

## Chapter 32 Electro-magnetic Waves

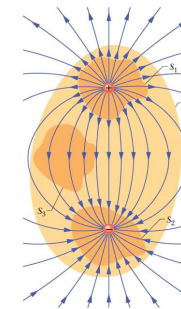
Maxwell Equations  
Electro-magnetic spectrum  
Speed of Light  
Energy and Momentum – Poynting Vector and Radiation Pressure  
Standing waves

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### Chapter 32 – El-mg Waves

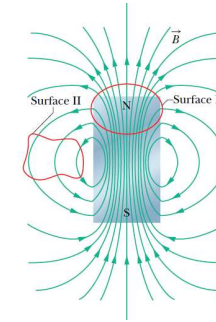
#### Maxwell's Equation – Gauss' Laws

Gauss' law for ELECTRIC field



$$\oiint \vec{E} \cdot d\vec{A} = \frac{1}{\epsilon_0} q_{enc}$$

Gauss' law for MAGNETIC field



$$\oiint \vec{B} \cdot d\vec{A} = 0$$

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### Chapter 32 – El-mg Waves

#### Maxwell's Equation – Integral form

Law	Equation	Relates ...
Gauss' law for electric field	$\oiint \vec{E} \cdot d\vec{A} = \frac{1}{\epsilon_0} q_{enc}$	net electric flux to net enclosed electric charge
Gauss' law for magnetic field	$\oiint \vec{B} \cdot d\vec{A} = 0$	net magnetic flux to net enclosed magnetic charge
Faraday/Lenz law of induction	$\oint \vec{E} \cdot d\vec{s} = -\frac{d\Phi_B}{dt}$	induced electric field to change in magnetic flux
Maxwell-Ampere law	$\oint \vec{B} \cdot d\vec{s} = \mu_0 i_{enc} + \epsilon_0 \mu_0 \frac{d\Phi_E}{dt}$	induced magnetic field to change in electric flux and to current

#### Conclusions:

Electric and magnetic fields can be separated for electrostatic charges and/or steady currents.  
Time-varying fields are no longer independent.  
Every accelerating charge radiates el.mg. energy, in particular a charge in SHM – see Fig 32.3 page 1053.

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### Chapter 32 – El-mg Waves

#### Maxwell's Equation – Integral and Differential Forms

Law	Equation	Relates ...
Gauss' law for electric field	$\oiint \vec{E} \cdot d\vec{A} = \frac{1}{\epsilon_0} q_{enc}$	net electric flux to net enclosed electric charge
	$\vec{\nabla} \cdot \vec{E} = \frac{1}{\epsilon_0} \rho_{enc}$	
Gauss' law for magnetic field	$\oiint \vec{B} \cdot d\vec{A} = 0$	net magnetic flux to net enclosed magnetic charge
	$\vec{\nabla} \cdot \vec{B} = 0$	
Faraday/Lenz law of induction	$\oint \vec{E} \cdot d\vec{s} = -\frac{d\Phi_B}{dt}$	induced electric field to change in magnetic flux
	$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$	
Maxwell-Ampere law	$\oint \vec{B} \cdot d\vec{s} = \epsilon_0 \mu_0 \frac{d\Phi_E}{dt} + \mu_0 i_{enc}$	induced magnetic field to change in electric flux and to current
	$\vec{\nabla} \times \vec{B} = \epsilon_0 \mu_0 \frac{\partial \vec{E}}{\partial t} + \mu_0 \vec{j}$	

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Chapter 15- Waves  
Mathematical Description - Review

**Sinusoidal wave form** (shape of the wave) described by a function

$$y(x, t) = A \sin(kx - \omega t + \phi) = A \sin\left(2\pi\left(\frac{x}{\lambda} - \frac{t}{T}\right) + \phi\right) = A \sin(2\pi k(x - vt) + \phi)$$

Wave equation  $\frac{\partial^2 y}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 y}{\partial t^2}$  where  $v = \frac{\omega}{k} = \frac{\lambda}{T} = \lambda f$

$y(x, t)$  displacement at time  $t$  and position  $x$   
 $A$  amplitude  $>0$

$\sin(kx - \omega t)$  magnitude of the maximum displacement of the elements from their eq  
 $(kx - \omega t)$  oscillating term  
 phase

$k$  angular wave (phase) number  
 related to wavelength  $\lambda$  by  $k = \frac{2\pi}{\lambda}$

$\omega$  angular frequency  
 related to time period  $T$  by  $\omega = \frac{2\pi}{T} = 2\pi f$

$\phi$  phase constant  
 the value of  $\phi$  is determined by the initial conditions ( $x=0, t=0$ )

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Chapter 32 - El-mg Waves  
Electro-magnetic Waves

**El-mg waves** = changing electric & magnetic fields traveling through vacuum and obeying Maxwell equations (w/out involving moving charges or currents)

Maxwell (1850s) showed this for light (el.mg waves traveling at the  $c$ ) and predicted radio waves  
 Heinrich Hertz (1888) - Helmholtz's student - experimental demo for radio waves (& sparks)

Postulate field configurations with wave-like behavior traveling in the  $+x$  direction in phase

$$\vec{E} = E_{max} \sin(kx - \omega t) \hat{j}$$

$$\vec{B} = B_{max} \sin(kx - \omega t) \hat{k}$$

E, B travel as package in  $x$  direction  
 with  $\lambda = \frac{2\pi}{k}$  and  $T = \frac{2\pi}{\omega}$

Test if they satisfy Maxwell's equations  
 Show they propagate at the speed of light  $c = 3 \cdot 10^8 \frac{m}{s}$   
 See YF derivation page 1055-1057

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Chapter 32 - El-mg Waves  
Electro-magnetic Waves

**Faraday's Law**  
 $\oint \vec{E} \cdot d\vec{s} = -\frac{d}{dt} \oint \vec{B} \cdot d\vec{A}$

**Ampere's law**  
 $\oint \vec{B} \cdot d\vec{s} = \epsilon_0 \mu_0 \frac{d}{dt} \oint \vec{E} \cdot d\vec{A}$

Harmonic wave solutions  
 $\vec{E} = E_{max} \cos(kx - \omega t) \hat{j}$   
 $\vec{B} = B_{max} \cos(kx - \omega t) \hat{k}$

$$((E + dE) - E) h = -\frac{d}{dt} B dx h \quad \therefore dx \quad dB h = -\epsilon_0 \mu_0 \frac{d}{dt} E dx h \quad \therefore dx$$

$$\frac{\partial E}{\partial x} = -\frac{\partial B}{\partial t} \quad \therefore \frac{\partial}{\partial x} \quad \frac{\partial B}{\partial x} = -\epsilon_0 \mu_0 \frac{\partial E}{\partial t} \quad \therefore \frac{\partial}{\partial t}$$

$$\frac{\partial^2 E}{\partial x^2} = -\frac{\partial}{\partial x} \frac{\partial B}{\partial t} = -\frac{\partial}{\partial t} \frac{\partial B}{\partial x} \quad \frac{\partial}{\partial t} \frac{\partial B}{\partial x} = -\epsilon_0 \mu_0 \frac{\partial^2 E}{\partial t^2}$$

$$\frac{\partial^2 E}{\partial x^2} = \epsilon_0 \mu_0 \frac{\partial^2 E}{\partial t^2} \quad \text{where} \quad \epsilon_0 \mu_0 = \frac{1}{v^2}$$

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Chapter 32 - El-mg Waves  
Speed of Light

Speed of the electromagnetic waves

$$c = \lambda f = \frac{\omega}{k} = \frac{1}{\sqrt{\epsilon_0 \mu_0}} = 3 \cdot 10^8 \frac{m}{s}$$

Compare to  $\omega = \sqrt{\frac{k}{m}} = \frac{2\pi}{T}$  SHM of mass on spring  
 $v = \sqrt{\frac{FT}{\mu}}$  Mechanical wave on taut string

Relationship between electric and magnetic field

$$\frac{E(t)}{B(t)} = \frac{E_{max}}{B_{max}} = c$$

Fields are in phase

Proof:  $E = E_{max} \sin(kx - \omega t) \quad \therefore \frac{\partial}{\partial x}$   
 $\frac{\partial E}{\partial x} = k E_{max} \cos(kx - \omega t)$   
 Use  $\frac{\partial E}{\partial x} = -\frac{\partial B}{\partial t}$  from Faraday's law  $\therefore \frac{\partial}{\partial t}$   
 $\int k E_{max} \cos(kx - \omega t) dt = \int -\frac{\partial B}{\partial t} dt$   
 $\frac{k}{\omega} E_{max} \sin(kx - \omega t) = \frac{E_{max}}{c} \sin(kx - \omega t) = \frac{E}{c} = B$

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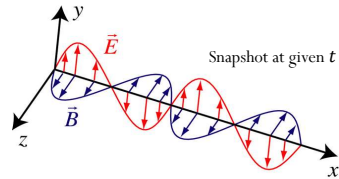
## Chapter 32 – El-mg Waves

## Key Properties of electro-magnetic Waves

$$\vec{E}(\vec{r}, t) = E_{\max} \cos(kx - \omega t) \hat{j}$$

$$\vec{B}(\vec{r}, t) = B_{\max} \cos(kx - \omega t) \hat{k}$$

E, B travel as package in x direction  
with  $\lambda = \frac{2\pi}{k}$  and  $T = \frac{2\pi}{\omega}$



## Properties of ALL el.mg. waves

1. Transverse wave
2. Fields are perpendicular to each other  
Fields are perpendicular to direction of propagation
3. Fields are in phase
4. Satisfy Maxwell equations under two conditions
  - a) field amplitudes/magnitudes are related
  - b) wave speed in vacuum
5. No medium is required for propagation  
In dielectric medium

$$\vec{E} \perp \hat{i} ; \vec{B} \perp \hat{i}$$

$$\vec{E} \perp \vec{B}$$

$$\vec{E} \times \vec{B} \sim \hat{i}$$

$$E = c B$$

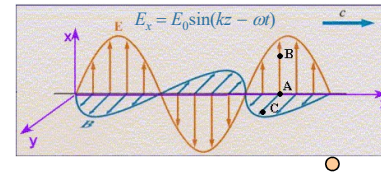
$$c = \frac{1}{\sqrt{\epsilon_0 \mu_0}} = 3 \cdot 10^8 \frac{m}{s}$$

$$v = \frac{1}{\sqrt{\epsilon \mu}} < c$$

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## Chapter 32 – El-mg Waves

## Key Properties of electro-magnetic Waves



An electro-magnetic wave is traveling in the z direction.

1. The points A, B, C have the same z coordinate. Compare the magnitude of the electric field at A and B:

- A)  $E_A > E_B$
- B)  $E_A = E_B$
- C)  $E_A < E_B$

2. Consider a point (x,y,z) at time t when  $E_x$  is negative and has its maximum value.

What is  $B_y$ ?

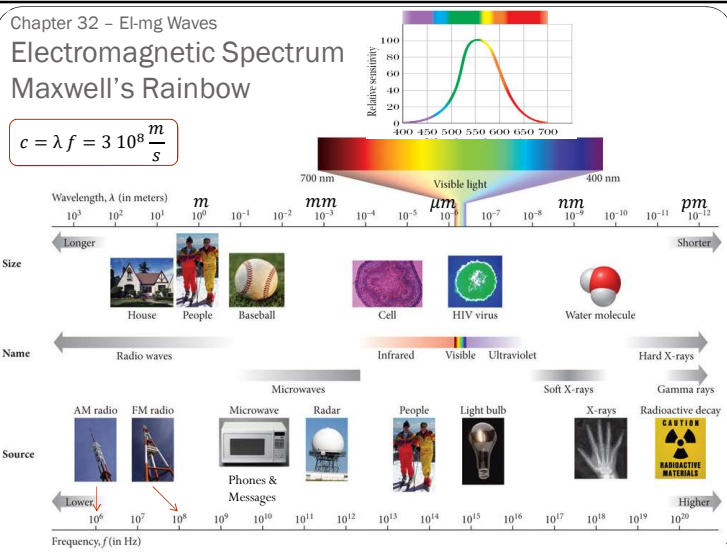
- A)  $B_y$  is positive and has its maximum value
- B)  $B_y$  is negative and has its maximum value
- C)  $B_y$  is zero
- D) We do not have enough information

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## Chapter 32 – El-mg Waves

Electromagnetic Spectrum  
Maxwell's Rainbow

$$c = \lambda f = 3 \cdot 10^8 \frac{m}{s}$$



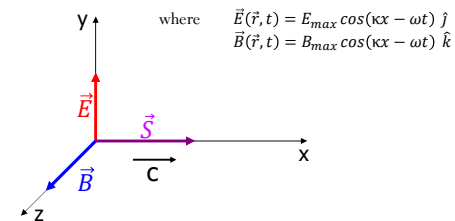
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## Chapter 32 – El-mg Waves

## Energy Transport

El-mg waves carry energy and its rate of transport per unit area is described by:

Poynting vector  $\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B}$  John Henry Poynting 1884



Direction: direction of traveling wave and energy transport

Magnitude:  $S(t) = \frac{1}{\mu_0} E B = \frac{1}{\mu_0} E \frac{E}{c} = \frac{E^2}{c \mu_0} = c \epsilon_0 E^2$

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## Chapter 32 – El-mg Waves

## Energy Transport

Relate magnitude of  $\vec{S}$  with energy contained in the electric and magnetic field of the el-mg. wave

Energies densities are the same

$$u_E = u_B \quad \text{using} \quad c \epsilon_0 = 1/c \mu_0$$

$$\text{Prove} \quad u_E = \frac{1}{2} \epsilon_0 E^2 \stackrel{c}{=} \frac{1}{2c} \frac{1}{c \mu_0} E^2 = \frac{1}{2 \mu_0} \left(\frac{E}{c}\right)^2 = \frac{1}{2 \mu_0} B^2 = u_B$$

Total energy density

$$u = u_E + u_B = \epsilon_0 E^2$$

$$S = c u$$

Take averages

$$S_{avg} = c u_{ave} = c \epsilon_0 (E^2)_{avg} = c \epsilon_0 \frac{1}{2} E_m^2 = c \epsilon_0 E_{rms}^2$$

Intensity of the wave

$$I = \frac{P_{ave}}{A} = \frac{1}{A} \frac{u_{ave} vol}{t} = \frac{1}{A} \frac{u_{ave}(c t) A}{t} = c u_{ave}$$

Conclusion

$$I = S_{avg} = \frac{1}{c \mu_0} E_{rms}^2$$

$$\text{Unit:} \quad [S] = \frac{W}{m^2}$$

Example: average intensity of sunlight on Earth  $\sim 1400 \frac{W}{m^2} \sim 140 \frac{mW}{cm^2}$

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## Chapter 32 – El-mg Waves

## Radiation Pressure

**Radiation pressure** of el-mag waves ( $p_r$ )

Assume el-mg wave incident on object and totally absorbed over  $\Delta t$

$U$  = energy carried by wave

$\Delta p = \frac{\Delta U}{c}$  (Maxwell) change in linear momentum of an object after absorbing  $\Delta U$

$$p_r = \frac{F}{A} = \frac{1}{A} \frac{\Delta p}{\Delta t} = \frac{1}{A} \frac{\frac{\Delta U}{c}}{\Delta t} = \frac{1}{c} \frac{1}{A} \frac{\Delta U}{\Delta t} = \frac{I}{c}$$

$$\text{Total absorption:} \quad p_r = \frac{I}{c} \quad \text{since} \quad \Delta p = p_i - 0 = p_i$$

Assume el-mg wave incident on object and totally reflected over  $\Delta t$

$\Delta p = \frac{2 \Delta U}{c}$  (Maxwell) change in linear momentum of an object after reflecting  $\Delta U$

$$\text{Perfect reflection} \quad p_r = \frac{2I}{c} \quad \text{since} \quad \Delta p = p_i - (-p_i) = 2p_i$$

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## Chapter 32 – El-mg Waves

## Radiation Pressure

Examples:

$$\text{Atmospheric pressure} \quad p_0 = 101 \text{ kPa} \approx 10^5 \text{ Pa}$$

$$\text{Pressure from sun} \quad p_r = \frac{I}{c} = \frac{1400 \text{ W/m}^2}{3 \cdot 10^8 \text{ m/s}} = 4.7 \text{ } \mu\text{Pa} = 4.6 \cdot 10^{-1} p_0$$

$$\text{Laser pointer} \quad p_r = \frac{2I}{c} = \frac{2 \frac{P}{\pi r^2}}{c} = \frac{2 \cdot 10^{-3}}{\pi \cdot 10^{-6} \cdot 3 \cdot 10^8} = 2.1 \text{ } \mu\text{Pa}$$

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## Chapter 32 – El-mg Waves

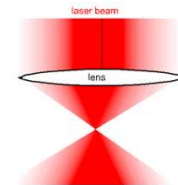
## Radiation Pressure : Applications

Optical (laser) tweezers

intense laser focused on small area

1987 Arthur Ashkin - used in cell biology

2018 ½ Nobel Prize



Solar sail

IKARUS – Japan, 2010-2013

(Interplanetary Kite-craft Accelerated by Radiation Of the Sun)

<https://www.youtube.com/watch?v=7Mb47w0vB04>

7.5  $\mu\text{m}$  over 200  $\text{m}^2$  (14  $\text{m}$  square); 2  $\text{kg}$

reflectance-adjustable LCD panel

Sunjammer – NASA, Jan 2015 (planned)

1200  $\text{m}^2$  (35  $\text{m}$  square); 32  $\text{kg}$

cancelled >4y, >\$21M

“It looks like it just wasn’t big enough for us to afford it,”

Dana Rohrabacher (R-Calif.)



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