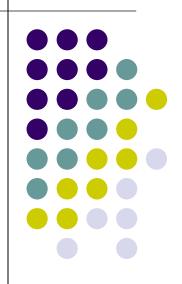
Lossy Medium

EE142

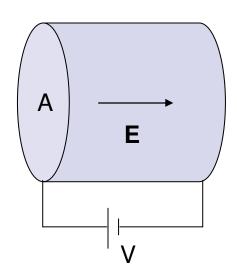
Dr. Ray Kwok



•reference:

Fundamentals of Engineering Electromagnetics, David K. Cheng (Addison-Wesley) Electromagnetics for Engineers, Fawwaz T. Ulaby (Prentice Hall)

Ohm's Law



$$\begin{split} V &= IR \\ E\ell &= \left(JA\right)\!\!\left(\rho\frac{\ell}{A}\right) \\ \vec{E} &= \vec{J}\rho \qquad \text{resisitivity} \\ \vec{J} &= \frac{1}{\rho}\vec{E} \\ \vec{J} &= \sigma\vec{E} \qquad \text{conductivity} \end{split}$$

- Low resistivity => "conductor" ~<10⁻⁵ Ω om (♥ T)
- High resistivity => "insulator" $\sim > 10^{10} \Omega$ m
- Intermediate resistivity => "semiconductor" typical ~10⁻³ to $10^5 \Omega$ om ($colone{color}$ e^{Eg/kT})
- unit of conductivity = S/m = $Siemens/meter = mho/m = (<math>\Omega$ m^{-1}



EM Wave through medium



$$\epsilon \nabla \cdot \vec{E} = \rho_f$$

$$\nabla \times \vec{E} = -\mu \frac{\partial \vec{H}}{\partial t}$$

$$\nabla \cdot \vec{H} = 0$$

$$\nabla \times \vec{H} = \vec{J}_f + \epsilon \frac{\partial \vec{E}}{\partial t}$$

$$\nabla \cdot \vec{E} = 0$$
 (homogeneous, linear, isotropic)
$$\nabla \times \vec{E} = -j\omega\mu \vec{H}$$

$$\rho \approx 0, \ J \neq 0$$

$$\nabla \cdot \vec{H} = 0$$

$$\nabla \times \vec{H} = \vec{J}_f + j\omega\epsilon\vec{E}$$

$$\nabla \times \vec{H} = (\sigma + j\omega \varepsilon)\vec{E} = j\omega \left(\varepsilon + \frac{\sigma}{j\omega}\right)\vec{E} = j\omega \varepsilon_c \vec{E}$$

$$\varepsilon_{c} \equiv \varepsilon - j \frac{\sigma}{\omega} \equiv \varepsilon' - j\varepsilon''$$

finite σ means complex ϵ



Loss Tangent

$$\tan \delta \equiv \frac{\epsilon''}{\epsilon'} = \frac{\sigma}{\omega \epsilon}$$

good conductor $\sigma >> \omega \epsilon$ good insulator $\sigma << \omega \epsilon$

Low $tan\delta \rightarrow low dielectric loss$

the smaller, the better !!!

DIELECTRIC CONSTANTS AND LOSS TANGENTS FOR SOME MATERIALS

Material	Frequency	ϵ_r	tan δ (25°C)	
Alumina (99.5%)	10 GHz	9.5–10.	0.0003	
Barium tetratitanate	6 GHz	$37 \pm 5\%$	0.0005	
Beeswax	10 GHz	2.35	0.005	
Beryllia	10 GHz	6.4	0.0003	
Ceramic (A-35)	3 GHz	5.60	0.0041	
Fused quartz	10 GHz	3.78	0.0001	
Gallium arsenide	10 GHz	13.	0.006	
Glass (pyrex)	3 GHz	4.82	0.0054	
Glazed ceramic	10 GHz	7.2	0.0034	
Lucite	10 GHz	2.56	0.005	
Nylon (610)	3 GHz	2.84	0.012	
Parafin	10 GHz	2.24	0.002	
Plexiglass	3 GHz	2.60	0.0002	
Polyethylene	10 GHz	2.25	0.0037	
Polystyrene	10 GHz	2.54	0.0004	
Porcelain (dry process)	100 MHz	5.04	0.00033	
Rexolite (1422)	3 GHz	2.54	0.0078	
Silicon	10 GHz	11.9	0.0048	
Styrofoam (103.7)	3 GHz	1.03	0.0001	
Teflon	10 GHz	2.08	0.0001	
litania (D-100)	6 GHz	$96 \pm 5\%$	0.0004	
Vaseline	10 GHz	2.16	0.001	
Water (distilled)	3 GHz	76.7	0.001	

Example



A sinusoidal E-field with amplitude of 250 V/m and frequency 1 GHz exists in a lossy dielectric medium that has a $\varepsilon_r = 2.5$ and loss tangent of 0.001. Find the average power dissipated in the medium per cubic meter.

$$\tan \delta = 0.001 = \frac{\sigma}{\omega \varepsilon} = \frac{\sigma}{\omega \varepsilon_o \varepsilon_r} = \frac{\sigma}{\left(2\pi \cdot 10^9 \right) \left(\frac{10^{-9}}{36\pi}\right) (2.5)}$$

$$\sigma = 1.39 \cdot 10^{-4}$$
 S/m

The average power dissipated per unit volume is

$$\frac{P_{ave}}{V} = \frac{1}{2}\vec{J} \cdot \vec{E} = \frac{1}{2}\sigma E^2 = \frac{1}{2}(1.39 \cdot 10^{-4})(250)^2$$

$$\frac{P_{ave}}{V} = 4.34 \quad \text{W/m}^2$$

$$P_{\text{ave}} = \frac{1}{2} \frac{V^2}{R} = \frac{1}{2} \frac{(E\ell)^2}{\rho \ell / A} = \frac{1}{2} \sigma E^2 (\ell A)$$

Wave Equation

$$\nabla \times \vec{E} = -\mu \frac{\partial \vec{H}}{\partial t}$$

$$\nabla \times \vec{H} = \varepsilon_{c} \frac{\partial \vec{E}}{\partial t}$$

$$\nabla^2 \vec{E} = \mu \varepsilon_c \frac{\partial^2 \vec{E}}{\partial t^2}$$

plane wave equation still holds with modification of $\boldsymbol{\epsilon}$

$$\vec{E}(r,t) = \vec{E}_o e^{j(\omega t - k_c r)} \equiv \vec{E}_o e^{j\omega t} e^{-\gamma r} \qquad \text{allow k be complex since ϵ is}$$

$$\gamma \equiv jk_c = j\omega\sqrt{\mu\epsilon_c} \equiv \alpha + j\beta$$

phase constant

propagation constant



Complex Propagation Constant

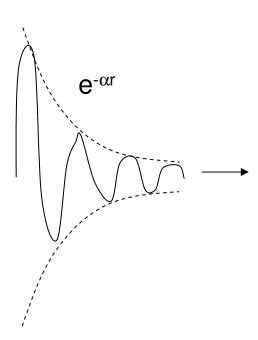


$$\gamma = j\omega\sqrt{\mu\epsilon_{\rm c}} = j\omega\sqrt{\mu\epsilon\left(1 - j\frac{\epsilon''}{\epsilon'}\right)} = j\omega\sqrt{\mu\epsilon\left(1 - j\frac{\sigma}{\omega\epsilon}\right)} = j\omega\sqrt{\mu\epsilon(1 - j\tan\delta)} = \alpha + j\beta$$

The phasor

$$\vec{E}(r) = \vec{E}_o e^{-\gamma r} = \vec{E}_o e^{-\alpha r} e^{-j\beta r}$$

attenuation



dB scale



power intensity ratio in log scale, not a unit!!

$$(dB) = 10 log \left(\frac{I}{I_o}\right) = 10 log \left(\frac{P}{P_o}\right) = 20 log \left(\frac{V}{V_o}\right) > 0 \text{ gain} < 0 \text{ loss}$$
 sound intensity power voltage

```
10 \log(2) \approx 3, 3 dB = double

10 \log(1/2) \approx -3, -3 dB = half

10 \log(10) = 10, 10 dB = 10x

10 \log(100) = 20, 20 dB = 100x

10 \log(0.1) = -10, -10 dB = 1/10
```

What is 6 dB?
$$-9$$
 dB? 7 dB? -44 dB? $4x$ $1/8$ $5x$ $4 x $10^{-5}$$

dBm & dBW



$$dBW \equiv 10 \log \left(\frac{P}{1W}\right)$$

$$dBm \equiv 10 \log \left(\frac{P}{1mW} \right)$$

become real units

$$0 \text{ dBm} = 1 \text{ mW}$$

 $30 \text{ dBW} = 1 \text{ kW}$
 $-30 \text{ dBm} = 1 \text{ }\mu\text{W}$

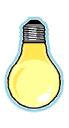
What is 40 dBW? -7 dBm? -26 dBm? 21 dBm?

10 kW 0.2 mW 2.5 μW 1/8 W

Example





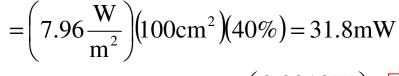


How much electricity generated by the solar cell? What if a 40 W bulb is used? 200 W bulb?

Intensity = power/area =
$$\frac{100}{4\pi R^2} = \frac{100}{4\pi (1)^2} = 7.96 \frac{W}{m^2}$$

1 m

Power generated in solar cell



solar cell 10 x 10 cm² 40% efficiency

In terms of dB =
$$10\log\left(\frac{0.0318W}{100W}\right) = -35dB$$

system "gain"

40 W bulb?
$$-35 = 10 \log \left(\frac{P}{40} \right)$$

Power of electricity generated = 12.6 mW

200 W bulb?
$$-35 = 10 \log \left(\frac{P}{200} \right)$$

Power of electricity generated = 63.2 mW

Attenuation

$$\vec{E}(r) = \vec{E}_o e^{-\gamma r} = \vec{E}_o e^{-\alpha r} e^{-j\beta r}$$

$$A(r)[dB] = 20\log\left|\frac{\vec{E}(r)}{\vec{E}(0)}\right| = 20\log\left|e^{-\alpha r}\right|$$

$$\log_{a}(b) = \frac{\log_{c}(b)}{\log_{c}(a)}$$

$$A(r)[dB] = \frac{-20\alpha r}{\ln(10)} = -8.686\alpha r$$



For example, if the electric field intensity going through a medium attenuates at a rate of 0.4 dB/m, what is α ?

-0.4 dB = -8.686
$$\alpha$$
 (1 m) α = 0.4/8.686 = 0.046 nepers/m

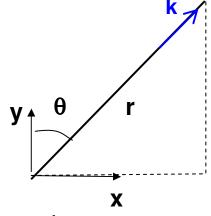
Note: nepers (np) is not a real unit. similar to radians !!!

Note also α is a positive number for attenuation.

$$\alpha$$
[dB/m] = 8.686 α [np/m]

Attenuation term





How to express $e^{-\alpha r}$ term??

$$\hat{\mathbf{k}} = \hat{\mathbf{x}}\sin\theta + \hat{\mathbf{y}}\cos\theta$$

$$e^{-\alpha r} = e^{-\alpha(x\sin\theta + y\cos\theta)} = e^{-\alpha\sqrt{x^2 + y^2}}$$
 Same?

$$x\sin\theta + y\cos\theta = x\left(\frac{x}{r}\right) + y\left(\frac{y}{r}\right) = \frac{x^2 + y^2}{\sqrt{x^2 + y^2}} = \sqrt{x^2 + y^2} \qquad \text{Yes, Q.E.D.}$$

Can think of :
$$\vec{\alpha} \equiv \alpha \hat{k}$$

$$\vec{\alpha} \cdot \vec{r} = \alpha (x \sin \theta + y \cos \theta)$$

$$e^{-\vec{\alpha}\cdot\vec{r}} = e^{-\alpha(x\sin\theta + y\cos\theta)} = e^{-\alpha\sqrt{x^2 + y^2}}$$

Low-loss dielectric (ε "<< ε ') or (σ << $\omega\varepsilon$)



$$\gamma = j\omega\sqrt{\mu\epsilon\left(1 - j\frac{\epsilon''}{\epsilon'}\right)} = j\omega\sqrt{\mu\epsilon\left(1 - j\frac{\sigma}{\omega\epsilon}\right)} = j\omega\sqrt{\mu\epsilon(1 - j\tan\delta)}$$

$$(1+x)^{n} = \sum_{k=0}^{n} \frac{n!}{k!(n-k)!} x^{n-k} = 1 + nx + \frac{n(n-1)}{2} x^{2} + \dots$$

$$(1+x)^n \approx 1+nx$$
 for small x

$$\gamma = j\omega\sqrt{\mu\epsilon} \left(1 - j\frac{\epsilon''}{\epsilon'}\right)^{1/2} \approx j\omega\sqrt{\mu\epsilon} \left(1 - \frac{1}{2}j\frac{\epsilon''}{\epsilon'}\right) \equiv \alpha + j\beta$$

$$\alpha = \frac{\omega\sqrt{\mu\epsilon}}{2} \frac{\epsilon''}{\epsilon'} = \frac{\omega\sqrt{\mu\epsilon}}{2} \frac{\sigma}{\omega\epsilon} = \frac{\sigma}{2} \sqrt{\frac{\mu}{\epsilon}} \quad \text{small}$$

$$\beta = \omega \sqrt{\mu \varepsilon} \approx \omega / v$$

$$\eta_{c} \equiv \sqrt{\frac{\mu}{\epsilon_{c}}} = \sqrt{\frac{\mu}{\epsilon(1 - j\epsilon''/\epsilon')}} = \sqrt{\frac{\mu}{\epsilon}} \left(1 - j\frac{\epsilon''}{\epsilon'}\right)^{-1/2} \approx \sqrt{\frac{\mu}{\epsilon}} \left(1 + \frac{j}{2}\frac{\epsilon''}{\epsilon'}\right) \approx \sqrt{\frac{\mu}{\epsilon}}$$

Good Conductor (σ>>ωε)

$$\gamma = j\omega \sqrt{\mu\epsilon \left(1 - j\frac{\sigma}{\omega\epsilon}\right)} \approx j\omega \sqrt{-j\frac{\sigma}{\omega\epsilon}\mu\epsilon}$$

$$\sqrt{-j} = \sqrt{e^{-j\pi/2}} = e^{-j\pi/4} = \frac{1-j}{\sqrt{2}}$$

$$\gamma = j\omega \frac{1-j}{\sqrt{2}} \sqrt{\frac{\mu\sigma}{\omega}} = (j+1)\sqrt{\frac{\mu\sigma\omega}{2}} \equiv \alpha + j\beta$$

$$\alpha = \beta = \sqrt{\frac{\mu\sigma\omega}{2}}$$

$$\eta_{c} \equiv \sqrt{\frac{\mu}{\epsilon_{c}}} = \sqrt{\frac{\mu}{\epsilon(1 - j\sigma/\omega\epsilon)}} \approx \sqrt{j\frac{\mu\omega\epsilon}{\epsilon\sigma}} \approx \frac{1 + j}{\sqrt{2}}\sqrt{\frac{\mu\omega}{\sigma}} = (1 + j)\sqrt{\frac{\mu\omega}{2\sigma}} = (1 + j)\frac{\alpha}{\sigma}$$

Skin Depth δ

$$\delta \equiv \frac{1}{\alpha} = \sqrt{\frac{2}{\mu \sigma \omega}} = \delta_s$$

(NOT loss tangent δ !!!!!)

$$\vec{E}(r) = \vec{E}_{o}e^{-\alpha r}e^{-j\beta r}$$

At $r = \delta$, |E| decreases to 1/e (or 63% drop).

$$A(r)[dB] = 20log \left| \frac{\vec{E}(r)}{\vec{E}(0)} \right| = 20log \left| e^{-\alpha r} \right| = -8.686\alpha r$$

At $r = \delta$, |E| decreases by -8.7 dB. At $r = 2\delta$, |E| decreases by -17.3 dB....



General Material

$$\gamma = j\omega\sqrt{\mu\epsilon_c} = j\omega\sqrt{\mu(\epsilon' - j\epsilon'')} \equiv \alpha + j\beta$$
$$\gamma^2 = -\omega^2\mu(\epsilon' - j\epsilon'') = \alpha^2 - \beta^2 + 2j\alpha\beta$$

$$-\omega^2\mu\epsilon' = \alpha^2 - \beta^2 \qquad \text{real}$$

$$j\omega^2\mu\epsilon''=j2\alpha\beta$$

imaginary

$$\alpha^{2} = \beta^{2} - \omega^{2} \mu \epsilon' = \left(\frac{\omega^{2} \mu \epsilon''}{2\alpha}\right)^{2} - \omega^{2} \mu \epsilon'$$

$$4\alpha^4 + 4\alpha^2\omega^2\mu\epsilon' - (\omega^2\mu\epsilon'')^2 = 0$$

$$\alpha^2 = \frac{-4\omega^2 \mu \epsilon' \pm \sqrt{(4\omega^2 \mu \epsilon')^2 + 16(\omega^2 \mu \epsilon'')^2}}{8}$$

$$\alpha^2 = \frac{\omega^2 \mu \epsilon'}{2} \left(-1 \pm \sqrt{1 + \tan^2 \delta} \right)$$

$$\alpha^2 = \frac{\omega^2 \mu \epsilon'}{2} \left(\sqrt{1 + \tan^2 \delta} - 1 \right)$$



$$\beta^{2} = \left(\frac{\omega^{2}\mu\epsilon''}{2\alpha}\right)^{2} = \frac{\omega^{2}\mu\epsilon'\tan^{2}\delta}{2(\sqrt{1+\tan^{2}\delta}-1)}$$

$$\beta^{2} = \frac{\omega^{2}\mu\epsilon'\tan^{2}\delta}{2(\sqrt{1+\tan^{2}\delta}-1)(\sqrt{1+\tan^{2}\delta}+1)}$$

$$\beta^{2} = \frac{\omega^{2}\mu\epsilon'\tan^{2}\delta}{2(1+\tan^{2}\delta-1)(\sqrt{1+\tan^{2}\delta}+1)}$$

$$\beta^{2} = \frac{\omega^{2}\mu\epsilon'\tan^{2}\delta(\sqrt{1+\tan^{2}\delta}+1)}{2(1+\tan^{2}\delta-1)}$$

$$\beta^2 = \frac{\omega^2 \mu \epsilon'}{2} \left(\sqrt{1 + \tan^2 \delta} + 1 \right)$$

$$\eta_{c} = \sqrt{\frac{\mu}{\varepsilon_{c}}} = \sqrt{\frac{\mu}{\varepsilon'(1-j\tan\delta)}}$$

$$\eta_{c} = \sqrt{\frac{\mu}{\epsilon'}} (1 - j \tan \delta)^{-1/2}$$

Summary



		Lossless	Low-loss	Good	
	Any Medium	Medium	Medium	Conductor	Units
		$(\sigma = 0)$	$(\varepsilon''/\varepsilon'\ll 1)$	$(\varepsilon''/\varepsilon'\gg 1)$	
$\alpha =$	$\omega \left[\frac{\mu \varepsilon'}{2} \left[\sqrt{1 + \left(\frac{\varepsilon''}{\varepsilon'} \right)^2} - 1 \right] \right]^{1/2}$	0	$\frac{\sigma}{2}\sqrt{\frac{\mu}{\varepsilon}}$.	$\sqrt{\pi f \mu \sigma}$	(Np/m)
$\beta =$	$\omega \left[\frac{\mu \varepsilon'}{2} \left[\sqrt{1 + \left(\frac{\varepsilon''}{\varepsilon'} \right)^2} + 1 \right] \right]^{1/2}$	$\omega\sqrt{\mu\varepsilon}$	$\omega\sqrt{\mu\varepsilon}$	$\sqrt{\pi f \mu \sigma}$	(rad/m)
$\eta_{\rm c} =$	$\sqrt{\frac{\mu}{\varepsilon'}} \left(1 - j \frac{\varepsilon''}{\varepsilon'} \right)^{-1/2}$	$\sqrt{\frac{\mu}{\varepsilon}}$	$\sqrt{\frac{\mu}{\varepsilon}}$	$(1+j)\frac{\alpha}{\sigma}$	(Ω)
$u_p =$	ω/β	$1/\sqrt{\mu\varepsilon}$	$1/\sqrt{\mu\varepsilon}$	$\sqrt{4\pi f/\mu\sigma}$	(m/s)
$\lambda =$	$2\pi/\beta = u_{\rm p}/f$	$u_{\rm p}/f$	u_p/f	$u_{\rm p}/f$	(m)

Notes: $\varepsilon' = \varepsilon$; $\varepsilon'' = \sigma/\omega$; in free space, $\varepsilon = \varepsilon_0$, $\mu = \mu_0$; in practice, a material is considered a low-loss medium if $\varepsilon''/\varepsilon' = \sigma/\omega\varepsilon < 0.01$ and a good conducting medium if $\varepsilon''/\varepsilon' > 100$.

Example - The skin depth of a certain nonmagnetic conducting material is 2μm at 5 GHz. Determine the phase velocity in the material. What is the attenuation (in dB) when the wave penetrates 10 μm into the material?



phase velocity $v = \omega/\beta$

for conductor, $\alpha = \beta = 1/\delta$

 $V = \omega \delta = (2\pi)(5 \times 10^9) (2 \times 10^{-6}) = 6.28 \times 10^4 \text{ m/s}$

$$A(r)[dB] = 20 \log \left| \frac{\vec{E}(r)}{\vec{E}(0)} \right| = 20 \log \left| e^{-\alpha r} \right| = -8.686 \alpha r$$

$$A(r)[dB] = -8.686r/\delta = -8.686(10/2) = -43.4dB$$

in just 5 skin depth.

(-43 dB = 1 / 20,000 !!!)

Only surface current on conductors.

Example — (a) Calculate the dielectric loss (in dB) of an EM wave propagating through 100 m of teflon at 1 MHz. (b) at 10 GHz?



Teflon: $\varepsilon_r = 2.08$, $\tan \delta = 0.0004$ at 25°C assuming frequency independence.

(a)
$$\tan \delta = \frac{\sigma}{\omega \varepsilon}$$

$$\sigma = \omega \varepsilon_o \varepsilon_r \tan \delta = \left(2\pi \cdot 10^6 \right) \left(\frac{10^{-9}}{36\pi}\right) (2.08)(0.0004) = 4.6 \cdot 10^{-8}$$
 S/m

$$\alpha = \frac{\sigma}{2} \sqrt{\frac{\mu}{\epsilon}} = \frac{\sigma \eta_o}{2\sqrt{\epsilon_r}} = \frac{(4.6 \cdot 10^{-8})(377)}{2\sqrt{2.08}} = 6.04 \cdot 10^{-6} \quad \text{np/m}$$

$$A(dB) = -8.686\alpha r = -8.686(6.04 \cdot 10^{-6})(100) = -0.005$$

(b)
$$\sigma = \omega \varepsilon_o \varepsilon_r \tan \delta = (2\pi \cdot 10^{10}) \left(\frac{10^{-9}}{36\pi}\right) (2.08)(0.0004) = 4.6 \cdot 10^{-4} \text{ S/m}$$

$$\alpha = \frac{\sigma}{2} \sqrt{\frac{\mu}{\epsilon}} = \frac{\sigma \eta_o}{2\sqrt{\epsilon_r}} = \frac{(4.6 \cdot 10^{-4})(377)}{2\sqrt{2.08}} = 6.04 \cdot 10^{-2} \text{ np/m}$$

$$A(dB) = -8.686\alpha r = -8.686(6.04 \cdot 10^{-2})(100) = -50 dB$$

Coaxial cable works well at low freq (TV to antenna) but not so well at high freq. !!

Example — In a nonmagnetic, lossy, dielectric medium, a 300-MHz plane wave is characterized by the magnetic field phasor $\vec{H} = (\hat{x} - j4\hat{z})e^{-2y}e^{-j9y}$ A/m. Obtain time-domain expressions for the electric and magnetic field vectors. What is the polarization state of this wave?



$$\vec{H}(\vec{r},t) = \Re e \{ (\hat{x} - j4\hat{z})e^{-2y}e^{j(\omega t - 9y)} \}$$

$$\vec{H}(\vec{r},t) = \hat{x}e^{-2y}\cos(\omega t - 9y) + \hat{z}4e^{-2y}\sin(\omega t - 9y)$$

$$\alpha = 2, \ \beta = 9 \qquad -\omega^{2}\mu\epsilon' = \alpha^{2} - \beta^{2}$$

$$\omega^{2}\mu\epsilon'' = 2\alpha\beta$$

$$\frac{\epsilon''}{\epsilon'} = \frac{2\alpha\beta}{\beta^{2} - \alpha^{2}} = \frac{2(2)(9)}{9^{2} - 2^{2}} = 0.468 = \tan\delta$$

$$\epsilon_{r} = \frac{\epsilon'}{\epsilon_{o}} = \frac{\beta^{2} - \alpha^{2}}{\omega^{2}\mu_{o}\epsilon_{o}} = \frac{77c^{2}}{\omega^{2}} = \frac{77(3 \cdot 10^{8})^{2}}{(2\pi \cdot 300 \cdot 10^{6})^{2}} = 1.95$$

$$\eta_{c} = \sqrt{\frac{\mu}{\epsilon'}} (1 - j\tan\delta)^{-1/2} = \frac{\eta_{o}}{\sqrt{\epsilon_{r}}} (1 - j\tan\delta)^{-1/2} = \frac{377}{\sqrt{1.95}} (1 - j0.468)^{-1/2}$$

$$\eta_{c} = 257\angle 12.5^{\circ} = 257e^{j0.22}$$

$$\vec{E} = \hat{z}\eta_c e^{-2y} \cos(\omega t - 9y) - \hat{x}4\eta_c e^{-2y} \sin(\omega t - 9y)$$

$$\vec{E} = \hat{z}257e^{-2y} \cos(\omega t - 9y + 0.22) - \hat{x}1028e^{-2y} \sin(\omega t - 9y + 0.22)$$

Exercise (how to write the attenuation?)



Given E_0 at the origin has a amplitude of 1 V/m along the y-axis in a non-magnetic medium, with the propagation given by:

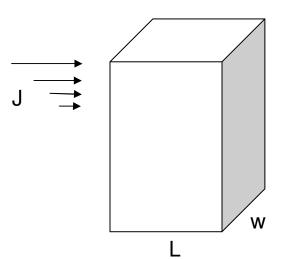
$$\sim \sin(\omega t + 30x - 15z)$$

$$\varepsilon = (4 - j0.02)\varepsilon_o$$

Write
$$\vec{H}(\vec{r},t) = ?$$

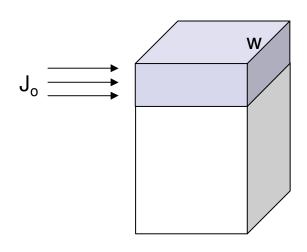
Surface resistance for conductors





$$J = \sigma E = J_o e^{-\alpha z} e^{-j\beta z} = J_o e^{-(1+j)z/\delta}$$

$$I = \int \vec{J} \cdot d\vec{a} = J_o w \int_0^\infty e^{-(1+j)z/\delta} dz = \frac{J_o w \delta}{(1+j)} = \frac{J_o w \delta}{\sqrt{2}} e^{-j\pi/4}$$



$$\delta/\sqrt{2}$$
 $V = E_o L = \frac{J_o}{\sigma} L$ on surface

$$Z = \frac{V}{I} = \frac{1+j}{\sigma\delta} \frac{L}{w} \equiv Z_s \frac{L}{w}$$
 similar to R & ρ

$$Z_{s} = \frac{1+j}{\sigma\delta}$$

 $Z_s = \frac{1+j}{2}$ surface impedance (Ω)

Homework



- Determine the frequency at which a time-harmonic electric field intensity causes a conduction current density and a displacement current density of equal magnitude in
 - (a) seawater with $\varepsilon_r = 72$ and $\sigma = 4$ S/m, and
 - (b) moist soil with $\varepsilon_r = 2.5$ and $\sigma = 10^{-3}$ S/m.
- Calculations concerning the electromagnetic effect of currents in a good conductor usually neglect the displacement current even at microwave frequencies.
 - (a) Assuming $\varepsilon_r = 1$ and $\sigma = 5.7 \times 10^7$ S/m for copper, compare the magnitude of the displacement current density with that of the conduction current density at 100 GHz.
 - (b) Write the differential equation in phasor form for magnetic field intensity H in a source-free good conductor.