

## Lecture 5: Poisson & Laplace equations, Part II

Reading: Stechmann Ch. 3.3-3.4

Last time:

$$\begin{aligned} -\nabla^2 u &= F \quad (\text{Poisson's equation}) \\ \nabla^2 u &= 0 \quad (\text{Laplace's equation}) \end{aligned}$$


Equilibrium heat / concentration ( $\overset{0}{u}_t = \kappa \nabla^2 u + F(\vec{x})$  or  $\overset{0}{\phi}_t = D \nabla^2 \phi + F(\vec{x})$ )

Electrostatics,  $\vec{E} = -\nabla \phi: -\nabla^2 \phi = \frac{\rho(\vec{x})}{\epsilon_0}$

Fundamental solution:  $\nabla^2 \Phi = \delta(\vec{x}) \quad \vec{x} \in \mathbb{R}^n$

$n=2$ . Polar coordinates  $(r, \theta)$ :  $\nabla^2 \Phi = \frac{1}{r} \partial_r (r \overset{\partial_r \Phi}{\Phi}_r) + \frac{1}{r^2} \Phi_{\theta\theta}$

Seek  $\Phi = \Phi(r)$ , find:  $\boxed{\Phi(r) = \frac{1}{2\pi} \log(r)}$  ( $r = |\vec{x}|$ )

$n=3$ . Spherical coordinates  $(r, \theta, \eta)$  

$$\nabla^2 \Phi = \frac{1}{r^2} \partial_r (r^2 \partial_r \Phi) + \frac{1}{r^2 \sin \theta} \partial_\theta (\sin \theta \partial_\theta \Phi) + \frac{1}{r^2 \sin^2 \theta} \partial_\eta^2 \Phi$$

Seek  $\Phi = \Phi(r)$ , find:  $\boxed{\Phi(r) = \frac{1}{4\pi r}}$  ( $r = |\vec{x}|$ )

(Check!  $\nabla^2 \left( \frac{1}{4\pi r} \right) = \frac{1}{r^2} \partial_r \left( r^2 \left( \frac{1}{4\pi r^2} \right) \right) = 0 \quad (r > 0)$ )

$\Phi$  solves: Poisson's equation with an impulsive force ( $\delta(\vec{x})$ )

and/or: Laplace's equation in  $\mathbb{R}^n \setminus \{0\}$ .

★ Solution to Poisson's equation,  $\nabla^2 u = f(\vec{x})$ :

$$u(\vec{x}) = f * \Phi = \int_{\mathbb{R}^n} f(\vec{y}) \Phi(\vec{x} - \vec{y}) dV_y$$

Why? Formally, 
$$\begin{aligned} \nabla^2 u &= \int_{\mathbb{R}^n} f(\vec{y}) \nabla^2 \Phi(\vec{x} - \vec{y}) dV_y \\ &= \int_{\mathbb{R}^n} f(\vec{y}) \delta(\vec{x} - \vec{y}) dV_y = f(\vec{x}) \end{aligned}$$

$$n=2: u(\vec{x}) = \frac{1}{2\pi} \int_{\mathbb{R}^2} f(\vec{y}) \log(|\vec{x} - \vec{y}|) dV_y \rightarrow dy_1 dy_2$$

$$n=3: u(\vec{x}) = \frac{1}{4\pi} \int_{\mathbb{R}^3} \frac{f(\vec{y})}{|\vec{x} - \vec{y}|} dV_y \rightarrow dy_1 dy_2 dy_3$$

Once again we are simply adding up solutions.  $(\int \leftrightarrow \sum)$

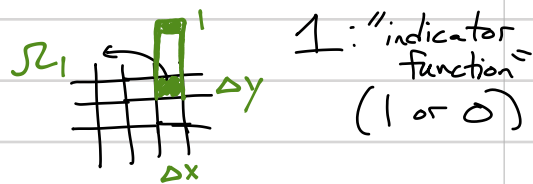
★ IF a PDE is linear and homogeneous (no forcing),

the sum of two solutions is a solution.

- Let  $\nabla^2 \Phi_1 = 0$  and  $\nabla^2 \Phi_2 = 0$ , then  $\nabla^2 (\Phi_1 + \Phi_2) = 0$

★ With forcing, it's almost as nice.

- $\vec{x} \in \mathbb{R}^2$ : Let  $\nabla^2 \Phi_1 = g_1 \mathbf{1}_{\{\vec{x} \in \Omega_1\}}$



or  $\nabla^2 \Phi_1 = g_1 \Delta x \Delta y \left( \frac{1}{\Delta x \Delta y} \mathbf{1}_{\{\vec{x} \in \Omega_1\}} \right)$

or  $\nabla^2 \Phi_1 = g_1 \Delta x \Delta y \delta_\Delta(\vec{x} - \vec{x}_1)$ , where  $\delta_\Delta(\vec{x}) = \begin{cases} \frac{1}{\Delta x \Delta y} & \vec{x} \in \Omega_1 \\ 0 & \vec{x} \notin \Omega_1 \end{cases}$

Also, let  $\nabla^2 \Phi_2 = g_2 \Delta x \Delta y \delta_\Delta(\vec{x} - \vec{x}_2)$

Then:  $\nabla^2 (\Phi_1 + \Phi_2) = (g_1 \delta_\Delta(\vec{x} - \vec{x}_1) + g_2 \delta_\Delta(\vec{x} - \vec{x}_2)) \Delta x \Delta y$

Adding more, as  $\Delta x, \Delta y \rightarrow 0$ ,  $\nabla^2 (\Phi_1 + \Phi_2 + \dots) = \int_{\mathbb{R}^2} g(\vec{y}) \delta(\vec{x} - \vec{y}) dx dy = g(\vec{x})$

Examples:

Electrostatic point charge.  $\rho(\vec{x}) = -\delta(\vec{x})$  (negatively charged particle)

$$\vec{x} \in \mathbb{R}^3. \quad \ominus \quad \phi?$$

$$\nabla^2 \phi = -\frac{\rho(\vec{x})}{\epsilon_0} = \frac{1}{\epsilon_0} \delta(\vec{x})$$

$$\text{So } \phi(\vec{x}) = \frac{1}{\epsilon_0} \Phi(\vec{x}) = \frac{-1}{4\pi\epsilon_0 |\vec{x}|} \text{ in 3D.}$$

$$(\text{use } \nabla^2 \Phi = \delta(\vec{x}) \text{ and linearity } (\frac{1}{\epsilon_0}) \times)$$

$$(\text{or } \phi(\vec{x}) = \int_{\mathbb{R}^3} (\frac{1}{\epsilon_0} \delta(\vec{y})) \frac{-1}{4\pi|\vec{x}-\vec{y}|} dV_y = \frac{-1}{4\pi\epsilon_0 |\vec{x}|})$$

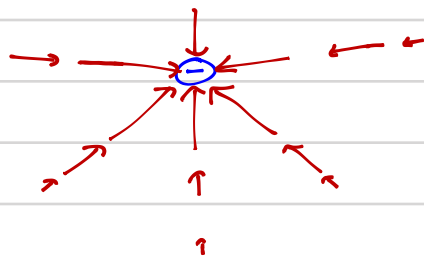
So  $\phi(\vec{x})$  grows as  $\frac{-1}{4\pi\epsilon_0 r}$  (potential energy)

$$\vec{E}\text{-field? } \vec{E} = -\nabla\phi$$

$$\textcircled{1} \quad \vec{E} = -\left(\hat{r}\partial_r + \frac{1}{r}\hat{\theta}\partial_\theta + \frac{1}{r\sin\theta}\hat{\varphi}\partial_\varphi\right)\left(\frac{-1}{4\pi\epsilon_0 r}\right)$$

$$= \hat{r} \left(\frac{-1}{4\pi\epsilon_0 r^2}\right)$$

$(\frac{1}{r^2} \text{ decay})$



$$\textcircled{2} \quad r = |\vec{x}| = \sqrt{x^2 + y^2 + z^2},$$

$$\vec{E} = -\nabla\phi = -(\hat{x}\partial_x + \hat{y}\partial_y + \hat{z}\partial_z)\left(\frac{-1}{4\pi\epsilon_0 |\vec{x}|}\right)$$

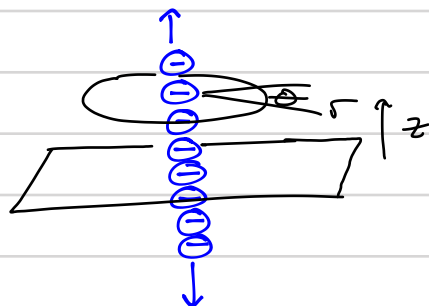
$$\text{Note: } \partial_x \left(\frac{1}{|\vec{x}|}\right) = \partial_x (x^2 + y^2 + z^2)^{-1/2} = -\frac{1}{2}(x^2 + y^2 + z^2)^{-3/2} \cdot 2x = \frac{-x}{|\vec{x}|^3}$$

$$\vec{E} = \hat{x} \left(\frac{-x}{4\pi\epsilon_0 |\vec{x}|^3}\right) + \hat{y} \left(\frac{-y}{4\pi\epsilon_0 |\vec{x}|^3}\right) + \hat{z} \left(\frac{-z}{4\pi\epsilon_0 |\vec{x}|^3}\right)$$

$$= \frac{-1}{4\pi\epsilon_0 |\vec{x}|^3} (x\hat{x} + y\hat{y} + z\hat{z})$$

$\underbrace{\hspace{1.5cm}}_{|\vec{x}|\hat{r}}$

• Line of charge:



Every cross-section is identical.

$$\nabla^2 \phi = -\frac{\rho(\vec{x})}{\epsilon_0}, \quad \rho(\vec{x}) = -\tilde{\rho}(r) \quad \left( \int_0^{2\pi} \int_0^R \tilde{\rho}(r) r \, dr \, d\theta = 1 \right)$$

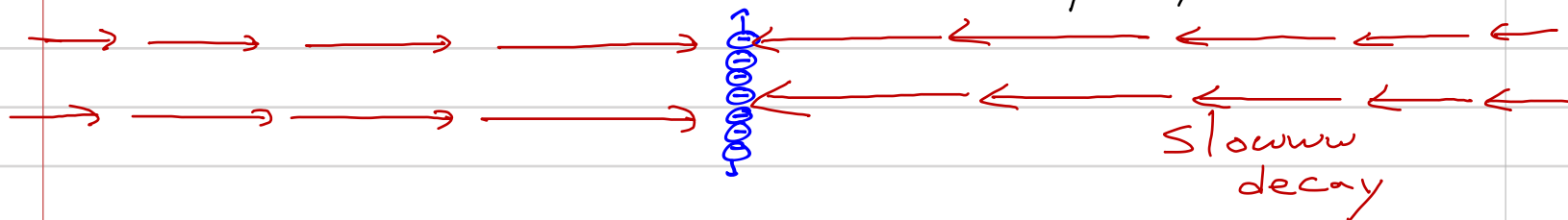
Seek a 2D solution,  $\phi(r, \theta, z) = \phi(r, \theta) = \phi(r)$

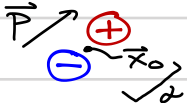
Polar coordinates:  $\frac{1}{r} \partial_r \left( \frac{1}{r} \partial_r \right) = \frac{1}{\epsilon_0} \tilde{\rho}(r)$

$$\phi = \frac{1}{2\pi\epsilon_0} \log(r)$$

$$\vec{E} = -\nabla \phi = -\left( \hat{r} \partial_r + \frac{1}{r} \hat{\theta} \partial_\theta + \hat{z} \partial_z \right) \phi$$

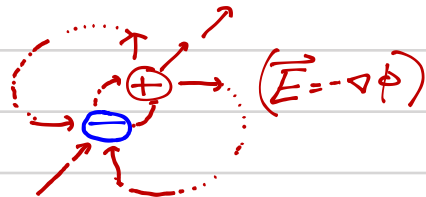
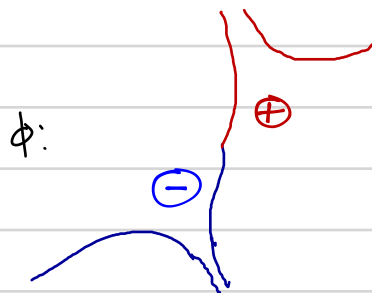
$$= -\hat{r} \left( \frac{1}{2\pi\epsilon_0 r} \right) \quad \text{Only decays like } \frac{1}{r}.$$



★ Dipoles  $\vec{p}$    $|\vec{p}| = 1$

$$\vec{x} \in \mathbb{R}^3: \quad \rho(\vec{x}) = \frac{1}{\alpha} \left( \delta(\vec{x} - [\vec{x}_0 + \frac{\alpha}{2} \vec{p}]) - \delta(\vec{x} - [\vec{x}_0 - \frac{\alpha}{2} \vec{p}]) \right)$$

$$\phi = \frac{1/\alpha}{4\pi\epsilon_0 |\vec{x} - (\vec{x}_0 + \frac{\alpha}{2} \vec{p})|} - \frac{1/\alpha}{4\pi\epsilon_0 |\vec{x} - (\vec{x}_0 - \frac{\alpha}{2} \vec{p})|}$$



as  $\alpha \rightarrow 0$ , Taylor expand:

$$\phi = \frac{1/\alpha}{4\pi\epsilon_0} \left[ \frac{1}{|\vec{x} - \vec{x}_0|} + \frac{\alpha}{2} \vec{p} \cdot \nabla_0 \left( \frac{1}{|\vec{x} - \vec{x}_0|} \right) + O(\alpha^2) \right]$$

$$- \frac{1}{4\pi\epsilon_0} \left[ \frac{1}{|\vec{x} - \vec{x}_0|} - \frac{\alpha}{2} \vec{p} \cdot \nabla_0 \left( \frac{1}{|\vec{x} - \vec{x}_0|} \right) + O(\alpha^2) \right]$$

$$\xrightarrow{\alpha \rightarrow 0} \frac{1}{4\pi\epsilon_0} \vec{p} \cdot \nabla_0 \left( \frac{1}{|\vec{x} - \vec{x}_0|} \right) = \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot (\vec{x} - \vec{x}_0)}{|\vec{x} - \vec{x}_0|^3}$$

$$(\sim \frac{1}{r^2})$$