

hw 3.

### Problem 1.

a) for  $\forall u \in X$ , and  $v \in X$  satisfying the strong formulation, suppose that  $w = u + v$

$$\begin{aligned} \text{then } \frac{1}{2} a(w, w) - l(w) &= \frac{1}{2} a(u, u) - l(u) && \text{(use linear and bilinear)} \\ &+ a(u, v) - l(v) \\ &+ \frac{1}{2} a(v, v) \end{aligned}$$

$$\begin{aligned} \text{where } a(u, v) - l(v) &= \int_0^{1/2} k^L u_x v_x - f^L v \, dx + \int_{1/2}^1 k^R u_x v_x - f^R v \, dx \\ &= k^L u_x^L v \Big|_0^{1/2} + \int_0^{1/2} -k^L u_{xx}^L v - f^L v \, dx \\ &\quad + k^R u_x^R v \Big|_{1/2}^1 + \int_{1/2}^1 (-k^R u_{xx}^R v - f^R v) \, dx \end{aligned}$$

$$\text{with } u(0) = v(1) = 0, \quad k^L u_x^L(1/2) = k^R u_x^R(1/2), \text{ we have } k^L u_x^L v \Big|_0^{1/2} + k^R u_x^R v \Big|_{1/2}^1 = 0$$

$$\text{and } -k^L u_{xx}^L - f^L = -k^R u_{xx}^R - f^R = 0, \text{ so we have}$$

$$a(u, v) - l(v) = 0 \quad (1)$$

$$\text{Furthermore, } a(v, v) = \int_0^{1/2} k^L v_x v_x \, dx + \int_{1/2}^1 k^R v_x v_x \, dx \geq 0$$

$$\begin{aligned} \text{so } \frac{1}{2} a(w, w) - l(w) &= \frac{1}{2} a(u, u) - l(u) + \frac{1}{2} a(v, v) \\ &\geq \frac{1}{2} a(u, u) - l(u) \end{aligned}$$

$$\text{which means that } u = \arg \min_{w \in X} \frac{1}{2} a(w, w) - l(w)$$

from (1), we know that

$$a(u, v) = l(v)$$

b)  $u^L(0) = 0, u^R(1) = 0, u^L(1/2) = u^R(1/2)$  are essential boundary/interface conditions  
 $-k^L u_x^L(1/2) = -k^R u_x^R(1/2)$  is natural boundary/interface condition.

c) because  $k^L u_x^L(1/2) = k^R u_x^R(1/2)$  leads to a delta function in  $x = 1/2$

$u_{xx}$  contains  $(u_x(1/2 + 0) - u_x(1/2 - 0)) \delta(x - 1/2)$ , because  
 $\int_0^1 u_{xx}^2 \, dx \geq \text{bounded terms} + \int_0^1 (\delta(x))^2 \, dx$  which is unbounded, so  $u \notin H^2(\Omega)$

However,  $\int_0^1 u_x^2 \, dx$  is bounded. Although there is a discontinuous point at  $1/2$ ,

but it will not cause any problem. so  $u \in H^1(\Omega)$

## Problem 2.

$$\begin{aligned}
 (a) \quad \int_V dV (\nabla^2 u + f) v &= \int_V dV (\nabla \cdot (\nabla u v) - (\nabla u) \cdot (\nabla v) + f v) \\
 &= \int_S dS (\vec{n} \cdot \nabla u) v + \int_V dV (-\nabla u \cdot \nabla v + f v) \\
 &= \int_S dS \frac{\partial u}{\partial n} v \Big|_{\Gamma} + \int_V dV (-\nabla u \cdot \nabla v + f v)
 \end{aligned}$$

because  $-\frac{\partial u}{\partial n} \Big|_{\Gamma_R} = h_c u \Big|_{\Gamma_R}$   $v \Big|_{\Gamma_D} = 0$

$$\int_V dV (\nabla^2 u + f) v = - \int_{\Gamma_R} dS h_c u v + \int_V dV (-\nabla u \cdot \nabla v + f v)$$

when  $u$  is the solution of strong formulation,  $\Delta u = 0$ , so  $RHS = 0$ .

suppose that  $a(u, v) = \int_{\Omega} dV \nabla u \cdot \nabla v + \int_{\Gamma_R} dS h_c u v$

$$l(v) = \int_{\Omega} dV f v$$

This satisfies that  $a(u, v) = l(v)$  for  $\forall v \in X$

functional  $J(w) = \frac{1}{2} a(w, w) - l(w)$

$$= \frac{1}{2} \int_{\Omega} dV (w)^2 + \frac{1}{2} \int_{\Gamma_R} dS h_c w^2 - \int_{\Omega} dV f w$$

(b)  $u \Big|_{\Gamma_D} = 0$  is essential,  $-\frac{\partial u}{\partial n} \Big|_{\Gamma_R} = h_c u \Big|_{\Gamma_R}$  is natural :

## Problem 3

$$\begin{aligned}
 (a) \quad \int_0^1 (u_{xxxx} - f) v dx &= u_{xxxx} v \Big|_0^1 - \int_0^1 u_{xxx} v_x dx - \int_0^1 f v dx \\
 &= u_{xxxx} v \Big|_0^1 - u_{xxx} v_x \Big|_0^1 + \int_0^1 u_{xxx} v_{xx} dx - \int_0^1 f v dx
 \end{aligned}$$

we can define  $a(u, v) = \int_0^1 u_{xx} v_{xx} dx$

$$\ell(v) = \int_0^1 f v \, dx$$

$$X = \{ v \in H^2(\Omega) \mid v(0) = v(1) = 0, v_x(0) = v_x(1) = 0 \}$$

$$\text{So that } \int_0^1 (u_{xxxx} - f) v \, dx = a(u, v) - \ell(v) \text{ for } \forall v \in X$$

if  $u$  is strong formulation solution, then  $a(u, v) = \ell(v)$  for  $\forall v \in X$

$$J(w) = \int_0^1 (w_{xx})^2 \, dx - \int_0^1 f w \, dx$$

$$(b) \quad X = \{ v \in H^2(\Omega) \mid v(0) = v(1) = 0, v_x(0) = v_x(1) = 0 \} :$$

here we use  $H^2(\Omega)$  because  $J(w)$  contains  $\int_0^1 w_{xx}^2 \, dx$ , which should be bounded

$$\begin{aligned} (c) \quad |\ell(v)| &= |v_x(\tfrac{1}{2})| \\ &= \left| \int_0^{\frac{1}{2}} u_{xx}(x) \, dx \right| \quad (\text{here we use } v_x(0) = 0 \text{ for } v \in X) \\ &\leq \sqrt{\int_0^{\frac{1}{2}} 1 \, dx} \sqrt{\int_0^{\frac{1}{2}} u_{xx}^2 \, dx} \\ &= \frac{1}{\sqrt{2}} \sqrt{\int_0^{\frac{1}{2}} u_{xx}^2 \, dx} \\ &\leq \frac{1}{\sqrt{2}} \sqrt{\int_0^1 u_{xx}^2 \, dx} \end{aligned}$$

$$\begin{aligned} \|v\|_X &= \sqrt{\int_0^1 v^2 + v_x^2 + v_{xx}^2 \, dx} \\ &\geq \sqrt{\int_0^1 v_{xx}^2 \, dx} \end{aligned}$$

$$\text{So we have } |\ell(v)| \leq \frac{1}{2} \|v\|_X, \text{ i.e. } C = \frac{1}{2}$$

Problem 4.

$$(a) \quad \forall v \in X$$

$$\begin{aligned} \int_0^1 (u_{xxxx} - f) v \, dx &= u_{xxx} v \Big|_0^1 - \int_0^1 u_{xxx} v_x \, dx - \int_0^1 f v \, dx \\ &= u_{xxx} v \Big|_0^1 - u_{xx} v_x \Big|_0^1 + \int_0^1 u_{xx} v_{xx} \, dx - \int_0^1 f v \, dx \end{aligned}$$

we can define  $a(u, v) = \int_0^1 u_{xxx} v_{xx} dx$

$$l(v) = \int_0^1 f v dx$$

because  $u_{xx}(0) = u_{xx}(1) = 0$ , so  $u_{xxx} v_{xx} \big|_0^1 = 0$

$v \in X$ , so  $v(0) = v(1) = 0$ , so  $u_{xxx} v \big|_0^1 = 0$

so

$$\int_0^1 (u_{xxx} - f) v dx = a(u, v) - l(v)$$

so we can obtain the same  $a(u, v)$  and  $l(v)$ , but

using only  $v(0) = v(1) = 0$

b)  $u(0) = 0, u(1) = 0$  are essential  $u_{xx}(0) = 0, u_{xx}(1) = 0$  are natural