{\nu}}] = 0,$$

Full conformal algebra

A comprehensive table of the **105 commutation** relations among the co-variant components of the Conformal generators is presented below:

	P_+	P_1	P_2	K^3	D^1	D^2	J^3	K^1	K^2	P_-	\mathcal{R}_+	\mathcal{R}_1	\mathcal{R}_2	\mathcal{R}_-	D
P_+	0	0	0	$i(CP_- - SP_+)$	iCP_1	iCP_2	0	iSP_1	iSP_2	0	$-2iCD$	$-2iD^1$	$-2iK^2$	$-2i(SD - K^3)$	$-iP_+$
P_1	0	0	0	0	iP_+	0	$-iP_2$	iP_-	0	0	$2iD^1$	$2iD$	$-2iJ^3$	$2iK^1$	$-iP_1$
P_2	0	0	0	0	0	iP_+	iP_1	0	iP_-	0	$2iD^2$	$2iJ^3$	$2iD$	$2iK^2$	$-iP_2$
K^3	$-i(CP_- - SP_+)$	0	0	0	$iSD^1 - iCK^1$	$iSD^2 - iCK^2$	0	$-iSK^1 - iCD^1$	$-iSK^2 - iCD^2$	$-i(SP_- + CP_+)$	$i(S\mathcal{R}_+ - C\mathcal{R}_-)$	0	0	$-i(C\mathcal{R}_+ + S\mathcal{R}_-)$	0
D^1	$-iCP_1$	$-iP_+$	0	$-iSD^1 + iCK^1$	0	$-iCJ^3$	$-iD^2$	$-iK^3$	$-iSJ^3$	$-iSP_1$	$-iC\mathcal{R}_1$	$-i\mathcal{R}_+$	0	$-iS\mathcal{R}_1$	0
D^2	$-iCP_2$	0	$-iP_+$	$-iSD^2 + iCK^2$	iCJ^3	0	iD^1	iSJ^3	$-iK^3$	$-iSP_2$	$-iC\mathcal{R}_2$	0	$-i\mathcal{R}_+$	$-iS\mathcal{R}_2$	0
J^3	0	iP_2	$-iP_1$	0	iD^2	$-iD^1$	0	iK^2	$-iK^1$	0	0	$i\mathcal{R}_2$	$-i\mathcal{R}_1$	0	0
K^1	$-iSP_1$	$-iP_-$	0	$iSK^1 + iCD^1$	iK^3	$-iSJ^3$	$-iK^2$	0	iCJ^3	iCP_1	$-iS\mathcal{R}_1$	$-i\mathcal{R}_-$	0	$iC\mathcal{R}_1$	0
K^2	$-iSP_2$	0	$-iP_-$	$iSK^2 + iCD^2$	iSJ^3	iK^3	iK^1	$-iCJ^3$	0	iCP_2	$-iS\mathcal{R}_2$	0	$-i\mathcal{R}_-$	$iC\mathcal{R}_2$	0
P_-	0	0	0	$i(SP_- + CP_+)$	iSP_1	iSP_2	0	$-iCP_1$	$-iCP_2$	0	$-2i(SD + K^3)$	$-2iK^1$	$-2iK^2$	$2iCD$	$-iP_-$
\mathcal{R}_+	$2iCD$	$-2iD^1$	$-2iD^2$	$-i(S\mathcal{R}_+ - C\mathcal{R}_-)$	$iC\mathcal{R}_1$	$iC\mathcal{R}_2$	0	$iS\mathcal{R}_1$	$iS\mathcal{R}_2$	$2i(SD + K^3)$	0	0	0	0	$i\mathcal{R}_+$
\mathcal{R}_1	$2iD^1$	$-2iD$	$-2iJ^3$	0	$i\mathcal{R}_+$	0	$-i\mathcal{R}_2$	$i\mathcal{R}_-$	0	$2iK^1$	0	0	0	0	$i\mathcal{R}_1$
\mathcal{R}_2	$2iK^2$	$2iJ^3$	$-2iD$	0	0	$i\mathcal{R}_+$	$i\mathcal{R}_1$	0	$i\mathcal{R}_-$	$2iK^2$	0	0	0	0	$i\mathcal{R}_2$
\mathcal{R}_-	$2i(SD - K^3)$	$-2iK^1$	$-2iK^2$	$i(C\mathcal{R}_+ + S\mathcal{R}_-)$	$iS\mathcal{R}_1$	$iS\mathcal{R}_2$	0	$-iC\mathcal{R}_1$	$-iC\mathcal{R}_2$	$-2iCD$	0	0	0	0	$i\mathcal{R}_-$
D	iP_+	iP_1	iP_2	0	0	0	0	0	0	iP_-	$-i\mathcal{R}_+$	$-i\mathcal{R}_1$	$-i\mathcal{R}_2$	$-i\mathcal{R}_-$	0

Conformal algebra in IFD

	P_0	P_1	P_2	K^3	$-K^1$	$-K^2$	J^3	$-J^2$	J^1	P_3	\mathfrak{K}_0	\mathfrak{K}_1	\mathfrak{K}_2	\mathfrak{K}_3	D
P_0	0	0	0	iP_3	iP_1	iP_2	0	0	0	0	$-2iD$	$2iK^1$	$-2iJ^1$	$-2iK^3$	$-iP_0$
P_1	0	0	0	0	iP_0	0	$-iP_2$	iP_3	0	0	$-2iK^1$	$2iD$	$-2iJ^3$	$-2iJ^2$	$-iP_1$
P_2	0	0	0	0	0	iP_0	iP_1	0	iP_3	0	$-2iK^2$	$2iJ^3$	$2iD$	$2iJ^1$	$-iP_2$
K^3	$-iP_3$	0	0	0	iJ^2	$-iJ^1$	0	iK^1	iK^2	$-iP_0$	$-i\mathfrak{K}_3$	0	0	$-i\mathfrak{K}_0$	0
$-K^1$	$-iP_1$	$-iP_0$	0	$-iJ^2$	0	$-iJ^3$	iK^2	$-iK^3$	0	0	$-i\mathfrak{K}_1$	$-i\mathfrak{K}_0$	0	0	0
$-K^2$	$-iP_2$	0	$-iP_0$	iJ^1	iJ^3	0	$-iK^1$	0	$-iK^3$	0	$-i\mathfrak{K}_2$	0	$-i\mathfrak{K}_0$	0	0
J^3	0	iP_2	$-iP_1$	0	$-iK^2$	iK^1	0	iJ^1	iJ^2	0	0	$i\mathfrak{K}_2$	$-i\mathfrak{K}_1$	0	0
$-J^2$	0	$-iP_3$	0	$-iK^1$	iK^3	0	$-iJ^1$	0	iJ^3	iP_1	0	$-i\mathfrak{K}_3$	0	$i\mathfrak{K}_1$	0
J^1	0	0	$-iP_3$	$-iK^2$	0	iK^3	$-iJ^2$	$-iJ^3$	0	iP_2	0	0	$-i\mathfrak{K}_3$	$i\mathfrak{K}_2$	0
P_3	0	0	0	iP_0	$i0$	$i0$	0	$-iP_1$	$-iP_2$	0	$-2iK^3$	$2iJ^2$	$-2iJ^1$	$2iD$	$-iP_3$
\mathfrak{K}_0	$2iD$	$2iK^1$	$2iK^2$	$i\mathfrak{K}_3$	$i\mathfrak{K}_1$	$i\mathfrak{K}_2$	0	0	0	$2iK^3$	0	0	0	0	$i\mathfrak{K}_0$
\mathfrak{K}_1	$-2iK^1$	$-2iD$	$-2iJ^3$	0	$i\mathfrak{K}_0$	0	$-i\mathfrak{K}_2$	$i\mathfrak{K}_3$	0	$-2iJ^2$	0	0	0	0	$i\mathfrak{K}_1$
\mathfrak{K}_2	$2iJ^1$	$2iJ^3$	$-2iD$	0	0	$i\mathfrak{K}_0$	$i\mathfrak{K}_1$	0	$i\mathfrak{K}_3$	$2iJ^1$	0	0	0	0	$i\mathfrak{K}_2$
\mathfrak{K}_3	$2iK^3$	$2iJ^2$	$-2iJ^1$	$i\mathfrak{K}_0$	0	0	0	$-i\mathfrak{K}_1$	$-i\mathfrak{K}_2$	$-2iD$	0	0	0	0	$i\mathfrak{K}_3$
D	iP_0	iP_1	iP_2	0	0	0	0	0	0	iP_3	$-i\mathfrak{K}_0$	$-i\mathfrak{K}_1$	$-i\mathfrak{K}_2$	$-i\mathfrak{K}_3$	0

Conformal algebra in LFD

	P_+	P_1	P_2	K^3	$-F^1$	$-F^2$	J^3	$-E^1$	$-E^2$	P_-	\mathfrak{R}_+	\mathfrak{R}_1	\mathfrak{R}_2	\mathfrak{R}_-	D
P_+	0	0	0	$-iP_+$	0	0	0	iP_1	iP_2	0	0	$2iF^1$	$2iE^2$	$-2i(D - K^3)$	$-iP_+$
P_1	0	0	0	0	iP_+	0	$-iP_2$	iP_-	0	0	$-2iF^1$	$2iD$	$-2iJ^3$	$-2iE^1$	$-iP_1$
P_2	0	0	0	0	0	iP_+	iP_1	0	iP_-	0	$-2iF^2$	$2iJ^3$	$2iD$	$-2iE^2$	$-iP_2$
K^3	iP_+	0	0	0	$-iF^1$	$-iF^2$	0	iE^1	iE^2	$-iP_-$	$i\mathfrak{R}_+$	0	0	$-i\mathfrak{R}_-$	0
$-F^1$	0	$-iP_+$	0	iF^1	0	0	iF^2	$-iK^3$	$-iJ^3$	$-iP_1$	0	$-i\mathfrak{R}_+$	0	$-i\mathfrak{R}_1$	0
$-F^2$	0	0	$-iP_+$	iF^2	0	0	$-iF^1$	iJ^3	$-iK^3$	$-iP_2$	0	0	$-i\mathfrak{R}_+$	$-i\mathfrak{R}_2$	0
J^3	0	iP_2	$-iP_1$	0	$-iF^2$	iF^1	0	$-iE^2$	iE^1	0	0	$i\mathfrak{R}_2$	$-i\mathfrak{R}_1$	0	0
$-E^1$	$-iP_1$	$-iP_-$	0	$-iE^1$	iK^3	$-iJ^3$	iE^2	0	0	0	$-i\mathfrak{R}_1$	$-i\mathfrak{R}_-$	0	0	0
$-E^2$	$-iP_2$	0	$-iP_-$	$-iE^2$	iJ^3	iK^3	$-iE^1$	0	0	0	$-i\mathfrak{R}_2$	0	$-i\mathfrak{R}_-$	0	0
P_-	0	0	0	iP_-	iP_1	iP_2	0	0	0	0	$-2i(D + K^3)$	$2iE^1$	$2iE^2$	0	$-iP_-$
\mathfrak{R}_+	0	$2iF^1$	$2iF^2$	$-i\mathfrak{R}_+$	0	0	0	$i\mathfrak{R}_1$	$i\mathfrak{R}_2$	$2i(D + K^3)$	0	0	0	0	$i\mathfrak{R}_+$
\mathfrak{R}_1	$-2iF^1$	$-2iD$	$-2iJ^3$	0	$i\mathfrak{R}_+$	0	$-i\mathfrak{R}_2$	$i\mathfrak{R}_-$	0	$-2iE^1$	0	0	0	0	$i\mathfrak{R}_1$
\mathfrak{R}_2	$-2iE^2$	$2iJ^3$	$-2iD$	0	0	$i\mathfrak{R}_+$	$i\mathfrak{R}_1$	0	$i\mathfrak{R}_-$	$-2iE^2$	0	0	0	0	$i\mathfrak{R}_2$
\mathfrak{R}_-	$2i(D - K^3)$	$2iE^1$	$2iE^2$	$i\mathfrak{R}_-$	$i\mathfrak{R}_1$	$i\mathfrak{R}_2$	0	0	0	0	0	0	0	0	$i\mathfrak{R}_-$
D	iP_+	iP_1	iP_2	0	0	0	0	0	0	iP_-	$-i\mathfrak{R}_+$	$-i\mathfrak{R}_1$	$-i\mathfrak{R}_2$	$-i\mathfrak{R}_-$	0

Kinematic and dynamic generators of the Conformal group

The generators of conformal transformations:

$$(\text{translation}) \quad P^{\hat{\mu}} = -i\partial^{\hat{\mu}} \quad (26)$$

$$(\text{dilation}) \quad D = -ix_{\hat{\mu}}\partial^{\hat{\mu}} \quad (27)$$

$$(\text{rotation}) \quad L^{\hat{\mu}\hat{\nu}} = i(x^{\hat{\mu}}\partial^{\hat{\nu}} - x^{\hat{\nu}}\partial^{\hat{\mu}}) \quad (28)$$

$$(\text{SCT}) \quad \mathfrak{K}^{\hat{\mu}} = -i(2x^{\hat{\mu}}x_{\hat{\nu}}\partial^{\hat{\nu}} - x^2\partial^{\hat{\mu}}) \quad (29)$$

Since $[\mathfrak{K}^{\hat{\tau}}, x^{\hat{\tau}}] = -i(2x^{\hat{\tau}}x^{\hat{\tau}} - (x^{\hat{\alpha}}.x_{\hat{\alpha}})\mathbb{C}) \rightarrow -i(x^0.x^0 + \vec{x}.\vec{x})$ as $\delta \rightarrow 0$, and $[\mathfrak{K}^{\hat{\tau}}, x^{\hat{\tau}}] = -i(2x^{\hat{\tau}}x^{\hat{\tau}} - (x^{\hat{\alpha}}.x_{\hat{\alpha}})\mathbb{C}) \rightarrow -i(2x^+.x^+)$ as $\delta \rightarrow \pi/4$, the conformal generator (LF time component) \mathfrak{K}_- is Kinematic in LFD, but Dynamic in IFD. And $[D, x^{\hat{\tau}}] = -ix^{\hat{\tau}}$, so D is always Kinematic in both IFD and LFD.

Interpolation angle	Kinematic	Dynamic
$\delta = 0$	$\mathcal{K}^{\hat{1}} = -J^2, \mathcal{K}^{\hat{2}} = J^1, J^3, P^1, P^2, P^3, D$	$\mathcal{D}^{\hat{1}} = -K^1, \mathcal{D}^{\hat{2}} = -K^2, K^3, P^0, \mathfrak{K}_0, \mathfrak{K}_1, \mathfrak{K}_2, \mathfrak{K}_3$
$0 \leq \delta < \pi/4$	$\mathcal{K}^{\hat{1}}, \mathcal{K}^{\hat{2}}, J^3, P^1, P^2, P_{\pm}, D$	$\mathcal{D}^{\hat{1}}, \mathcal{D}^{\hat{2}}, K^3, P_{\pm}, \mathfrak{K}_{\pm}, \mathfrak{K}_1, \mathfrak{K}_2, \mathfrak{K}_{\pm}$
$\delta = \pi/4$	$\mathcal{K}^{\hat{1}} = -E^1, \mathcal{K}^{\hat{2}} = -E^2, J^3, K^3, P^1, P^2, P_{\pm}, D, \mathfrak{K}_{\pm}$	$\mathcal{D}^{\hat{1}} = -F^1, \mathcal{D}^{\hat{2}} = -F^2, P_{\pm}, \mathfrak{K}_{\pm}, \mathfrak{K}_1, \mathfrak{K}_2$