Homework sheet 2 - Due 02/16/2025

Comment: Part of this exercise includes functions of matrices. These are defined by the Taylor series of the corresponding function, e.g. $e^M = \sum_{k=0}^{M^k} \frac{M^k}{k!}$ for any matrix M.

Problem 1: Matrix Operations [1 + 2 + 1 + 2 + 2 + 1 + 1 = 10 points]

In this exercise we prove some useful matrix identities.

a) For matrices A, B, C, prove

$$[A, BC] = B[A, C] + [A, B]C.$$
 (1)

b) Prove the Bianchi identity for matrices A, B, C

$$[A, [B, C]] + [B, [C, A]] + [C, [A, B]] = 0.$$
(2)

c) Prove that

$$[A, B]^{\dagger} = -[A^{\dagger}, B^{\dagger}] \text{ and } [A, B]^{T} = -[A^{T}, B^{T}].$$
 (3)

d) Consider two matrices A, B such that [A, B] = C where [A, C] = 0 = [B, C]Prove

$$e^{\alpha A}Be^{-\alpha A} = B + \alpha C \tag{4}$$

for arbitrary $\alpha \in \mathbb{C}$.

e) Campbell-Baker-Hausdorff formula. Consider matrices A, B, C with the same properties as in the previous problem. Show

$$e^A e^B = e^{A+B+C/2}. (5)$$

Hint: Define a function $T(\alpha) = e^{\alpha A}e^{\alpha B}$ and first study its α -derivative. Use the result from part d).

f) For an involutory matrix A (i.e. $A^2 = 1$), prove

$$e^{i\alpha A} = \cos(\alpha)\mathbf{1} + i\sin(\alpha)A$$
 (6)

g) Prove that, for any diagonalizable matrix M

$$\ln(\det(M)) = \operatorname{tr}(\ln(M)). \tag{7}$$

Problem 2: Single-qubit gates [1 + 1 + 1 + 2 + 2 + 2 + 1 = 10 points]

In quantum information theory it is common practice to denote the Pauli gates as

$$X = \sigma_x, Y = \sigma_y, Z = \sigma_z. \tag{8}$$

Rotations about the x-axis are denoted $R_x(\theta_x) = e^{-i\theta_x X/2}$ (and analogously for y and z).

- a) Show that, up to a phase, the $\pi/8$ gate $T = R_z(\pi/4)$.
- b) Show that, up to a phase, the Hadamard gate is a concatenation of $R_y(\theta_y)$ and $R_z(\theta_z)$. Determine the angles θ_y, θ_z .
- c) Show that, since YXY = -X,

$$YR_X(\theta_x)Y = R_X(-\theta_x),$$

- (i.e. the direction of rotation is reversed by Y).
- d) Show the following identities for single-qubit gates

$$HXH = Z, HYH = -Y, HZH = X. (9)$$

e) Show that (up to a phase)

$$HTH = R_x(\pi/4). \tag{10}$$

- f) Calculate eigenstates of X, Y along with corresponding eigenvalues.
- g) Show that the eigenstate of X with eigenvalue +1 can be obtained by applying $R_y(\pi/2)$ on $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$. Illustrate this statement on the Bloch sphere.

Problem 3: Higher spin systems. Spin-1 systems [2 + 2 + 1 + 3 + 2 = 10 points]

a) In the lectures we found out that $\hat{J}_{\pm} | j, m \rangle = \hbar c_{\pm} | j, m \pm 1 \rangle$, but did not determine the constant c_{\pm} . Calculate c_{\pm} (we assume $c_{\pm} > 0$) when all $|j, m\rangle$ are orthonormalized.

- b) For spin-j systems consider the normalized magnetizations $\hat{m}_i = \hat{J}_i/[\hbar j]$ and calculate the Heisenberg-bound on the combined uncertainty of \hat{m}_x and \hat{m}_y . Why is $j \to \infty$ sometimes called the semiclassical limit?
- c) Explicitly present \hat{J}_x , \hat{J}_y , \hat{J}_z for spin-1 systems (in the basis where \hat{J}_z is diagonal).

Comment: Exemplary spin-1 systems of relevance in atomic physics are ⁸⁷Rb and ²³Na (their groundstate forms a "hyperfine" triplet), which were the first cold atomic gases to display Bose-Einstein condensation (Nobel prize 2001).

d) Convince yourself that spin nematicity operators (i.e. magnetic quadrupole operators)

$$\hat{N}_{ij} = \frac{1}{2} \{ \hat{J}_i, \hat{J}_j \} - \frac{1}{3} \delta_{ij} \hat{\vec{J}}^2, \quad i, j = x, y, z.$$
 (11)

vanish for spin-1/2 systems but they do exist for spin-1 systems. How many non-trivial \hat{N}_{ij} are there for spin-1? Calculate them explicitly.

e) Which \hat{N}_{jk} and \hat{J}_i are compatible? Calculate $[\hat{J}_i, \hat{N}_{jk}]$ to find out.

Problem 4: Lie Algebra for special unitary group SU(N)[2+3+2+3+2+3=10 points + 5 bonus points.]

For a Lie group G, the Lie algebra \mathfrak{g} is given by the d_G -dimensional real vector space of generators λ_a of the group supplemented with the Lie-bracket

$$[\lambda_a, \lambda_b] = i \sum_c f_{abc} \lambda_c. \tag{12}$$

The exponential map relates Lie algebra and Lie group

$$\exp: \mathfrak{g} \to G, \alpha \mapsto e^{i\alpha}, \tag{13}$$

In this exercise we consider the special unitary group G = SU(N) of $N \times N$ unitary matrices U with unit determinant $\det(U) = 1$. Elements of the Lie algebra $\alpha \in \mathfrak{su}(N)$ are $N \times N$ matrices, the Lie bracket is just the matrix commutator and the exponential map is just the matrix exponential. This is called the "fundamental representation" of the Lie algebra.

a) Prove that $\mathfrak{su}(N)$ is spanned by traceless, Hermitian matrices.

Hint: *Eq.* (7).

The Lie algebra can be equipped with an inner product

$$\langle \alpha, \beta \rangle = \frac{1}{2} \operatorname{tr} \left[\alpha \beta \right],$$
 (14)

and we assume the $\{\lambda_a\}_{a=1}^{d_G}$ to be orthonormal with respect to this inner product, hence elements $\alpha \in \mathfrak{su}(N)$ can be expanded as $\alpha = \sum_{a=1}^{d_{SU(N)}} \alpha_a \lambda_a$.

- b) Prove that $d_{SU(N)} = N^2 1$ and use the results from homework sheet 1 to convince yourself that the Pauli matrices form an orthonormal basis for the fundamental representation of $\mathfrak{su}(2)$. Which physical spin does the fundamental representation of SU(2) correspond to?
- c) For general N, use the orthonormal basis of the fundamental representation to show that the "structure factors" f_{abc} are real, totally antisymmetric tensors.

A "faithful representation" of a Lie algebra is an injective map $D: \alpha \mapsto D(\alpha)$, where $\alpha \in \mathfrak{g}$ and $D(\alpha)$ is a $d_{D(\mathfrak{g})} \times d_{D(\mathfrak{g})}$ dimensional matrix and the matrices $\{D(\lambda_a)\}_{a=1}^{d_G}$ fulfill the same Lie algebra as $\{\lambda_a\}_{a=1}^{d_G}$

$$[D(\lambda_a), D(\lambda_b)] = i \sum_c f_{abc} D(\lambda_c). \tag{15}$$

- d) Use the Bianchi identity to show that the $d_G \times d_G$ matrices $[T_a]_{bc} = -if_{abc}$ fulfill the Lie algebra (they form the "adjoint representation" $D(\lambda_a) = T_a$).
- e) Write down the matrices of the adjoint representation of SU(2). Diagonalize one of the matrices. Which spin does this representation correspond to?
- f) Write down the matrices of the orthonormal basis for the fundamental representation of SU(3) and determine a set of matrices which form an SU(2) sub-algebra.
- g) In the energy window $E \lesssim 900 MeV$ only three quarks are relevant for quantum chromodynamics. They are distinguished by their flavor quantum number: up $(|u\rangle)$, down $(|d\rangle)$, strange $(|s\rangle)$ with an approximate SU(3) symmetry between them.

Comment: Before this energy range was reached, aspects of particle physics could be understood by means of Heisenberg's SU(2) isospin, essentially acting in $|u\rangle$, $|d\rangle$ space. Once experiments surpassed the energy of the strange-quark rest mass $\sim 95 MeV/c^2$, SU(2) isospin flavor symmetry had to be extended to SU(3).

- i) Discuss the dimension of fundamental and adjoint SU(3) representations
- ii) Based on the newly acquired knowledge on Lie algebras, explain the appearance of an octet of mesons (=quark-antiquark boundstates) for $E \lesssim 900 MeV$. Why is there only a triplet for $E \lesssim 200 MeV$?