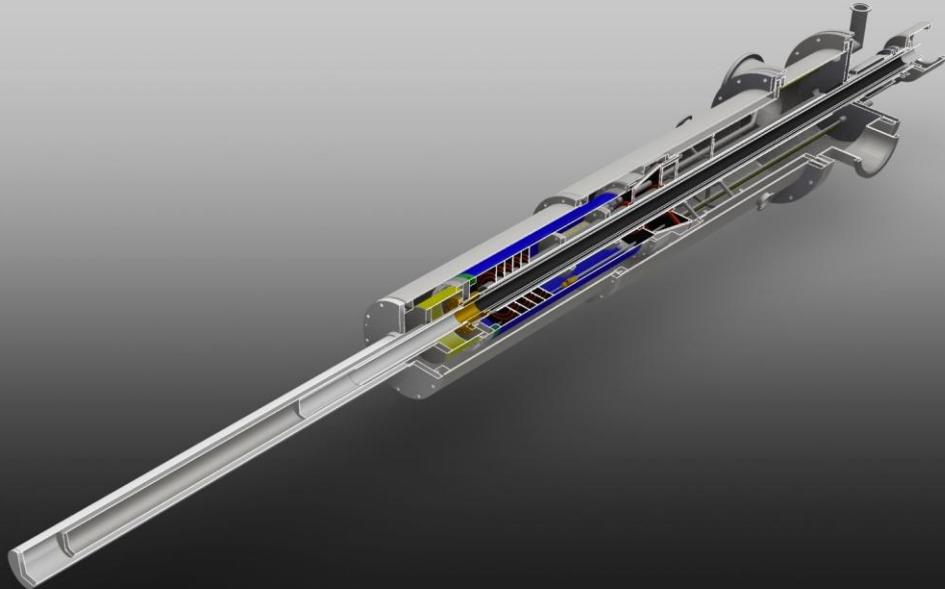
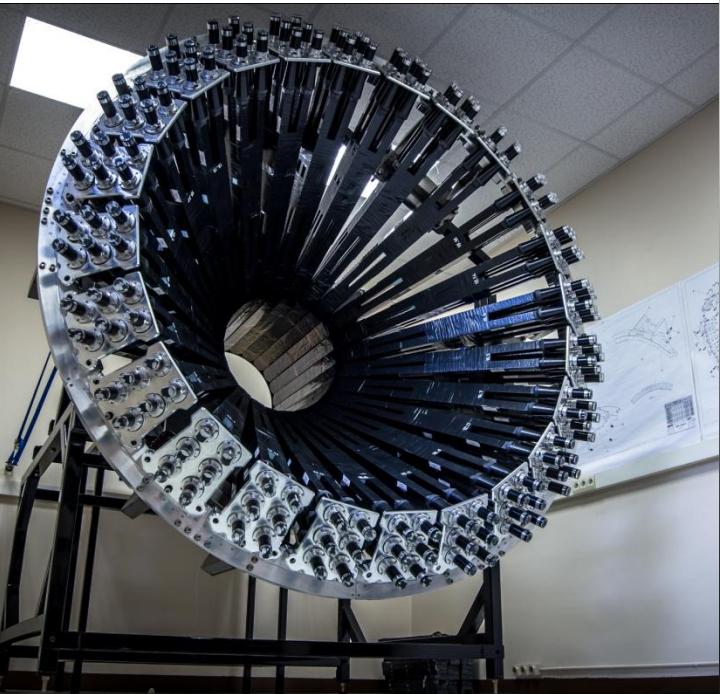
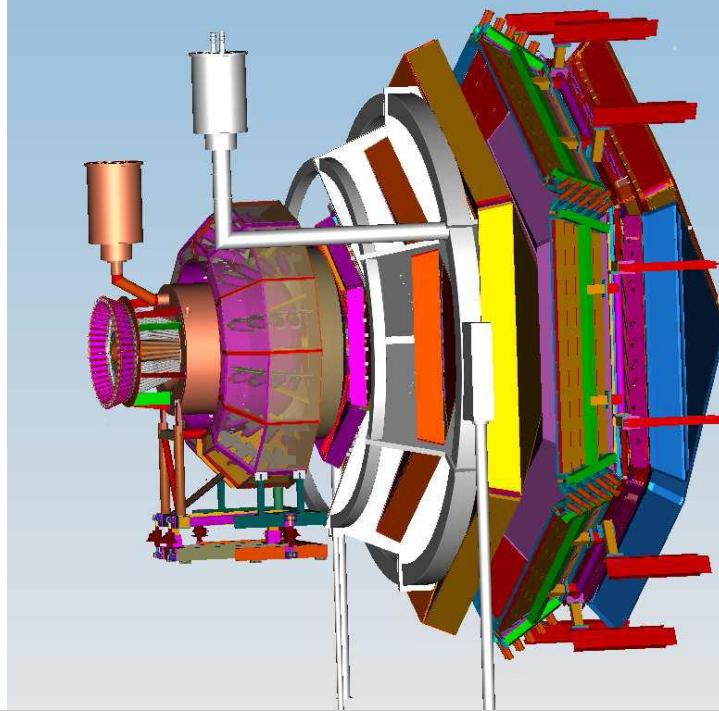


C12-15-004: DVCS, DIS, and SIDIS with a longitudinally polarized deuterium target

Silvia Niccolai



PAC44, JLab, 7/26/2016



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Co-spokespersons:

- A. Biselli (Fairfield U),
- C. Keith (JLab),
- S. Kuhn (ODU)
- S. Niccolai (IPN Orsay),
- S. Pisano (INFN Frascati),
- D. Sokhan (Glasgow U)



A CLAS collaboration proposal

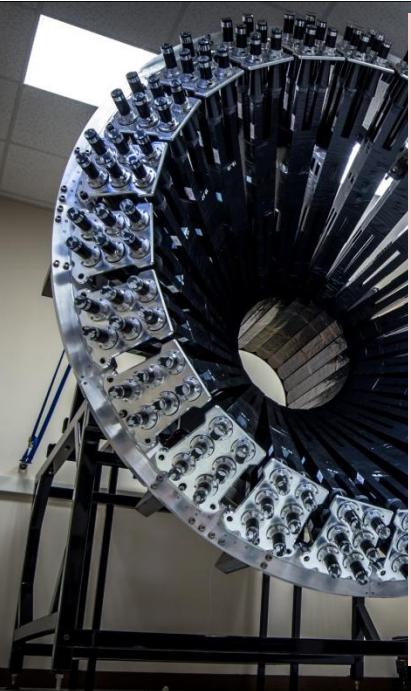
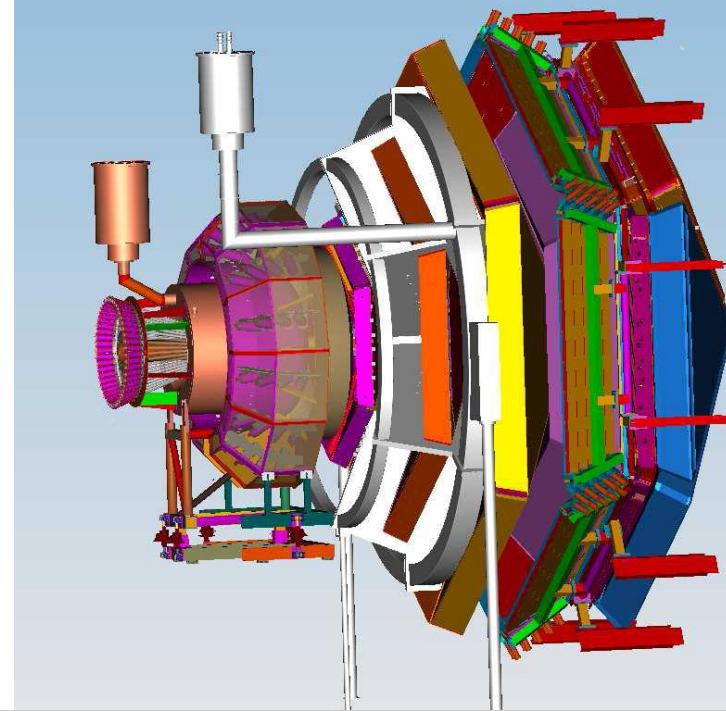


C12-15-004: DVCS, DIS, and SIDIS with a longitudinally polarized deuterium target

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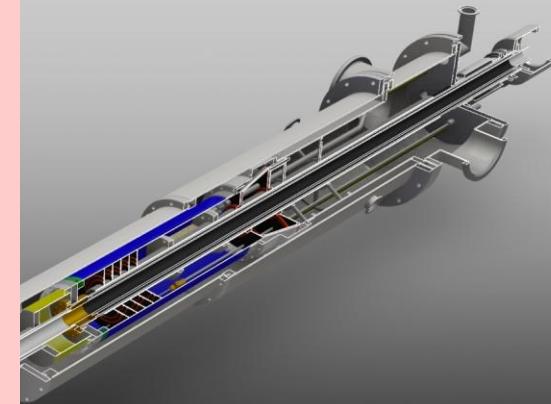


PAC44, JLab, 7/26/2016

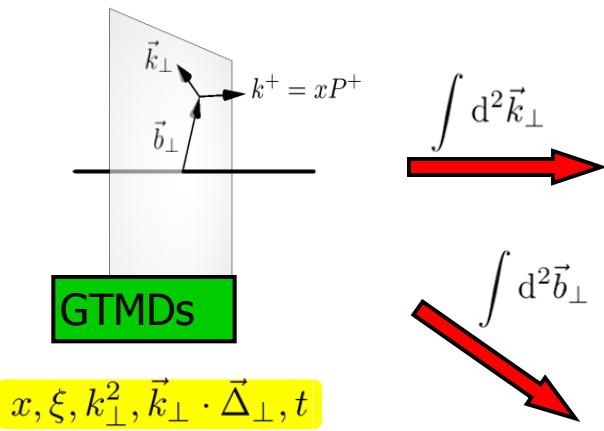
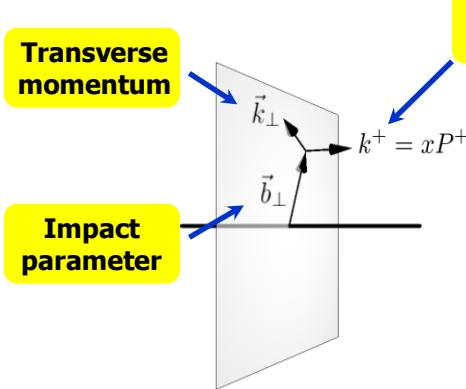


Outline:

- Multi-dimensional mapping of the nucleon
- Flavor separation of parton distributions
- Polarized target experiments with CLAS12
- Run group M: beam time request
- Experimental setup
- Physics cases:
 - ✓ Polarized nDVCS
 - ✓ SIDIS
 - ✓ DIS
 - ✓ Additional physics topics
- Conclusions/summary



Multi-dimensional mapping of the nucleon



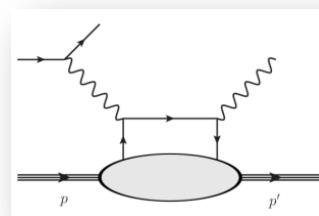
A complete picture of nucleon structure requires the measurement of all these distributions

DVCS

GPDs

TMDs

SIDIS

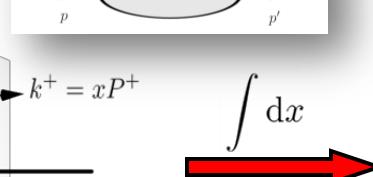


Elastic Scattering

FFs

PDFs

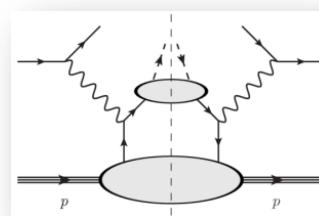
DIS



$$\int dx$$

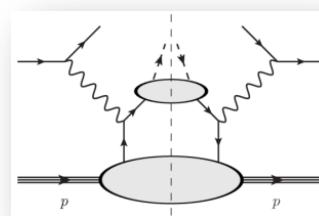
$$\int d^2 \vec{b}_\perp$$

$$\int d^2 \vec{k}_\perp$$



$$x, \xi, t$$

$$x, k_\perp^2$$



Flavor decomposition of parton distributions

- Proton and neutron have **different flavor compositions** → observables on these two target types are linked to different mixtures of **quark structure functions**

$$A_l^{p,n}(x) = \frac{\sum e_i^2 \Delta q_i(x, Q^2)}{\sum e_i^2 q_i(x, Q^2)}$$

Virtual-photon asymmetry from DIS

$$A_p^{\pi^+ - \pi^-}(x) = \frac{4\Delta u_v(x) - \Delta d_v(x)}{4u_v(x) - d_v(x)} \quad A_d^{\pi^+ - \pi^-}(x) = \frac{\Delta u_v(x) + \Delta d_v(x)}{u_v(x) - d_v(x)}$$

SIDIS pion asymmetry
(integrated over k_T)

$$\begin{aligned} \sigma_{KM}^{K^+}(p) &= 4h_{1L}^{\perp u} H_1^{\perp(1/2)u/K^+} + h_{1L}^{\perp d} H_1^{\perp(1/2)d/K^+} + h_1^{\perp \bar{s}} H_1^{\perp(1/2)\bar{s}/K^+} \\ \sigma_{KM}^{K^+}(n) &= 4h_{1L}^{\perp d} H_1^{\perp(1/2)u/K^+} + h_{1L}^{\perp u} H_1^{\perp(1/2)d/K^+} + h_1^{\perp \bar{s}} H_1^{\perp(1/2)\bar{s}/K^+} \end{aligned}$$

SIDIS $A_{UL}^{\sin 2\phi}$ for K^+

$$\mathcal{H}_p(\xi, t) = \frac{4}{9} \mathcal{H}_u(\xi, t) + \frac{1}{9} \mathcal{H}_d(\xi, t) \quad \mathcal{H}_n(\xi, t) = \frac{1}{9} \mathcal{H}_u(\xi, t) + \frac{4}{9} \mathcal{H}_d(\xi, t)$$

Compton form factors
from DVCS

- Experiments on **proton and deuterium targets** must be performed to extract the flavor dependence of the various kinds of parton distributions
- **Longitudinally polarized observables**: nuclear targets, such as **NH₃** and **ND₃** are used
- ND₃ typically has ½ the polarization of NH₃ (~40%)

ND₃ experiments with CLAS12: history

- E12-06-109 (Polarized DIS): **30 days on NH₃** (+overhead), **50 days on ND₃** (+overhead) (**run group C**)
- E12-06-119b (DVCS on longitudinally polarized protons): **120 days on NH₃** (**run group C**)
- E12-07-107 (SIDIS with long. pol. target), E12-09-007b, E12-09-009 (K-SIDIS with long. pol. target)
→ beam time for **ND₃** ~40% of that for **NH₃**

PR12-15-004 (« Deeply virtual Compton scattering on a longitudinally polarized deuterium target »), requesting **100 days** + overhead on **ND₃** was **conditionally approved (C2)** by **PAC 43**

*“...To obtain full approval, the collaboration needs to fulfill two conditions. One would be the submission of a **Run Group proposal**, connected to RG-C, that has been fully vetted according to standard procedures in the CLAS12 collaboration. The second would be the **submission of a new proposal, defining a new run group**, for the extended running time, optimized for this measurement (for example with increased neutron detection efficiency), and possibly incorporating other experiments. The PAC encourages the collaboration to consider the opportunities, and looks forward to understanding the full physics potential of a new run group.”*

- ✓ PR12-06-109a, polarized nDVCS using the **existing 50 days of Run group Cb**; submitted to PAC44, after CLAS review and approval
- ✓ **C12-15-004: New run group M: three experiments, inclusion of FT to optimize nDVCS**

CLAS12 Run group M

Production data taking at $10^{35}\text{cm}^{-2}\text{s}^{-1}$ on ND ₃ (10 nA)	50 days
Production data taking at $0.5 \cdot 10^{35}\text{cm}^{-2}\text{s}^{-1}$ on ND ₃ (5 nA)	10 days – with Forward Tagger
Target work	4 days
Production data taking on C ₁₂ target	5 days
Moeller polarimeter runs	1 day
Configuration change	3 days
Total beam time request	73 days

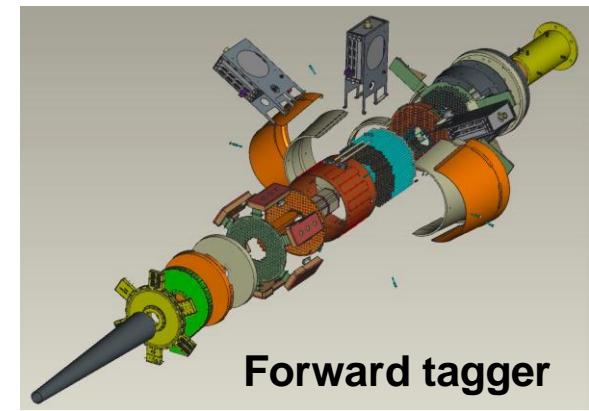
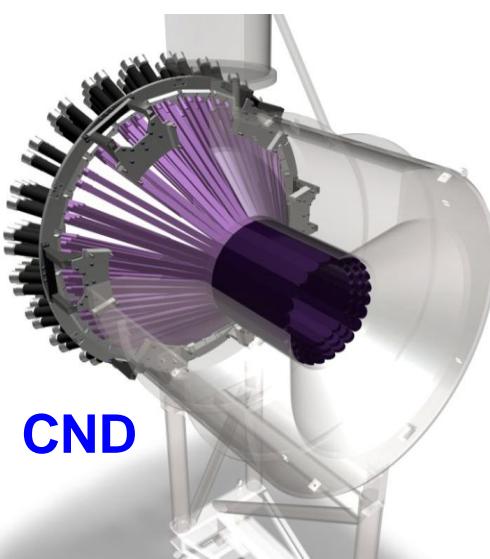
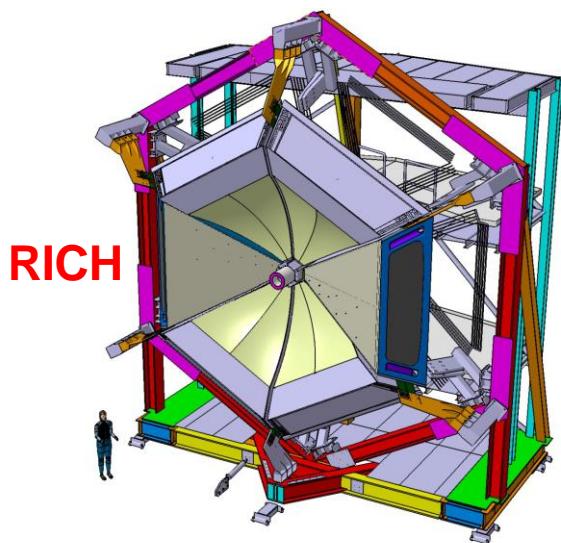
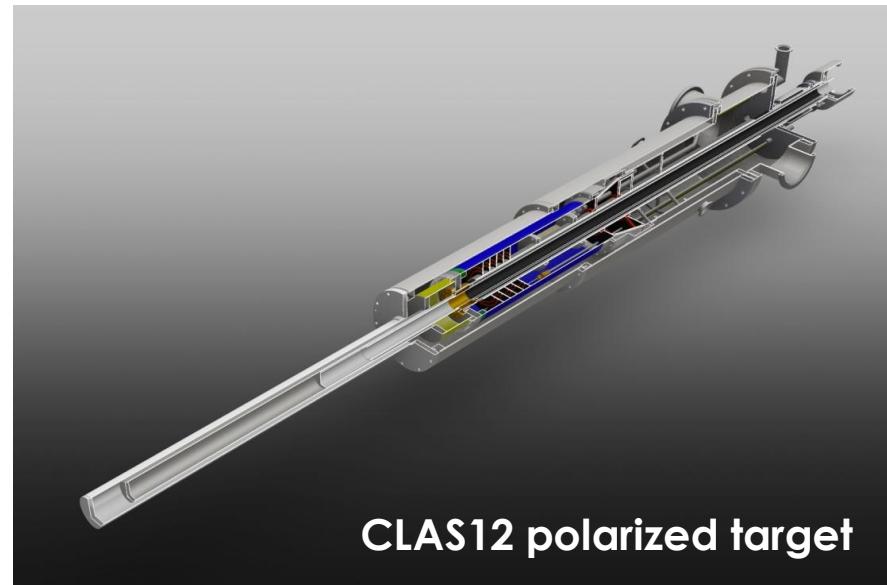
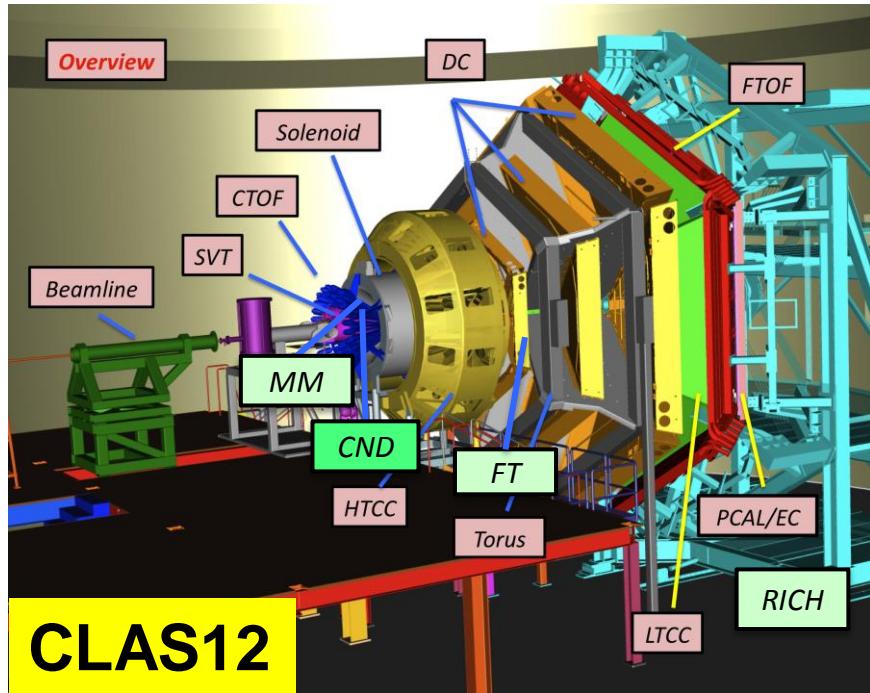
Beam energy: 11 GeV - Beam polarization: 85%

Assuming a total of 60 days (50 on ND₃ + 10 of overhead) for RG Cb
→ 133 days of beam time for polarized deuterium (110 on ND₃ + 23 of overhead)

Physics topics:

- Polarized nDVCS (GPDs)
 - Polarized SIDIS (TMDs)
 - Polarized DIS (polarized PDFs)
- Additional topics: polarized nTCS, polarized nDVMP, di-hadron SIDIS

Experimental setup of Run-group M



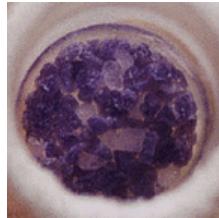
Added to optimize
nDVCS measurement

The CLAS12 longitudinally polarized target



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of VIRGINIA

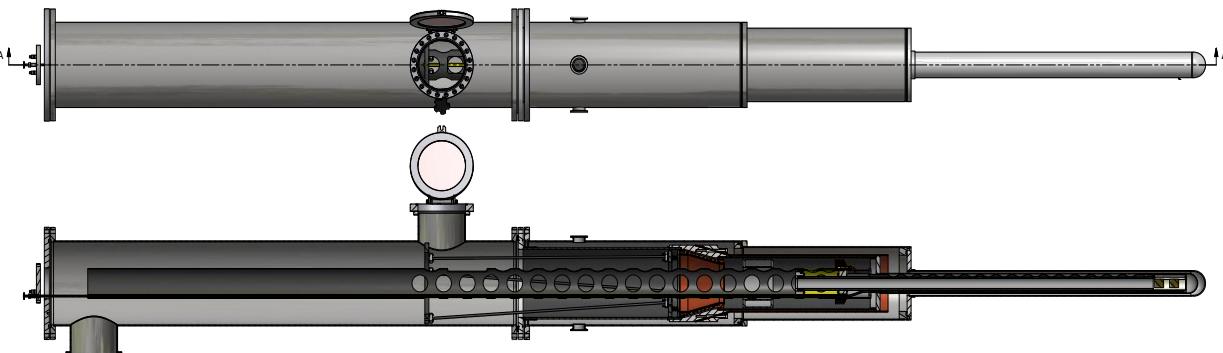
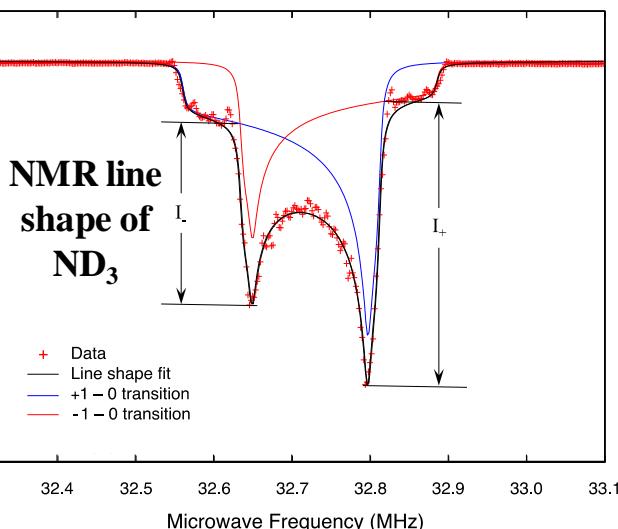
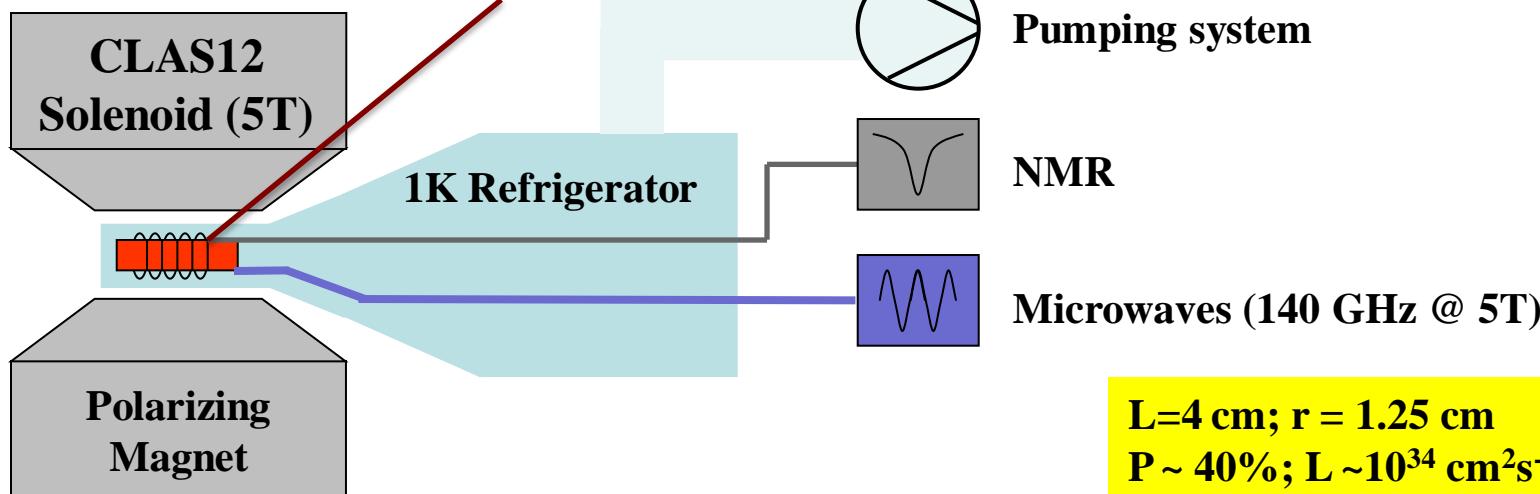
Irradiated
deuterated
ammonia
(ND₃)



Jefferson Lab
Thomas Jefferson National Accelerator Facility

OLD DOMINION
UNIVERSITY
IDEA FUSION

CHRISTOPHER
NEWPORT
UNIVERSITY

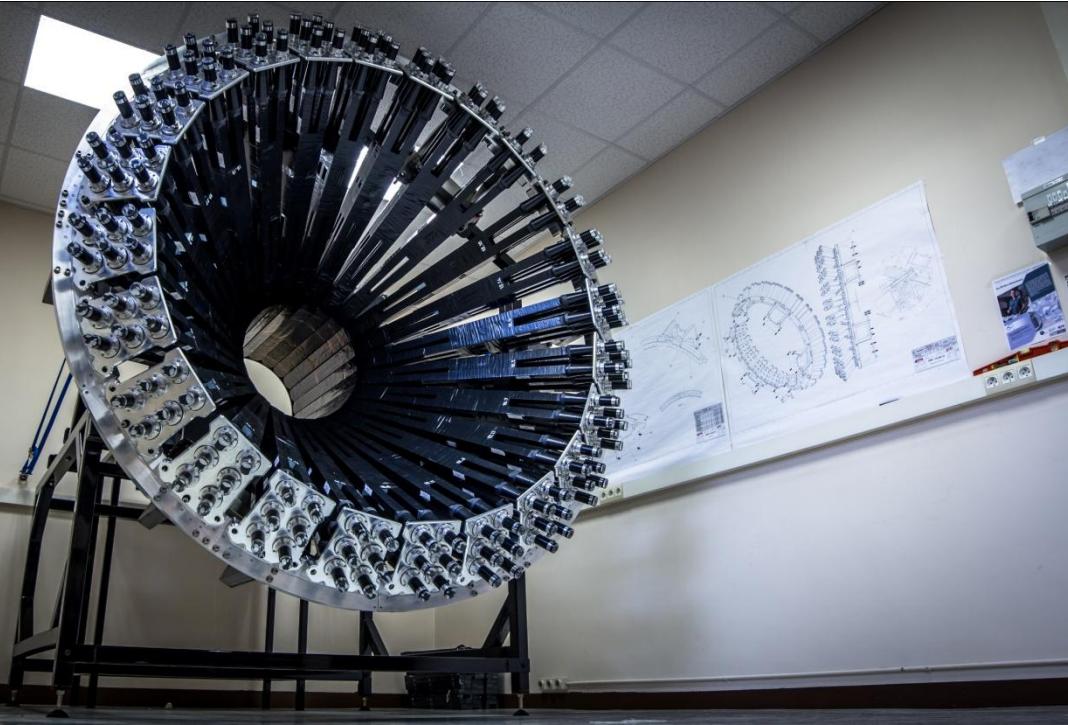


Central Neutron Detector (CND)

Main physics goal: detect the **recoiling neutron in nDVCS**

Requirements/performances:

- good **neutron/photon separation** for $0.2 < p_n < 1 \text{ GeV}/c \rightarrow \sim 150 \text{ ps time resolution}$
- momentum resolution $\delta p/p < 10\%$
- **neutron detection efficiency** $\sim 10\%$



Project status:

- Construction completed
- Detector shipped to JLab in June 2015
- Cosmic data analysis: $\sigma_t \sim 150 \text{ ps}$ for all blocks, confirmed with tests at JLab
- Readiness Review passed in June
→ CND included in CLAS12 baseline equipment

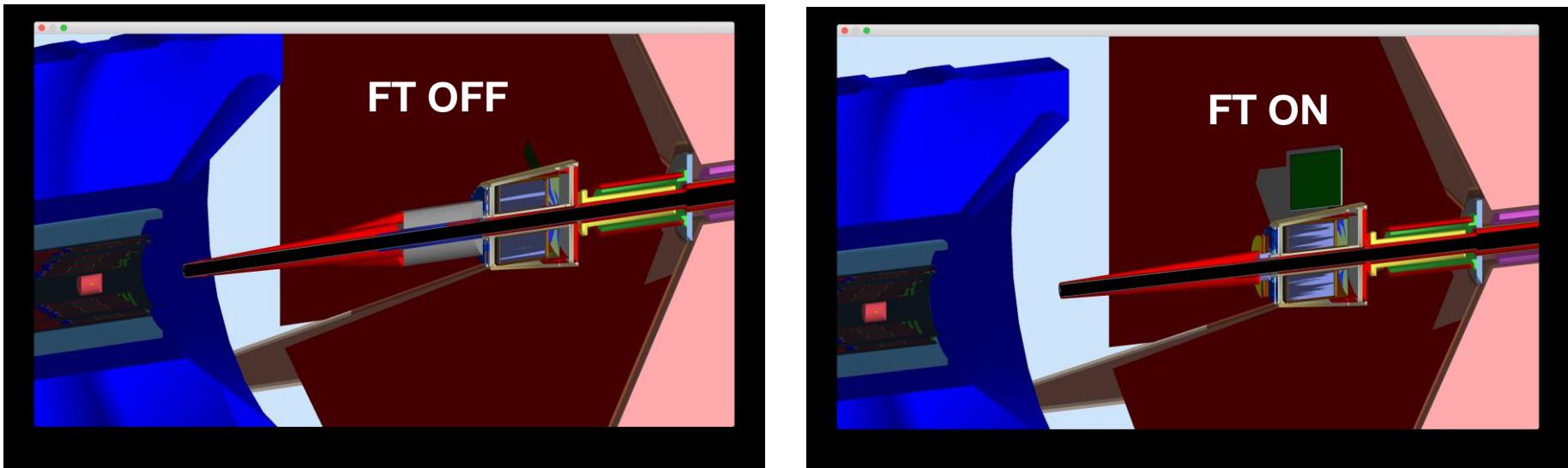


CND design: **scintillator barrel** - 3 radial layers, 48 bars per layer **coupled two-by-two** downstream by a “**u-turn**” **lightguide**, 144 long light guides with **PMTs** upstream



Forward tagger

- FT (calorimeter + hodoscope + tracker): polar angles between **2.5° and 4.5°**
→ Important for DVCS, low-angle **photons**, covering holes at $\phi \sim 0^\circ\text{-}360^\circ$
- GEMC simulations on **ND₃** target, **rastered beam**, including **background**
- With the current design of the shielding, **high occupancy** (5%) in DC-R1 at **10 nA**
→ FT on for only a **subset** of the experiment (**10 days**), with **lower current (5 nA)**



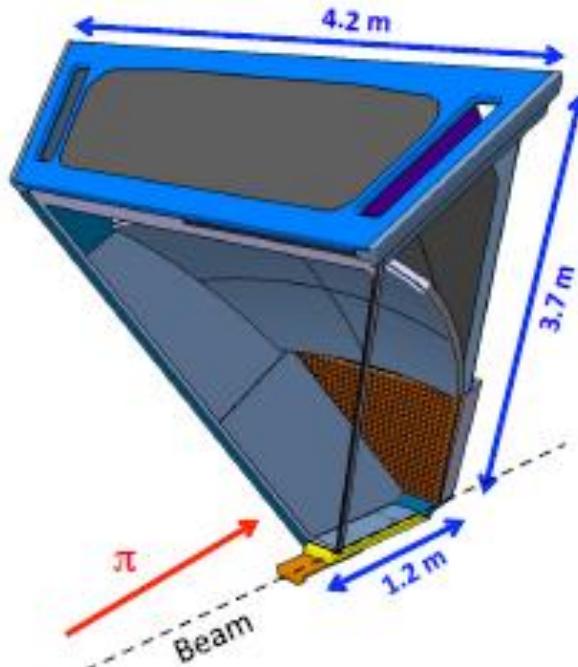
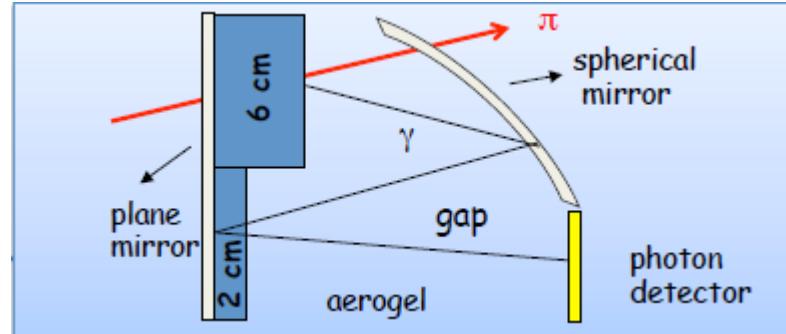
Configuration change (from Bob Miller and Chris Keith): **3 PAC days**

- 1 (calendar) day to move the SVT/solenoid/HTCC upstream
- 1 day to remove the Moller shielding
- 1 day to install the lead shield, Moller shield, outer shielding cone, and nose cone
- 1 day to move back the SVT/solenoid/HTCC into position
- 2 days to reinstall the polarized target and recover the polarization

CLAS12 RICH

Design goals:

- $\pi/K/p$ identification from 3 up to 8 GeV/c and 25°
- $\sim 4\sigma$ pion-kaon separation
- pion rejection factor $\sim 1:500$
- will cover 1 sector (2 afterwards) of CLAS12, replacing LTCC



Project overview:

- 2010: Concept of Design and Technology
- 2011: Tests of components and small prototype
- 2012: Tests of large scale prototype
- 2013: June: RICH Technical Review
September: Project Review with DOE
→ **Start Construction Phase**
- 2014: RICH Mechanical Review
- 2015: June: RICH Internal Review
October: Project Mid-term Review DOE
- 2016: **RICH Readiness Review**
- 2017: September: **Ready for Installation (1 sector)**



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FEDERICO SANTA MARIA



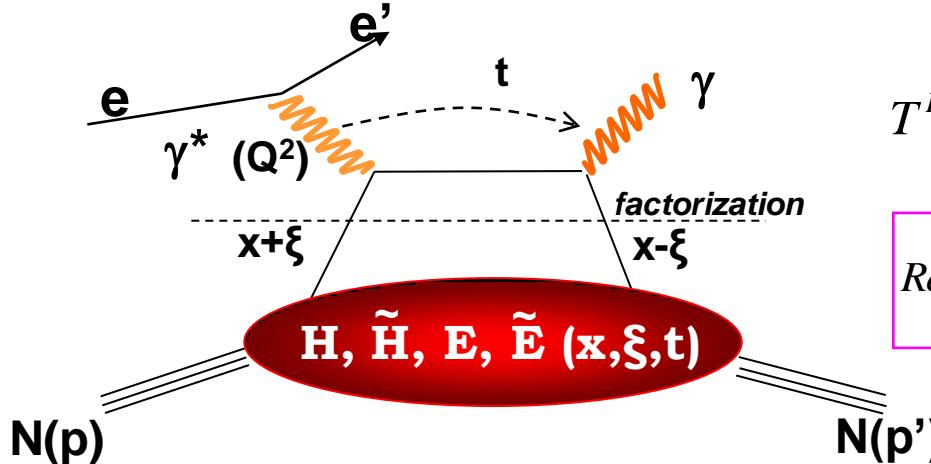
University of
Connecticut



Deeply-virtual Compton scattering on the neutron with a longitudinally polarized deuteron target

*A. Biselli (Fairfield U), C. Keith (JLab), S. Niccolai (IPN Orsay) (contact person),
S. Pisano (INFN Frascati), D. Sokhan (Glasgow U)*

Deeply Virtual Compton Scattering and GPDs



$$T^{DVCS} \sim P \int_{-1}^{+1} \frac{GPDs(x, \xi, t)}{x \pm \xi} dx \pm i\pi GPDs(\pm \xi, \xi, t) + \dots$$

$$Re \mathcal{H}_q = e_q^2 P \int_0^{+1} (H^q(x, \xi, t) - H^q(-x, \xi, t)) \left[\frac{1}{\xi - x} + \frac{1}{\xi + x} \right] dx$$

$$Im \mathcal{H}_q = \pi e_q^2 [H^q(\xi, \xi, t) - H^q(-\xi, \xi, t)]$$

At LO QCD, twist 2, chiral-even, quark sector \rightarrow **4 GPDs for each quark flavor**

conserve nucleon spin

Vector: $\mathbf{H} (\mathbf{x}, \xi, t)$

Axial-Vector: $\tilde{\mathbf{H}} (\mathbf{x}, \xi, t)$

flip nucleon spin

Tensor: $\mathbf{E} (\mathbf{x}, \xi, t)$

Pseudoscalar: $\tilde{\mathbf{E}} (\mathbf{x}, \xi, t)$

Nucleon tomography

$$q(x, \delta_\perp) = \int_0^\infty \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{i \Delta_\perp \delta_\perp} H(x, 0, -\Delta_\perp^2)$$

$$\Delta q(x, \delta_\perp) = \int_0^\infty \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{i \Delta_\perp \delta_\perp} \tilde{H}(x, 0, -\Delta_\perp^2)$$

Quark angular momentum (Ji's sum rule)

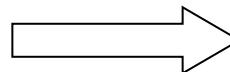
$$\frac{1}{2} \int_{-1}^1 x dx (H_q(x, \xi, t=0) + E_q(x, \xi, t=0)) = J_q = \frac{1}{2} \Delta \Sigma + \Delta L$$

X. Ji, Phy.Rev.Lett.78,610(1997)

DVCS observables and flavor separation of CFFs

Polarized beam, unpolarized target:

$$\Delta\sigma_{LU} \sim \sin\phi \operatorname{Im}\{F_1\mathcal{H} + \xi(F_1+F_2)\tilde{\mathcal{H}} - kF_2\mathcal{E}\}$$



$$\operatorname{Im}\{\mathcal{H}_p, \mathcal{H}_p, \tilde{\mathcal{E}}_p\}$$

Proton
Neutron

Unpolarized beam, longitudinal target:

$$\Delta\sigma_{UL} \sim \sin\phi \operatorname{Im}\{F_1\tilde{\mathcal{H}} + \xi(F_1+F_2)(\mathcal{H} + x_B/2\mathcal{E}) - \xi k F_2 \tilde{\mathcal{E}} + \dots\}$$



$$\begin{aligned} &\operatorname{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\} \\ &\operatorname{Im}\{\mathcal{H}_n, \tilde{\mathcal{E}}_n\} \end{aligned}$$

Polarized beam, longitudinal target:

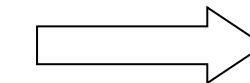
$$\Delta\sigma_{LL} \sim (A+B\cos\phi) \operatorname{Re}\{F_1\tilde{\mathcal{H}} + \xi(F_1+F_2)(\mathcal{H} + x_B/2\mathcal{E}) \dots\}$$



$$\begin{aligned} &\operatorname{Re}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\} \\ &\operatorname{Re}\{\mathcal{H}_n, \mathcal{E}_n\} \end{aligned}$$

Unpolarized beam, transverse target:

$$\Delta\sigma_{UT} \sim \cos\phi \sin(\phi_s - \phi) \operatorname{Im}\{k(F_2\mathcal{H} - F_1\mathcal{E}) + \dots\}$$



$$\begin{aligned} &\operatorname{Im}\{\mathcal{H}_p, \mathcal{E}_p\} \\ &\operatorname{Im}\{\mathcal{H}_n\} \end{aligned}$$

$$\begin{aligned} \xi &= x_B/(2-x_B) \\ k &= -t/4M^2 \end{aligned}$$

$$\mathcal{H}_p(\xi, t) = \frac{4}{9} \mathcal{H}_u(\xi, t) + \frac{1}{9} \mathcal{H}_d(\xi, t)$$

$$\mathcal{H}_n(\xi, t) = \frac{1}{9} \mathcal{H}_u(\xi, t) + \frac{4}{9} \mathcal{H}_d(\xi, t)$$



$$\begin{aligned} \mathcal{H}_u(\xi, t) &= \frac{9}{15} (4\mathcal{H}_p(\xi, t) - \mathcal{H}_n(\xi, t)) \\ \mathcal{H}_d(\xi, t) &= \frac{9}{15} (4\mathcal{H}_n(\xi, t) - \mathcal{H}_p(\xi, t)) \end{aligned}$$

Proton and neutron GPDs (and CFFs) are linear combinations of quark GPDs

A combined analysis of DVCS observables for proton and neutron targets is necessary to perform the flavor separation of the CFFs

What we learned from CLAS pDVCS asymmetries

Extraction of CFFs from *combined analysis* of **CLAS data** (TSA, BSA, DSA – *eg1dvcs*)
CFFs fitting code by M. Guidal
M. Guidal, Eur. Phys. J. A 37 (2008) 319, etc...

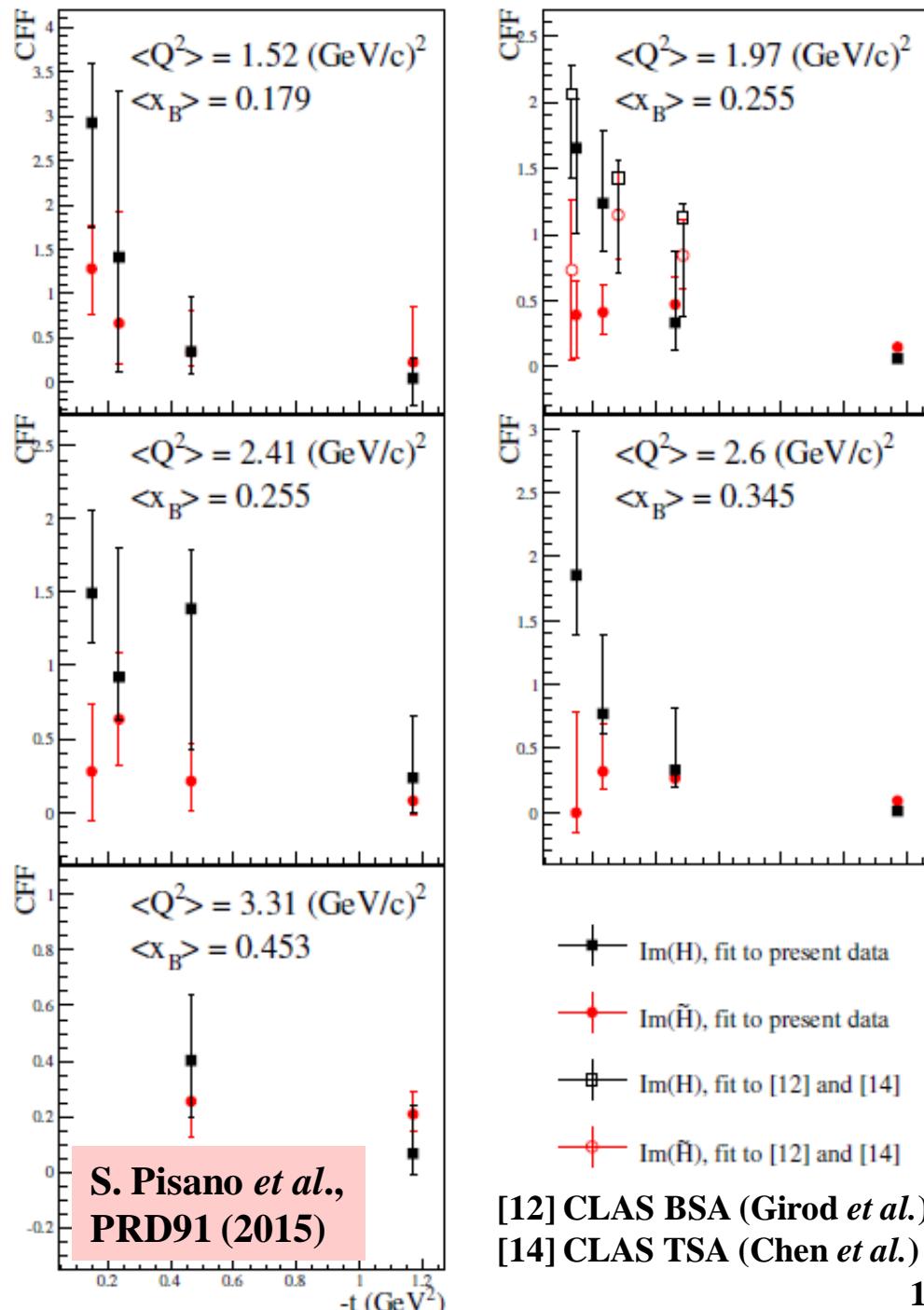
PROTON TOMOGRAPHY:

- $\text{Im}\mathcal{H}$ has steeper t -slope than $\text{Im}\tilde{\mathcal{H}}$: axial charge more “concentrated” than the electric charge
- $\text{Im}\mathcal{H}$, flatter t -slope at high x_B : faster quarks (valence) at the core of the nucleon, slower quarks (sea) at its periphery

$$\Delta q(x, b_\perp) = \int_0^\infty \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{i \Delta_\perp b_\perp} \tilde{H}(x, 0, -\Delta_\perp^2)$$

$$\int H(x, \xi, t) dx = F_1(t)$$

$$\int \tilde{H}(x, \xi, t) dx = G_A(t)$$

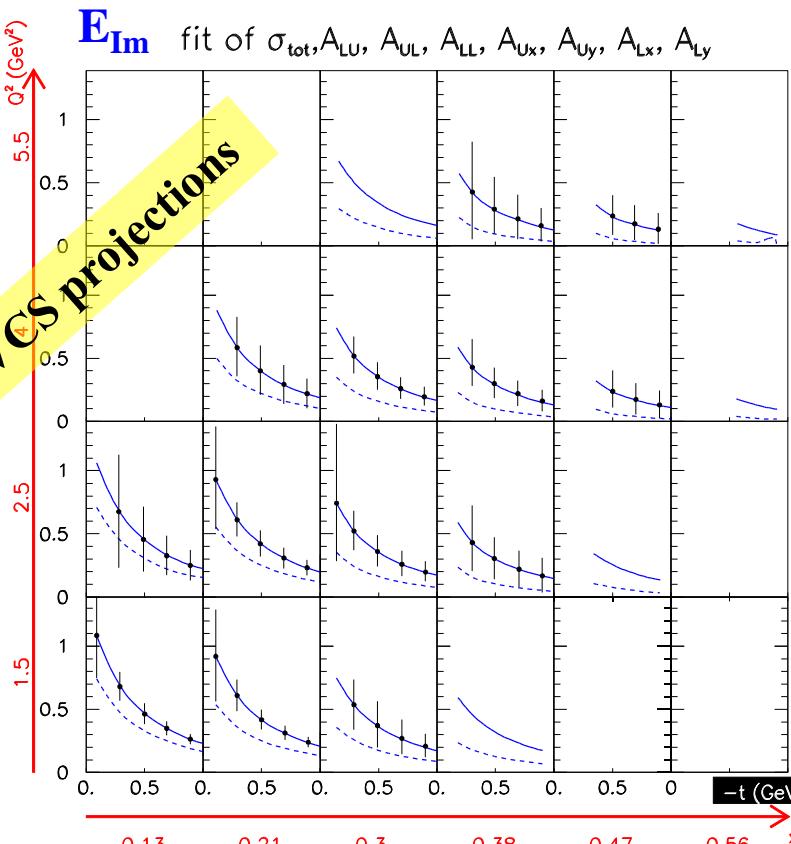
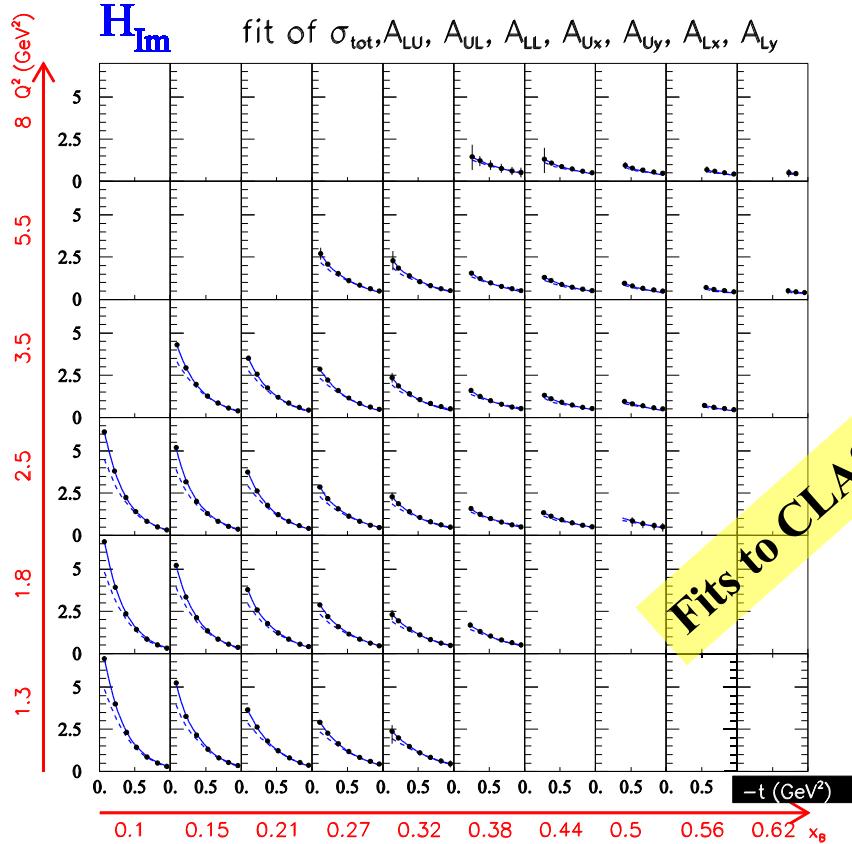


DVCS program for JLab12

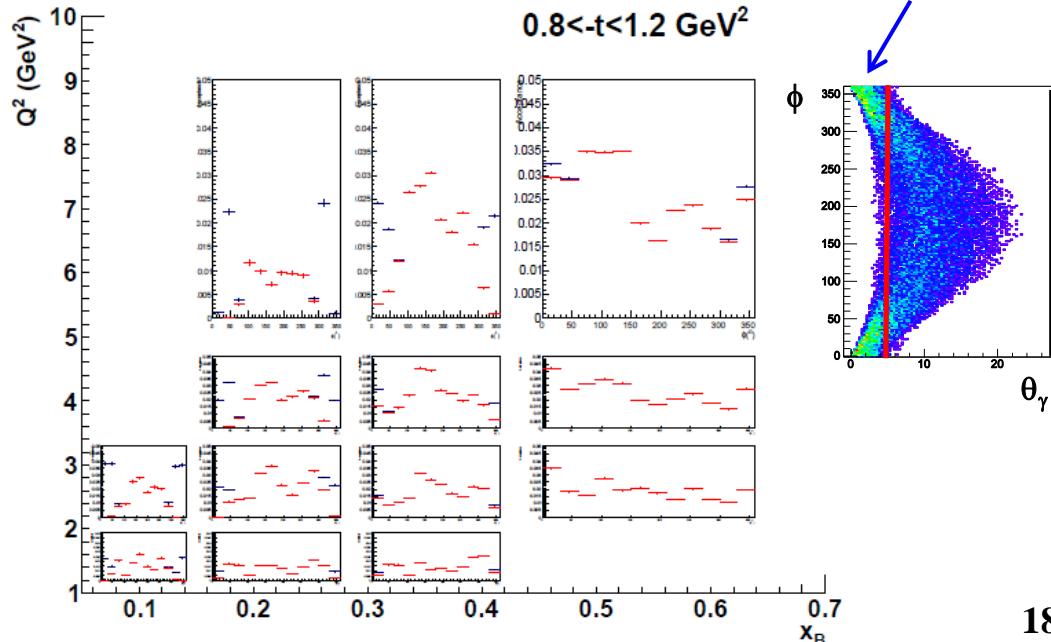
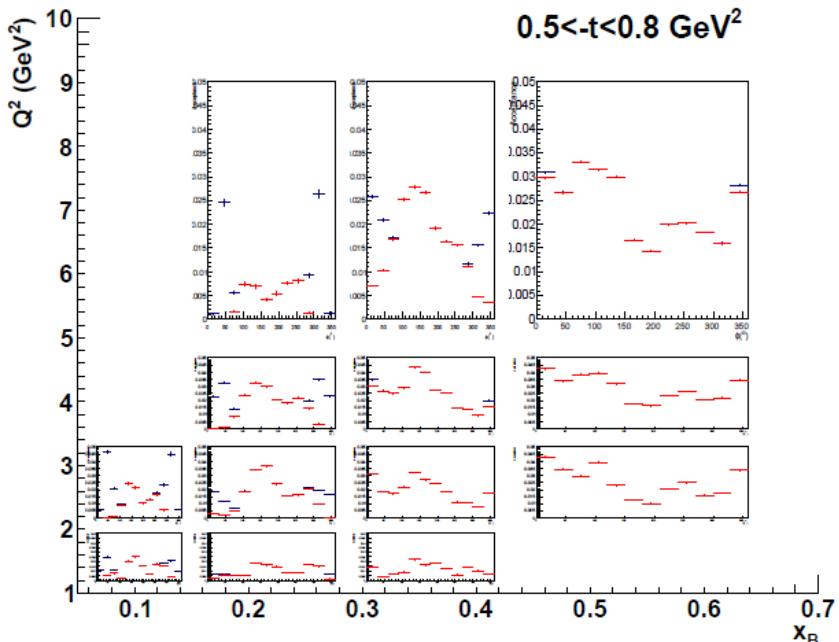
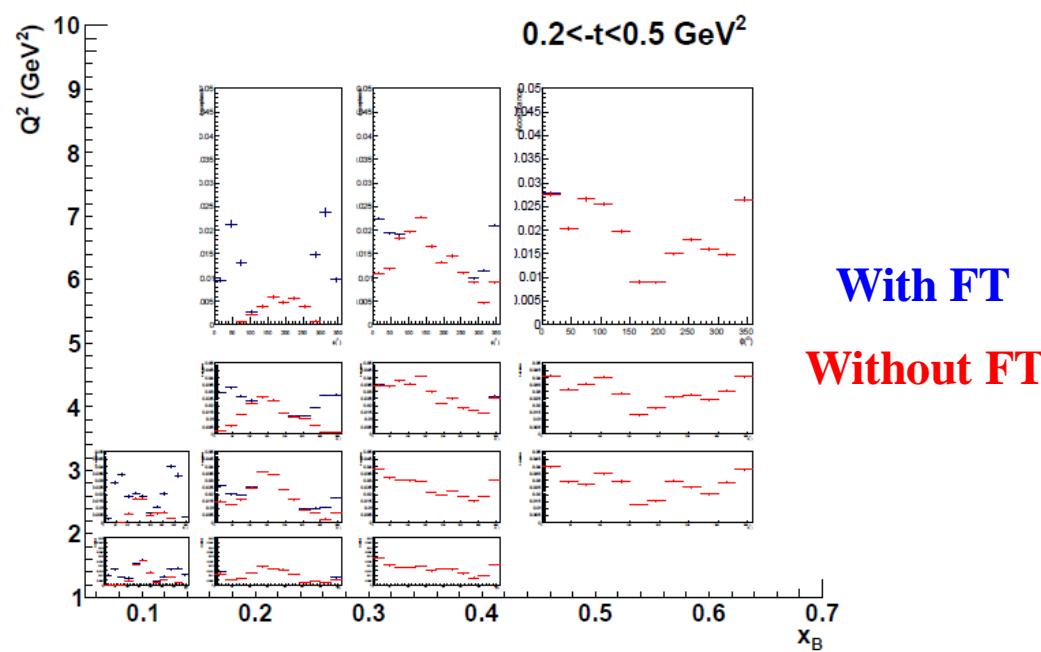
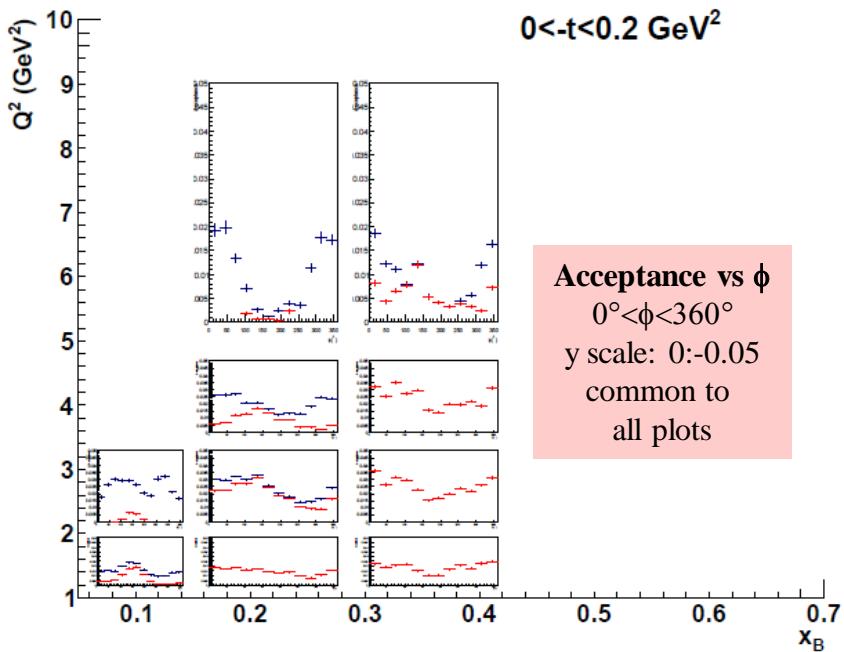
Proton
Neutron

Observable (target)	Target	Sensitivity to CFFs	Completed experiments	12-GeV experiments
$\Delta\sigma_{beam}(p)$	Unpolarized hydrogen	$\Im m \mathcal{H}_p$	Hall A, CLAS	Hall A, CLAS12, Hall C
BSA(p)	Unpolarized hydrogen	$\Im m \mathcal{H}_p$	HERMES, CLAS	CLAS12
TSA(p)	Long. pol. NH3	$\Im m \tilde{\mathcal{H}}_p, \Im m \mathcal{H}_p$	HERMES, CLAS	CLAS12
DSA(p)	Long. pol. NH3	$\Re e \mathcal{H}_p, \Re e \tilde{\mathcal{H}}_p$	HERMES, CLAS	CLAS12
tTSA(p)	Transv. pol. protons	$\Im m \mathcal{H}_p, \Im m \mathcal{E}_p$	HERMES	CLAS12
$\Delta\sigma_{beam}(n)$	Unpolarized deuterium	$\Im m \mathcal{E}_n$	Hall A	
BSA(n)	Unpolarized deuterium	$\Im m \mathcal{E}_n$		CLAS12
TSA(n)	Long. pol. ND3	$\Im m \mathcal{H}_n$		
DSA(n)	Long. pol. ND3	$\Re e \mathcal{H}_n$		

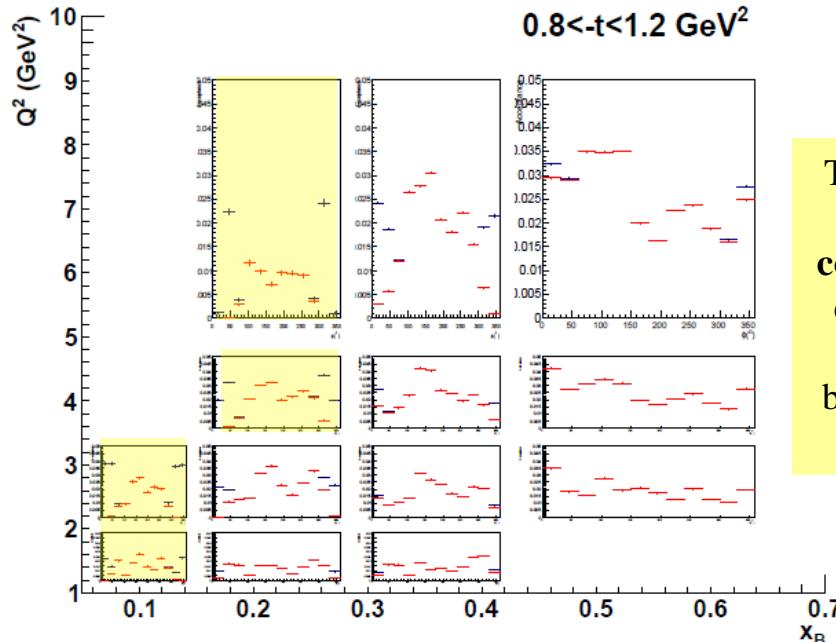
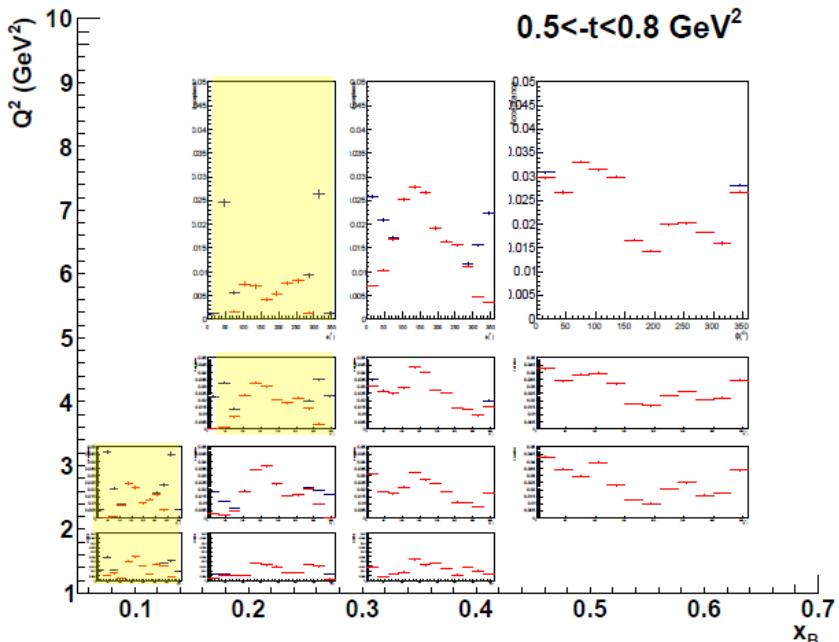
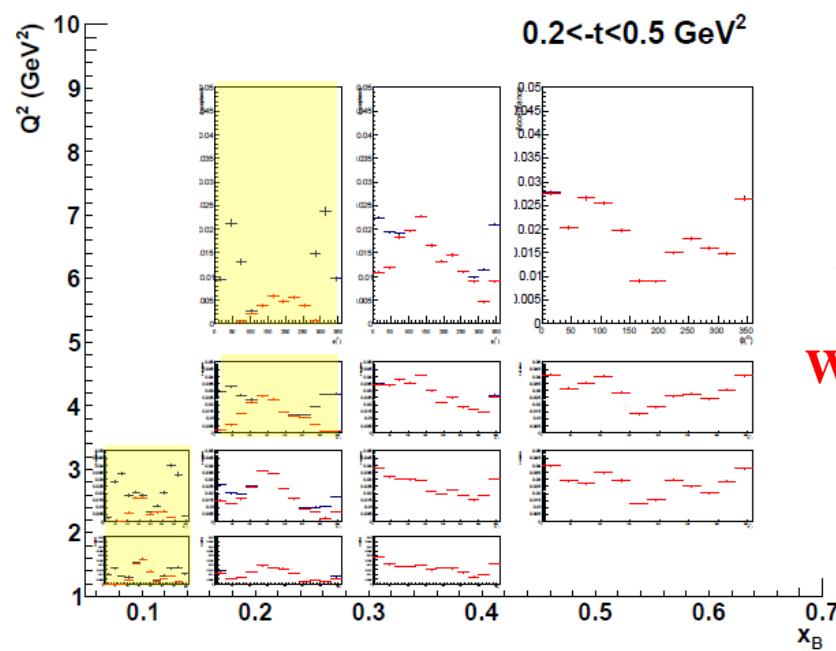
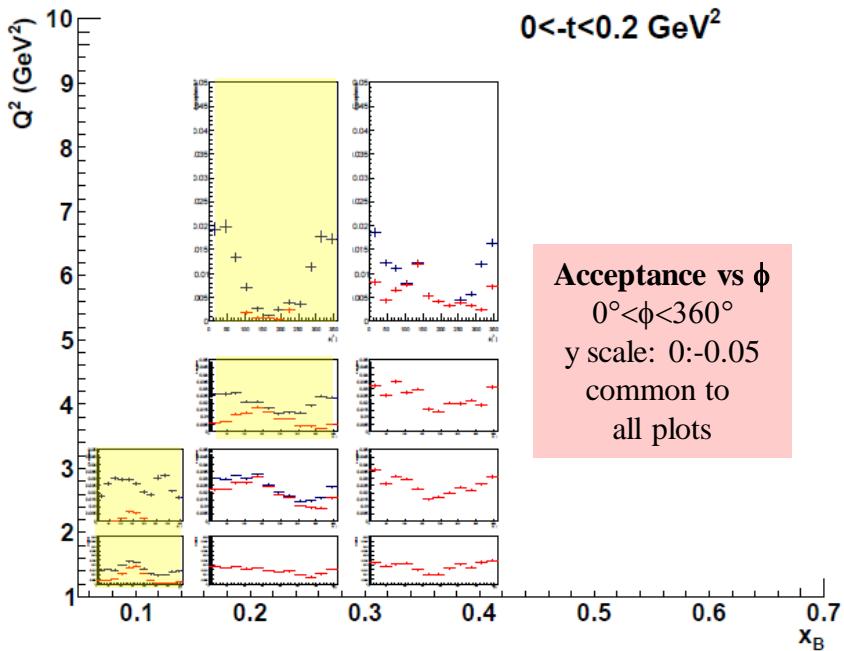
C12-15-004



Effect of the Forward Tagger on the acceptance for nDVCS



Effect of the Forward Tagger on the acceptance for nDVCS



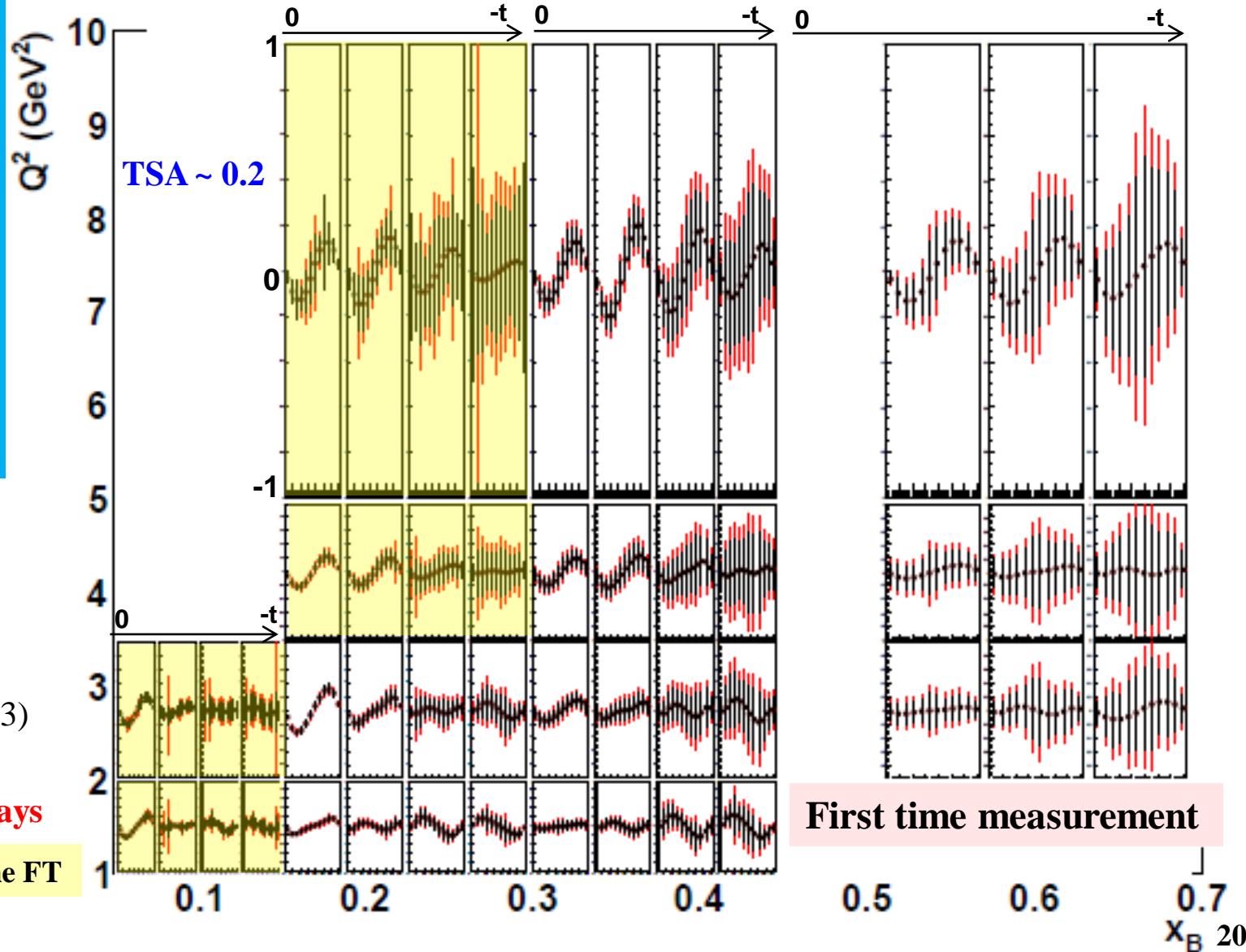
The inclusion of the FT completes the ϕ coverage, affecting 16 bins out of 49 (30%),

Projected results: target-spin asymmetry

- ✓ Count rates computed with nDVCS+BH event generator + CLAS12 + FT acceptance from FastMC + CND efficiency from GEANT4 simulation
- ✓ Asymmetries computed with VGG model

$$\sigma_A = \frac{1}{P_t} \cdot \frac{\sqrt{(1 - P_t \cdot A)^2}}{\sqrt{N}}$$

- $L = 3/20 \cdot 10^{35} \text{ cm}^{-2} \text{s}^{-1}$, Time = 100 days@10 nA + 10 days@5 nA with Forward tagger, $P_t = 0.4$

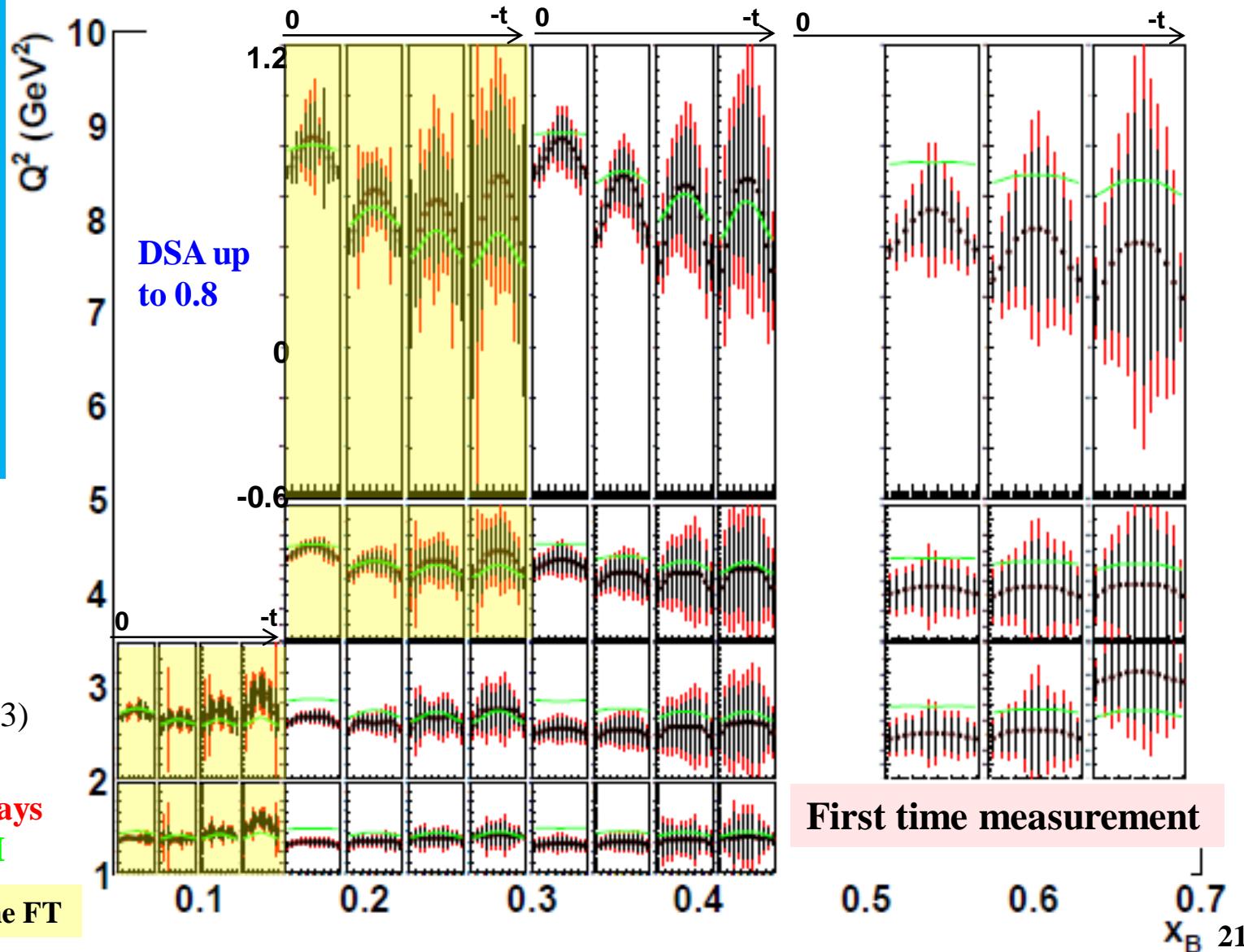


Projected results: double spin asymmetry

- ✓ Count rates computed with nDVCS+BH event generator + CLAS12 + FT acceptance from FastMC
- + CND efficiency from GEANT4 simulation
- ✓ Asymmetries computed with VGG model

$$\sigma_A = \frac{1}{P_b P_t} \cdot \frac{\sqrt{(1 - P_b P_t \cdot A)^2}}{\sqrt{N}}$$

- $L = 3/20 \cdot 10^{35} \text{ cm}^{-2} \text{s}^{-1}$, Time = 100 days@10 nA + 10 days@5 nA with Forward tagger, $P_t = 0.4$, $P_b = 0.85$



Projected beam spin asymmetry from C12-11-003

✓ Count rates computed with nDVCS+BH event generator + CLAS12 + FT acceptance from FastMC + CND efficiency from GEANT4 simulation
 ✓ Asymmetries computed with VGG model

- 4 bins in Q^2
- 4 bins in $-t$
- 4 bins in x_B
- 12 bins in ϕ

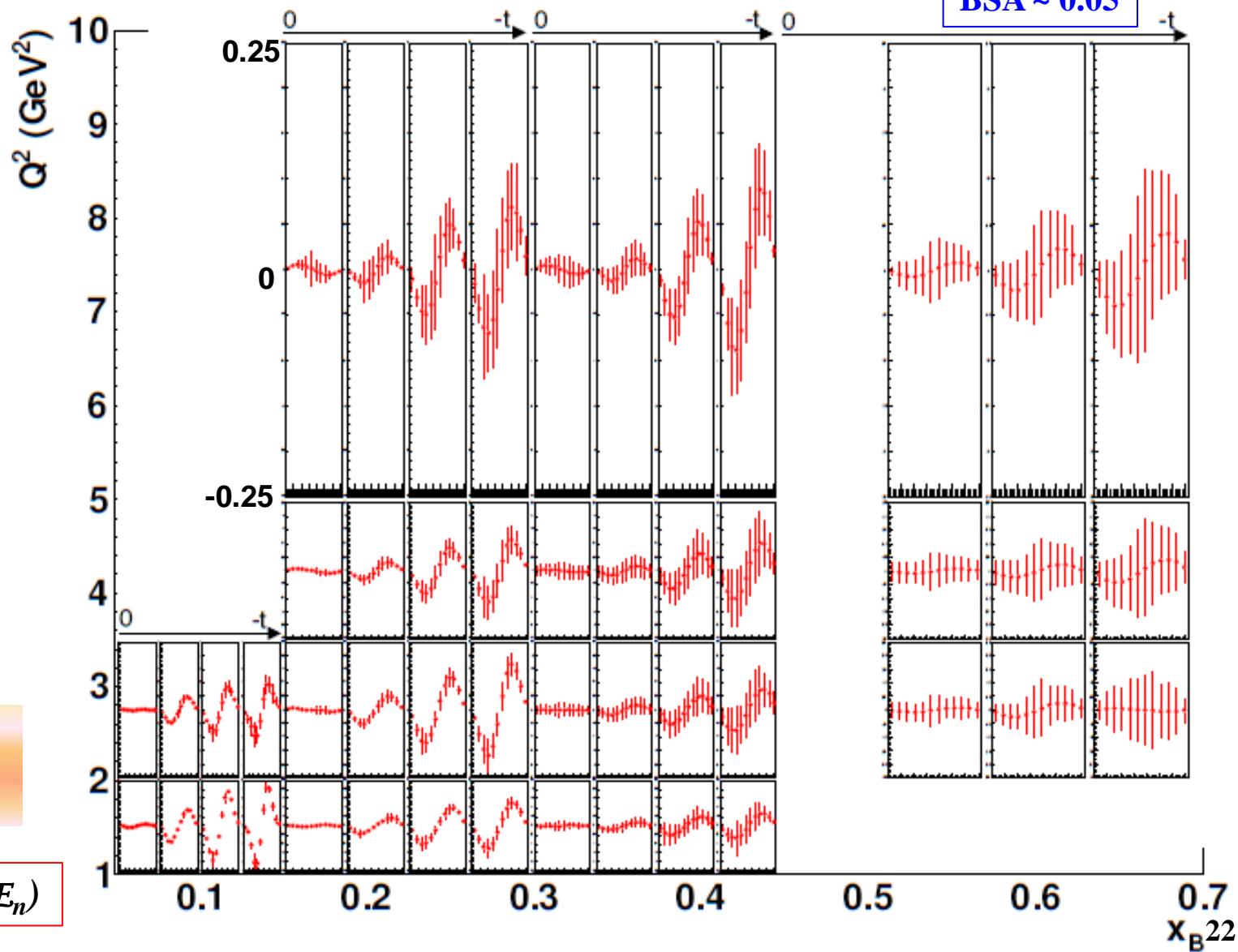
High-impact experiment

Sensitive to $Im(\mathcal{E}_n)$

$$\sigma_A = \frac{1}{P_b} \cdot \frac{\sqrt{(1 - P_b \cdot A)^2}}{\sqrt{N}}$$

- $L = 10^{35} \text{ cm}^{-2}\text{s}^{-1}$ for LD, Time = 80 days; $P_b = 0.85$

BSA ~ 0.05



Combined analysis of C12-15-004 and E12-11-003

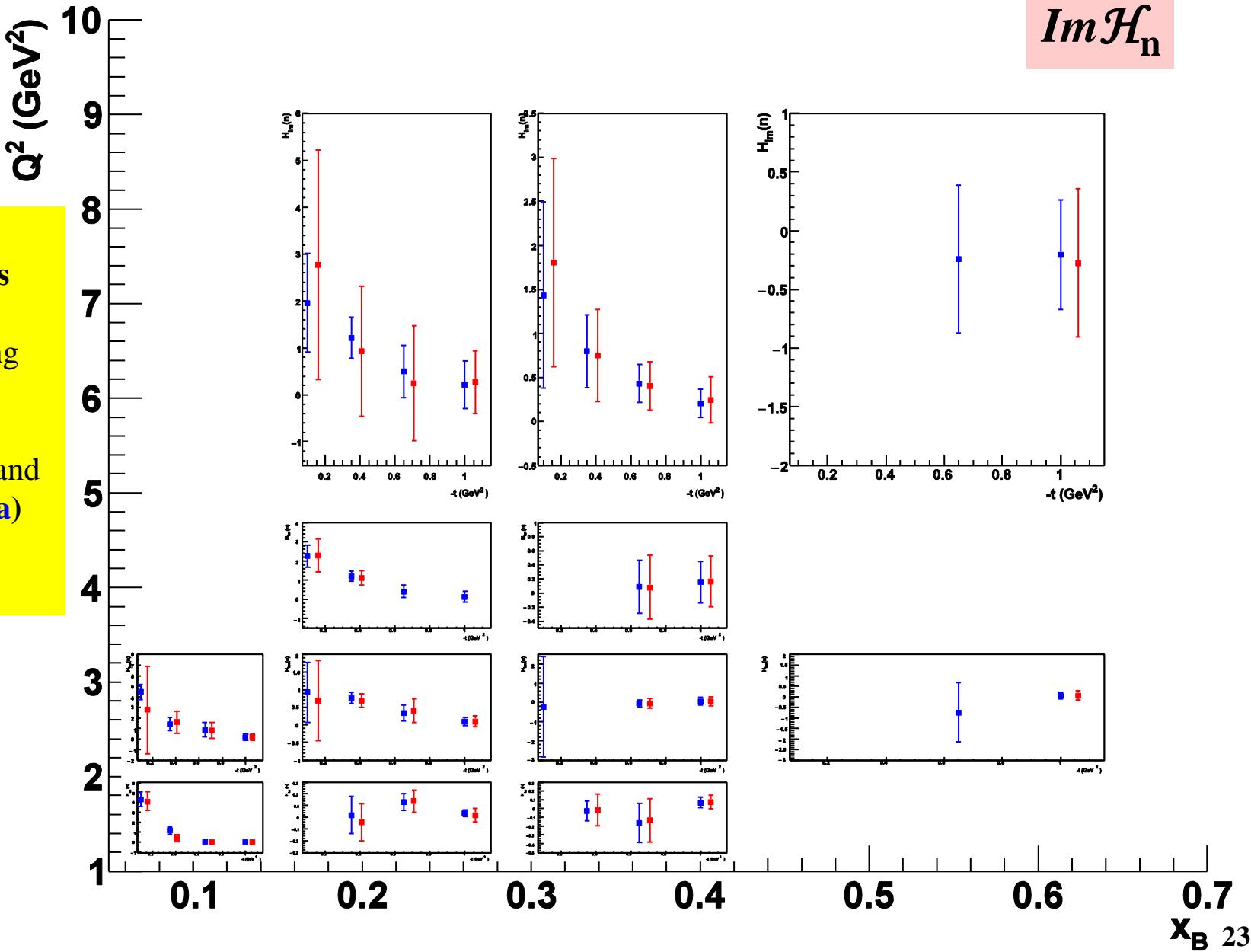
110 days
50 days

$Im\mathcal{H}_n$

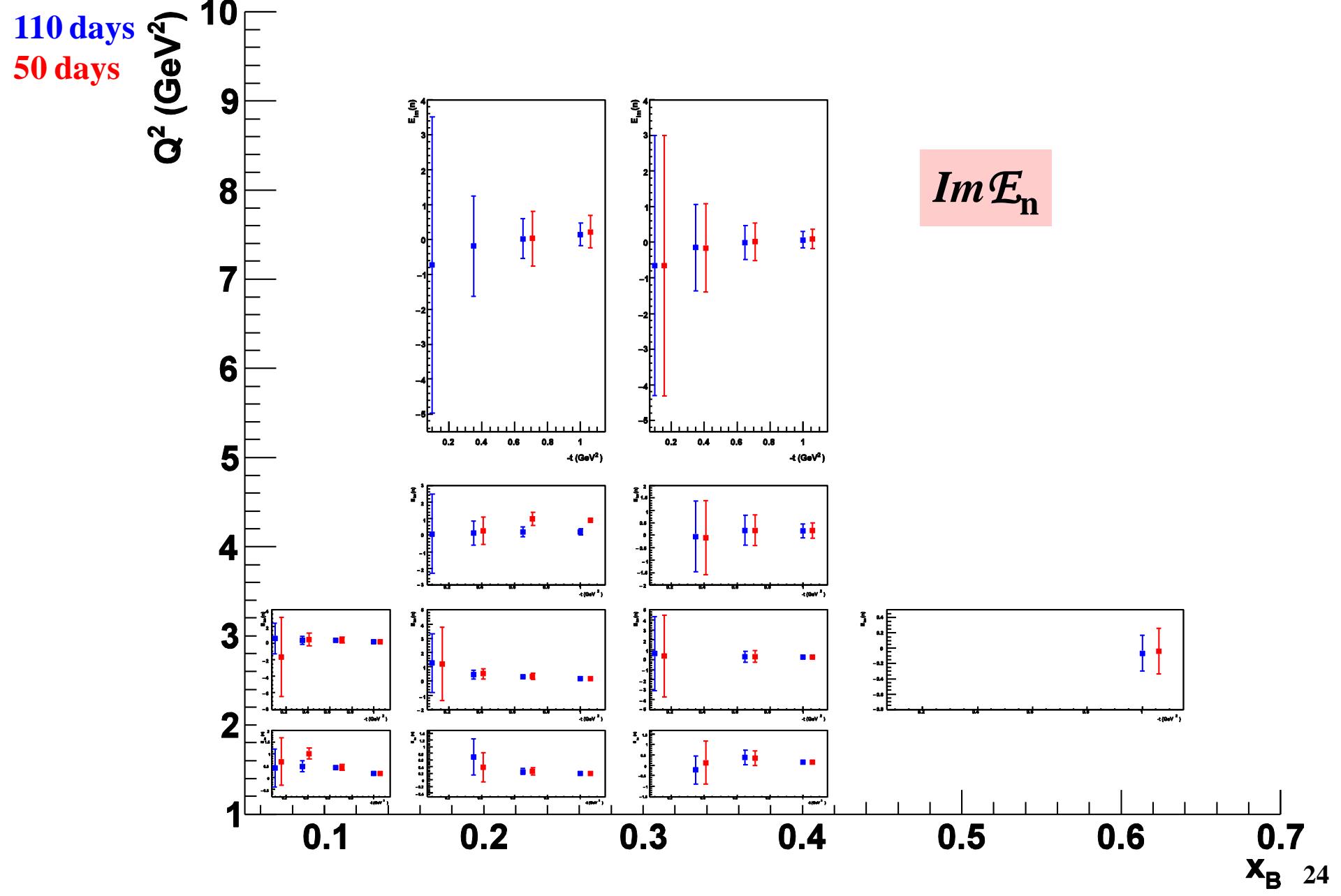
Extraction of neutron CFFs using M. Guidal's fitting code. Fit of TSA, DSA (C12-15-004 and PR12-06-109a) and BSA (E12-11-003)

7 CFFs as free fit parameters

$Im\tilde{\mathcal{E}} = 0$



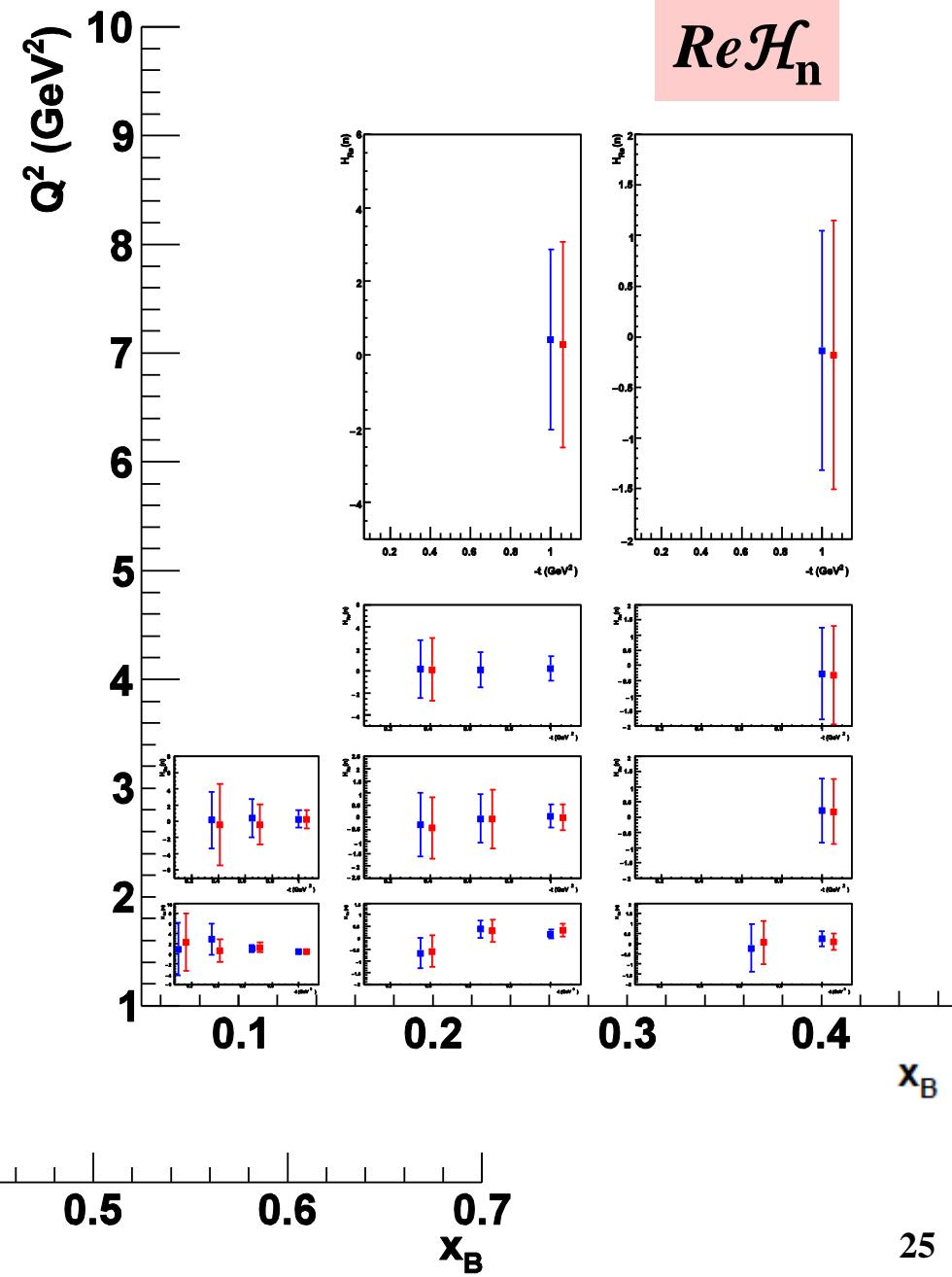
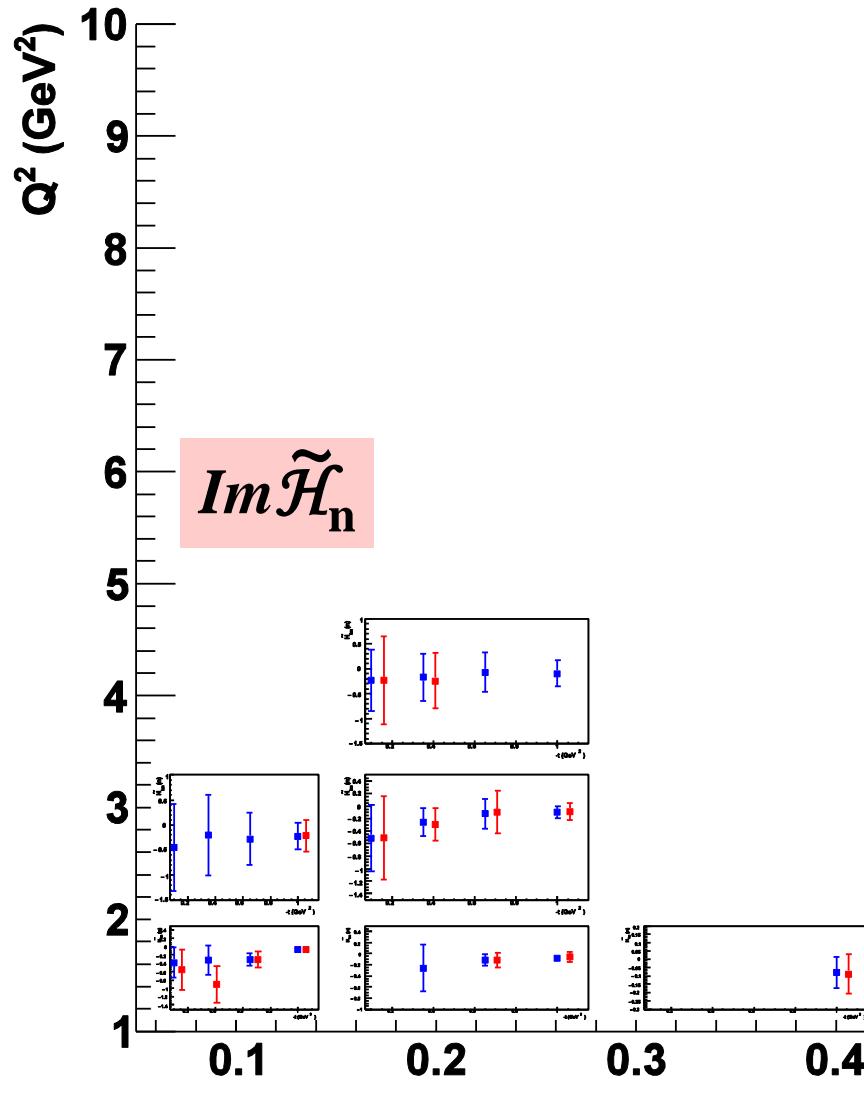
Combined analysis of C12-15-004 and E12-11-003



Combined analysis of C12-15-004 and E12-11-003

110 days

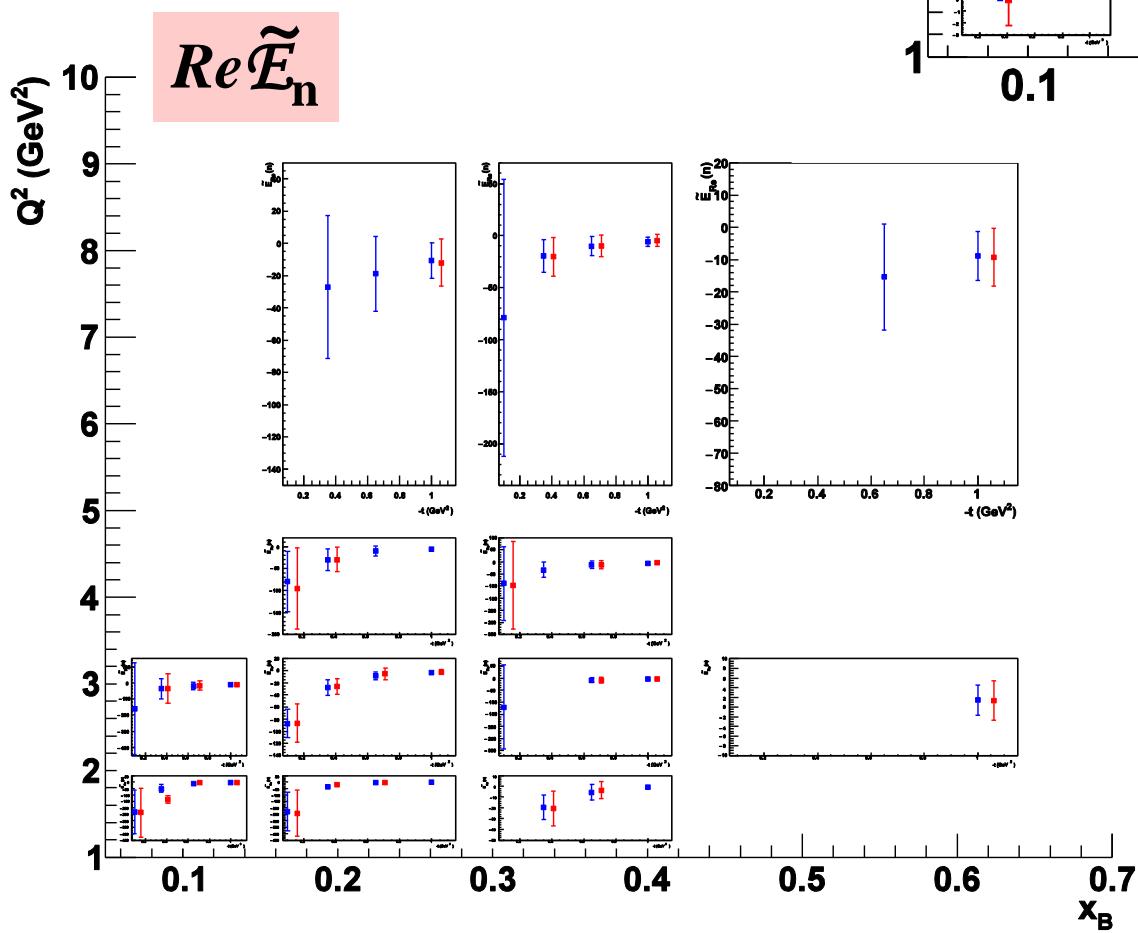
50 days



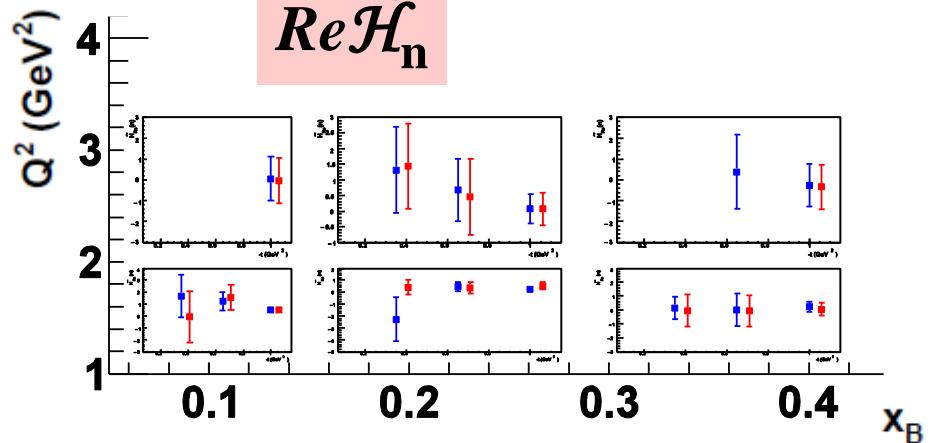
Combined analysis of C12-15-004 and E12-11-003

110 days

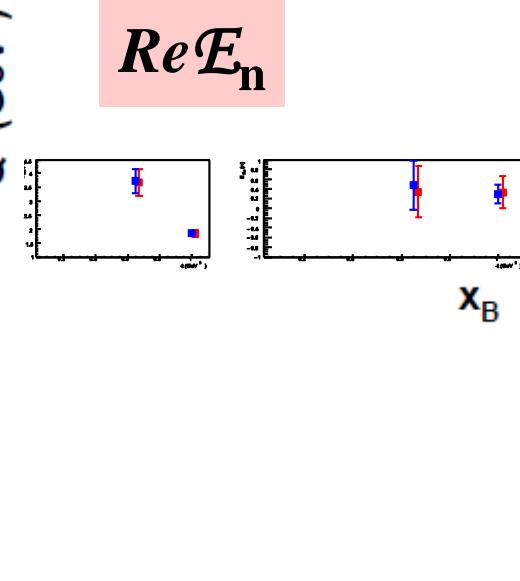
50 days



$Re \tilde{\mathcal{H}}_n$

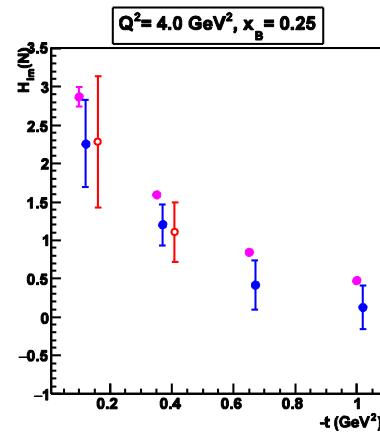
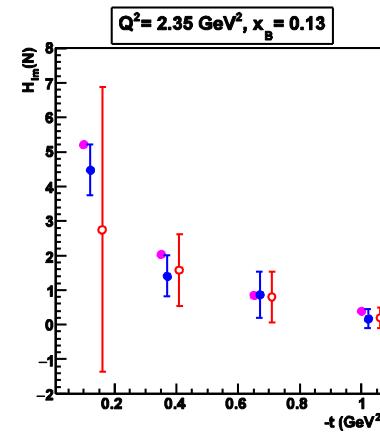
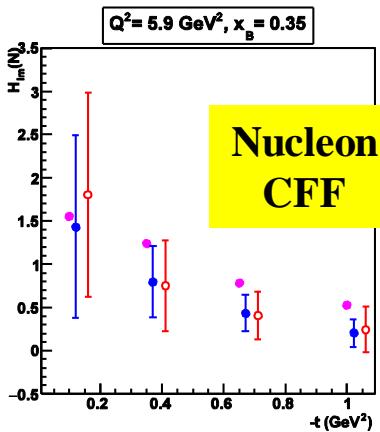
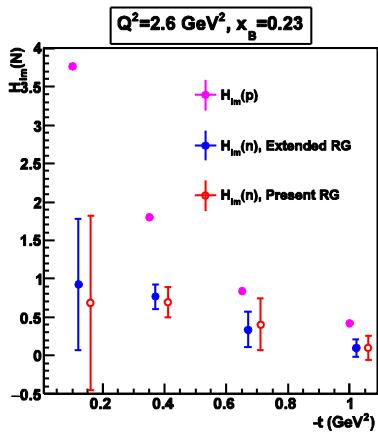


$Re \mathcal{E}_n$

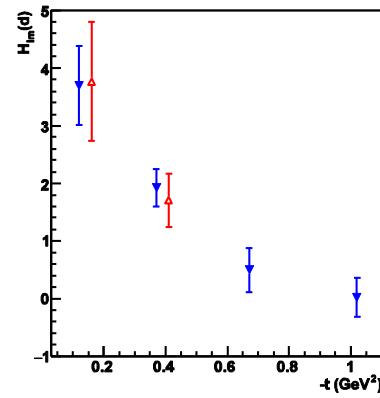
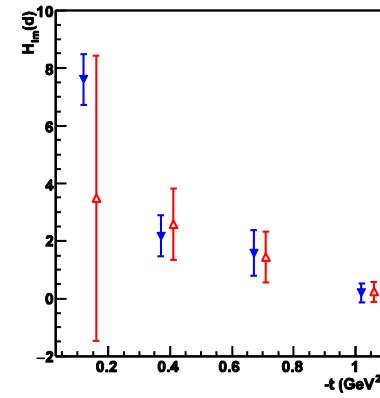
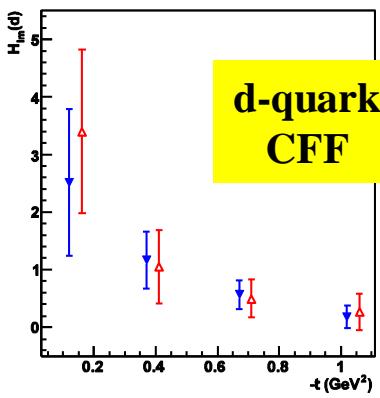
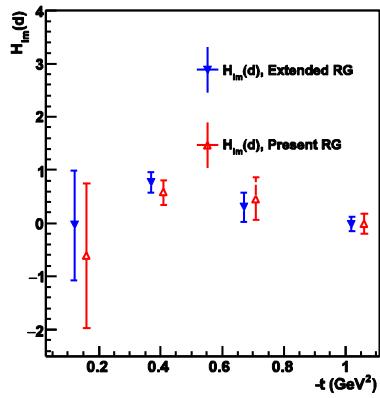


Projections for flavor separation ($Im\mathcal{H}$)

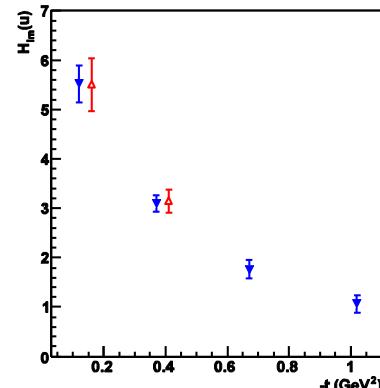
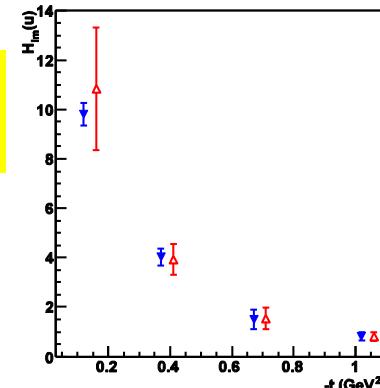
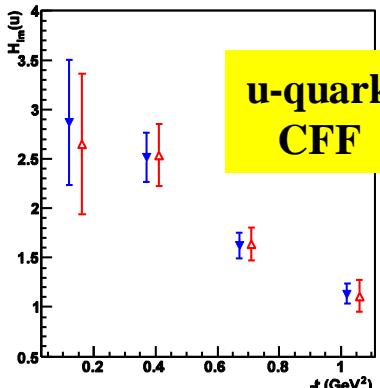
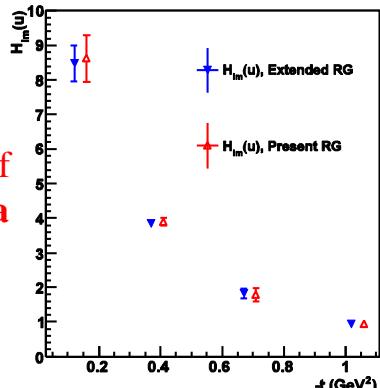
Fit using all the projected pDVCS asymmetries of the CLAS12 program



Fit using the projected nDVCS asymmetries of C12-15-004 and E12-11-003 (110 days)

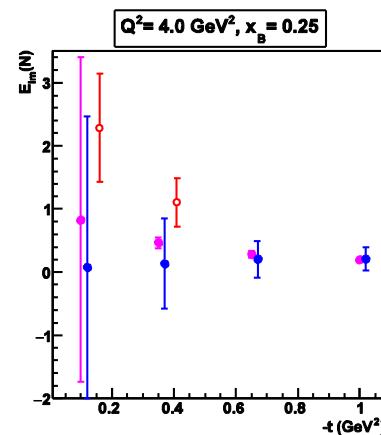
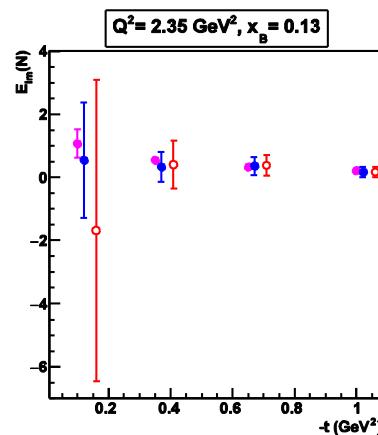
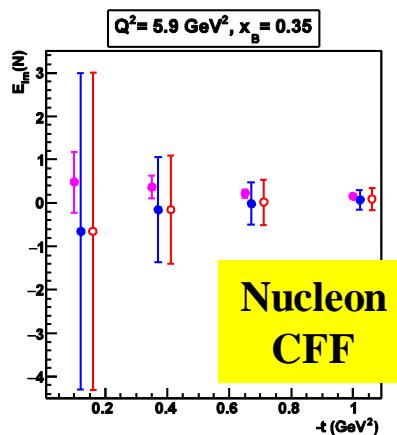
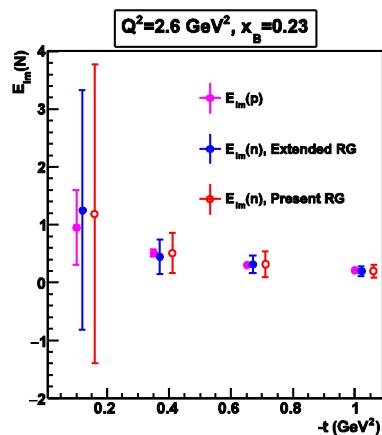


Fit using the projected nDVCS asymmetries of PR12-06-109a and E12-11-003 (50 days)

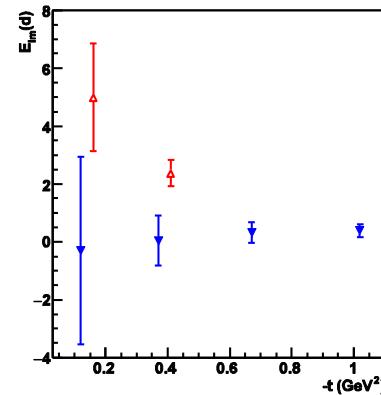
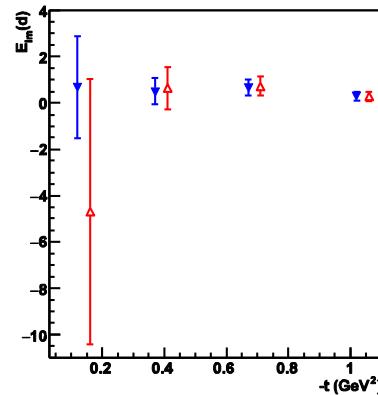
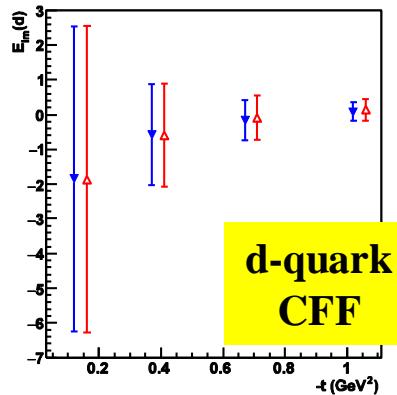
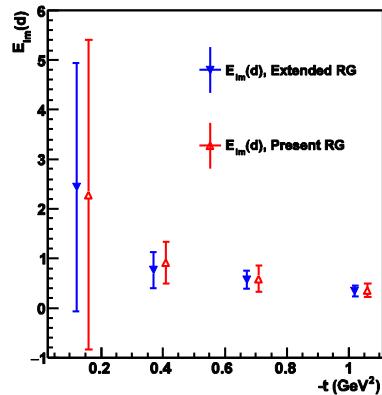


Projections for flavor separation ($\text{Im}E$)

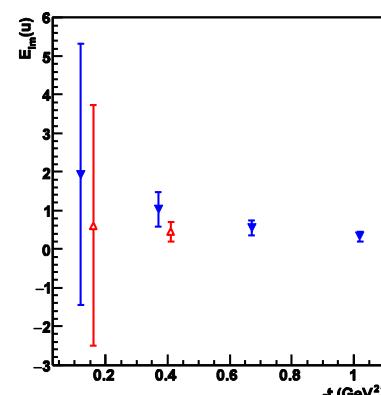
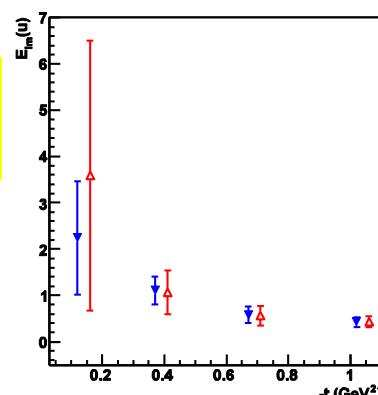
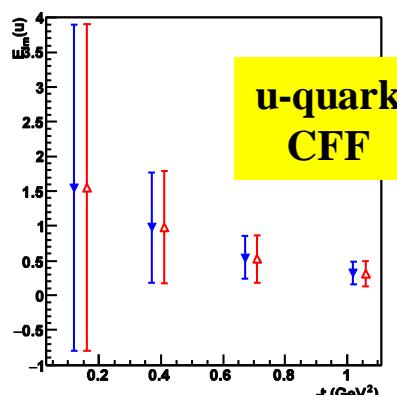
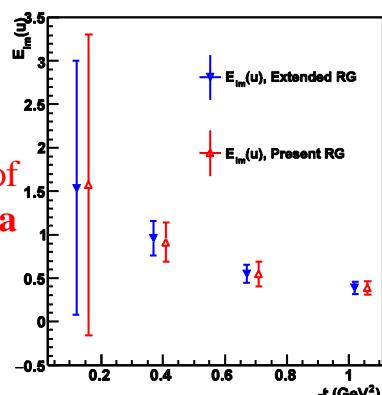
Fit using all
the projected
pDVCS
asymmetries
of the
CLAS12
program



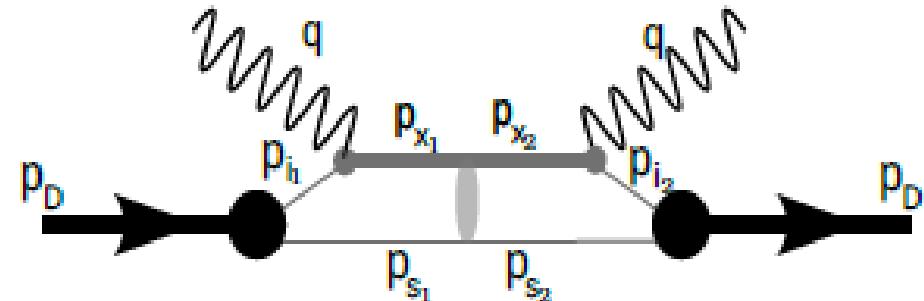
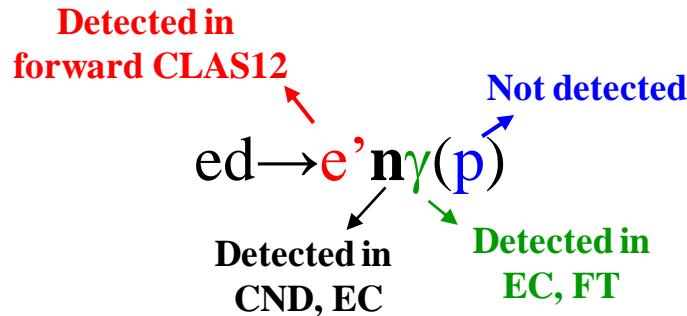
Fit using the
projected
nDVCS
asymmetries
of
C12-15-004
and
E12-11-003
(110 days)



Fit using the
projected
nDVCS
asymmetries of
PR12-06-109a
and
E12-11-003
(50 days)



Final-state interactions for nDVCS on the deuteron



FSI effects without charge exchange ($pn \rightarrow pn$) will be estimated empirically measuring **en γ and e γp on deuteron** in this same experiment and comparing with **proton data**

The effect of **charge-exchange FSI** can be **calculated** and corrected for.

M. Sargsian estimated this effect for the extraction of the **neutron form factor in polarized electron scattering from polarized He³** (S.Riordan *et al.*, PRL 105, 262302 (2010), W. Cosyn *et al.*, PRC 89, 014612 (2014)):

- $pn \rightarrow np$ interaction is due to **pion exchange** → amplitude has a $1/\sqrt{s}$ suppression factor as compared to non-charge exchange $pn \rightarrow pn$ scattering (W. R. Gibbs and B. Loiseau, PRC 50, 2742 (1994))
- 10% at the **cross-section level**
- **negligible impact on asymmetry, as pion-exchange is a spin-independent effect**
- $pn \rightarrow np$ has a **larger slope factor** in the amplitude → it decreases faster with t than $pn \rightarrow np$

A measurement of **fully exclusive nDVCS**, using **BoNuS** for the spectator proton and the **CND** for the active neutron could also be done, even if it will have limited statistics:
experimental check of FSI calculations, possibility to select kinematics minimizing FSI effects

Semi-Inclusive Deep Inelastic Scattering (SIDIS) on deuterium: The case for doubling the statistics

Silvia Pisano, INFN Frascati, contact person

Transverse Momentum Dependent PDFs&FFs

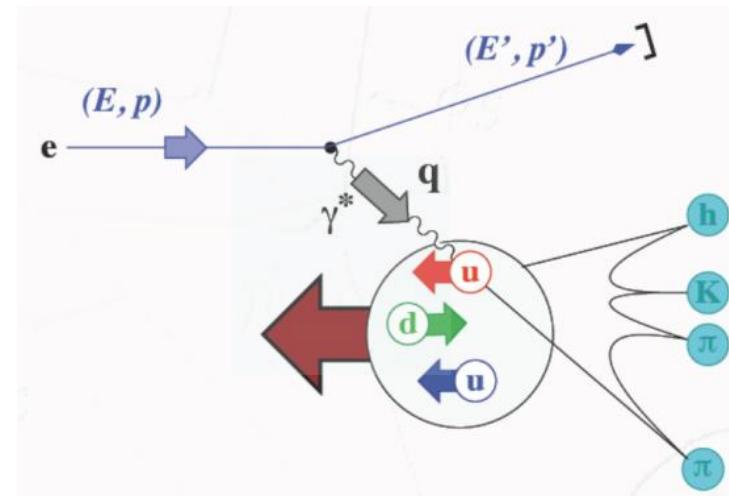
8 leading-twist TMDs

- They depend on the parton **longitudinal momentum fraction x** and on its **transverse momentum k_T**
→ **3D dynamics of the nucleon**
- Diagonal elements: **PDFs**
- Off-diagonal elements require **OAM**

Leading Twist TMDs

Quark Polarization			
	Un-Polarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	L	T
U	$f_1 = \bullet$		$h_1^\perp = \bullet - \bullet$ Boer-Mulders
L		$g_{1L} = \bullet \rightarrow - \rightarrow \bullet$ Helicity	$h_{1L}^\perp = \bullet \rightarrow - \bullet \rightarrow$ Kotzinian-Mulders
T	$f_{1T}^\perp = \bullet \uparrow - \bullet \downarrow$ Sivers	$g_{1T}^\perp = \bullet \uparrow - \bullet \uparrow$	$h_1 = \bullet \uparrow - \bullet \uparrow$ Transversity $h_{1T}^\perp = \bullet \uparrow - \bullet \uparrow$

Approved CLAS12 SIDIS experiments, with 50 days on longitudinally polarized ND₃:
E12-07-107 (pion SIDIS), E12-09-009 (kaon SIDIS)



Fragmentation Functions

Hadronization of the active quark → transition from **partonic to hadronic** degrees of freedom

q/H	U	L	T
U	D_1		H_1^\perp
L		G_{1L}	H_{1L}^\perp
T	H_1^\perp	G_{1T}	H_1, H_{1T}^\perp

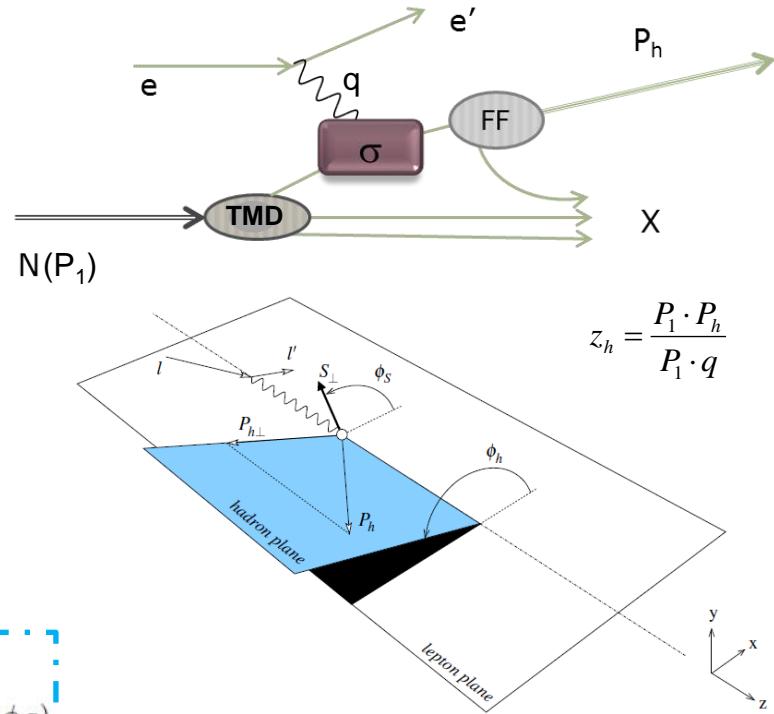
Single-hadron SIDIS cross-section

Depending on the polarizations of beam and target, different TMD and FF can be accessed:

Unpolarized target
Longitudinally pol. target
Transversely pol. target

$$\frac{d\sigma^h}{dx dy d\phi_S dz d\phi d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right)$$

$$\left\{ \begin{array}{l} \text{---} [F_{UU,T} + \epsilon F_{UU,L} \\ \quad + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)}] \\ + \lambda_l [\sqrt{2\epsilon(1-\epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)}] \\ + S_L [\sqrt{2\epsilon(1+\epsilon)} \sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)}] \\ + S_L \lambda_l [\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)}] \\ + S_T [\sin(\phi - \phi_S) (F_{UT,T}^{\sin(\phi-\phi_S)} + \epsilon F_{UT,L}^{\sin(\phi-\phi_S)}) \\ \quad + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi+\phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi-\phi_S)} \\ \quad + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} \\ \quad + \sqrt{2\epsilon(1+\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi-\phi_S)}] \\ + S_T \lambda_l [\sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi-\phi_S)} \\ \quad + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \\ \quad + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi-\phi_S)}] \end{array} \right\}$$



$$F_{ij,k} \propto TMD(x, k_T) \otimes FF(z_h, p_{\perp})$$

Depending on the produced hadron and on the **target type**, different **flavor combinations** of TMDs and FFs can be singled out

SIDIS on polarized ND_3 : single target spin asymmetry

N/q	U	L	T
U	f_1		h_1^\perp
L		g_1	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1 h_{1T}^\perp

$$\frac{d\sigma}{dxdy d\phi_S d\phi_h dP_{h\perp}^2} \propto S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \epsilon \sin(2\phi_h) F_{UL}^{\sin(2\phi_h)} \right]$$

h.t.

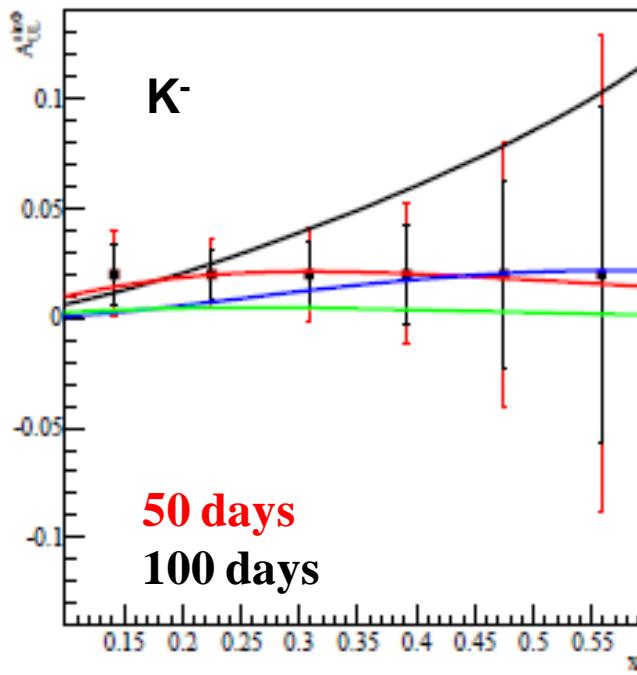
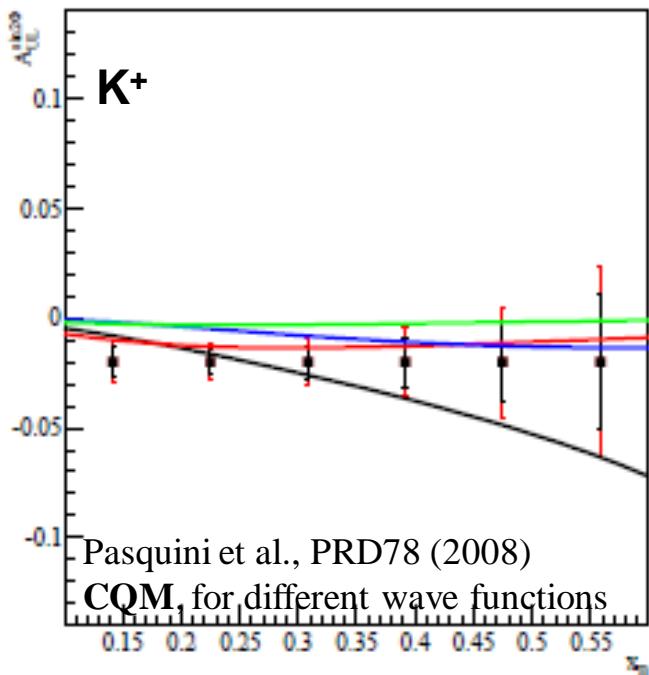
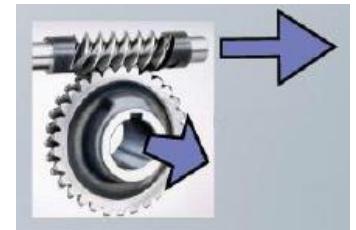
$$+ S_L \lambda_e \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_h) F_{LL}^{\cos(\phi_h)} \right]$$

$g_{1L} \otimes D_1$ h.t.

$h_{1L}^\perp \otimes H_1^\perp$

A_{UL} : $\sin 2\phi$ is the only **twist-2** term

- Kotzinian-Mulders function (worm gear)
- transversely polarized quark in a longitudinally polarized proton
- **quark spin-orbit correlations**
- provides independent information on **Collins function H_1^\perp**



Doubling the statistics helps at **high x_B** , where models differ the most, and it is crucial for the **K- channel**, which has lower statistics

SIDIS on polarized ND₃: double spin asymmetry

N/q	U	L	T
U	f_1		h_1^\perp
L		g_1	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	$h_1 h_{1T}^\perp$

$$\frac{d\sigma}{dxdy d\phi_S d\phi_h dP_{h\perp}^2} \propto S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \epsilon \sin(2\phi_h) F_{UL}^{\sin(2\phi_h)} \right]$$

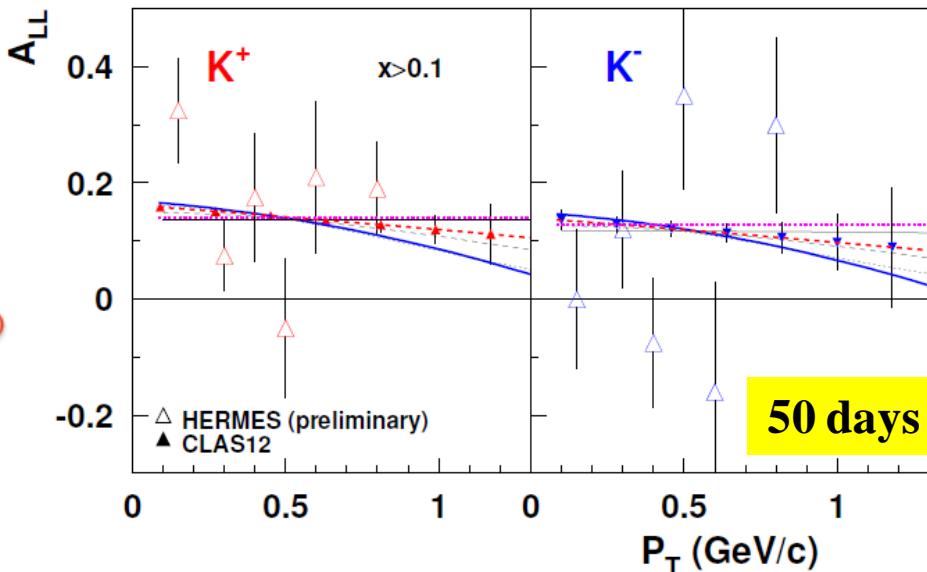
$$+ S_L \lambda_e \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_h) F_{LL}^{\cos(\phi_h)} \right]$$

h.t. $h_{1L}^\perp \otimes H_1^\perp$

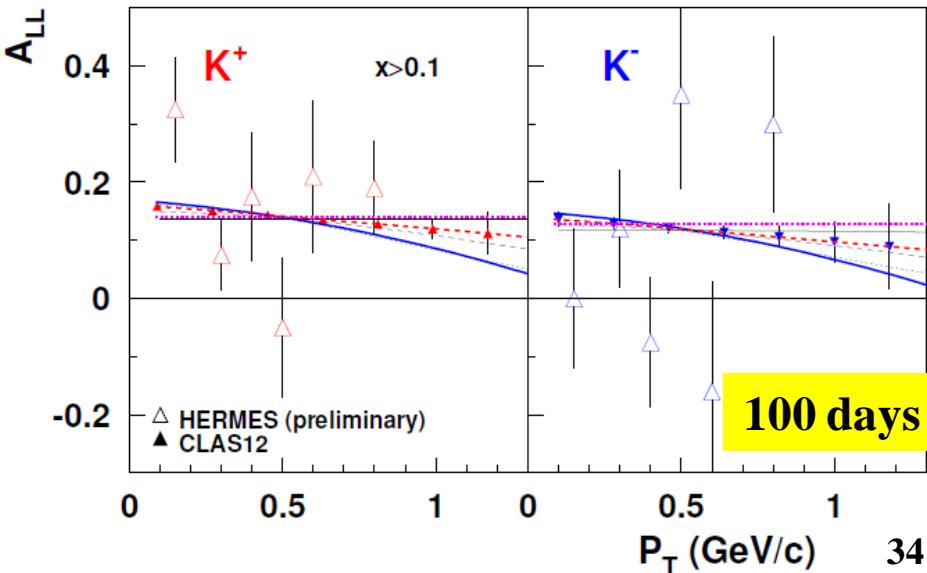
$g_{1L} \otimes D_1$ h.t.

- p_T dependence of **constant term** of A_{LL} (LT term): access to $g_{1L} \rightarrow$ **helicity dependence of k_T distributions of quarks**
- p & d data for p_T -dependent flavor decomposition

Doubling the statistics helps at **high p_T** , where models differ the most, and it is crucial for the **K⁻ channel**, which has lower statistics



Models: Anselmino et al., PRD74 (2006)
Different k_T widths for helicity distributions



SIDIS on polarized ND₃: double spin asymmetry

$$\frac{d\sigma}{dxdy d\phi_S d\phi_h dP_{h\perp}^2} \propto S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \epsilon \sin(2\phi_h) F_{UL}^{\sin(2\phi_h)} \right]$$

$$+ S_L \lambda_e \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_h) F_{LL}^{\cos(\phi_h)} \right]$$

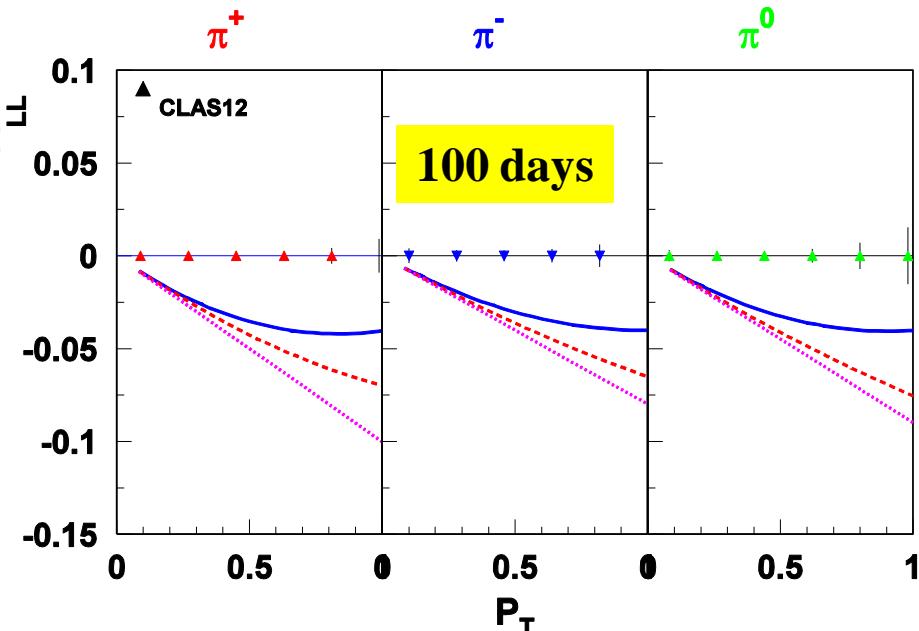
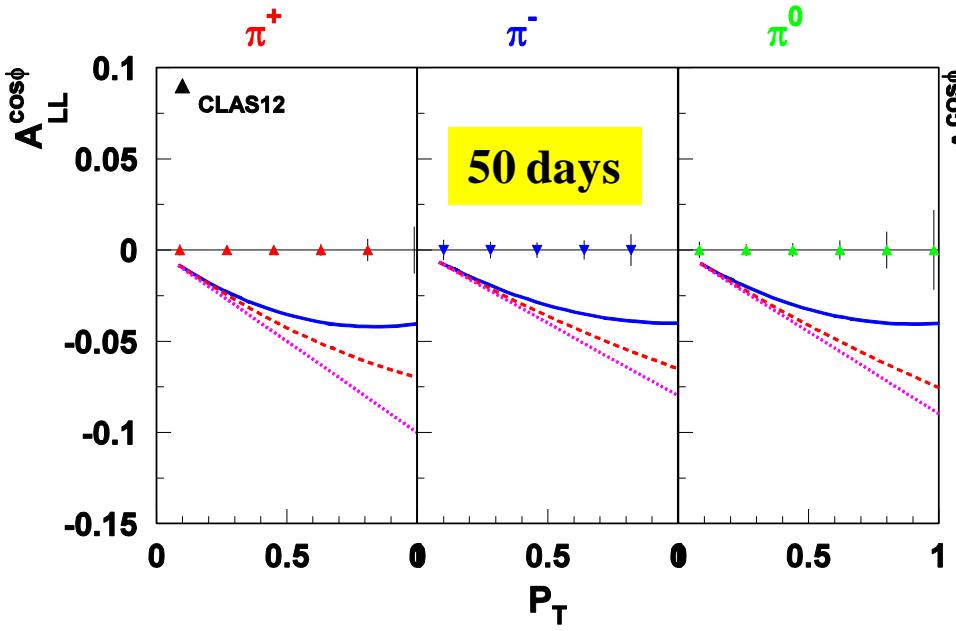
h.t. h_{LL}[†] ⊗ H_L[†]

g_{LL} ⊗ D_L h.t.

cos ϕ term of A_{LL} is linked to higher-twist TMDs

Higher-twist TMDs

N/q	U	L	T
U	f [†]	g [†]	h,e
L	f _L [†]	g _L [†]	h _L ,e _L
T	f _T , f _T [†]	g _T , g _T [†]	h _T ,e _T ,h _T [†] ,e _T [†]



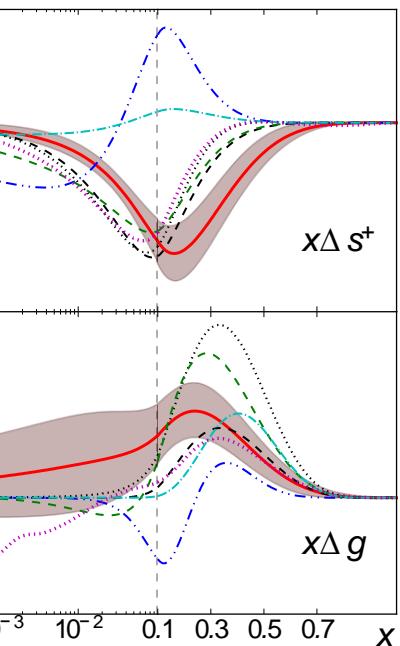
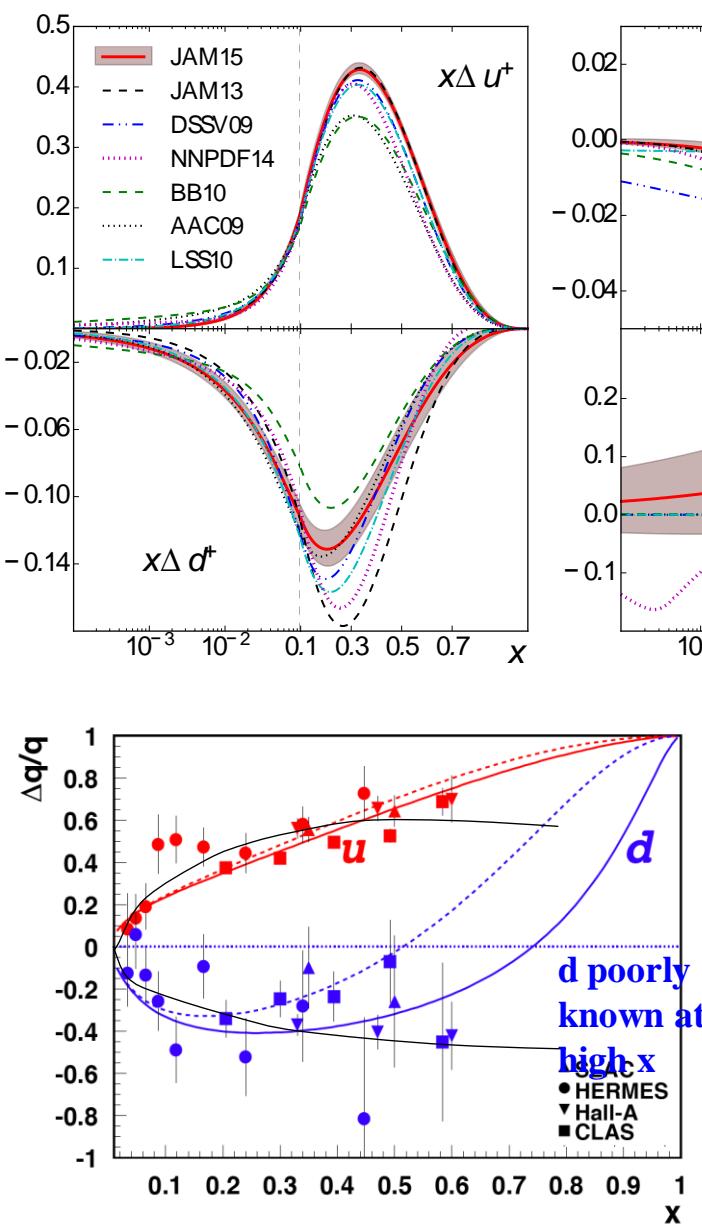
Models: Anselmino et al., PRD74 (2006)
Different k_T widths for helicity distributions

Also for pions, doubling the statistics helps at high p_T, where models differ the most

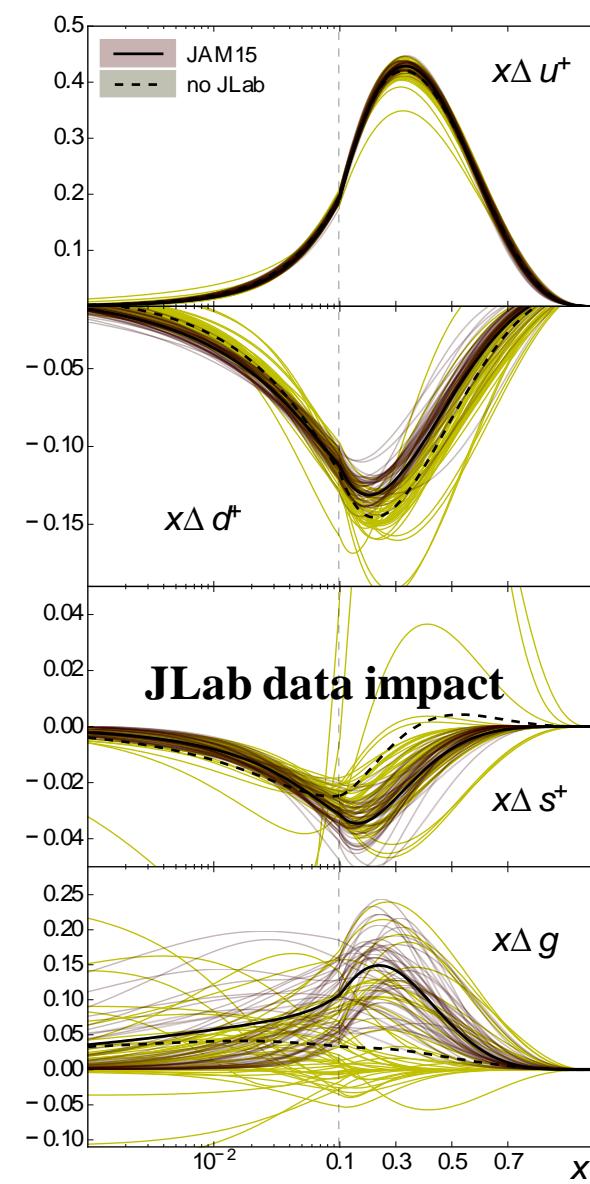
Deep Inclusive Scattering (DIS) on Deuterium: The case for doubling the statistics

Sebastian Kuhn, ODU, contact person

Present Status: What do we know about polarized PDFs?



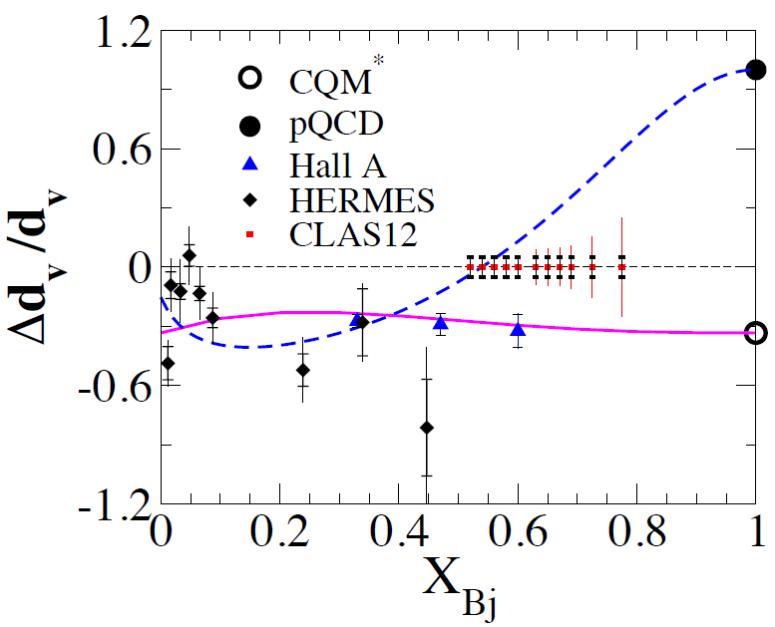
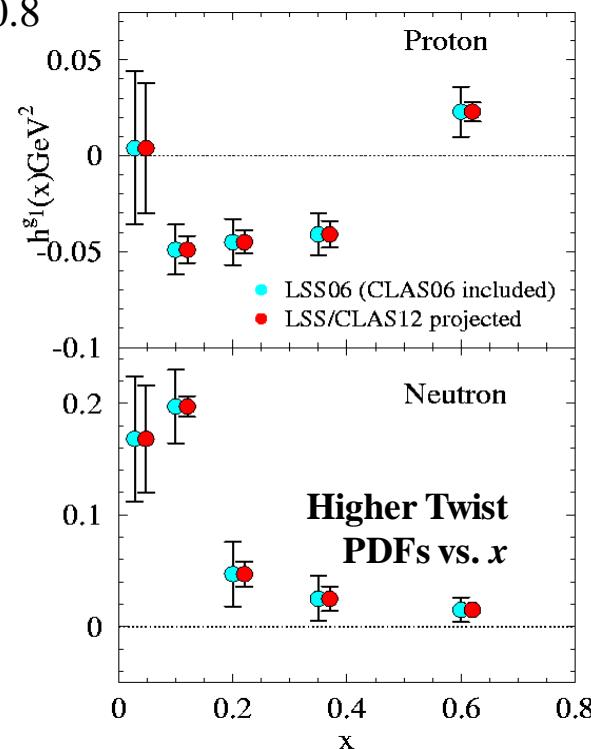
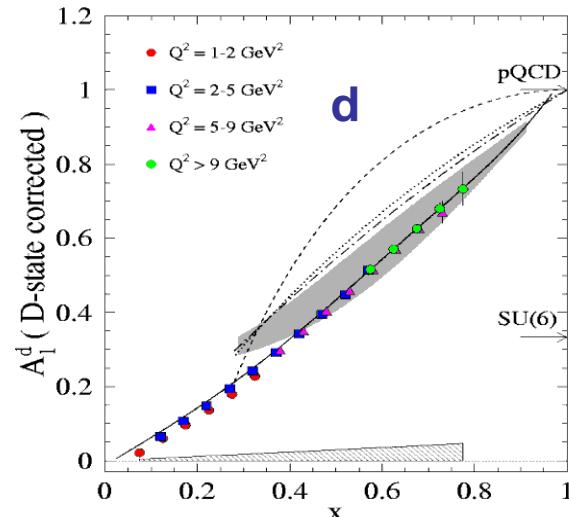
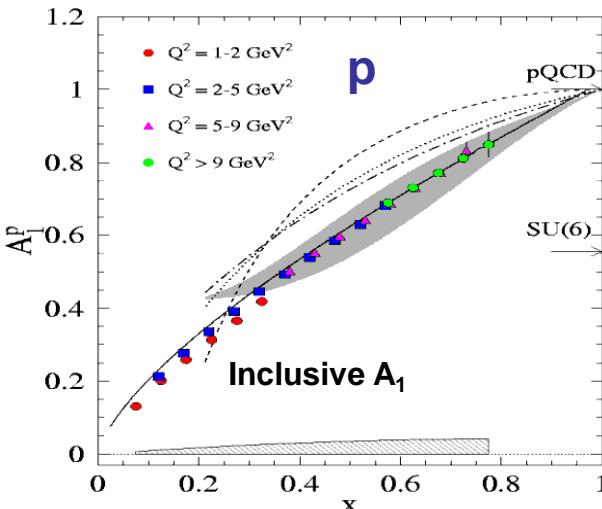
**JAM
Fit**



JLab data impact

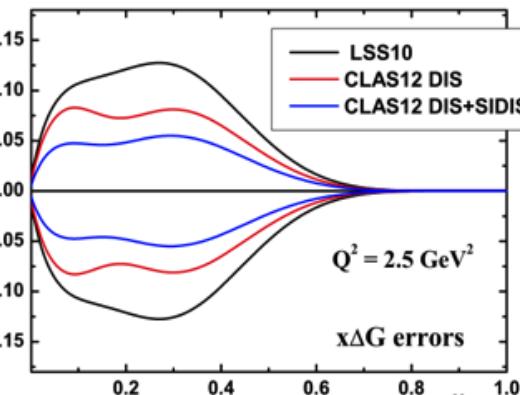
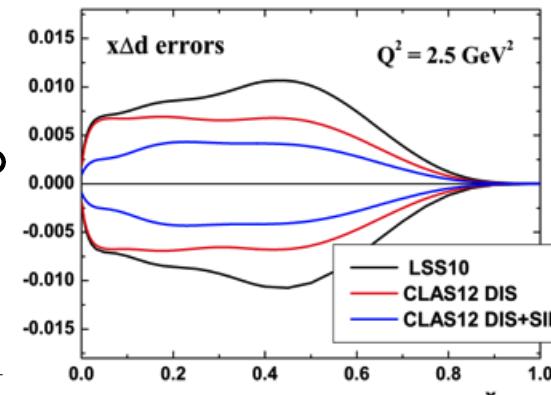
CLAS12 Approved (RG C: 30 days p, 50 days d)

- Goal of E12-06-109: **x-dependence of each PDF (q or g) for $0.06 < x < 0.8$**
- Measurement of **DIS and SIDIS** (no p_T dependence)

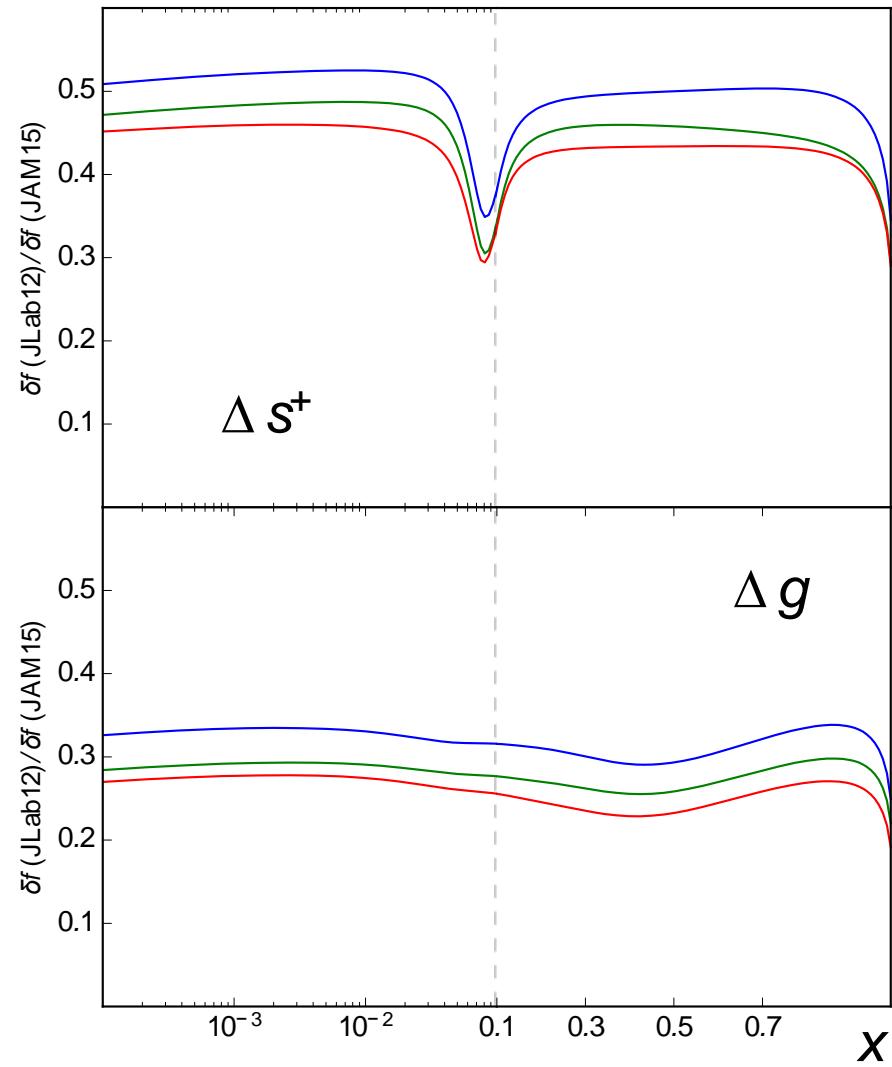
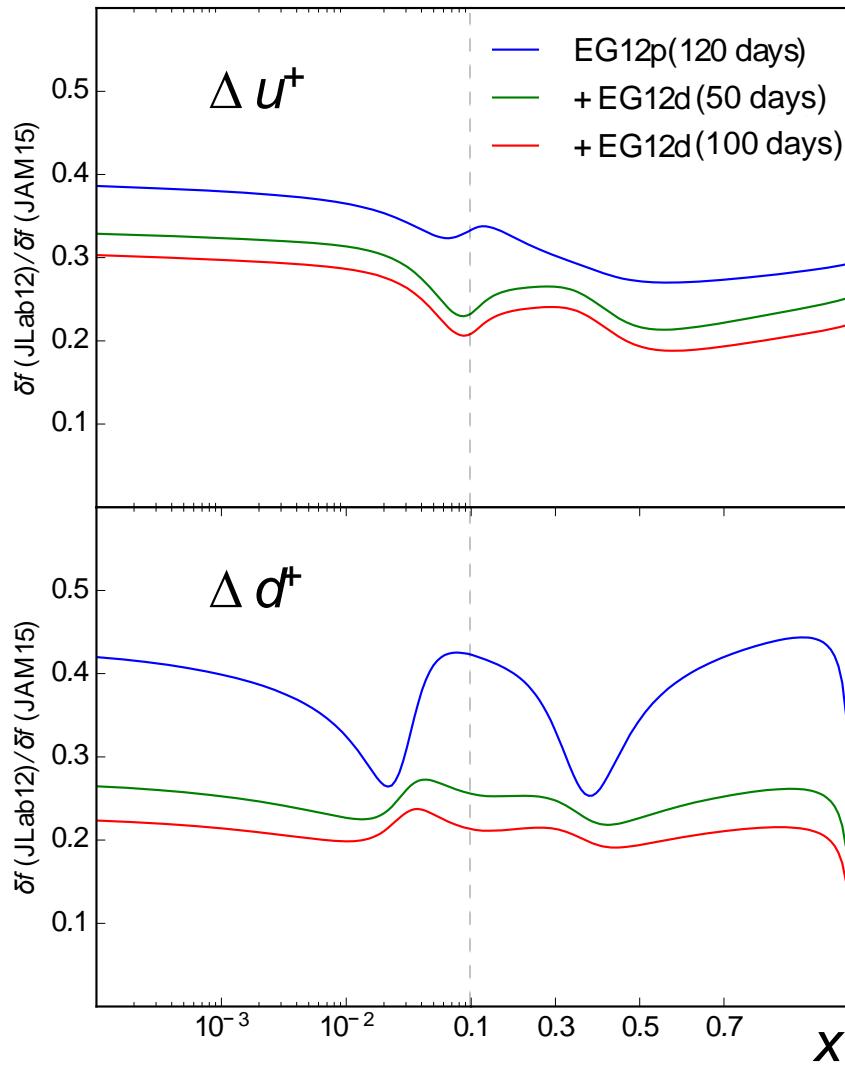


Additional information
on Δq from SIDIS

Improved PDFs from NLO analyses

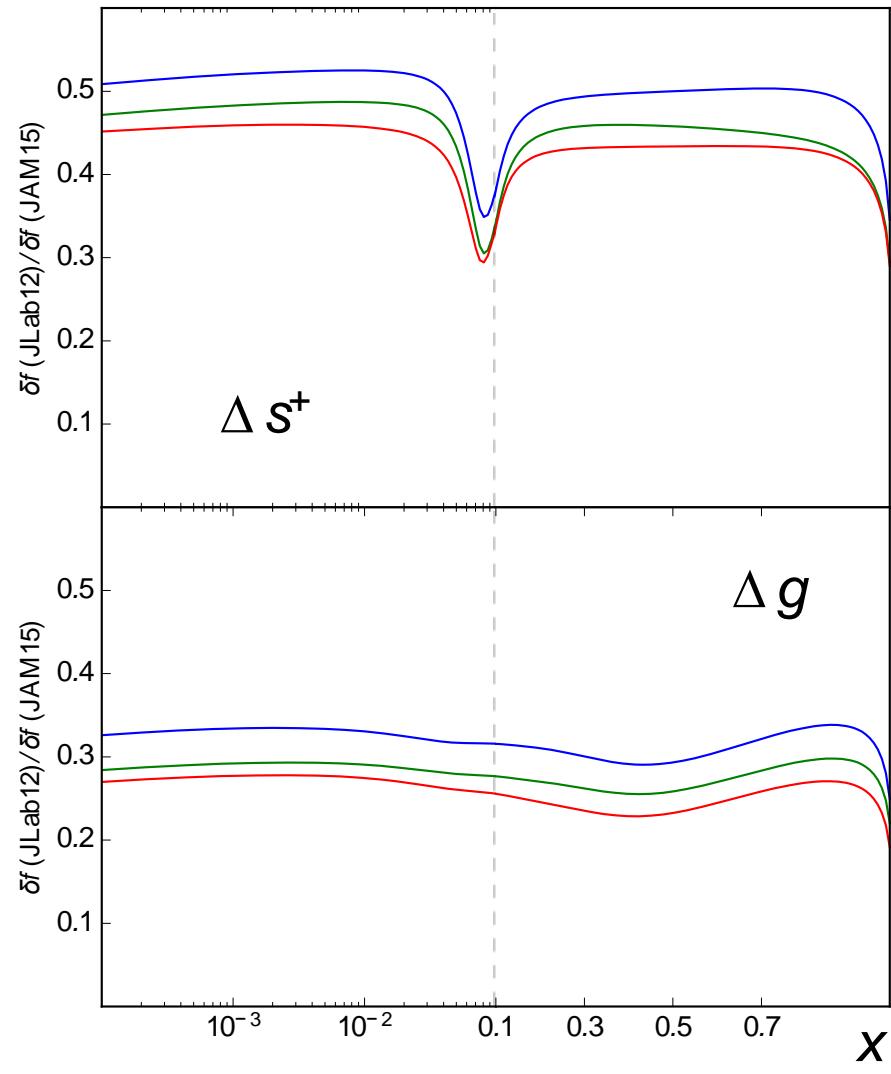
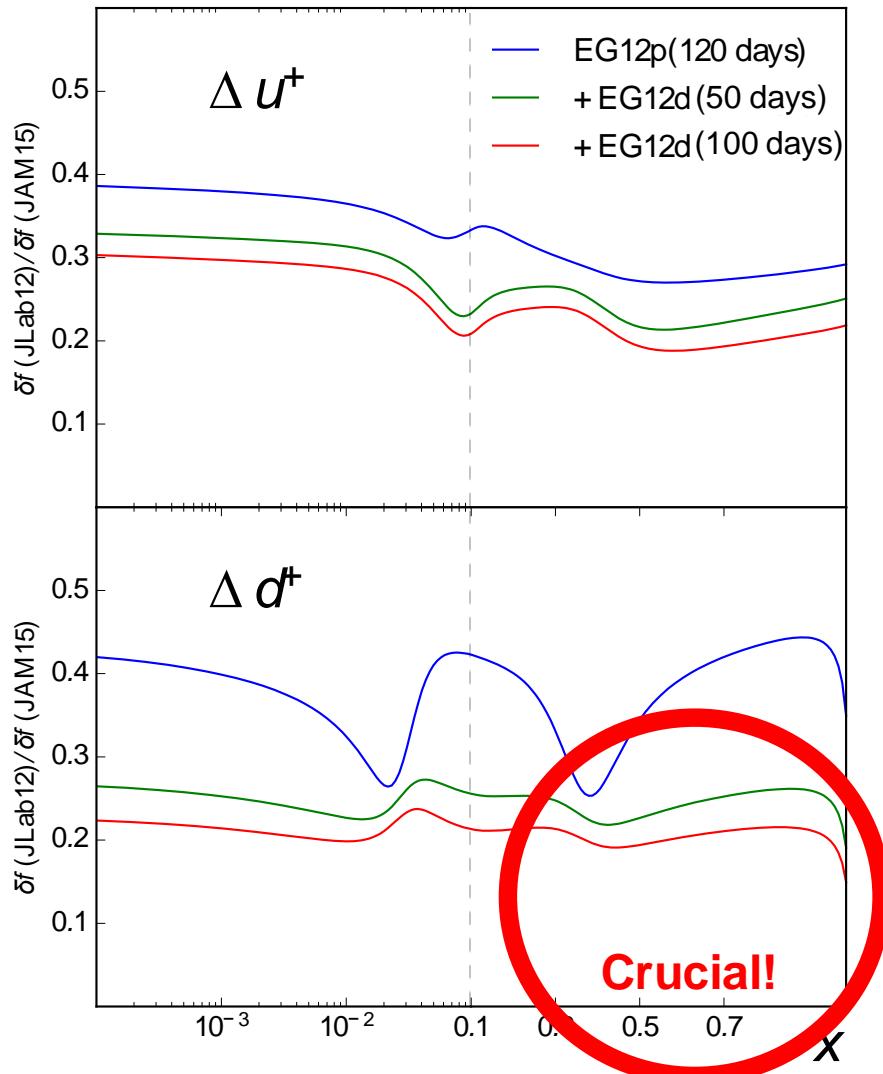


What will 50 more days on D bring us?



Relative uncertainties on polarized PDFs with various data sets (Courtesy N. Sato)

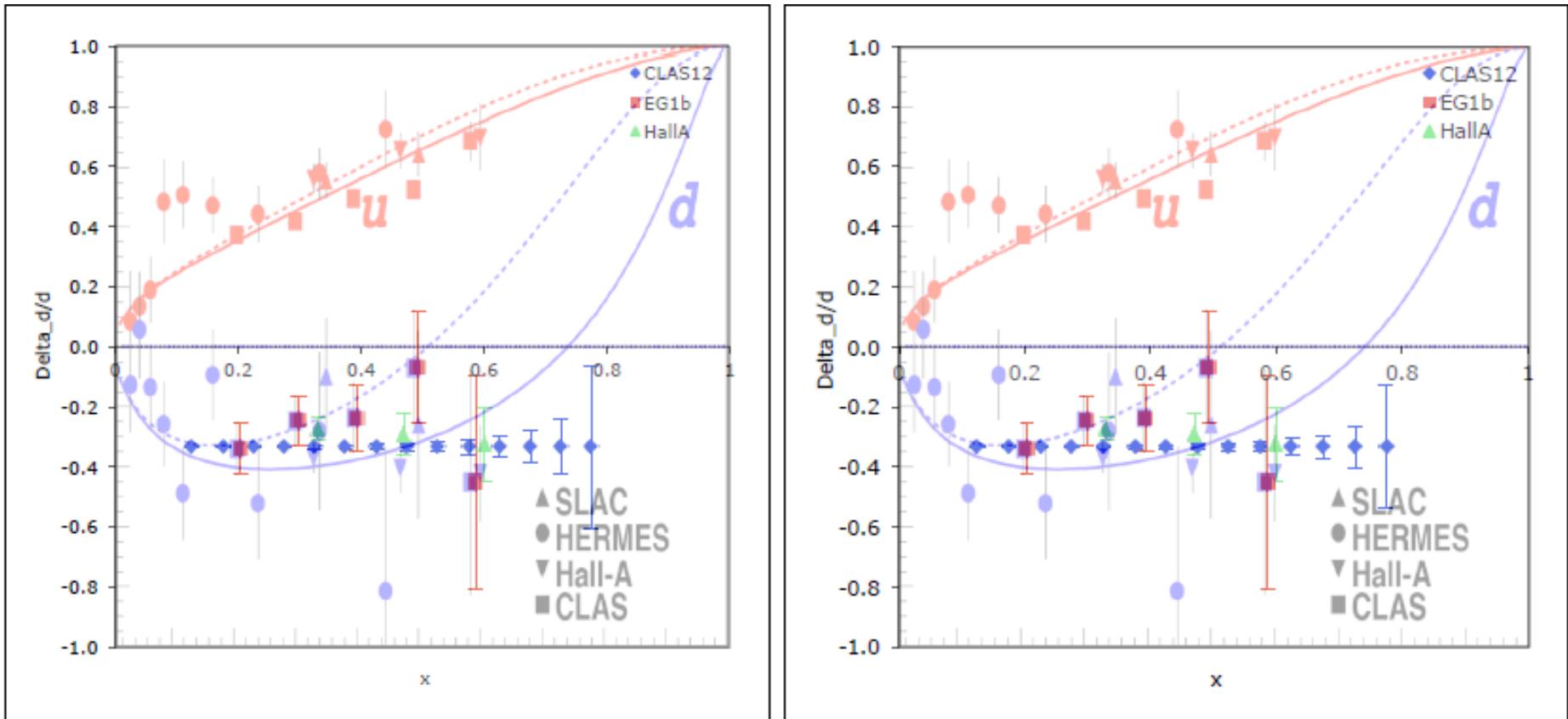
What will 50 more days on D bring us?



Relative uncertainties on polarized PDFs with various data sets (Courtesy N. Sato)

What will 50 more days on D bring us?

$\Delta d/d$ for $x \rightarrow 1$



Curves: pQCD without and with OAM effects (Brodsky et al. Nucl. Phys. B441 (1995) – Avakian et al. PRL99 (2007))

Projections: hyperfine perturbed model (N. Isgur et al., PRD59 (1999))

As approved

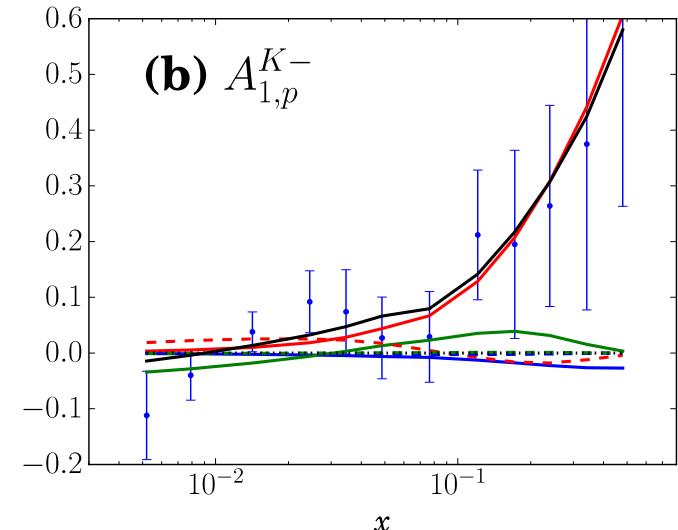
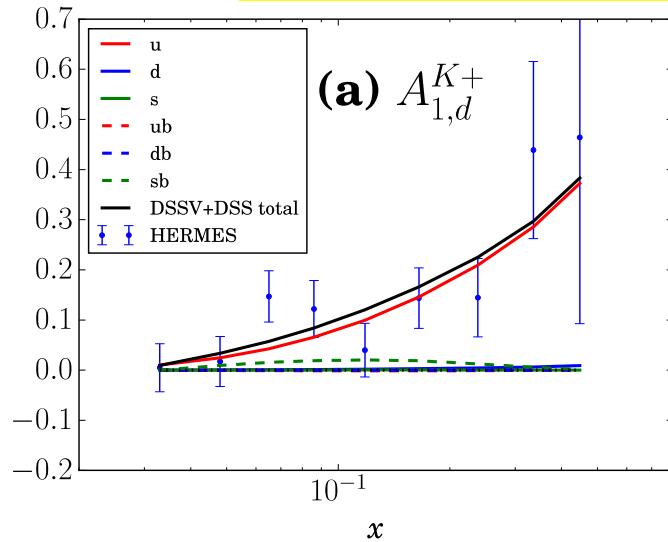
With “unlucky” statistical fluctuations, cannot exclude $\Delta d/d > 0$ at highest x ($\chi^2/\text{dof}=2.5, p=8\%$)

With doubled beam time

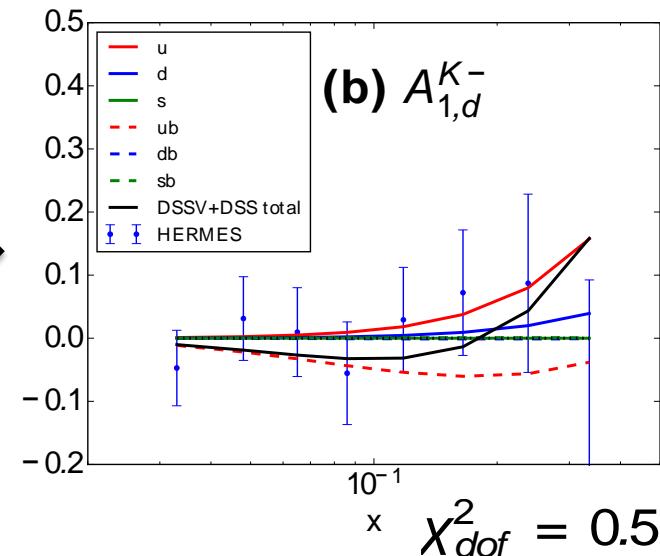
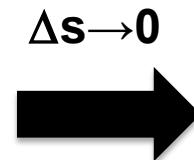
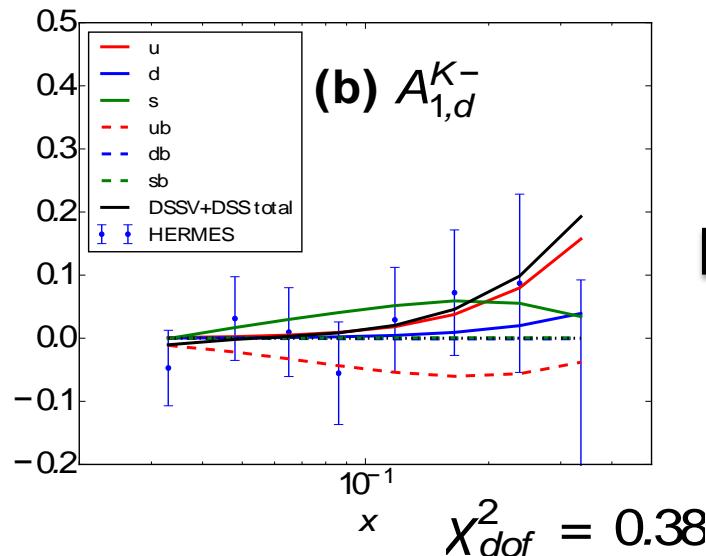
Even with “unlucky” statistical fluctuations, can exclude $\Delta d/d > 0$ at highest x ($\chi^2/\text{dof}=6.6, p=0.4\%$)

What will 50 more days on D bring us?

A₁ from SIDIS (kaon channels) – HERMES data

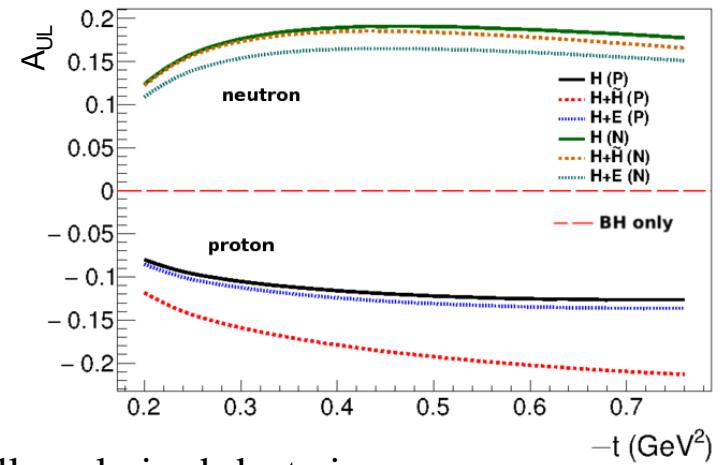
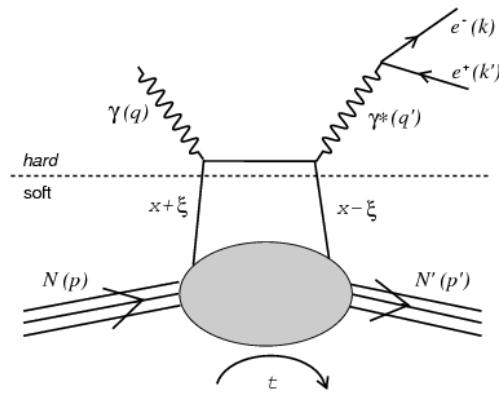


Only K⁻ production on D is uniquely sensitive to Δs (Δu and Δd largely cancel)



Additional topics

- First time measurement of **Timelike Compton scattering**, $\gamma n \rightarrow n e^+ e^-$, on longitudinally polarized neutron:
 - ✓ universality of GPDs
 - ✓ A_{UL} and A_{LL} have sensitivity to H_n and E_n
 - ✓ A_{UL} expected to be fairly big (0.15)
 - ✓ nTCS cross section only a factor of 2 less than the pTCS one
 - ✓ 120 days approved for pTCS with unpolarized target



- First-time measurement of **di-hadron SIDIS** on longitudinally polarized deuterium:
 - ✓ cleanest access to chiral-odd PDFs
 - ✓ promising preliminary results for A_{UL} on NH_3 with CLAS at 6 GeV
 - ✓ ND_3 needed for flavor dependence of the various PDFs and FFs, high stat for mapping
- **DVMP** on longitudinally polarized neutrons:
 - ✓ transversity neutron-GPDs in pseudoscalar channels (π^0, η, π^-)
 - ✓ link to transversity GPDs shown by 6-GeV CLAS results on π^0, η
 - ✓ A_{UL} proportional to LT interference term of the cross section $\rightarrow H_T, E_T, \bar{E}_T$

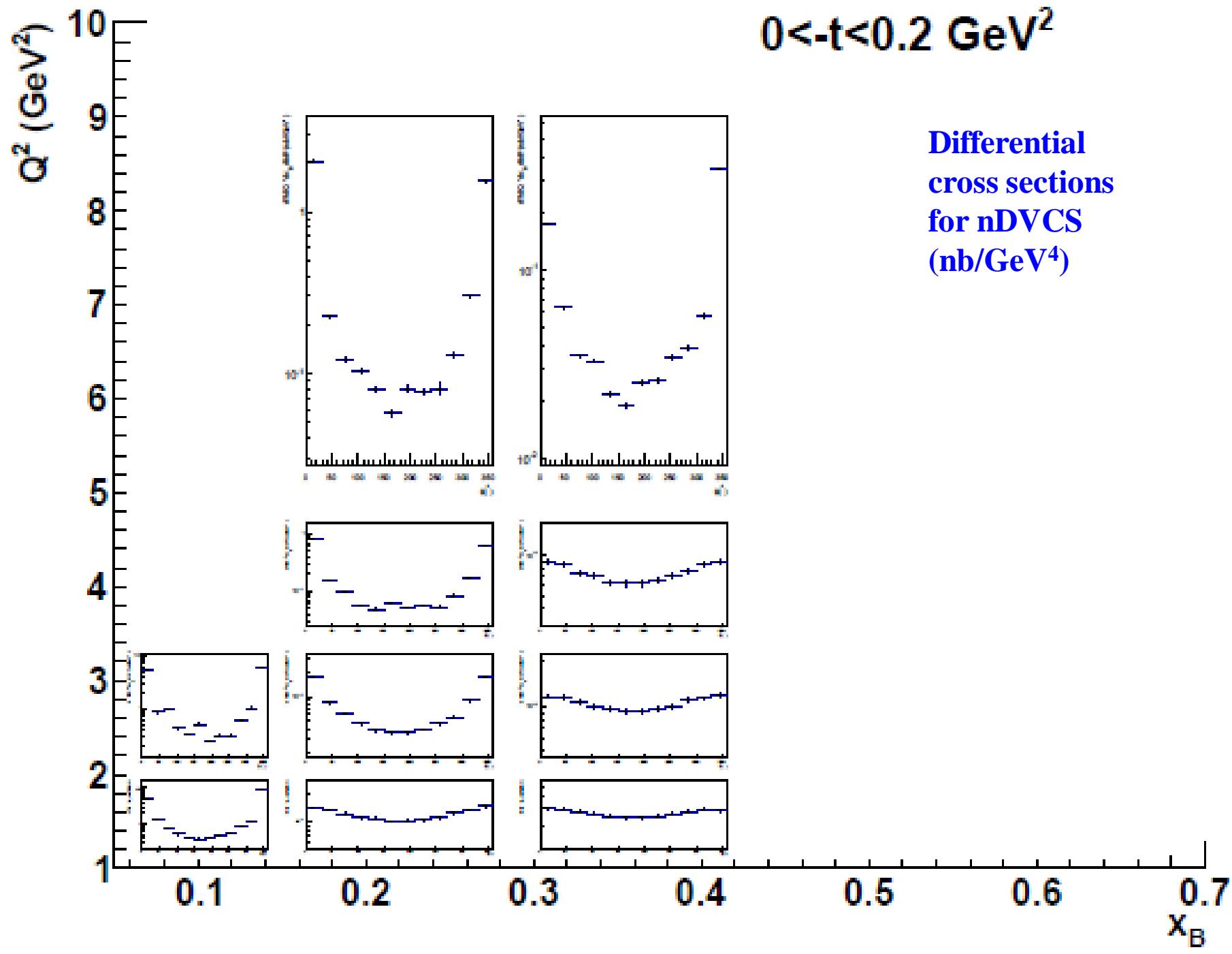
Conclusions/summary: CLAS12 Run group M

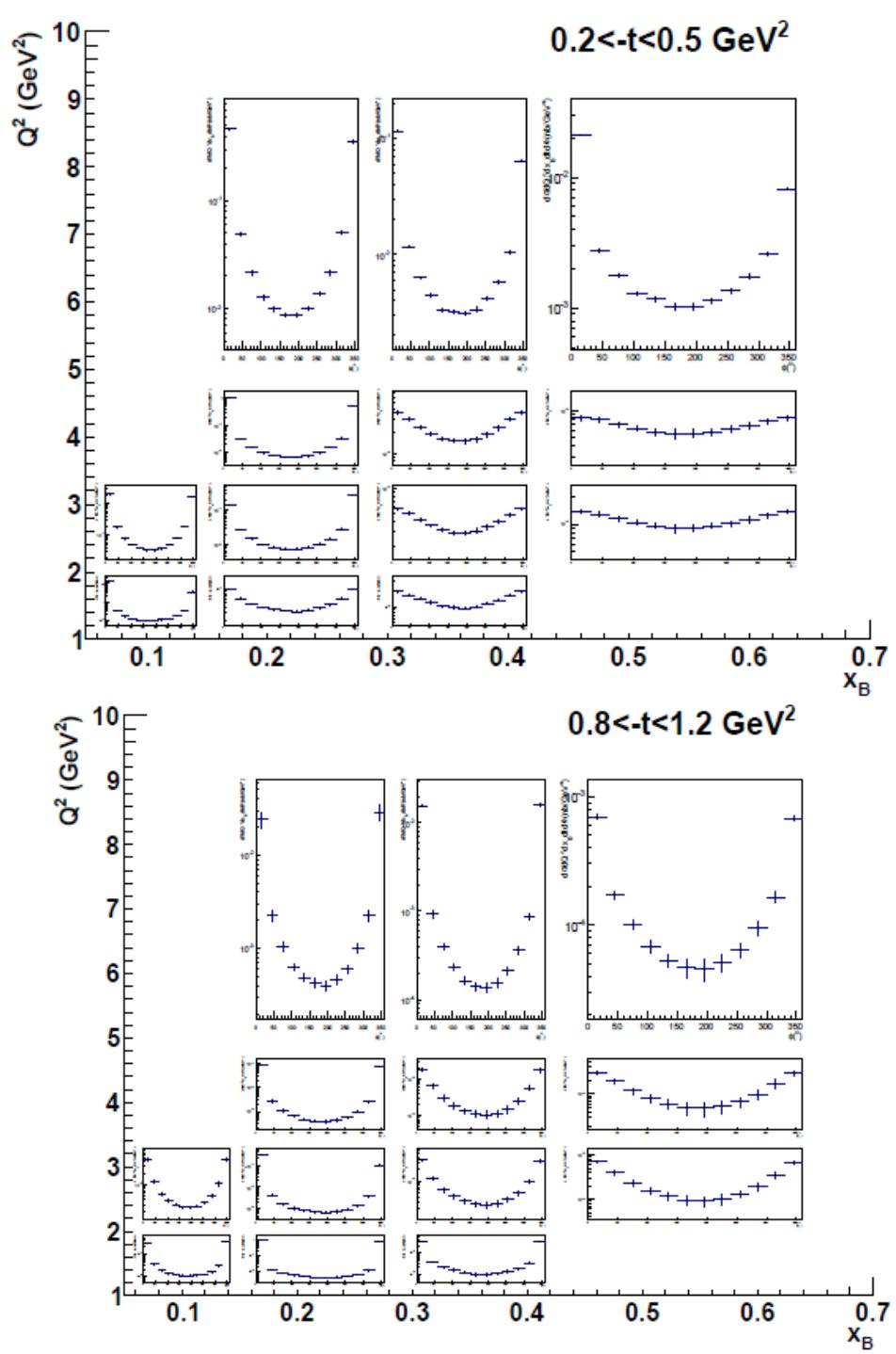
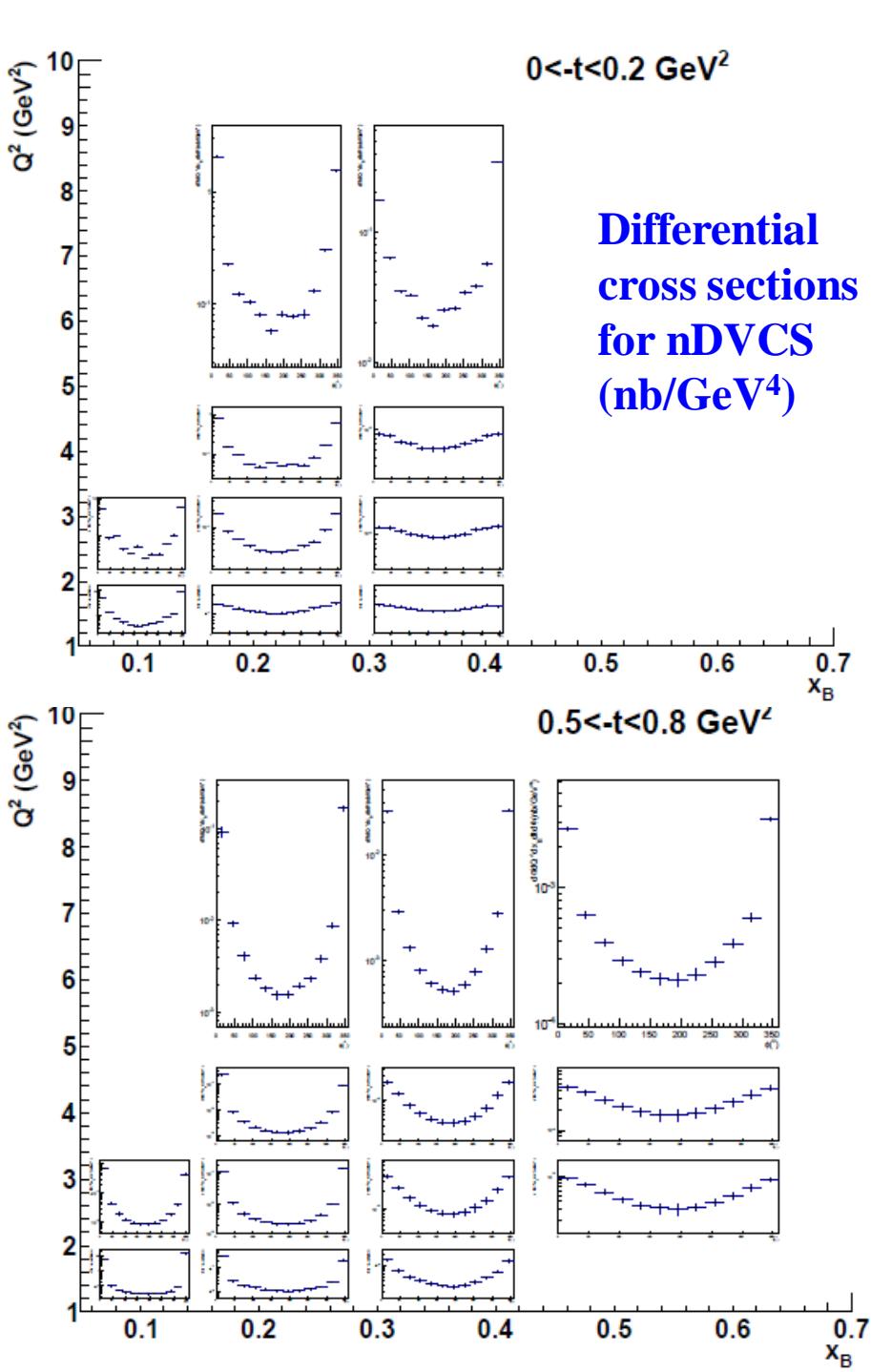
- The multi-dimensional structure of the nucleon is one of the primary goals of the 12-GeV upgrade of JLab, and it is the core of the CLAS12 program
- Parton distributions depend on quark flavors, measurements on proton and deuterium are necessary
- Current approved beam time on ND_3 for CLAS12 is $\sim 40\%$ of the NH_3 one
- We request 73 PAC days of new beam time (11 GeV, 85% polarization):
 - 60 of production running on ND_3 , 13 of ancillary measurements/configuration change
- Setup: CLAS12 + ND_3 polarized target, CND, RICH, FT (for 10 days at 5 nA)
- The inclusion of the FT completes the ϕ coverage for nDVCS, improvement for CFF fits beyond statistical effects
- This will bring the total running time on ND_3 for CLAS12 to 133 days – same as NH_3
- First time measurement of polarized nDVCS
- Increased statistics for SIDIS, important especially for kaon channels at high x, p_T
- High precision measurement of collinear structure functions
- Exploratory measurements of DVMP, nTCS, di-hadron SIDIS

Back-up slides

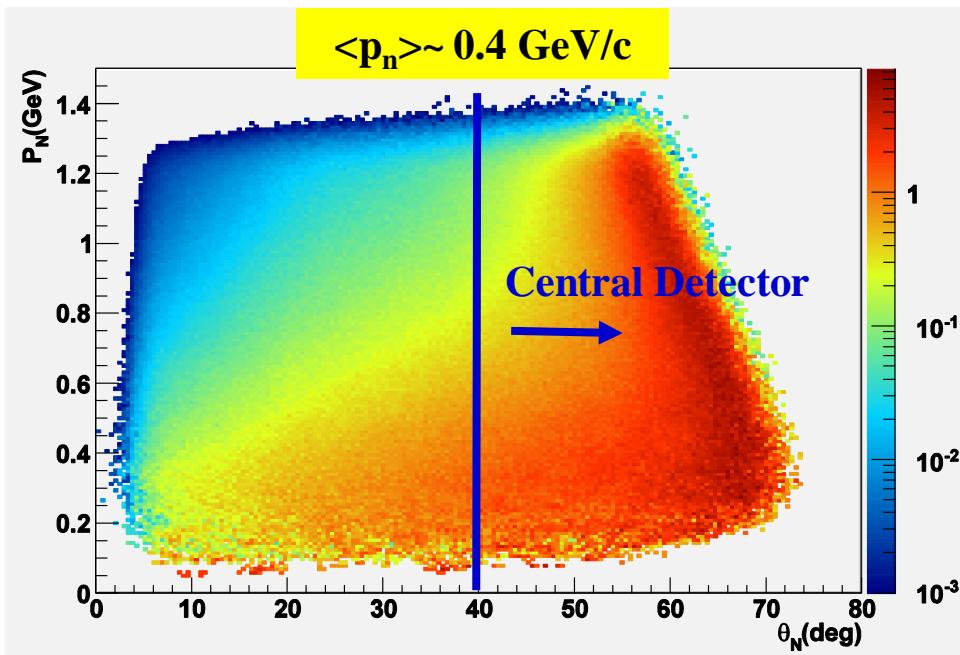
$0 < -t < 0.2 \text{ GeV}^2$

Differential
cross sections
for nDVCS
(nb/GeV⁴)



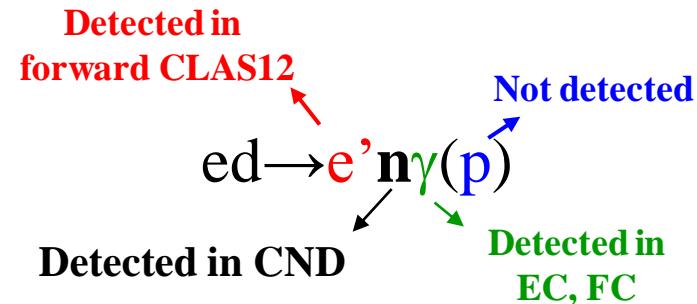


CND: requirements



More than **80%** of the neutrons have $\theta > 40^\circ$
→ Neutron detector in the CD

Resolution on **MM(en γ)** studied with nDVCS event generator + electron and photon resolutions obtained from CLAS12 FastMC + design specs for Forward Calorimeter
→ dominated by photon resolutions



$$p^\mu_e + p^\mu_n + p^\mu_p = p^\mu_{e'} + p^\mu_{n'} + p^\mu_{p'} + p^\mu_\gamma$$

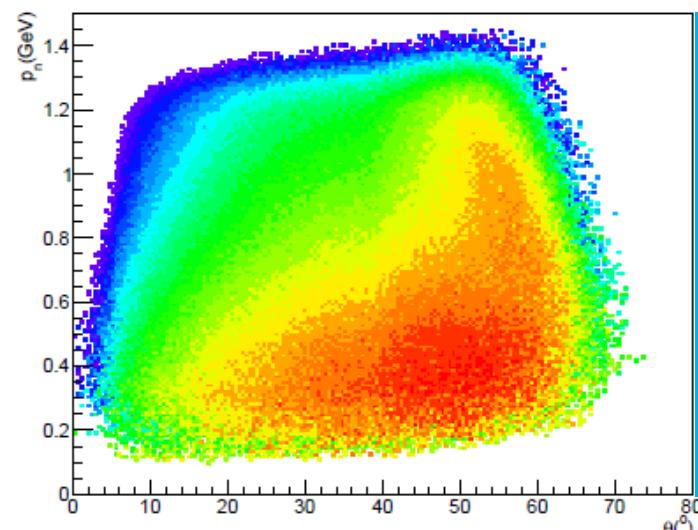
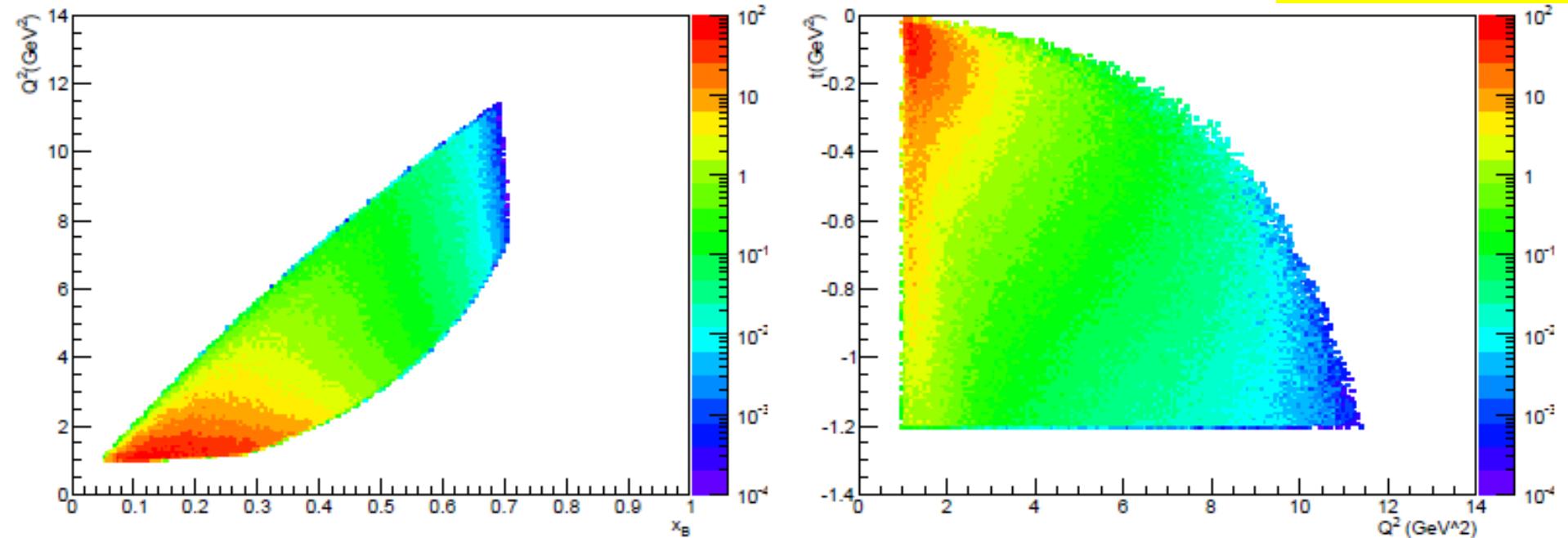
In the hypothesis of **absence of FSI**:
 $p^\mu_p = p^\mu_{p'}$, → **kinematics are complete**
 detecting **e', n (p, θ, ϕ), γ**

FSI effects will be estimated measuring
en γ , ep γ , on deuteron
 in this same experiment
 and compare with free-proton data

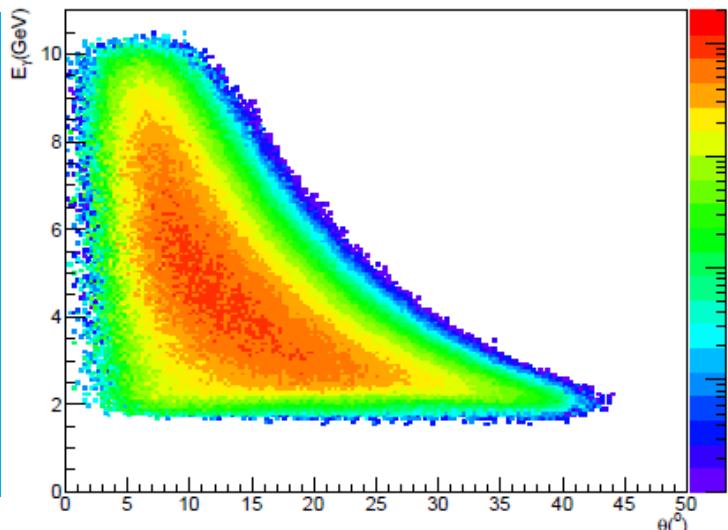
- The CND must ensure:
- **good neutron identification** for $0.2 < p_n \leq 1 \text{ GeV}/c$
 - $\sigma(\text{TOF}) \sim 150 \text{ ps}$ for n/γ β-separation
 - momentum resolution **up to 10%**
 - no stringent requirements for angular resolutions

nDVCS@CLAS12: kinematics and acceptances

$ed \rightarrow e'ny(p)$



nDVCS+BH event generator on a FREE deuteron target
+ CLAS12 acceptance from FastMC ($e\gamma$)
+ CND efficiency from GEANT4 simulation

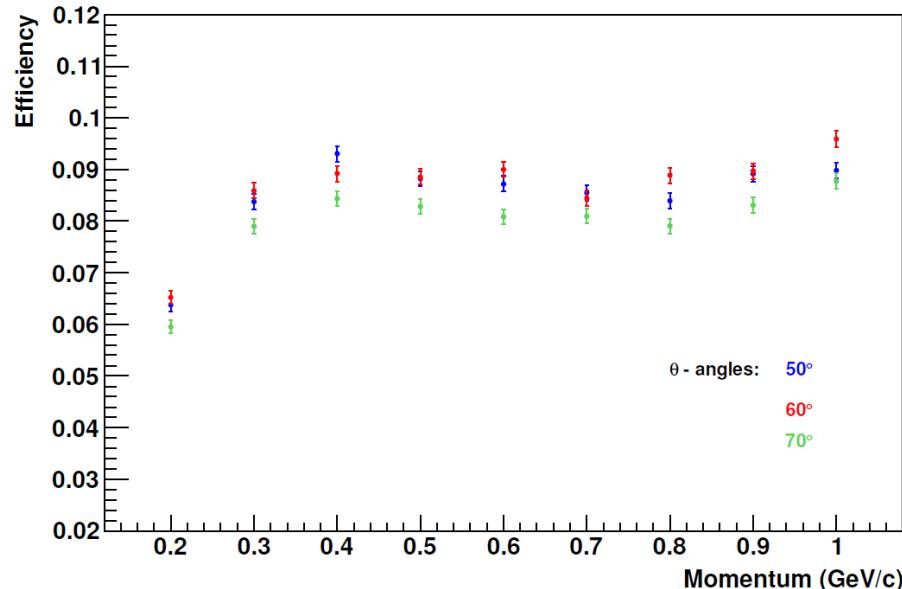


80% of neutrons
in CND ($\theta > 40^\circ$)

$E_e = 11$ GeV, $Q^2 > 1$ GeV 2 , $W > 2$ GeV 2 , $t < -1.2$ GeV 2

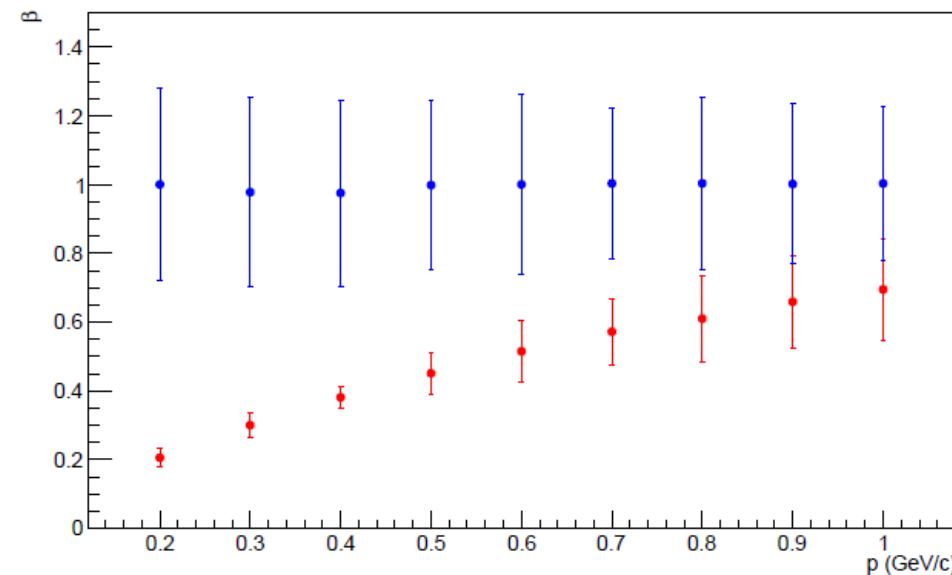
nDVCS-BH photons
have $E > 2$ GeV

CND: expected performances (Monte Carlo)



GEANT4 simulations used to evaluate:

- efficiency
- PID (neutron/photon separation)
- momentum and angular resolutions
- definition of reconstruction algorithms
- background studies



Measured σ_t and light loss due to u-turn implemented in the simulation

Efficiency ~ 8-9%
for a threshold of 3 MeV, TOF < 8 ns
and $p_n = 0.2 - 1 \text{ GeV}/c$



Definitions of nDVCS observables

For each 4-D. bin in $(Q^2, x_B, -t, \phi)$

$$A_{UL} = \frac{N^{++} + N^{-+} - N^{+-} - N^{--}}{D_f \left(P_t^- (N^{++} + N^{-+}) + P_t^+ (N^{+-} + N^{--}) \right)}$$

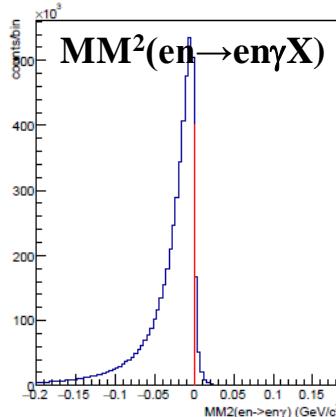
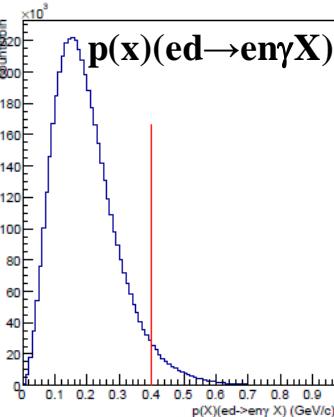
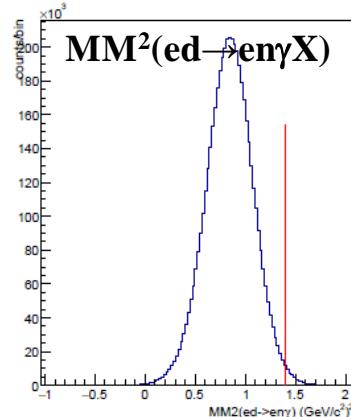
Target-spin asymmetry

$Im\{\mathcal{H}_n, \mathcal{E}_n\}$

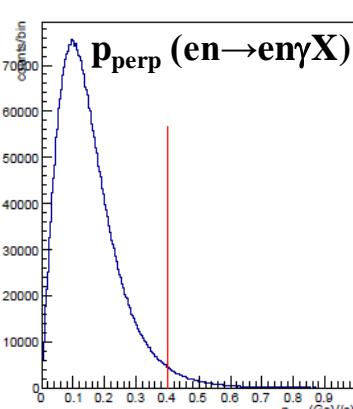
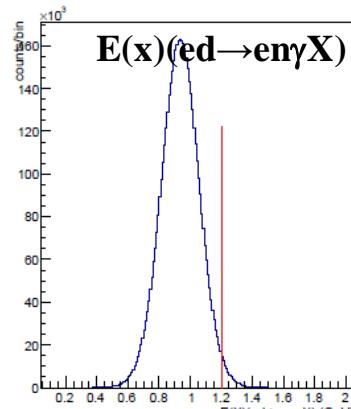
$$A_{LL} = \frac{N^{++} + N^{--} - N^{+-} - N^{-+}}{D_f P_b \left(P_t^- (N^{++} + N^{-+}) + P_t^+ (N^{+-} + N^{--}) \right)}$$

Double spin asymmetry

$Re\{\mathcal{H}_n, \mathcal{E}_n\}$



Yields extracted applying **exclusivity cuts** to $en\gamma X$ events



Exclusivity cuts suppress:

- nuclear background (^{14}N)
- π^0 background
- accidentals in the CND

Dilution factor (needs ^{12}C data)

$$D_f(en\gamma) = 1 - c \frac{N_{^{12}\text{C}}}{N_{ND_3}} \approx 0.7$$

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Target-spin asymmetry

$Im\{\mathcal{H}_n, \mathcal{E}_n\}$

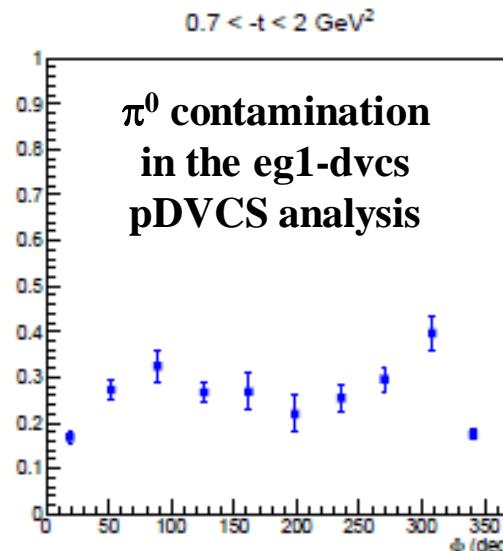
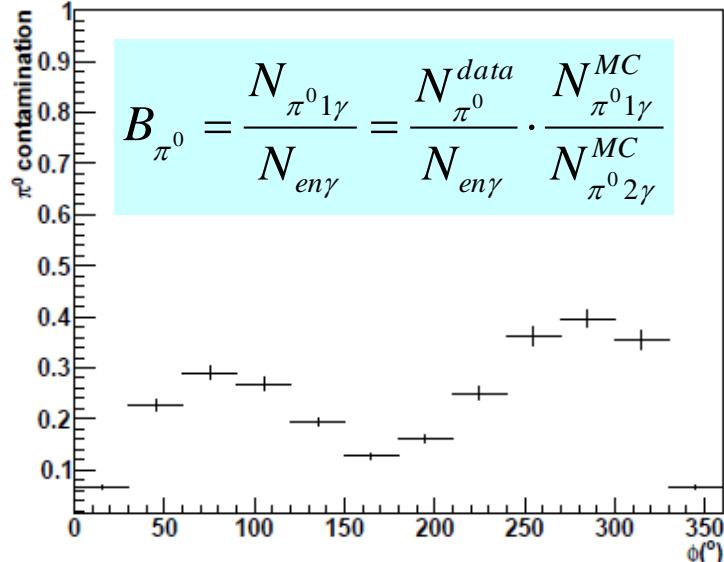
$$A_{LL} = \frac{N^{++} + N^{--} - N^{+-} - N^{-+}}{D_f P_b \left(P_t^- (N^{++} + N^{-+}) + P_t^+ (N^{+-} + N^{--}) \right)}$$

Double spin asymmetry

$$N^{bt} = (1 - B_{\pi^0}^{bt}) \cdot \frac{N_{eny}^{bt}}{FC^{bt}}$$

Charge-normalized DVCS/BH yield

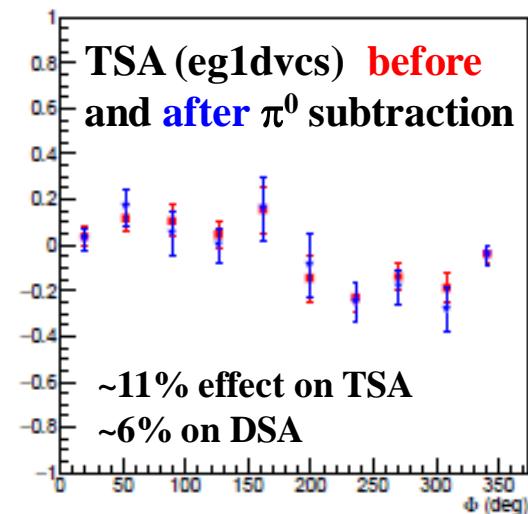
Projected π^0 contamination vs ϕ



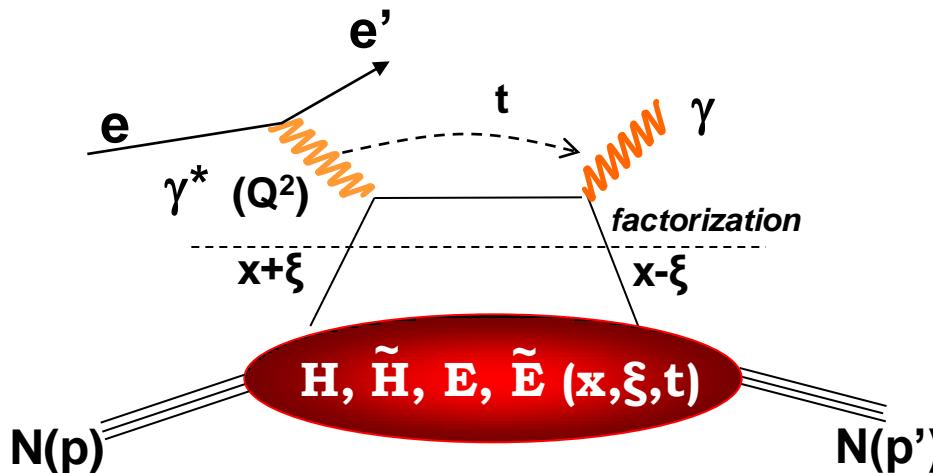
$Re\{\mathcal{H}_n, \mathcal{E}_n\}$

b : beam; t : target
 FC^{bt} : charge, (Faraday Cup)
 B_{π^0} : π^0 contamination
 P_t : target pol.; P_b : beam pol.
 D_f : dilution factor

0.7 < -t < 2 GeV²



Deeply Virtual Compton Scattering and GPDs



- $Q^2 = -(\mathbf{e} - \mathbf{e}')^2$
- $x_B = Q^2/2Mv$ $v = E_e - E_{e'}$
- $x + \xi, x - \xi$ longitudinal momentum fractions
- $t = \Delta^2 = (\mathbf{p} - \mathbf{p}')^2$
- $x \cong x_B/(2 - x_B)$

« Handbag » factorization valid
in the Bjorken regime:
high Q^2 , v (fixed x_B), $t \ll Q^2$

At LO QCD, twist 2, chiral-even, quark sector \rightarrow 4 GPDs for each quark flavor

conserve nucleon spin

Vector: $H(x, \xi, t)$ Axial-Vector: $\tilde{H}(x, \xi, t)$

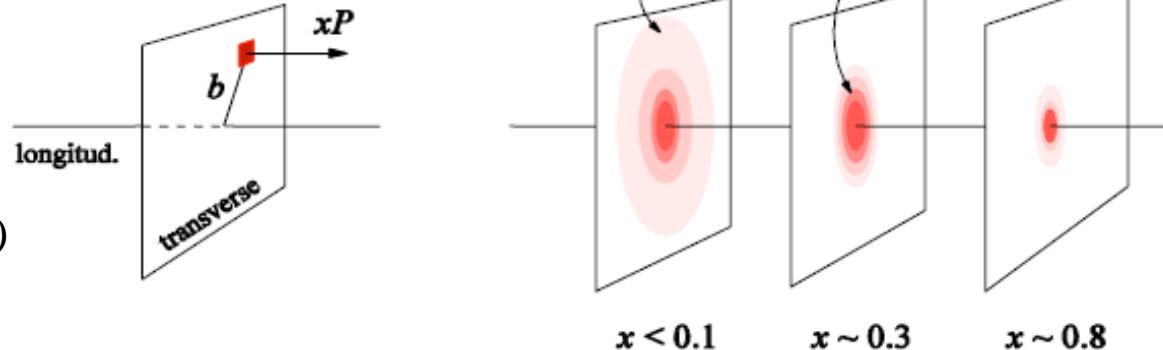
flip nucleon spin

Tensor: $E(x, \xi, t)$ Pseudoscalar: $\tilde{E}(x, \xi, t)$

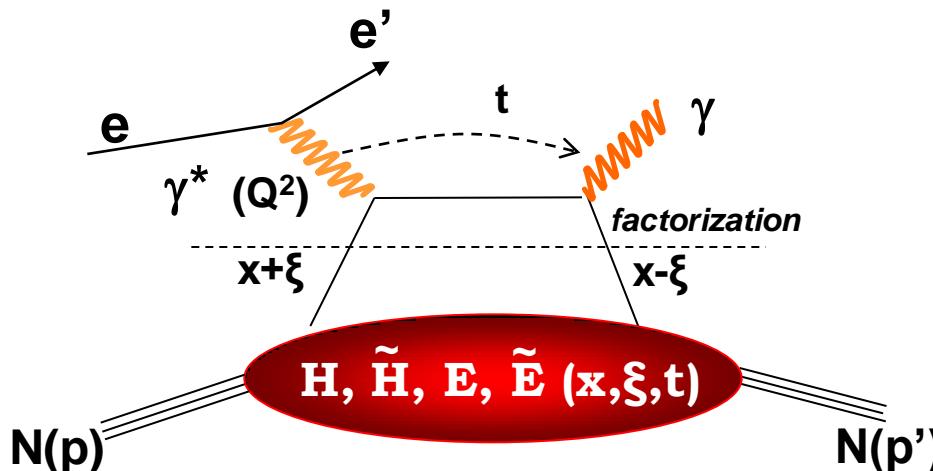
Nucleon tomography

$$q(x, b_\perp) = \int_0^\infty \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{i \Delta_\perp b_\perp} H(x, 0, -\Delta_\perp^2)$$

$$\Delta q(x, b_\perp) = \int_0^\infty \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{i \Delta_\perp b_\perp} \tilde{H}(x, 0, -\Delta_\perp^2)$$



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Nucleon tomography

$$q(x, \beta_\perp) = \int_0^\infty \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{i \Delta_\perp \beta_\perp} H(x, 0, -\Delta_\perp^2)$$

$$\Delta q(x, \beta_\perp) = \int_0^\infty \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{i \Delta_\perp \beta_\perp} \tilde{H}(x, 0, -\Delta_\perp^2)$$

Quark angular momentum (Ji's sum rule)

$$\frac{1}{2} \int_{-1}^1 x dx (H_q(x, \xi, t=0) + E_q(x, \xi, t=0)) = J_q = \frac{1}{2} \Delta \Sigma + \Delta L$$

X. Ji, Phy.Rev.Lett.78,610(1997)

Accessing GPDs through DVCS

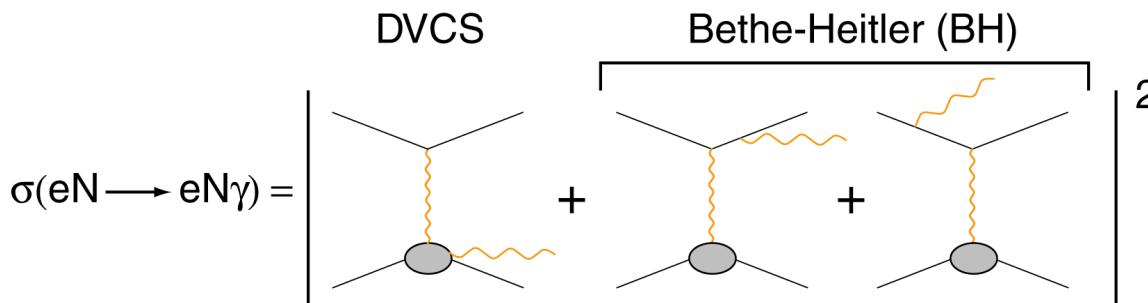
DVCS allows access to 4 complex GPDs-related quantities: Compton Form Factors (ξ, t)

$$T^{DVCS} \sim P \int_{-1}^{+1} \frac{GPDs(x, \xi, t)}{x \pm \xi} dx \pm i\pi GPDs(\pm \xi, \xi, t) + \dots$$

Only ξ and t are accessible experimentally

$$Re \mathcal{H}_q = e_q^2 P \int_0^{+1} \left(H^q(x, \xi, t) - H^q(-x, \xi, t) \right) \left[\frac{1}{\xi - x} + \frac{1}{\xi + x} \right] dx$$

$$Im \mathcal{H}_q = \pi e_q^2 [H^q(\xi, \xi, t) - H^q(-\xi, \xi, t)]$$

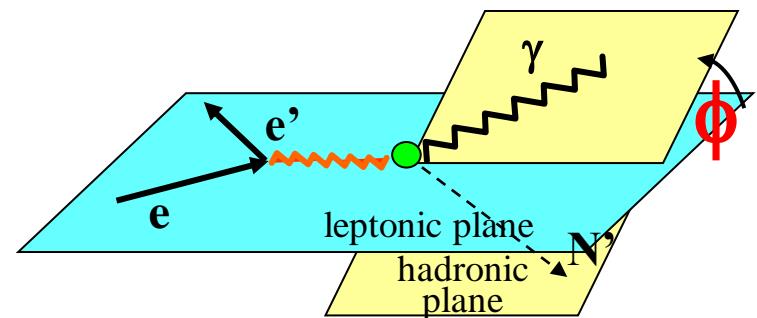


BH is calculable
(electromagnetic FFs)

$$\sigma \sim |T^{DVCS} + T^{BH}|^2 \propto \text{Re}(CFFs) \quad (\text{also DSA})$$

$$\Delta\sigma = \sigma^+ - \sigma^- \propto I(DVCS \cdot BH) \propto \text{Im}(CFFs)$$

$$A = \frac{\Delta\sigma}{2\sigma} \propto \frac{I(DVCS \cdot BH)}{|BH|^2 + |DVCS|^2 + I}$$



Sensitivity to CFFs of DVCS spin observables

$$A_{LU(UL)} = \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-} \propto \frac{s_{1,unp(UL)}^I \sin \phi}{c_{0,unp}^{BH} + c_{0,unp}^I + (c_{1,unp}^{BH} + c_{1,unp}^I) \cos \phi}$$

$$A_{LL} = \frac{\sigma^{++} + \sigma^{+-} - \sigma^{+-} - \sigma^{+-}}{\sigma^{++} + \sigma^{+-} + \sigma^{+-} + \sigma^{+-}} \propto \frac{c_{0,LP}^{BH} + c_{0,LP}^I + (c_{1,LP}^{BH} + c_{1,LP}^I) \cos \phi}{c_{0,unp}^{BH} + c_{0,unp}^I + (c_{1,unp}^{BH} + c_{1,unp}^I) \cos \phi}$$

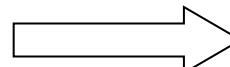
Twist 2
approximation
(-t << Q²)

$$(\xi = x_B/(2-x_B) \quad k = -t/4M^2)$$

Proton Neutron

Polarized beam, unpolarized target:

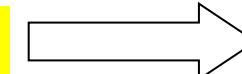
$$s_{1,unp}^I \sim \sin \phi \text{ Im}\{F_1 \mathcal{H} + \xi(F_1+F_2)\tilde{\mathcal{H}} - kF_2 \mathcal{E}\}$$



$$\begin{aligned} &\text{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p, \mathcal{E}_p\} \\ &\text{Im}\{\mathcal{H}_n, \tilde{\mathcal{H}}_n, \mathcal{E}_n\} \end{aligned}$$

Unpolarized beam, longitudinal target:

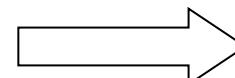
$$s_{1,UL}^I \sim \sin \phi \text{ Im}\{F_1 \tilde{\mathcal{H}} + \xi(F_1+F_2)(\mathcal{H} + x_B/2\mathcal{E}) - \xi k F_2 \tilde{\mathcal{E}} + \dots\}$$



$$\begin{aligned} &\text{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\} \\ &\text{Im}\{\mathcal{H}_n, \mathcal{E}_n\} \end{aligned}$$

Polarized beam, longitudinal target:

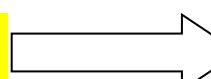
$$c_{1,Lp}^I \sim (A+B\cos\phi) \text{ Re}\{F_1 \tilde{\mathcal{H}} + \xi(F_1+F_2)(\mathcal{H} + x_B/2\mathcal{E}) \dots\}$$



$$\begin{aligned} &\text{Re}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\} \\ &\text{Re}\{\mathcal{H}_n, \mathcal{E}_n\} \end{aligned}$$

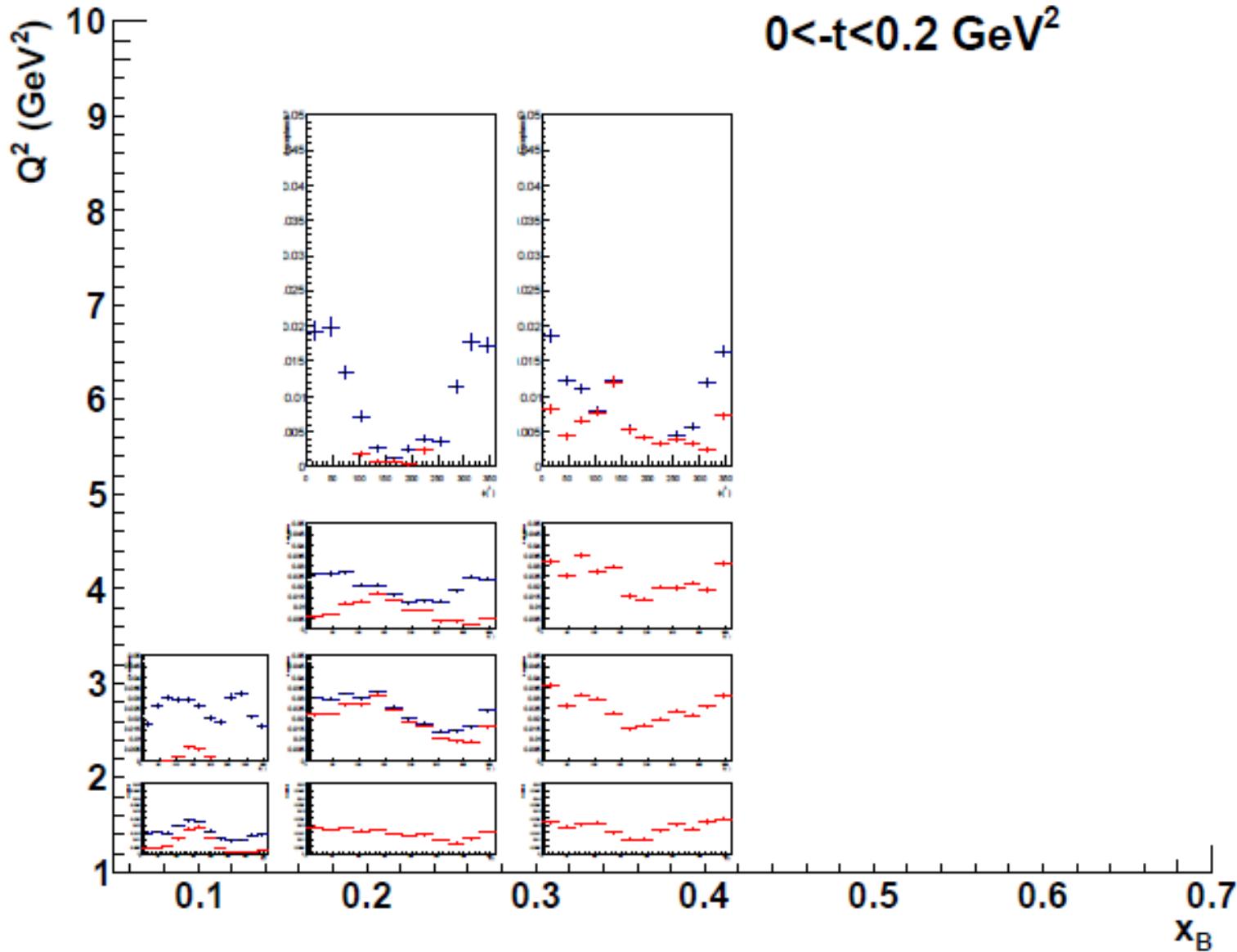
Unpolarized beam, transverse target:

$$\Delta\sigma_{UT} \sim \cos\phi \sin(\phi_s - \phi) \text{ Im}\{k(F_2 \mathcal{H} - F_1 \mathcal{E}) + \dots\}$$

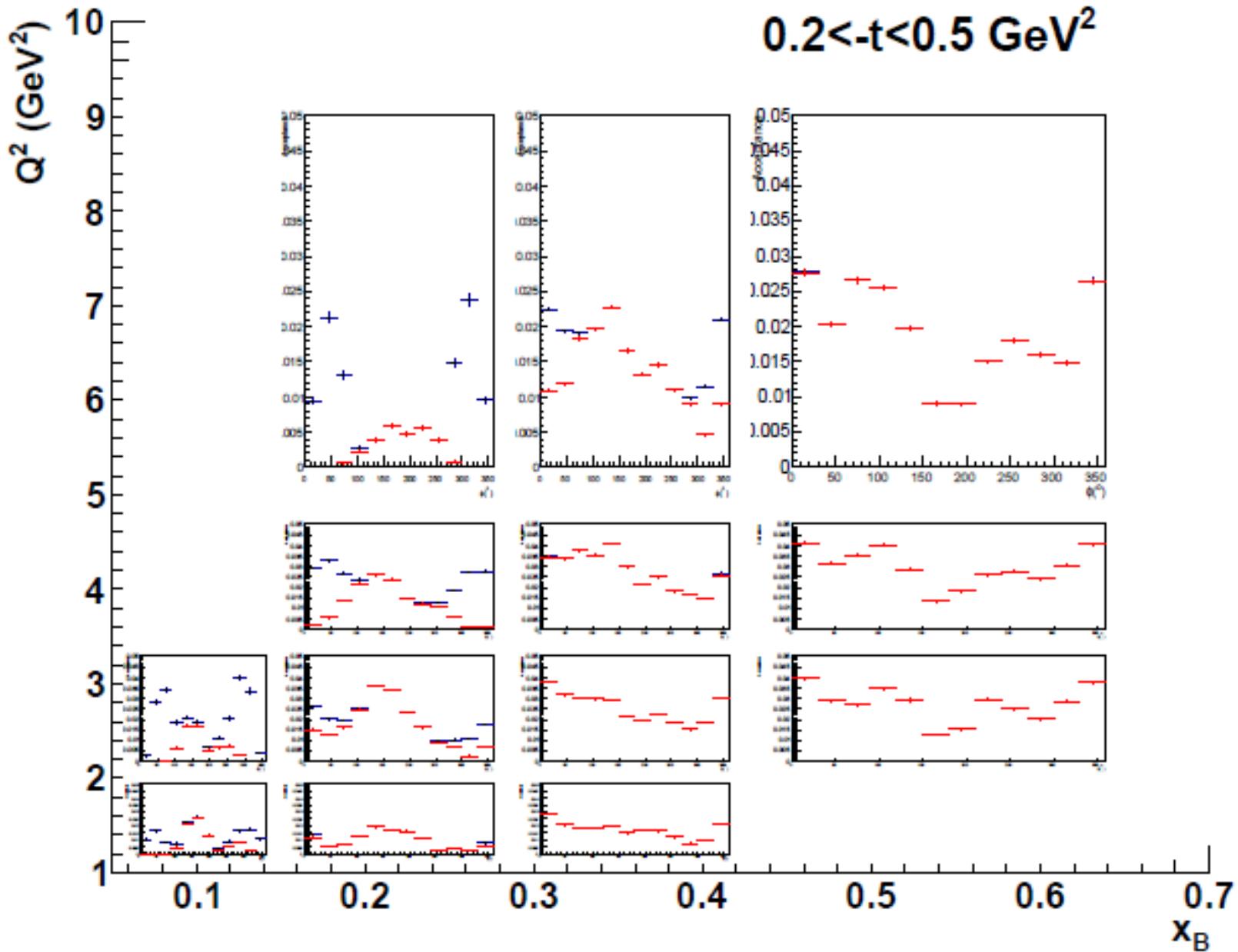


$$\begin{aligned} &\text{Im}\{\mathcal{H}_p, \mathcal{E}_p\} \\ &\text{Im}\{\mathcal{H}_n\} \end{aligned}$$

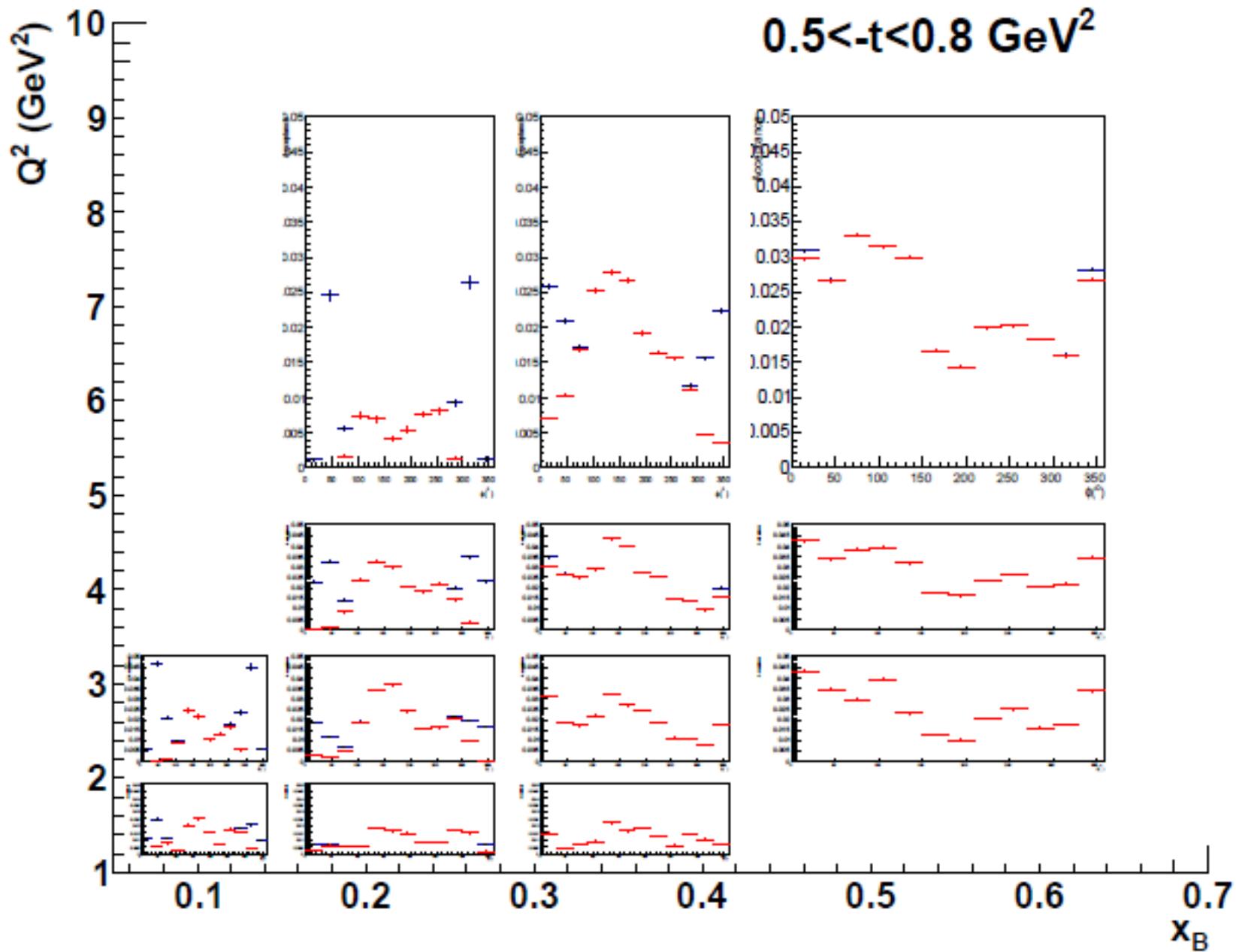
Effect of the Forward Tagger on the acceptance for nDVCS



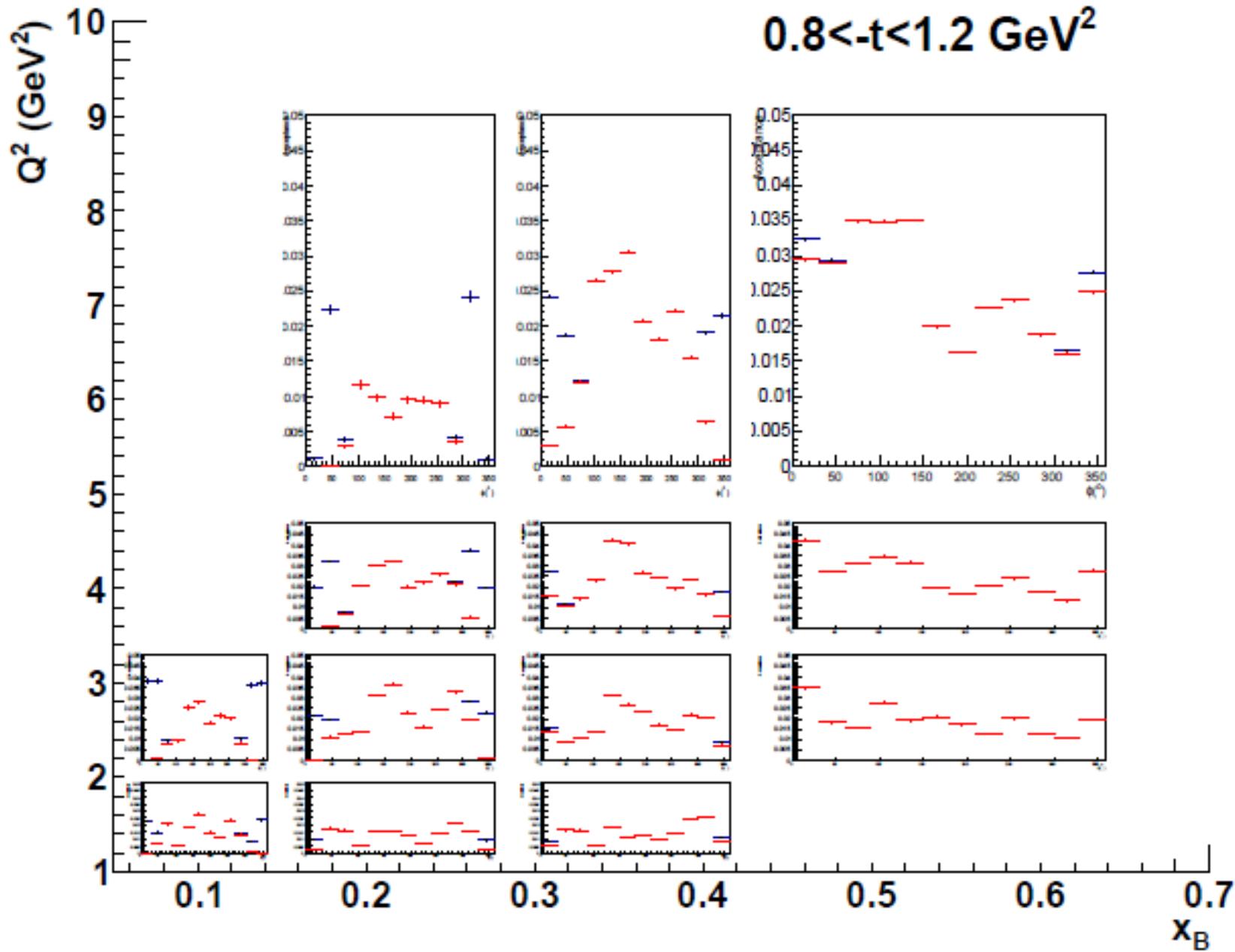
Effect of the Forward Tagger on the acceptance for nDVCS



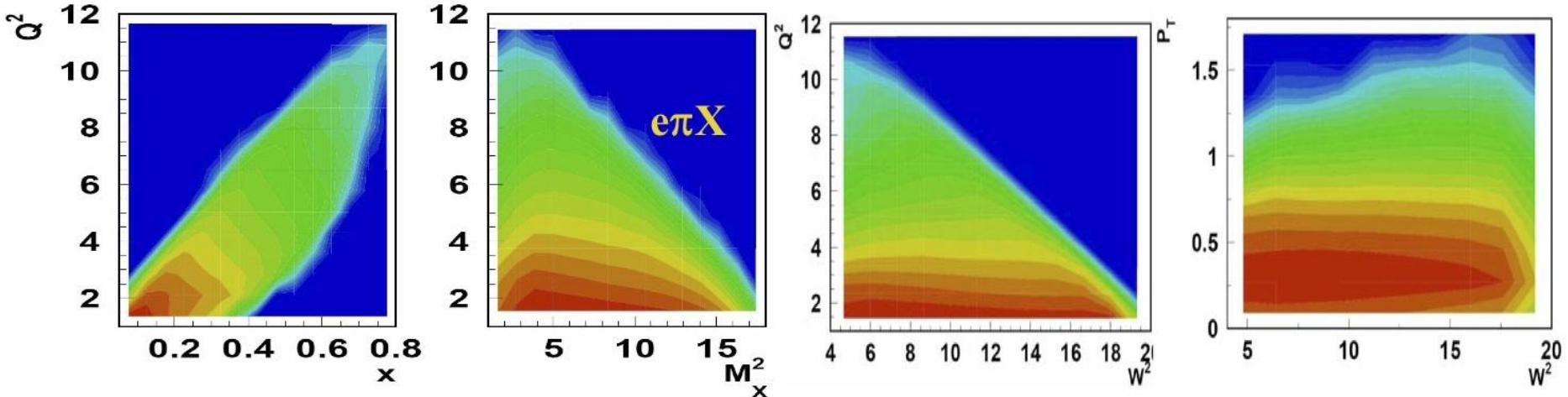
Effect of the Forward Tagger on the acceptance for nDVCS



Effect of the Forward Tagger on the acceptance for nDVCS



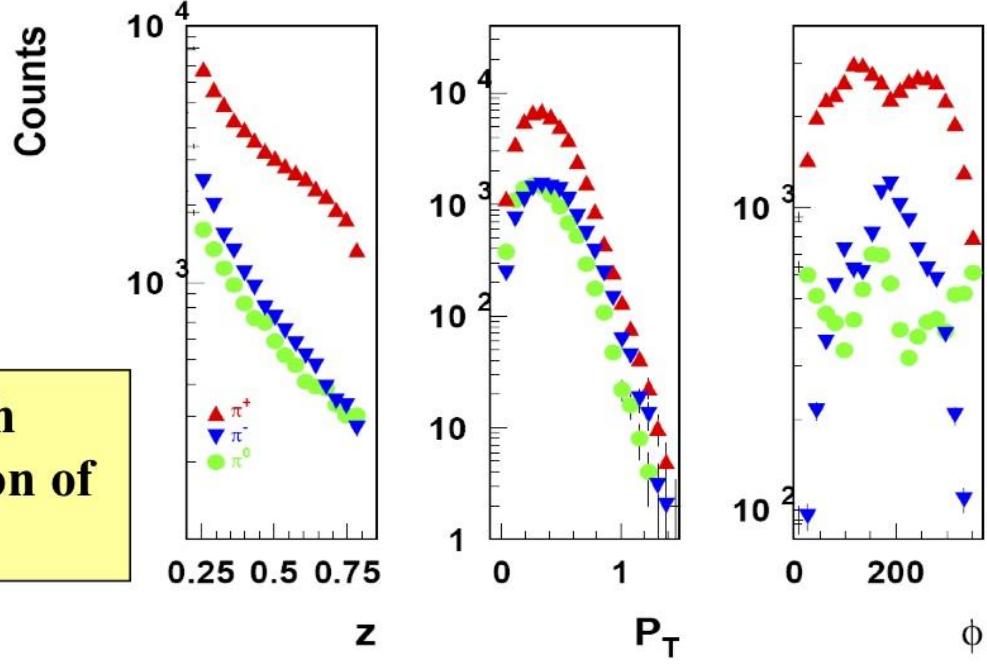
Pion distributions in $e p \rightarrow e' \pi X$



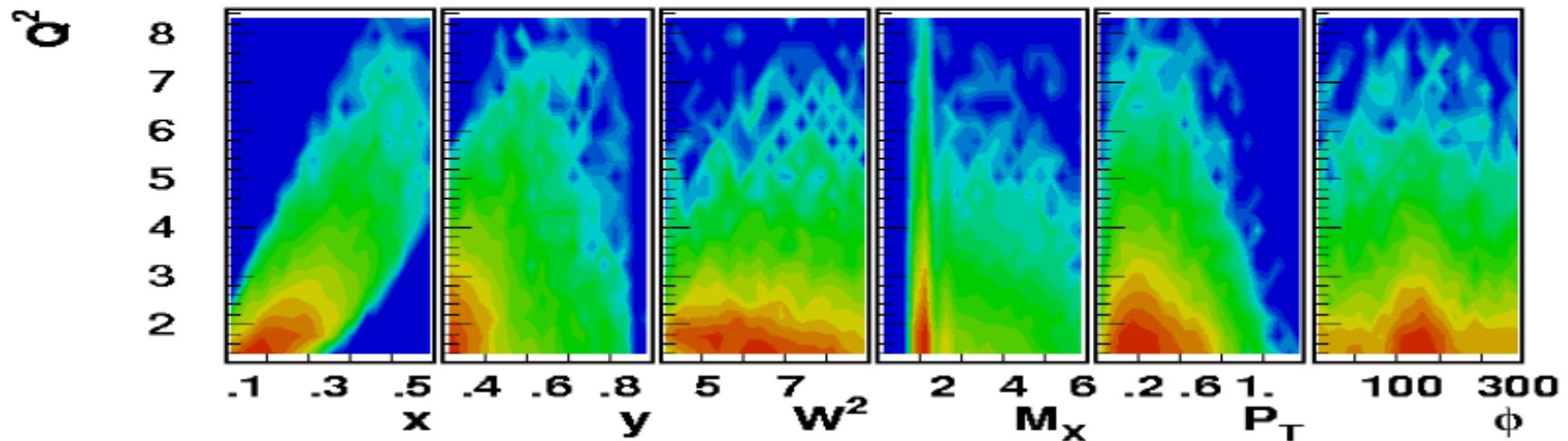
SIDIS
kinematics

$Q^2 > 1 \text{ GeV}^2$
 $W^2 > 4 \text{ GeV}^2(10)$
 $y < 0.85$
 $M_X > 2 \text{ GeV}$

Large Q^2 , M_X , and P_T accessible with CLAS12 are important for separation of Higher Twist contributions

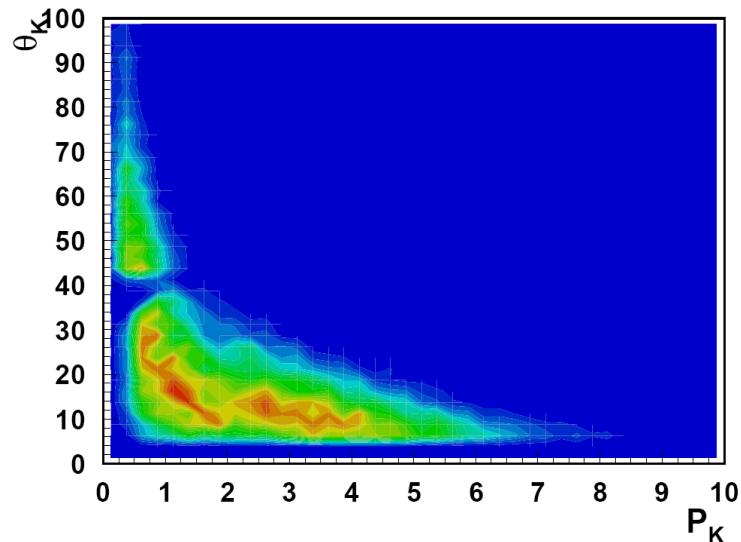


Kaon distributions in $e p \rightarrow e' K X$

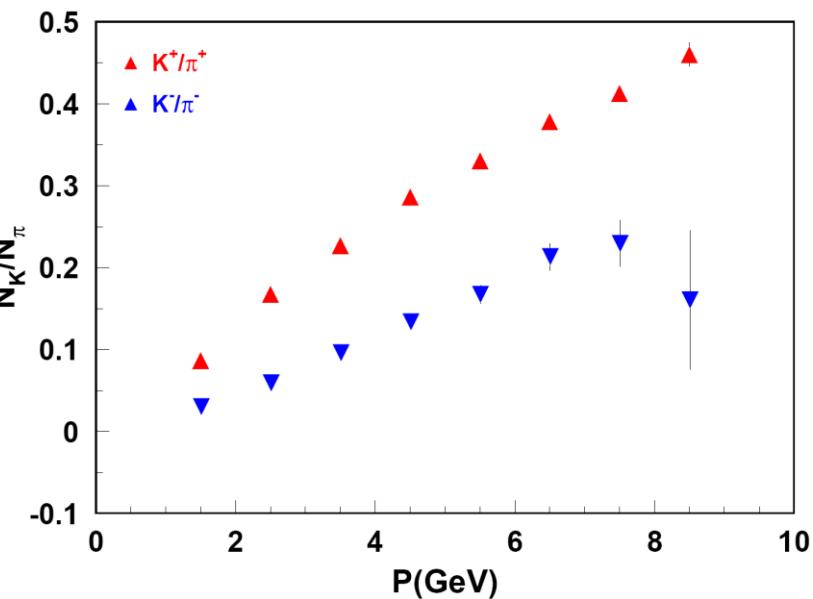


SIDIS kinematic

S
 $Q^2 > 1 \text{ GeV}^2$
 $W^2 > 4 \text{ GeV}^2(10)$
 $y < 0.85$
 $M_X > 2 \text{ GeV}$

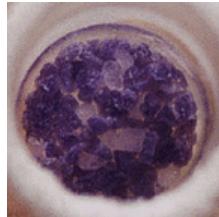


**High energy kaons are at small angles
($\theta < 30^\circ$)**



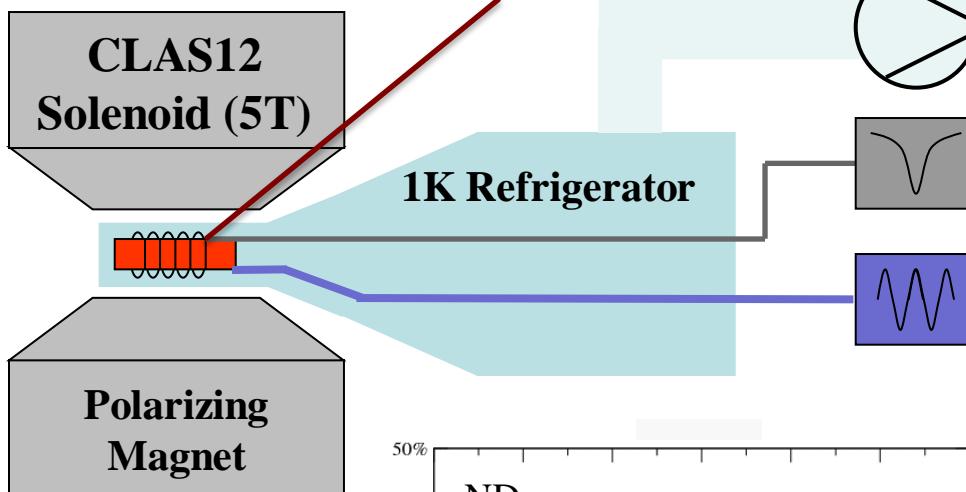
The CLAS12 longitudinally polarized target

Irradiated
deuterated
ammonia
(ND_3)

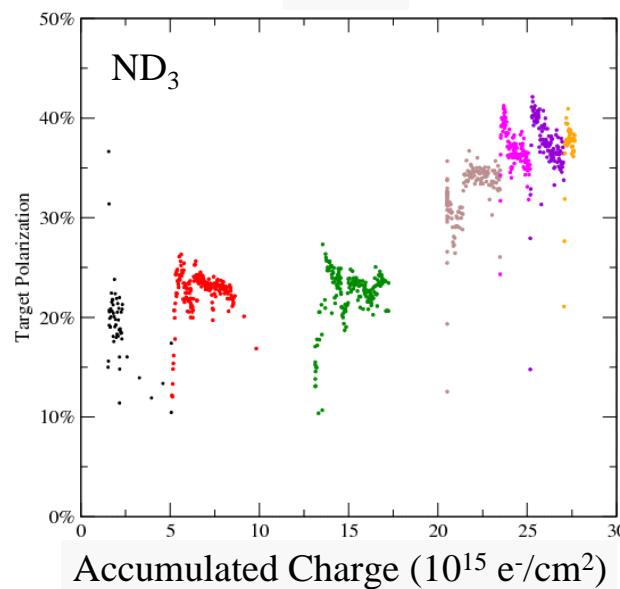


UVA, ODU,
CNU and JLab

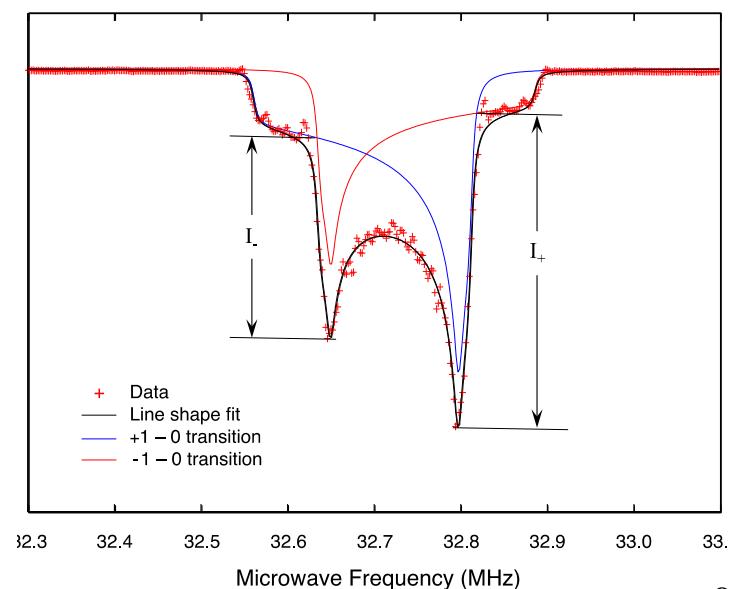
$L=4 \text{ cm}$; $r = 1.25 \text{ cm}$
 $P \sim 40\%$
 $L \sim 10^{34} \text{ cm}^2 \text{s}^{-1} / \text{nA}$

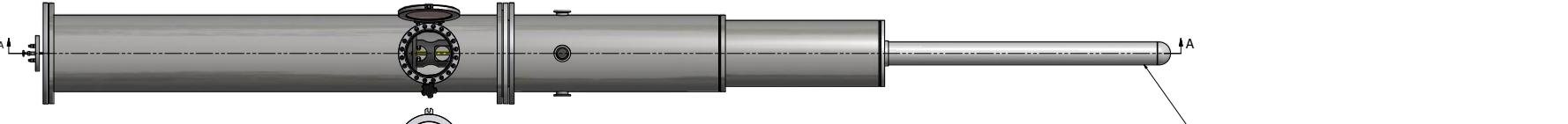


ND_3 polarization
improves with
“cold dose”

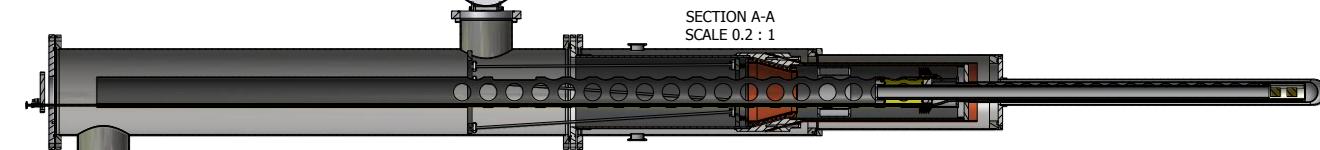


NMR line shape of ND₃

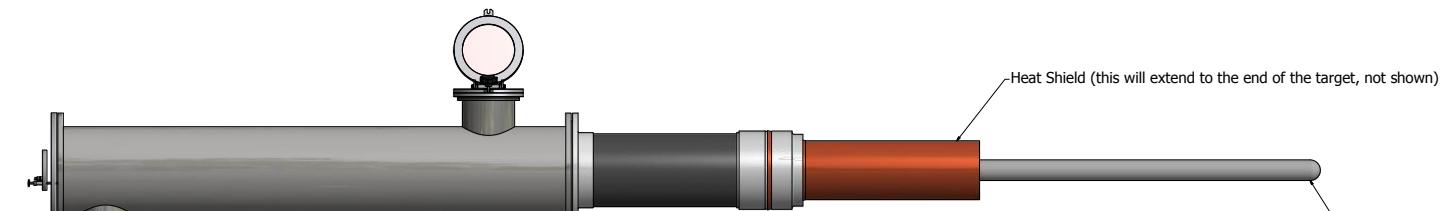




Insulating Vacuum Can (Thin scattering nose piece not shown)

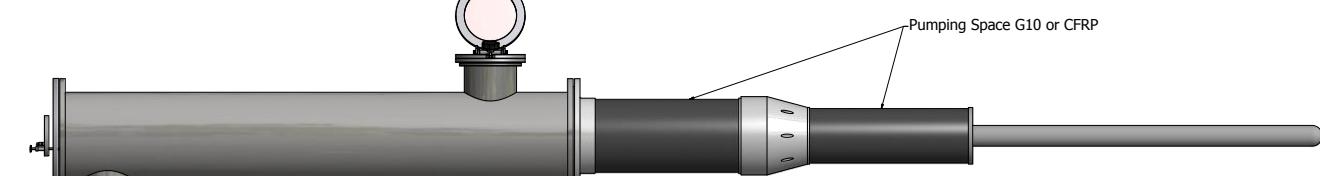


SECTION A-A
SCALE 0.2 : 1

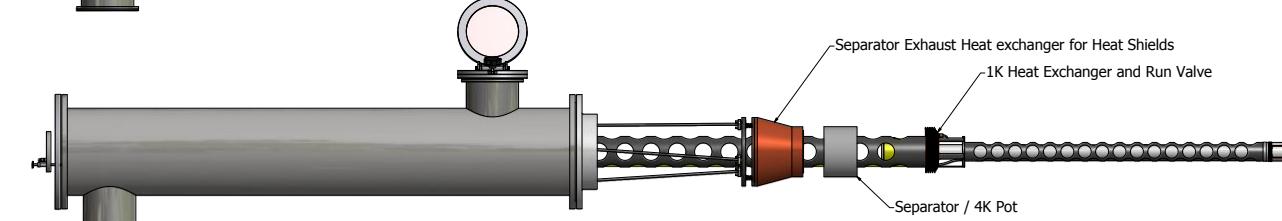


Heat Shield (this will extend to the end of the target, not shown)

1 K Pumping Space Aluminum (Filled with low pressure 1K He)



Pumping Space G10 or CFRP



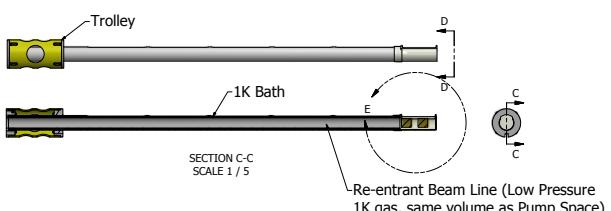
Separator / 4K Pot

1K Heat Exchanger and Run Valve

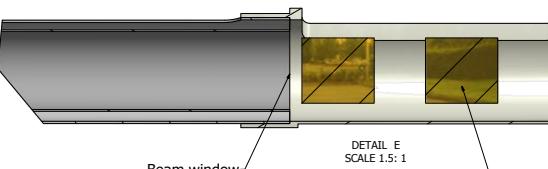
Separator Exhaust Heat exchanger for Heat Shields



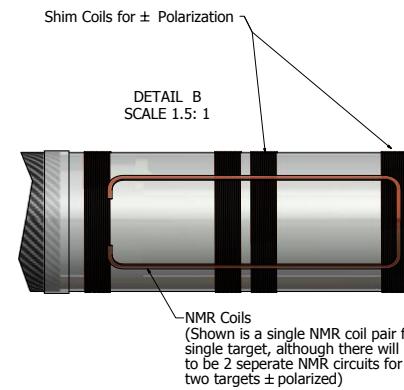
DETAIL E
SCALE 1.5 : 1



SECTION F-F
SCALE 1 / 5



DETAIL E
SCALE 1.5 : 1



DETAIL B
SCALE 1.5 : 1

NMR Coils
(Shown is a single NMR coil pair for a single target, although there will need to be 2 separate NMR circuits for the two targets ± polarized)