Nuclear Shell Model

2, 8, 20, 28, 50, 82, and 126.

Simple Shell Model

$$V(r) = -\frac{V_o}{1 + \exp[(r - R)/a]}$$
 $V_o \sim 57 \text{ MeV}, R \sim 1.25 A^{1/3} \text{ F, a} \sim 0.65 \text{ F.}$

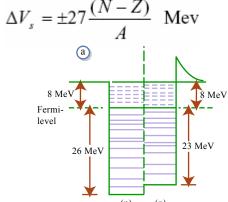
harmonic oscillator potential, $V(r) = m\omega^2 r^2 / 2$ $E_{\nu} = \hbar\omega(\nu + 3/2) = \hbar\omega(n_x + n_y + n_z + 3/2)$

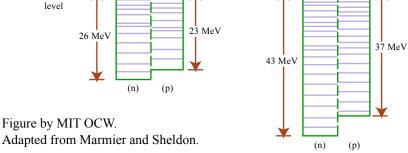
Shell Model with Spin-Orbit Coupling

$$H = \frac{p^2}{2m} + V(r) + V_{so}(r)\underline{s} \cdot \underline{L} \quad \underline{j} = \underline{S} + \underline{L} \quad \underline{S} \cdot \underline{L} = (j^2 - S^2 - L^2)/2$$

$$|\ell - s| \le j \le \ell + s \quad -j \le m_j \le j$$

Potential Wells for Neutrons and Protons





(b)

8 MeV

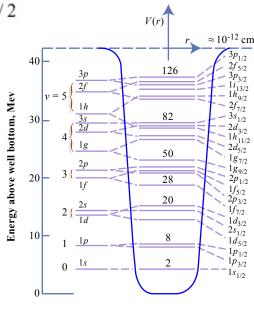


Figure by MIT OCW. Adapted from Meyerhof.

Nuclear Binding Energy and Stability

$$B(A,Z) \equiv [ZM_{H} + NM_{n} - M(A,Z)]c^{2}$$

$$i + I \to f + F + Q \qquad Q = T_{f} + T_{F} - (T_{i} + T_{I})$$

$$Q \equiv [(M_{i} + M_{I}) - (M_{f} + M_{F})]c^{2}$$

$$O = B(f) + B(F) - B(i) - B(I)$$

Separation Energy

$$S_n = [M_n + M(A-1, Z) - M(A, Z)]c^2$$

= B(A,Z) - B(A-1, Z)

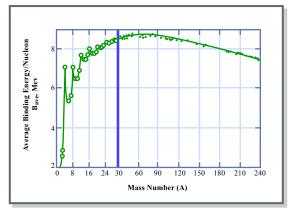


Figure by MIT OCW. Adapted from Meyerhof.

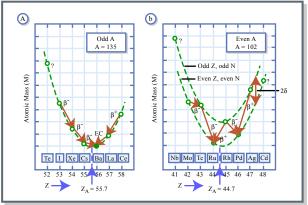
- ~ 1 Mev difference between the neutron absorbed being an even neutron or an odd neutron N~ Z for low A, but N > Z at high A.
- (i) In the case of odd A, only one stable isobar exists, except A = 113, 123.
- (ii) In the case of even A, only even-even nuclides exist, except A = 2, 6, 10,

Empirical Binding Energy Formula and Mass Parabolas

$$B(A,Z) = a_v A - a_s A^{2/3} - a_c \frac{Z(Z-1)}{A^{1/3}} - a_a \frac{(N-Z)^2}{A} + \delta$$

$$\delta = a_p / \sqrt{A}$$
 even-even nuclei
= 0 even-odd, odd-even nuclei
= $-a_p / \sqrt{A}$ odd-odd nuclei

$$\frac{a_{\rm v}}{16}$$
 $\frac{a_{\rm s}}{18}$ $\frac{a_{\rm c}}{0.72}$ $\frac{a_{\rm a}}{23.5}$ $\frac{a_{\rm p}}{11}$ Mev



$$M(A,Z)c^2 \cong A\Big[M_nc^2 - a_v + a_a + a_s/A^{1/3}\Big] + xZ + yZ^2 - \delta$$
 Figure by MIT OCW. Adapted From Meyerhof.
$$x = -4a_s - (M_\gamma - M_H)c^2 \cong -4a_a$$

$$y = \frac{4a_a}{A} + \frac{a_c}{A^{1/3}}$$

proton-rich.
$$_{9}F^{16} \rightarrow_{8}O^{16} + \beta^{+} + \nu$$
 $M(A, Z+1) > M(A,Z) + 2 m_{e}$ β^{+} - decay too many neutrons, $_{7}N^{16} \rightarrow_{8}O^{16} + \beta^{-} + \overline{\nu}$ $M(A,Z) > M(A,Z+1)$ β^{-} - decay

A competing process with positron decay is electron capture (EC). M(A, Z+1) > M(A,Z)twice the electron rest mass (1.02 Mev).

two exceptions, A = 113 123. Yet there are several exceptions, H^2 , Li^6 , B^{10} and N^{14} .

Radioactive-Series Decay

probability of a decay during a small time interval Δt . $P(\Delta t) = \lambda \Delta t$ the survival probability $S(t) \rightarrow e^{-\lambda t}$ $S(t_{1/2}) = 1/2 \rightarrow t_{1/2} = \ln 2/\lambda = 0.693/\lambda$ $\tau = \frac{1}{\lambda}$

 λ N(t), where N is the number of radioisotope atoms at time t, is called *activity*.

 $1 \text{ Ci} = 3.7 \text{ x } 10^{10} \text{ disintegrations/sec}, 1 \text{ Bq} = 2.7 \text{ x } 10^{-11} \text{ Ci}$

rate of radioactive decay.

Radioisotope Production by Bombardment

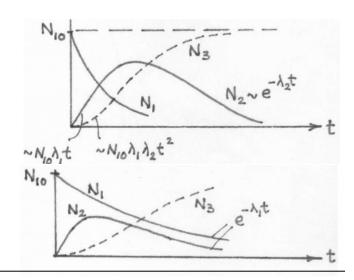
$$\frac{dN(t)}{dt} = Q_o - \lambda N(t) \qquad N(t) = \frac{Q_o}{\lambda} (1 - e^{-\lambda T}) e^{-\lambda (t - T)} \qquad N(t) = \frac{Q_o}{\lambda} (1 - e^{-\lambda t}), \quad t < T$$

Radioisotope Production in Series Decay

Series Decay with Short-Lived Parent

Series Decay with Short-Lived Parent

$$N_2(t) \approx N_{10} \frac{\lambda_1}{\lambda_2} e^{-\lambda t}$$
 $\lambda_2 N_2(t) \approx \lambda_1 N_1(t)$ secular equilibrium.



Charged-Particle Interactions: Stopping Power, Collisions and Ionization

Inelastic Collision with Atomic Electrons. ionization excitation of the atomic electrons

Inelastic Collision with a Nucleus. radiate (bremsstrahlung)

Elastic Collision with a Nucleus. Rutherfordscattering.

Elastic Collision with Atomic Electrons

Stopping Power: Energy Loss of Charged Particles in Matter

$$-\frac{dT}{dx}$$

$$\sim 5\infty \bar{I} \qquad 3Mc^{2}$$

$$-\frac{dT}{dx} = \frac{4\pi z^2 e^4 nZ}{m_e v^2} \ell n \left(\frac{2m_e v^2}{\overline{I}}\right)$$
 nonrelativistic
$$\ell n \left(\frac{2m_e v^2}{\overline{I}}\right) - \ell n \left(1 - \frac{v^2}{c^2}\right) - \frac{v^2}{c^2}$$
 The result is valid
$$\frac{ze^2}{\hbar v} = \left(\frac{e^2}{\hbar c}\right) \frac{z}{v/c} = \frac{z}{137(v/c)} << 1$$

$$-\left(\frac{dT}{dx}\right)_{class} = \frac{4\pi z^2 e^4 nZ}{m_e v^2} \ell n \left[\frac{M\hbar v}{2ze^2(m_e+M)} \frac{2m_e v^2}{\bar{I}}\right] \text{ if } \frac{ze^2}{\hbar v} >> 1$$

Charged-Particle Interactions: Radiation Loss, Range $\lambda_{min} = hc/T$ $v_{max} = T/h$

$$\sigma_{rad} \sim \frac{Z^2}{137} \left(\frac{e^2}{m_e c^2}\right)^2 \text{ cm}^2/\text{nucleus} \quad z^2 / m_e c^2 = r_e = 2.818 \text{ x } 10^{-13} \text{ cm}$$

$$\left[\frac{d\sigma}{d(h\nu)} \right]_{rad} = \sigma_o B Z^2 \frac{T + m_e c^2}{T} \frac{1}{h\nu}$$

$$-\left(\frac{dT}{dx} \right)_{rad} = n(T + m_e c^2) \sigma_{rad} \quad \text{ergs/cm}$$

$$nZ = (\rho N_o / A) Z = \rho N_o (Z / A)$$

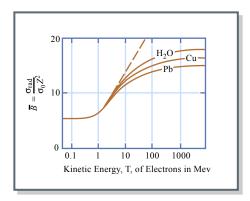


Figure by MIT OCW. Adapted from Evans.

Constant	Value	Unit	
		mks	cgs
Speed of light in vacuum c	2.997925(1)	x 10 ⁸ m s ⁻¹	x 10 ¹⁰ cm s ⁻¹
Elementary charge e	1.60210(2) 4.80298(7)	10 ⁻¹⁹	C 0 ⁻²⁰ emu 1 10 ⁻¹⁰ esu
Avogadro's number N	6.02252(9)	10 ²⁶ kmole ⁻¹	10 ²³ mole ⁻¹
Mass unit	1.66043(2)	10 ⁻²⁷ kg	10 ⁻²⁴ g
Electron rest mass m ₀	9.10908(13) 5.48597(3)	10 ⁻³¹ kg 10 ⁻⁴ u	10 ⁻²⁸ g 10 ⁻⁴ u
Proton rest mass M _P	1.67252(3) 1.00727663(8)	10 ⁻²⁷ kg u	10 ⁻²⁴ g
Neutron rest mass M _n	1.67482(3) 1.0086654(4)	10 ⁻²⁷ kg u	10 ⁻²⁴ g
Faraday constant Ne	9.64870(5) 2.89261(2)	10 ⁴ C mole ⁻¹	10 ³ emu 10 ¹⁴ esu
Planck constant $h = h/2$	6.62559(16) 1.054494(25)	10 ⁻³⁴ J s 10 ⁻³⁴ J s	10 ⁻²⁷ erg s 10 ⁻²⁷ erg s
Charge-to-mass ratio for electron e/m ₀	1.758796(6) 5.27274(2)	10 ¹¹ C kg ⁻¹	10 ⁷ emu 10 ¹⁷ esu
Rydberg constant 2 2m ₀ e ⁴ /h ³ c	1.0973731(1)	10 ⁷ m ⁻¹	10 ⁵ cm ⁻¹
Bohr radius ħ²/m ₀ e²	5.29167(2)	10 ⁻¹¹ m	10 ⁻⁹ cm
Compton wavelength of electron $\frac{h/m_0c}{\hbar/m_0c}$	2.42621(2) 3.86144(3)	10 ⁻¹² m 10 ⁻¹³ m	10 ⁻¹⁰ cm 10 ⁻¹¹ cm
Compton wavelength of proton $\frac{h/M_pc}{\hbar/M_pc}$	1.321398(13) 2.10307(2)	10 ⁻¹⁵ m 10 ⁻¹⁶ m	10 ⁻¹³ cm 10 ⁻¹⁴ cm

Figure by MIT OCW. Adapted from Meyerhof, Appendix D.

Conversion Factor	Value
l electron volt	1.60210(2) x 10 ⁻¹⁹ J 1.60210(2) x 10 ⁻¹² erg 8065.73(8) cm ⁻¹ 2.41804(2) x 10 ⁻¹⁴ s ⁻¹
$E_r \lambda_r$	12398.10(13) x 10 ⁻⁸ ev cm
1 u	931.478(5) Mev
Proton mass M _p c ²	938.256(5) Mev
Neutron mass M _n c ²	939.550(5) Mev
Electron mass m ₀ c ²	511006(2) ev
Rydberg $2\pi^2 m_0 e^4/h^2$	2.17971(5) x 10 ⁻¹¹ erg 13.60535(13) ev
Gas constant	8.31434 x 10 ⁷ erg mole ⁻¹ deg ⁻¹ 0.082053 liter atm mole ⁻¹ deg ⁻¹ 82.055 cm ³ atm mole ⁻¹ deg ⁻¹ 1.9872 cal _{th} mole ⁻¹ deg ⁻¹
Standard volume of ideal gas at NTP	22413.6 cm ³ mole ⁻¹
Mass on physical scale ($O^{16} = 16$) Mass on unified scale ($C^{12} = 12$)	1.000317917(17)
Mass on chemical scale (O = 16) Mass on unified scale ($C^{12} = 12$)	1.000043(5)

Figure by MIT OCW. Adapted from Evans

Lecture 9:

$$\begin{split} S_n &= [M(A-1,Z) + M_n - M(A,Z)]c^2 \\ &- \frac{\hbar}{2m} \frac{d^2 u_\ell}{dr^2} + \left[\frac{\ell(\ell+1)\hbar^2}{2mr^2} + V(r) \right] u_\ell(r) = E u_\ell(r) \\ E_\nu &= \hbar \omega (\nu + 3/2) = \hbar \omega (n_x + n_y + n_z + 3/2) \quad H = \frac{p^2}{2m} + V(r) + V_{so}(r) \underline{s} \cdot \underline{L} \\ &\left[\ell, m_\ell, s, m_s \right\rangle \equiv Y_\ell^{m_\ell} \chi_s^{m_s} \quad S^2 \chi_s^{m_s} = s(s+1)\hbar^2 \chi_s^{m_s}, \qquad s = 1/2 \\ S_z \chi_s^{m_s} &= m_s \hbar \chi_s^{m_s}, \qquad -s \leq m_s \leq s \quad \underline{j} = \underline{S} + \underline{L} \quad \underline{S} \cdot \underline{L} = (j^2 - S^2 - L^2)/2 \end{split}$$

$$j^{2} |jm_{j}\ell s\rangle = j(j+1)\hbar^{2} |jm_{j}\ell s\rangle, \quad |\ell-s| \leq j \leq \ell+s$$
Nucleon occupation = 2j+1
$$j_{z} |jm_{j}\ell s\rangle = m_{j}\hbar |jm_{j}\ell s\rangle, \quad -j \leq m_{j} \leq j$$

$$L^{2} |jm_{j}\ell s\rangle = \ell(\ell+1)\hbar^{2} |jm_{j}\ell s\rangle, \quad \ell = 0, 1, 2, ...$$

$$S^{2} |jm_{j}\ell s\rangle = s(s+1)\hbar^{2} |jm_{j}\ell s\rangle, \quad s = \frac{1}{2}$$

Lecture 10

$$B(A,Z) = [ZM_H + NM_n - M(A,Z)]c^2 \qquad Q = [(M_i + M_I) - (M_f + M_F)]c^2$$

$$T_i + M_i c^2 + T_I + M_I c^2 \rightarrow T_f + M_f c^2 + T_F + M_F c^2 \qquad Q = T_f + T_F - (T_i + T_I)$$

$$Q = B(f) + B(F) - B(i) - B(I)$$

$$S_a = [M_a(A', Z') + M(A - A', Z - Z') - M(A, Z)]c^2$$

$$S_a = [M_a + M(A - 1, Z) - M(A, Z)]c^2 = B(A, Z) - B(A - 1, Z)$$

Lecture 11

$$\delta = a_p / \sqrt{A}$$
 even-even nuclei
$$\underline{a_v} \quad \underline{a_s} \quad \underline{a_c} \quad \underline{a_a} \quad \underline{a_p}$$

$$= 0$$
 even-odd, odd-even nuclei
$$16 \quad 18 \quad 0.72 \quad 23.5 \quad 11 \quad \text{Mev}$$

$$= -a_p / \sqrt{A} \quad \text{odd-odd nuclei}$$

$$B(A,Z) = a_v A - a_s A^{2/3} - a_c \frac{Z(Z-1)}{A^{1/3}} - a_a \frac{(N-Z)^2}{A} + \delta$$

Lecture 12

$$P(dt) = \lambda dt$$

$$1 - P(dt)$$

$$[1 - P(dt)]^{2}$$

$$[1 - P(dt)]^{n}$$

$$T = \frac{\ln 2}{\lambda} = \frac{0.693}{\lambda}.$$

Formulas for Energy Release (Q), in Terms of Mass Differences, Δ_P and Δ_D , of Parent and Daughter Atoms

Type of Decay	Formula	
$egin{array}{c} lpha \ eta^- \ \gamma \ EC \ eta^+ \end{array}$	$Q_{\alpha} = \Delta_{P} - \Delta_{D} - \Delta_{Hc}$ $Q_{\beta} = \Delta_{P} - \Delta_{D}$ $Q_{IT} = \Delta_{P} - \Delta_{D}$ $Q_{EC} = \Delta_{P} - \Delta_{D} - E_{B}$ $Q_{\beta} = \Delta_{P} - \Delta_{D} - 2mc^{2}$	
-	νр -г -υ	

Figure by MIT OCW. Adapted from Meyerhof, Appendix D.

$$f = \frac{dt}{|dN|} = \frac{1}{\lambda N}$$

$$\frac{dN_1}{dt} = -\lambda_1 N_1$$

$$\frac{dN}{dt} = \lambda_2 N_2$$

$$\frac{dN_2}{dt} = \lambda_2 N_2$$

$$\frac{dN_3}{dt} = \lambda_2 N_2$$

$$\frac{dN_4}{dt} = \lambda_2 N = dt$$

$$\frac{dN_2}{dt} = \lambda_2 N_2$$

$$\frac{dN_3}{dt} = \lambda_2 N_2$$

$$\frac{dN_4}{dt} = \lambda_2 N_2$$

$$\frac{dN_5}{dt} = \lambda_2 N_5$$

$$\frac{dN_6}{dt} = \lambda_2 N_6$$

$$\frac{dN_6}{dt} = \lambda_2 N_$$

Lecture 13

$$\frac{p_e^2}{2m_e} = \frac{2(ze^2)^2}{m_e b^2 v^2} - \frac{dT}{dx} = \frac{4\pi z^2 e^4 nZ}{m_e v^2} \ell n \left(\frac{2m_e v^2}{\bar{I}}\right) - \frac{dT}{dx} = \frac{2\pi e^4 nZ}{m_e v^2} \left[\ell n \left(\frac{m_e v^2 T}{\bar{I}^2 (1-\beta^2)}\right) - \beta^2\right]$$

Lecture 14

$$\sigma_{ion} = \frac{2\alpha Z}{\beta^4} \ell n \left(\frac{\sqrt{2}T}{\bar{I}} \right) \qquad \text{ionization}$$

$$\sigma_{mic} = \frac{\alpha Z^2}{4\beta^4} \qquad \qquad \text{backscattering by nuclei}$$

$$\sigma_{el} = \frac{2\alpha Z}{\beta^4} \qquad \qquad \text{elastic scattering by atomic electrons}$$

$$\sigma_{rod}^{\prime} = \frac{8\alpha}{3\pi} \frac{1}{137} \frac{Z^2}{\beta^2} \qquad \qquad bremsstrahlung$$

$$nZ = (\rho N_o / A)Z = \rho N_o (Z / A)$$

$$nZ = (\rho N_o / A)Z = \rho N_o (Z / A) \qquad \frac{\left(dT / dx\right)_{rad}}{\left(dT / dx\right)_{ion}} \approx Z \left(\frac{m_e}{M}\right)^2 \left(\frac{T}{1400 m_e c^2}\right)$$

$$i = \frac{1}{W} \left(-\frac{dT}{dx} \right)$$

$$R \propto \int_{0}^{T_o} T dT = T_o^2$$
 $R \propto \int_{0}^{T_o} dT = T_o$ $\frac{R}{R_1} = \frac{\rho_1 \sqrt{A}}{\rho \sqrt{A_1}}$

$$R = 3.2x10^{-4} \frac{\sqrt{A}}{\rho} \times R_{air}$$