

Let

$$X = a_0 \mathbb{I} + \sum_{k=1}^3 \sigma_k a_k, \quad (1)$$

where  $\sigma_k$  are pauli matrices and  $a_0, a_1$  are numbers.

**1.  $\text{Tr } X$  and  $\text{Tr}(\sigma_k X)$ .**

Noticing pauli matrices are traceless:

$$\text{Tr}(X) = \text{Tr}(a_0 \mathbb{I}) + \text{Tr}\left(\sum_k \sigma_k a_k\right) = 2a_0. \quad (2)$$

Then consider

$$\begin{aligned} \text{Tr}(\sigma_k X) &= \text{Tr}(a_0 \sigma_k \mathbb{I}) + \text{Tr}\left(\sum_j \sigma_k \sigma_j a_j\right) \\ &= 0 + \sum_j a_j \text{Tr}(\sigma_j \sigma_k). \end{aligned} \quad (3)$$

Using the fact that  $\text{Tr}(\sigma_j \sigma_k) = 2\delta_{jk}$  :

$$\text{Tr}(\sigma_k X) = \sum_{k=1}^3 a_k \cdot 2\delta_{jk} = 2a_k. \quad (4)$$

**2. Find  $a_0, a_k$  w.r.t.  $X_{ij}$ .**

Write

$$X_{ij} = a_0 \delta_{ij} + \sum_{k=1}^3 \left( (\sigma_k)_{ij} a_k \right). \quad (5)$$

Then for diagonal elements:

$$\begin{aligned} X_{ii} &= a_0 + (\sigma_3)_{ii} a_3, \\ \Rightarrow X_{11} &= a_0 + a_3, \quad X_{22} = a_0 - a_3. \end{aligned} \quad (6)$$

Off diagonal elements: for  $m \neq n$  :

$$\begin{aligned} X_{mn} &= \sigma_{1_{mn}} a_1 + \sigma_{2_{mn}} a_2 \\ \Rightarrow X_{12} &= a_1 - ia_2, \quad X_{21} = a_1 + ia_2. \end{aligned} \quad (7)$$

## P2.

Consider a ket space spanned by  $\{|a'\rangle\}$  of Hermitian operator  $\mathbf{A}$ . No degeneracy.

### 1. $\Pi_{a'}(\mathbf{A} - a')$ is the null operator.

Proof: Since  $\{|a'\rangle\}$  spans the space, it's sufficient to show that  $\Pi_{a'}(\mathbf{A} - a')$  annihilates all basisket. To show, consider arbitrary basis ket  $|a''\rangle$ .

$$\begin{aligned} \left[ \prod_{a'} (\mathbf{A} - a') \right] |a''\rangle &= \dots \times (\mathbf{A} - a'') |a''\rangle \times \dots \\ &= \mathbf{A} |a''\rangle - a'' |a''\rangle \\ &= a'' |a''\rangle - a'' |a''\rangle = 0, \end{aligned} \quad (8)$$

and thus  $\Pi_{a'}(\mathbf{A} - a')$  annihilates all basis kets, and therefore nullify all vectors in this space.

### 2.

Let

$$\hat{P} \equiv \prod_{a'' \neq a'} \frac{\mathbf{A} - a''}{a' - a''}. \quad (9)$$

Notice that  $\hat{P}|a'\rangle = |a'\rangle$ ,  $\hat{P}|a_k\rangle = 0$  (for any  $a_k \neq a'$ ). Explicitly:

$$\hat{P}|a'\rangle = \frac{\mathbf{A} - a''}{a' - a''} |a'\rangle = \frac{a' - a''}{a' - a''} |a'\rangle = |a'\rangle; \quad (10)$$

and for any  $a_k \neq a'$ ,

$$\hat{P}|a_k\rangle = \dots \times \frac{\mathbf{A} - a_k}{a' - a_k} \times \dots |a_k\rangle = 0. \quad (11)$$

It's clear that  $\hat{P}$  is a projection operator onto the eigenspace corresponding to the eigenvalue  $a'$ .

### 3. Illustrate both results using $\mathbf{A} = S_z$ for spin 1/2 system.

Recall that the eigenkets  $\{|+\rangle, |-\rangle\}$  of the Hermitian operator  $S_z$  form an orthonormal basis, just like  $\{|a'\rangle\}$  in the assumption.

$$S_z |\pm\rangle = \pm \frac{\hbar}{2} |\pm\rangle, \quad (12)$$

and from which we can observe:

$$\prod_{a'} (S_z - a') |\pm\rangle = \left( S_z - \frac{\hbar}{2} \right) \left( S_z + \frac{\hbar}{2} \right) |\pm\rangle = 0. \quad (13)$$

Further,

$$\hat{P} = \prod_{a'' \neq a'} \frac{\mathbf{A} - a''}{a' - a''} = \begin{cases} \frac{S_z + \frac{\hbar}{2}}{\frac{\hbar}{2} + \frac{\hbar}{2}} & \text{for } a' = -\frac{\hbar}{2} \\ \frac{S_z - \frac{\hbar}{2}}{-\frac{\hbar}{2} - \frac{\hbar}{2}} & \text{for } a' = \frac{\hbar}{2} \end{cases}, \quad (14)$$

and thus for  $a' = -\frac{\hbar}{2}$ :  $\hat{P}|+\rangle = 0$ ,  $\hat{P}|-\rangle = 1$ ; for  $a' = \frac{\hbar}{2}$ :  $\hat{P}|-\rangle = 0$ ,  $\hat{P}|+\rangle = 1$ , as expected from part 2.

### P3

Construct  $|\mathbf{S} \cdot \hat{\mathbf{n}}; +\rangle$  as a linear combination of  $|+\rangle, |-\rangle$ , s.t.

$$\mathbf{S} \cdot \hat{\mathbf{n}} |\mathbf{S} \cdot \hat{\mathbf{n}}; +\rangle = \frac{\hbar}{2} |\mathbf{S} \cdot \hat{\mathbf{n}}; +\rangle, \quad (15)$$

We recall that, in spherical coordinates,

$$\hat{\mathbf{n}} = (\sin \beta \cos \alpha, \sin \beta \sin \alpha, \cos \beta); \quad (16)$$

and the definition of  $\mathbf{S}$  :

$$S_x = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, S_y = \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, S_z = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}. \quad (17)$$

Then,

$$\begin{aligned} \mathbf{S} \cdot \hat{\mathbf{n}} &= S_x n_x + S_y n_y + S_z n_z \\ &= \frac{\hbar}{2} \begin{pmatrix} \cos \beta & \sin \beta \cos \alpha - i \sin \beta \sin \alpha \\ \sin \beta \cos \alpha + i \sin \beta \sin \alpha & -\cos \beta \end{pmatrix} \\ &= \frac{\hbar}{2} \begin{pmatrix} \cos \beta & \sin \beta e^{-i\alpha} \\ \sin \beta e^{i\alpha} & -\cos \beta \end{pmatrix}. \end{aligned} \quad (18)$$

Let eigenket  $|\mathbf{S} \cdot \hat{\mathbf{n}}; +\rangle = \begin{pmatrix} c_1 \\ c_2 \end{pmatrix}$ . The eigenvalue equation then reads:

$$\begin{aligned} \frac{\hbar}{2} \begin{pmatrix} \cos \beta & \sin \beta e^{-i\alpha} \\ \sin \beta e^{i\alpha} & -\cos \beta \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} &= \frac{\hbar}{2} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} \\ \Rightarrow \begin{pmatrix} c_1 \cos \beta + c_2 \sin \beta e^{-i\alpha} \\ c_1 \sin \beta e^{i\alpha} - c_2 \cos \beta \end{pmatrix} &= \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} \end{aligned} \quad (19)$$

This yields two dependent equations

$$\begin{aligned} c_1(\cos \beta - 1) + c_2 \sin \beta e^{-i\alpha} &= 0, \\ c_1 \sin \beta e^{i\alpha} - c_2(\cos \beta + 1) &= 0. \end{aligned} \quad (20)$$

from which,

$$c_1 = \frac{\sin \beta e^{-i\alpha}}{1 - \cos \beta} c_2 \Rightarrow c_1 = \frac{\cos\left(\frac{\beta}{2}\right) e^{-i\alpha}}{\sin\left(\frac{\beta}{2}\right)} c_2. \quad (21)$$

Normalization condition for the eig solution requires

$$\begin{aligned} |c_1|^2 + |c_2|^2 &= 1 \\ \Rightarrow |c_2|^2 \left( \frac{\cos^2\left(\frac{\beta}{2}\right)}{\sin^2\left(\frac{\beta}{2}\right)} + 1 \right) &= 1 \\ \Rightarrow |c_2|^2 &= \sin^2\left(\frac{\beta}{2}\right). \end{aligned} \quad (22)$$

Then for an arbitrary azimuthal angle  $\alpha$ , we have

$$c_2 = \sin\left(\frac{\beta}{2}\right) e^{i\alpha}, c_1 = \cos\left(\frac{\beta}{2}\right). \quad (23)$$

Thus the eigenket is found to be

$$\cos\left(\frac{\beta}{2}\right) |+\rangle + \sin\left(\frac{\beta}{2}\right) e^{i\alpha} |-\rangle. \quad (24)$$

**P4.**

A two level system with Hamiltonian

$$H = H_{11}|1\rangle\langle 1| + H_{22}|2\rangle\langle 2| + H_{12}(|1\rangle\langle 2| + |2\rangle\langle 1|). \quad (25)$$

Find energy eigenvalues and eigenkets.

**1. method 1, solve eigenvalue problem Explicitely.**

**2. method 2, use result from P3.**