

# Engineering Optics

## Lecture 33

31/05/2023

*by*

**Debolina Misra**

Assistant Professor in Physics  
IIITDM Kancheepuram, Chennai, India

# The Michelson interferometer

$S \rightarrow$  a light source (may be a sodium lamp)

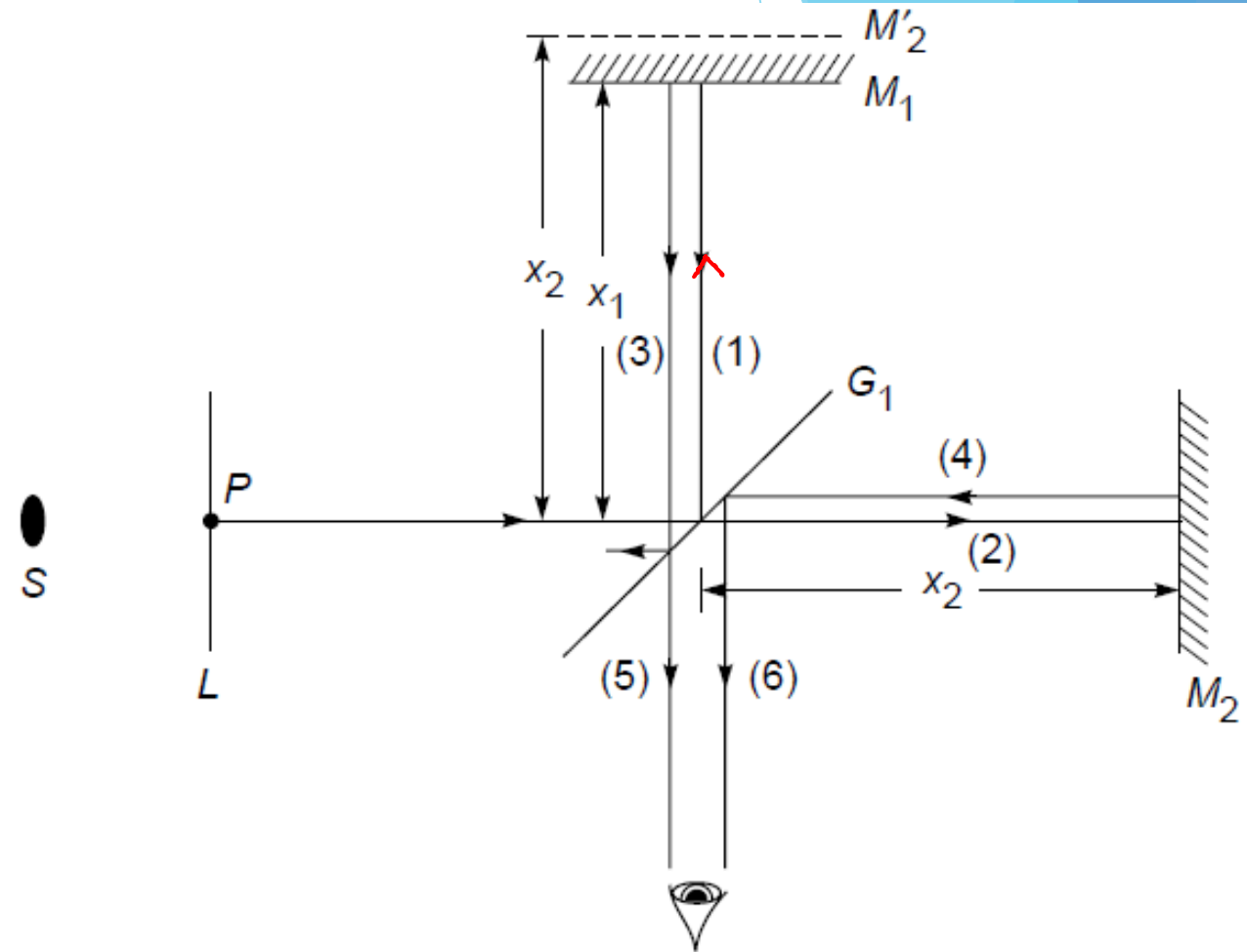
$L \rightarrow$  glass plate so that an extended source of almost uniform intensity is formed.

$G_1 \rightarrow$  a beam splitter  
a beam incident on  $G_1$  gets partially reflected and partially transmitted

$M_1$  and  $M_2 \rightarrow$  good-quality plane mirrors having very high reflectivity

One of the mirrors ( $M_2$ ) is fixed and the other (usually  $M_1$ ) is capable of moving away from or toward the glass plate  $G_1$  along an accurately machined track by means of a screw.

Usually mirrors  $M_1$  and  $M_2$  are perpendicular to each other and  $G_1$  is at  $45^\circ$  to the mirror.



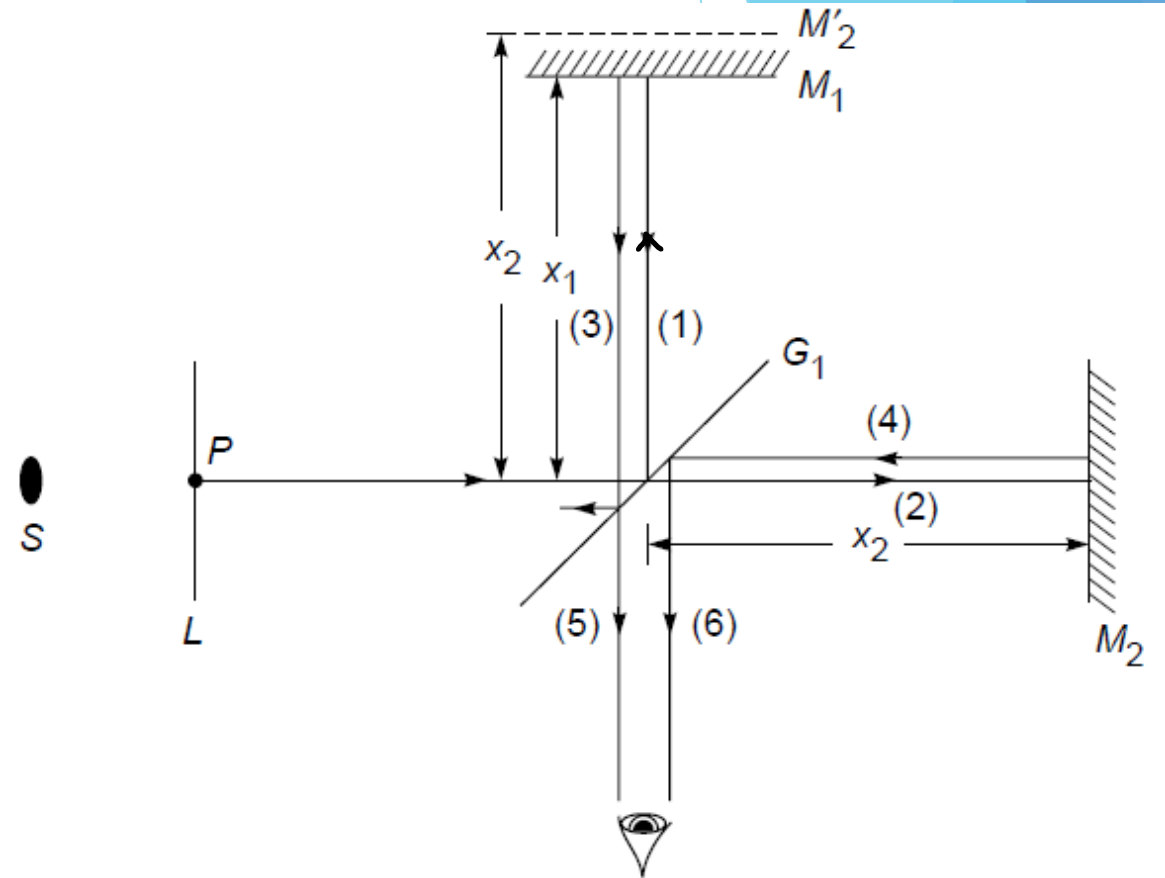
Schematic of the Michelson interferometer.

# The Michelson interferometer

if  $x_1$  and  $x_2$  are the distances of mirrors  $M_1$  and  $M_2$  from the plate  $G_1$ ,  $d = x_1 \sim x_2$

To the eye the waves emanating from point P will appear to get reflected by two parallel mirrors ( $M_1$  and  $M_2'$  — separated by a distance  $(x_1 \sim x_2)$ ).

if we use an extended source  $\rightarrow$  if we have a camera, then on the focal plane we will obtain circular fringes, each circle corresponding to a definite value of  $\theta$



Schematic of the Michelson interferometer.

# The Michelson interferometer

Now, if the beam splitter is just a simple glass plate, the beam reflected from mirror  $M_2$  will undergo an abrupt phase change of  $\pi$  (when getting reflected by the beam splitter), and since the extra path that one of the beams will traverse will be  $2(x_1 \sim x_2)$ , the condition for destructive interference will be

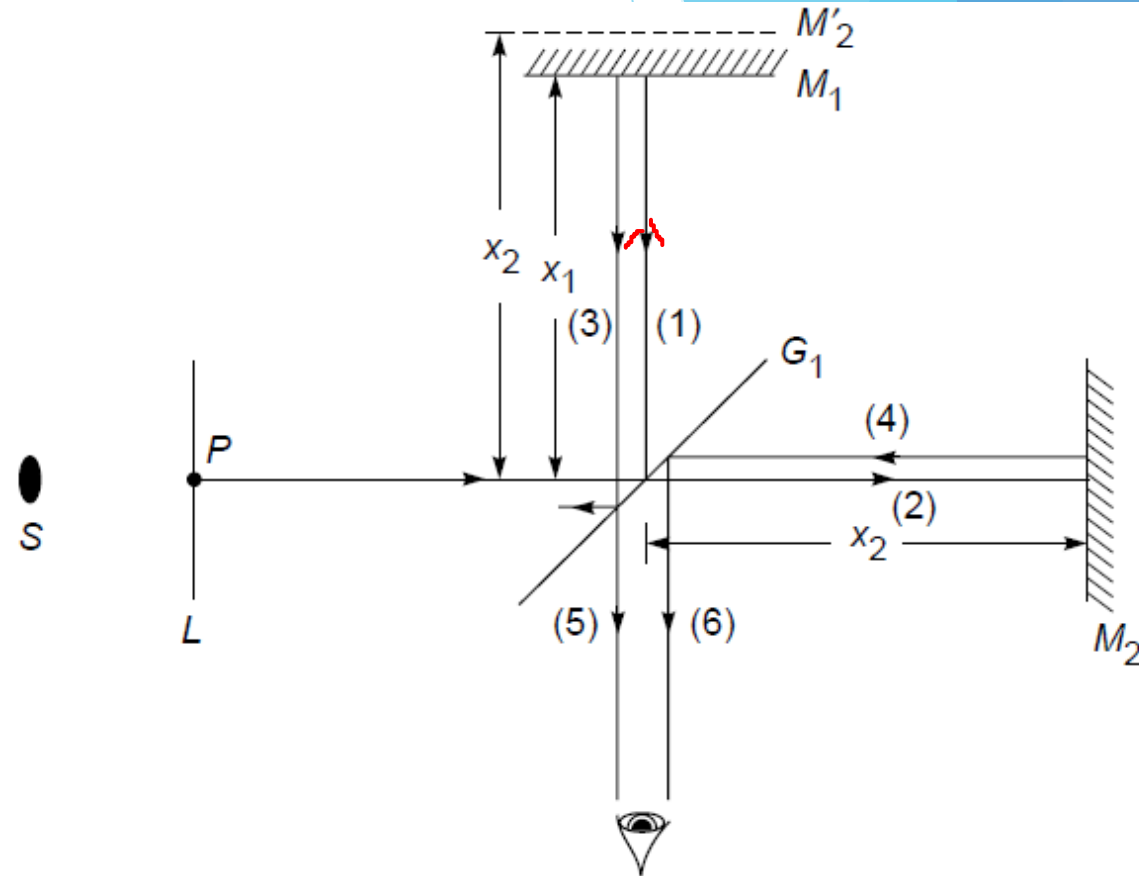
$$2d \cos \theta = m\lambda$$

where  $m = 0, 1, 2, 3, \dots$  and

$$d = x_1 \sim x_2$$

and the angle  $\theta$  represents the angle that the rays make with the axis (which is normal to the mirrors as shown in Fig. 15.35). Similarly, the condition for a bright ring is

$$2d \cos \theta = \left(m + \frac{1}{2}\right) \lambda$$



Schematic of the Michelson interferometer.

# The Michelson interferometer

Thus as we start reducing the value of  $d$ , the fringes will tend to collapse at the center

Conversely, if  $d$  is increased, the fringe pattern will expand.

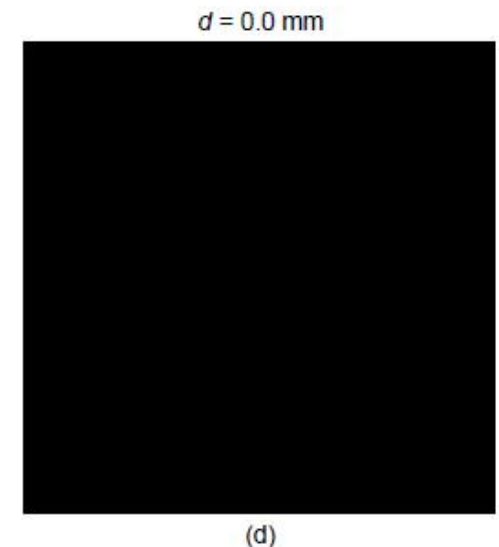
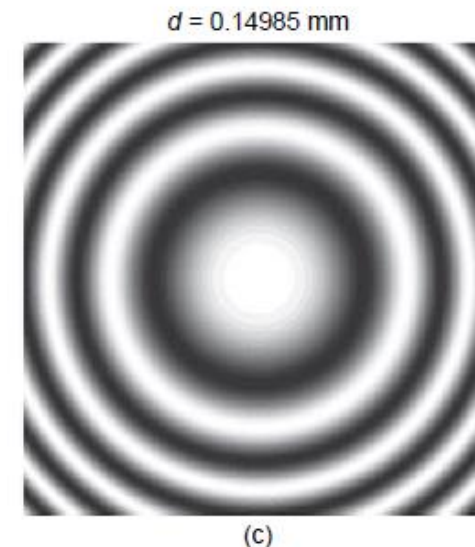
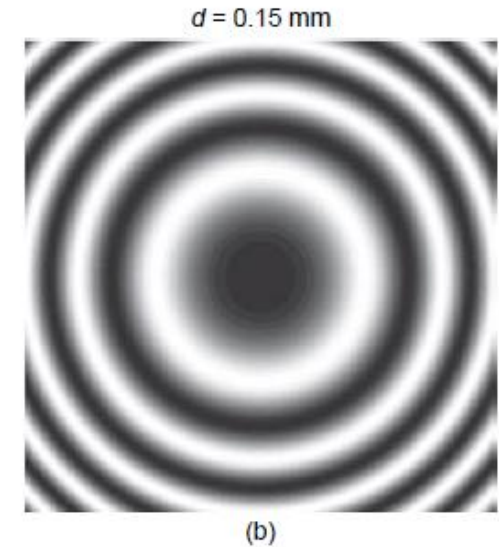
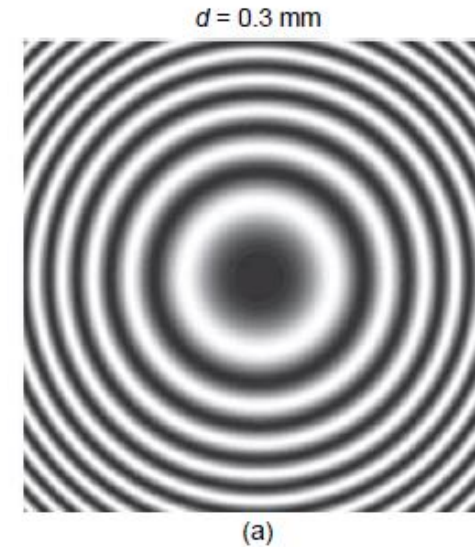
if  $N$  fringes collapse to the center as mirror  $M_1$  moves by a distance  $d_0$ , then we must have

$$2d = m\lambda$$

$$2(d - d_0) = (m - N)\lambda$$

where we have set  $\theta' = 0$  because we are looking at the central fringe. Thus

$$\lambda = \frac{2d_0}{N}$$



Computer-generated interference pattern produced by a Michelson interferometer.

# Wavelength measurement

This provides us with a method for the measurement of the wavelength. For example, in a typical experiment, if 1000 fringes collapse to the center as the mirror is moved through a distance of  $2.90 \times 10^{-2}$  cm, then

$$\lambda = 5800 \text{ \AA}$$

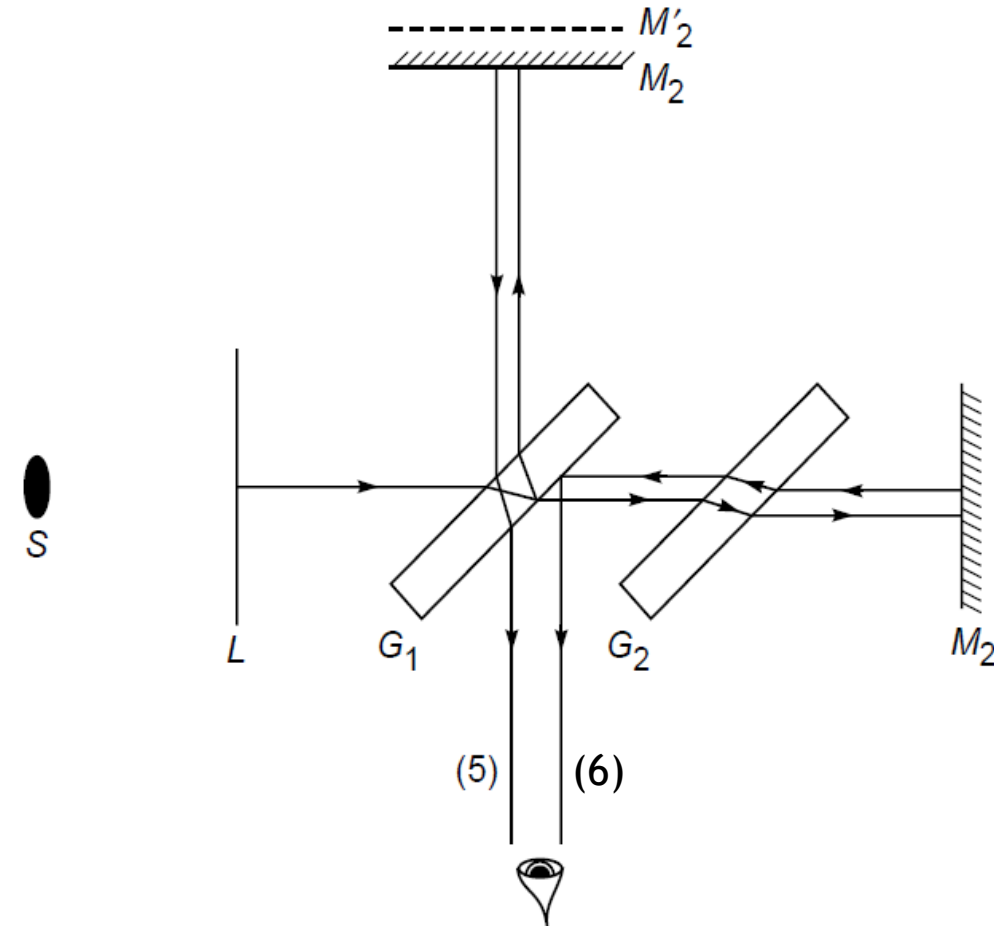
The above method was used by Michelson for the standardization of the meter. He found that the red cadmium line ( $\lambda = 6438.4696 \text{ \AA}$ ) is one of the ideal monochromatic sources, and as such this wavelength was used as a reference for the standardization of the meter.



In an actual Michelson interferometer, the beam splitter  $G_1$  consists of a plate (which may be about  $\frac{1}{2}$  cm thick), the back surface of which is partially silvered, and the reflections occur at the back surface as shown in Fig. 15.37. It is immediately obvious that beam 5 traverses the glass plate three times, and to compensate for this additional path, one introduces a “compensating plate”  $G_2$  which is exactly of the same thickness as  $G_1$ . The compensating plate is not really necessary for a monochromatic source because the additional path  $2(n - 1)t$  introduced by  $G_1$  can be compensated by moving mirror  $M_1$  by a distance  $(n - 1)t$ , where  $n$  is the refractive index of the material of the glass plate  $G_1$ .

# A point to note

- ▶ In an actual Michelson interferometer, the beam splitter  $G_1$  consists of a plate (which may be about 1/2 cm thick),
- ▶ The back surface of which is partially silvered, and the reflections occur at the back surface
- ▶ The compensating plate is not really necessary for a monochromatic source because the additional path introduced by  $G_1$  can be compensated by moving mirror  $M_1$
- ▶ *Q: How many time rays have crossed  $G_1$ ?*
- ▶ **What about white light??**
- ▶ Difficult to adjust  $M_1$  for each  $\lambda$
- ▶ Hence, compensating plate  $G_2$





# Application

- ▶ Can be used in the measurement of two closely spaced wavelengths:
- ▶ Sodium lamp → emits two closely spaced wavelengths 5890 and 5896 Å

If mirror  $M_1$  is moved away from (or toward) plate  $G_1$  through a distance  $d$ , then the maxima corresponding to the wavelength  $\lambda_1$  will not, in general, occur at the same angle as  $\lambda_2$ . Indeed, if the distance  $d$  is such that

$$\frac{2d}{\lambda_1} - \frac{2d}{\lambda_2} = \frac{1}{2}$$

and if  $2d \cos \theta' = m\lambda_1$ , then  $2d \cos \theta' = \left(m + \frac{1}{2}\right)\lambda_2$ . Thus, the maxima of  $\lambda_1$  will fall on the minima of  $\lambda_2$ , and conversely, and the fringe system will disappear.

# Application continued

$$\text{if } \frac{2d}{\lambda_1} - \frac{2d}{\lambda_2} = 1$$

then interference pattern will again reappear. In general, if

$$\frac{2d}{\lambda_1} - \frac{2d}{\lambda_2}$$

is  $1/2, 3/2, 5/2, \dots$ , we will have disappearance of the fringe pattern; and if it is equal to  $1, 2, 3, \dots$ , then the interference pattern will appear.

# Few points to note

- ▶ When the mirrors of the interferometer are inclined with respect to each other, making a small angle (i.e., when  $M_1$  and  $M_2$  are not quite perpendicular), ***Fizeau fringes*** are observed. The resultant wedge-shaped air film between  $M_2$  and  $M_1$  creates a pattern of straight parallel fringes.
- ▶ by appropriate adjustment of the orientation of the mirrors- $M_1$  and  $-M_2$ , fringes can be produced that are straight, circular, elliptical, parabolic, or hyperbolic—this holds as well for the real and virtual fringes.

# Experiment in brief

- ▶ <https://www.youtube.com/watch?v=j-u3IEgcTiQ>

## Problem:1

For a sodium lamp, the distance traversed by the mirror between two successive disappearances is 0.289 mm. Calculate the difference in the wavelengths of the  $D_1$  and  $D_2$  lines. Assume  $\lambda = 5890 \text{ \AA}$

Answer:

$$\frac{2d}{\lambda_1} - \frac{2d}{\lambda_2}$$

is  $1/2, 3/2, 5/2, \dots$ , we will have disappearance of the fringe pattern; and if it is equal to  $1, 2, 3, \dots$ , then the interference pattern will appear.

When the mirror moves through a distance 0.289 mm, the additional path introduced is 0.578 mm. Thus

$$\frac{0.578}{\lambda} - \frac{0.578}{\lambda + \Delta\lambda} = 1$$

$$\Delta\lambda = \frac{\lambda^2}{0.578} = \frac{(5890 \times 10^{-7})^2}{0.578} \text{ mm}$$
$$= 6 \text{ \AA}$$

Assume  $\Delta\lambda \times \lambda \ll \lambda^2$



## Problem:2

In the Michelson interferometer arrangement, if one of the mirrors is moved by a distance 0.08 mm, 250 fringes cross the field of view. Calculate the wavelength

## Answer:

The wavelength  $\lambda$  in Michelson interferometer is given by following equation.

$$\lambda = \frac{2d_0}{N}$$

Here,  $d_0$  is the distance moved by the mirror, and  $N$  is the number of fringes.

$$\begin{aligned}\lambda &= \frac{2(0.08 \text{ mm}) \left( \frac{1 \text{ cm}}{10 \text{ mm}} \right)}{250} \\ &= 6.4 \times 10^{-5} \text{ cm} \left( \frac{10^8 \text{ \AA}}{1 \text{ cm}} \right) \\ &= 6400 \text{ \AA}\end{aligned}$$

**Thank You**