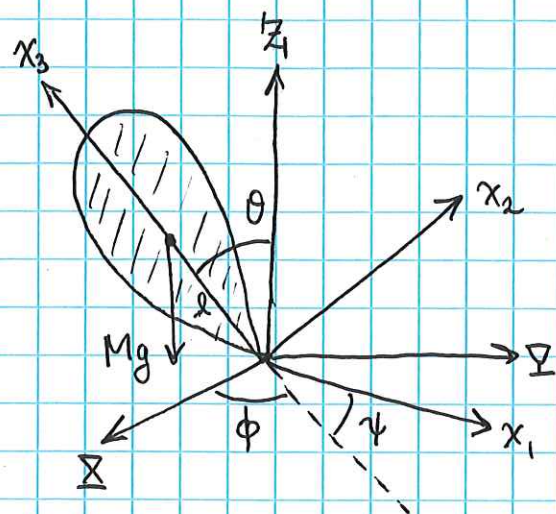


Summary

①
04/08/24

- Heavy symmetric top - sol.ⁿ w/ Euler angles.



$$L = T - U$$

$$= \frac{1}{2} I_1' (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta) + \frac{1}{2} I_3 (\dot{\psi} + \dot{\phi} \cos \theta)^2 - Mgl \cos \theta.$$

E-L eqn.^s:

$$\cdot \phi: \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{\phi}} = \frac{\partial L}{\partial \phi} \Rightarrow L_2 = (I_1' \sin^2 \theta + I_3 \cos^2 \theta) \dot{\phi} + I_3 \cos \theta \dot{\psi} = \text{const.}$$

$$\cdot \psi: \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{\psi}} = \frac{\partial L}{\partial \psi} \Rightarrow L_3 = I_3 (\dot{\psi} + \dot{\phi} \cos \theta) = \text{const.}$$

• Instead of using θ EOM, recall that time-translation symmetry of $L \Rightarrow$ energy = third conserved quantity.

$$E = T + U = \frac{1}{2} I_1' (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta) + \frac{1}{2} I_3 (\dot{\psi} + \dot{\phi} \cos \theta)^2 + Mgl \cos \theta.$$

$\dot{\phi}$ & $\dot{\psi}$ eliminated using L_2 & L_3 :

$$\begin{cases} L_2 = I_1' \sin^2 \theta \dot{\phi} + L_3 \cos \theta & \Rightarrow \dot{\phi} = \frac{L_2 - L_3 \cos \theta}{I_1' \sin^2 \theta} \\ L_3 = I_3 (\dot{\psi} + \dot{\phi} \cos \theta) & \Rightarrow \dot{\psi} = \frac{L_3}{I_3} - \frac{\cos \theta (L_2 - L_3 \cos \theta)}{I_1' \sin^2 \theta} \end{cases}$$

plugging into E:

②
04/08/24

$$E = \frac{1}{2} I_1' \dot{\theta}^2 + \frac{1}{2 I_1'} \frac{(L_2 - L_3 \cos \theta)^2}{\sin^2 \theta} + \frac{1}{2 I_3} L_3^2 + Mgl \cos \theta.$$

$$\Rightarrow E - \underbrace{\frac{L_3^2}{2 I_3} - Mgl}_{E'} = \frac{1}{2} I_1' \dot{\theta}^2 + \underbrace{\frac{1}{2 I_1'} \frac{(L_2 - L_3 \cos \theta)^2}{\sin^2 \theta} - Mgl (1 - \cos \theta)}_{U_{\text{eff}}(\theta)}.$$

$$\Rightarrow E' = \frac{1}{2} I_1' \dot{\theta}^2 + U_{\text{eff}}(\theta).$$

→ effective 1 DOF problem (in θ).

as before, this can be integrated.

$$\dot{\theta} = \frac{d\theta}{dt} \Rightarrow t = \int \frac{d\theta}{\sqrt{2[E - U_{\text{eff}}(\theta)]/I_1'}} \quad (0 \leq \theta \leq \pi).$$

→ full sol'n of problem.

invert $t(\theta) \rightarrow \theta(t)$ & plug into
(*) + (**) for $\phi(t)$ & $\psi(t)$.

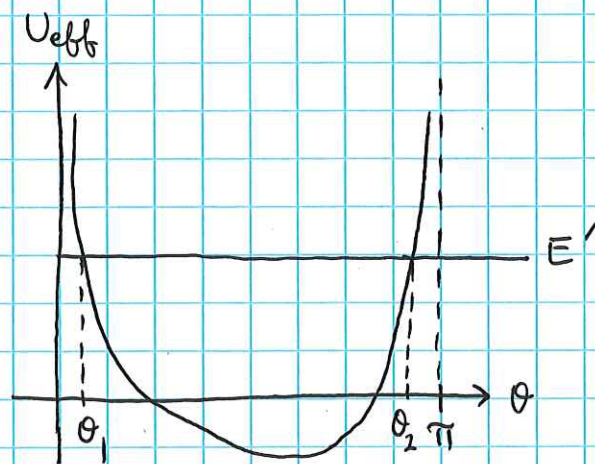
• we can learn much about the motion by considering prop.'s of $U_{\text{eff}}(\theta)$.

• First, note $U_{\text{eff}} \rightarrow +\infty$ as $\theta \rightarrow 0$ or π .

• It can also be shown $U_{\text{eff}}(\theta)$ has a single minimum for $0 \leq \theta \leq \pi$ (see end of notes for proof).

⇒ U_{eff} has the following generic form:

③
04/08/24



• For given value of E' ,
there are two turning pt.'s:

$$U_{\text{eff}}(\theta_{1,2}) = E' \quad (\dot{\theta} = 0).$$

⇒ motion restricted:

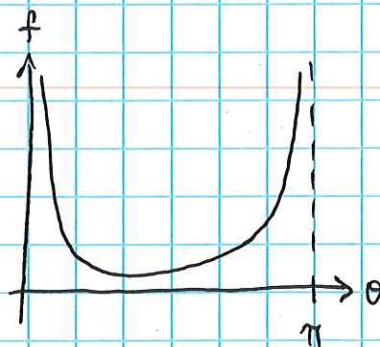
$$\theta_1 \leq \theta \leq \theta_2.$$

what else can we say about the motion in general?

→ consider just motion of the symmetry axis x_3 , which is described by θ & ϕ .

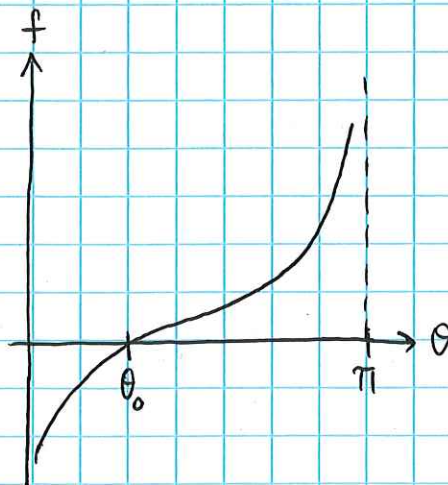
$$\rightarrow \dot{\phi} = \frac{L_2 - L_3 \cos \theta}{I_1' \sin^2 \theta} = \frac{L_3}{I_1'} \frac{(L_2/L_3 - \cos \theta)}{\sin^2 \theta} \equiv \frac{L_3}{I_1'} f(\theta).$$

(i) if $L_2/L_3 > 1$, then:



$$\Rightarrow \dot{\phi} > 0.$$

(ii) if $L_2/L_3 < 1$, then:



$$\cos \theta_0 = L_2/L_3$$

⇒ $\dot{\phi}$ changes sign.

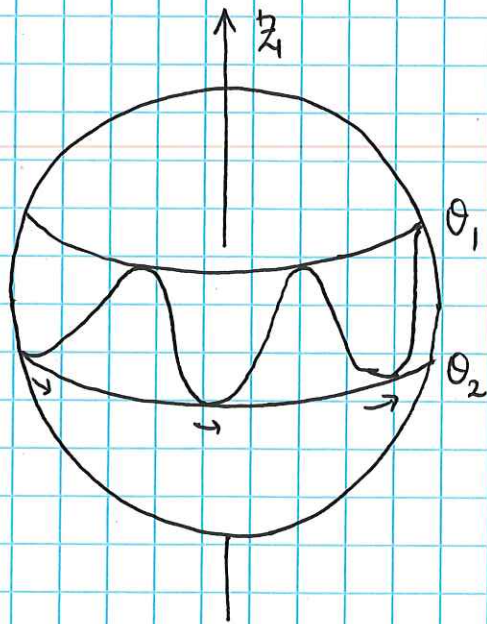
$$\Rightarrow \begin{cases} \dot{\phi} < 0 & \text{for } 0 \leq \theta \leq \theta_0 \\ \dot{\phi} > 0 & \text{for } \theta_0 < \theta \leq \pi \end{cases}$$

Now, since we know motion restricted to $\theta_1 \leq \theta \leq \theta_2$, we may distinguish different cases:

$$(1) \quad \begin{array}{c} \text{---} \theta_0 \quad \theta_1 \quad \theta_2 \text{---} \rightarrow \theta \Rightarrow \dot{\phi} > 0 \text{ for } \theta_1 \leq \theta \leq \theta_2. \\ \text{---} \theta_1 \quad \theta_2 \quad \theta_0 \text{---} \rightarrow \theta \Rightarrow \dot{\phi} < 0 \quad " \quad " \end{array}$$

(note: in case (i) above $\dot{\phi} > 0$).

In this situation where $\dot{\phi}$ has single sign throughout motion, x_3 precesses monotonically about \hat{z} axis (ϕ motion), while also oscillating up & down. (θ motion).



"nutration".

← trace of x_3 axis
on surface of sphere
centered at the origin.

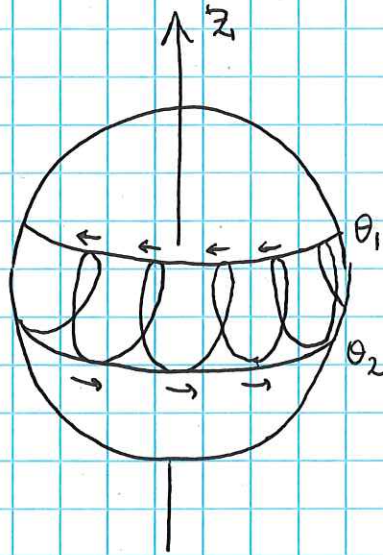
(2)



$$\Rightarrow \dot{\phi} < 0 \text{ for } \theta_1 \leq \theta < \theta_0$$

$$\dot{\phi} > 0 \text{ for } \theta_0 < \theta \leq \theta_2$$

in this case, ϕ changes sign throughout the motion. & signs are opposite on limiting circles θ_1 & θ_2

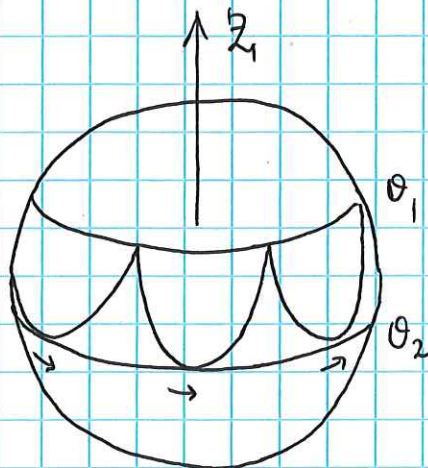


(3)



$$\theta_1 = \theta_0 \text{ or } \theta_2 = \theta_0$$

$$\Rightarrow \dot{\phi} \text{ \& \; } \dot{\theta} \text{ vanish simultaneously.}$$



$$(\theta_1 = \theta_0)$$

(6)

04/08/24

Ex: near vertical oscillationQ: under what conditions is $\theta=0$ a pt. of stable equilibrium?

$$U_{\text{eff}}(\theta) = \frac{(L_z - L_3 \cos\theta)^2}{2I_1' \sin^2\theta} - Mgl(1 - \cos\theta).$$

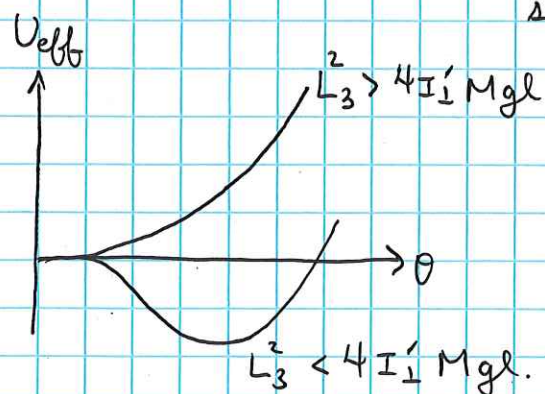
First, observe that when $\theta=0$ $L_z = L_3$, since z_1 & z_3 axes coincide. More explicitly:

$$L_z = (I_1' \sin^2\theta + I_3 \cos^2\theta) \dot{\phi} + I_3 \dot{\psi} \cos\theta \xrightarrow{\theta=0} I_3 (\dot{\phi} + \dot{\psi})$$

$$L_3 = I_3 (\dot{\phi} \cos\theta + \dot{\psi}) \xrightarrow{\theta=0} I_3 (\dot{\phi} + \dot{\psi}) \quad \leftarrow =$$

$$L_z = L_3 \Rightarrow U_{\text{eff}}(\theta) = \frac{L_3^2}{2I_1'} \frac{(1 - \cos\theta)^2}{\sin^2\theta} - Mgl(1 - \cos\theta).$$

$$\theta \text{ near zero: } \left. \begin{array}{l} \cos\theta \approx 1 - \theta^2/2 \\ \sin\theta \approx \theta \end{array} \right\} U_{\text{eff}}(\theta) \approx \underbrace{\left(\frac{L_3^2}{8I_1'} - \frac{Mgl}{2} \right)}_{\text{stable if } > 0} \theta^2.$$



\Rightarrow motion about $\theta=0$ stable if $L_3^2 > 4I_1' Mgl$.

or $\Omega_3^2 > 4I_1' Mgl / I_3^2$ ($L_3 = I_3 \Omega_3$), that is, ~~motion~~ $\theta=0$ is stable if spin about symm. axis is fast enough.

Appendix

04/08/24

(7)

• proof that $U_{\text{eff}}(\theta)$ in problem of heavy symmetric top has at most two turning pt.'s in range $0 \leq \theta \leq \pi$.

$$E' = \frac{1}{2} I_1' \dot{\theta}^2 + U_{\text{eff}}(\theta).$$

$$U_{\text{eff}}(\theta) = \frac{1}{2I_1'} \frac{(L_2 - L_3 \cos \theta)^2}{\sin^2 \theta} - Mgl(1 - \cos \theta).$$

• let $u = \cos \theta \rightarrow \dot{u} = -\dot{\theta} \sin \theta$. ($-1 \leq u \leq 1$).

$$\Rightarrow \dot{u}^2 = \frac{2E'}{I_1'} (1 - u^2) - \left(\frac{L_2 - L_3 u}{I_1'} \right)^2 + \frac{2Mgl}{I_1'} (1 - u)(1 - u^2) \equiv f(u)$$

\Rightarrow turning pt.'s $\dot{u} = 0$ correspond to zeros of polynomial $f(u) = 0$.

• now, $f(u)$ = cubic poly. s.t. $f(u) \xrightarrow{u \rightarrow \infty} \frac{2Mgl}{I_1'} u^3$.

$$\& f(\pm 1) = - \left(\frac{L_2 \mp L_3}{I_1'} \right)^2 < 0$$

\Rightarrow at least one root for $u > 1$. Since $f(u)$ = cubic, there are at most two roots for $-1 \leq u \leq 1$.

(& these roots correspond to turning pt.'s).

