Lab Notes

Haran Jackson

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Chapter 1

To Do

1.1 Coding

- Test split-WENO in 2D
- Implement in CUDA
- Solid EOSs
- Implement Barton's damage model
- Test compressible Euler vs GPR - use narrowing domain from Toro's book
- Try conservative formulation [Peshkov,Grmela,Romenski]

1.2 Theoretical

- Investiage approximate Riemann solvers (e.g. Dumbser's HLLEM)
- Convergence of conservation of mass in RGFM

1.3 Papers

- Isobaric cookoff (make cookoff the focus)
- Split solver vs ADER-WENO
- Application of ADER-WENO to Barton's solid model (with Tome)
- HPR-RGFM paper (solve stationarity of interface solver under some conditions)
- Application of ADER-WENO to equations in other areas of Physics, Biology, & Economics

Chapter 2

Faster Solvers

2.1 Improved Multidimensional WENO Oscillation Indicator

Given 2D reconstruction on stencil S_j:

$$w_j(x,y) = \tilde{w}_{pq}\psi_p(x)\psi_q(y) \tag{2.1}$$

the oscillation indicator for this reconstruction is [1]:

$$\begin{aligned}
o_j &= \sum_{|\alpha|=1}^{N} \int_{C_i} (D^\alpha w_j)^2 \, d\boldsymbol{x} \qquad (2.2)\\
&= \int_0^1 \left\{ \left(\frac{\partial w_j}{\partial x}\right)^2 + \left(\frac{\partial w_j}{\partial y}\right)^2 + \left(\frac{\partial^2 w_j}{\partial x^2}\right)^2 + \left(\frac{\partial^2 w_j}{\partial y^2}\right)^2 + \left(\frac{\partial^2 w_j}{\partial x \partial y}\right)^2 \right\} d\boldsymbol{\chi}\\
&= \int_0^1 \left\{ \left(\tilde{w}_{pq}\psi_p^{'}(x)\psi_q(y)\right)^2 + \left(\tilde{w}_{pq}\psi_p(x)\psi_q^{'}(y)\right)^2 + \left(\tilde{w}_{pq}\psi_p^{''}(x)\psi_q(y)\right)^2 + \left(\tilde{w}_{pq}\psi_p(x)\psi_q^{'}(y)\right)^2 + \left(\tilde{w}_{pq}\psi_p^{'}(x)\psi_q^{'}(y)\right)^2 \right\} d\boldsymbol{\chi}\\
&= \int_0^1 \left\{ \tilde{w}_{p_1 q_1}\psi_{p_1}^{'}(x)\psi_{q_1}(y)\tilde{w}_{p_2 q_2}\psi_{p_2}^{'}(x)\psi_{q_2}(y) + \cdots \right\} d\boldsymbol{\chi}\\
&= \tilde{w}_{p_1 q_1}\tilde{w}_{p_2 q_2} \int_{C_i} \left\{ \psi_{p_1}^{'}(x)\psi_{q_1}(y)\psi_{p_2}^{'}(x)\psi_{q_2}(y) + \psi_{p_1}(x)\psi_{q_1}(y)\psi_{p_2}(x)\psi_{q_2}(y) + \cdots \right\} d\boldsymbol{\chi}\\
&= \tilde{w}_{p_1 q_1}\tilde{w}_{p_2 q_2}\Sigma_{p_1 p_2 q_1 q_2}
\end{aligned}$$

where

$$\Sigma = \Sigma^{(1)} \otimes \Sigma^{(0)} + \Sigma^{(0)} \otimes \Sigma^{(1)} + \Sigma^{(2)} \otimes \Sigma^{(0)} + \Sigma^{(0)} \otimes \Sigma^{(2)} + \Sigma^{(1)} \otimes \Sigma^{(1)} \tag{2.3}$$

$$\Sigma_{ij}^{(\beta)} = \int_0^1 \psi_i^{(\beta)}\psi_j^{(\beta)} dx \tag{2.4}$$

2.2 Fast WENO Oscillation Indicator Calculation

The WENO oscillation indicator is defined as:

$$o = \Sigma_{mn} w_m w_n \tag{2.5}$$

where w_i are the WENO coefficients calculated for a particular stencil, and:

$$\Sigma_{mn} = \sum_{\alpha=1}^{N} \int_0^1 \psi_m^{(\alpha)} \psi_n^{(\alpha)} \tag{2.6}$$

By considering that:

$$o = \sum_{\alpha=1}^{N} \int_0^1 \left(\frac{d^\alpha w}{dx^\alpha} \right)^2 > 0 \tag{2.7}$$

we have that Σ is symmetric positive definite. Thus, it has a Cholesky decomposition $\Sigma = LL^T$. Thus:

$$o = \left\| w^T L \right\|^2 \tag{2.8}$$

L can be precalculated, and o calculated quickly as:

```
o = 0
for j = 1...n do
    tmp = 0
    for i = j...n do
        tmp = tmp + w(i) * L(i,j)
    end for
    o = o + tmp * tmp
end for
```

2.3 Approximating Interface Terms in FV

Instead of calculating the following:

$$\int D \left(q^- (x_0, t), q^+ (x_0, t) \right) dt \tag{2.9}$$

I propose calculating the following:

$$D \left(\frac{1}{\Delta t} \int q^- (x_0, t) \, dt, \frac{1}{\Delta t} \int q^+ (x_0, t) \, dt \right) \tag{2.10}$$

This obtains a large speedup with no discernable difference in the results of Stokes' First Problem.

2.4 Analytical Results for Basis Vectors

For $N = 1$, the Gauss-Legendre nodes on $[0, 1]$ are $\left\{\frac{1}{2}\left(1-\frac{1}{\sqrt{3}}\right), \frac{1}{2}\left(1+\frac{1}{\sqrt{3}}\right)\right\}$. Thus:

$$\psi_1(x) = -\sqrt{3}x + \frac{1+\sqrt{3}}{2} \tag{2.11a}$$
$$\psi_2(x) = \sqrt{3}x + \frac{1-\sqrt{3}}{2} \tag{2.11b}$$

$$\psi_1(1) = \frac{1-\sqrt{3}}{2} \tag{2.12a}$$
$$\psi_2(1) = \frac{1+\sqrt{3}}{2} \tag{2.12b}$$

$$\psi_1(1)\,\psi_1(1) = 1 - \frac{\sqrt{3}}{2} \tag{2.13a}$$
$$\psi_1(1)\,\psi_2(1) = -\frac{1}{2} \tag{2.13b}$$
$$\psi_2(1)\,\psi_2(1) = 1 + \frac{\sqrt{3}}{2} \tag{2.13c}$$

$$\int_m^{m+1} \psi_1(x)\,dx = \frac{-\sqrt{3}}{2}(2m+1) + \frac{1+\sqrt{3}}{2} = \frac{1}{2} - m\sqrt{3} \tag{2.14a}$$
$$\int_m^{m+1} \psi_2(x)\,dx = \frac{\sqrt{3}}{2}(2m+1) + \frac{1-\sqrt{3}}{2} = \frac{1}{2} + m\sqrt{3} \tag{2.14b}$$

The WENO matrices are:

$$\begin{pmatrix} \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2}-\sqrt{3} & \frac{1}{2}+\sqrt{3} \end{pmatrix}, \begin{pmatrix} \frac{1}{2}+\sqrt{3} & \frac{1}{2}-\sqrt{3} \\ \frac{1}{2} & \frac{1}{2} \end{pmatrix} \tag{2.15}$$

The inverses are:

$$\frac{1}{\sqrt{3}}\begin{pmatrix} \sqrt{3}+\frac{1}{2} & -\frac{1}{2} \\ \sqrt{3}-\frac{1}{2} & \frac{1}{2} \end{pmatrix}, \frac{1}{\sqrt{3}}\begin{pmatrix} \frac{1}{2} & \sqrt{3}-\frac{1}{2} \\ -\frac{1}{2} & \sqrt{3}+\frac{1}{2} \end{pmatrix} \tag{2.16}$$

The weights for both nodes are 0.5 so $\int_0^1 \psi_i\psi_j dx = \frac{\delta_{ij}}{2}$ and $\int_0^1 \psi_i\psi_j' dx = (-1)^j \frac{\sqrt{3}}{2}$.

$$\begin{aligned} I_{11} - I_2^T &= \frac{1}{2}\begin{pmatrix} 2-\sqrt{3} & -1 \\ -1 & 2+\sqrt{3} \end{pmatrix} - \frac{1}{2}\begin{pmatrix} -\sqrt{3} & -\sqrt{3} \\ \sqrt{3} & \sqrt{3} \end{pmatrix} \\ &= \frac{1}{2}\begin{pmatrix} 2 & -\left(1-\sqrt{3}\right) \\ -\left(1+\sqrt{3}\right) & 2 \end{pmatrix} \end{aligned} \tag{2.17}$$

$$\left(I_{11} - I_2^T\right)^{-1} = \frac{1}{3}\begin{pmatrix} 2 & 1-\sqrt{3} \\ 1+\sqrt{3} & 2 \end{pmatrix} \tag{2.18}$$

Using a precalculated, analytical form of U in the DG predictor for $N = 1$ obtains a ~30% speedup on Stokes' First Problem.

2.5 Distortion ODEs

2.5.1 Linearized Distortion ODEs Solver

Note that $A^* = \left(\frac{\rho}{\rho_0}\right)^{\frac{1}{3}} I$ is a stationary point of the ODE for A. Linearizing the ODE around A^* gives:

$$\begin{aligned}
\frac{dA}{dt} &\approx J_A\left(A^*\right)\left(A - A^*\right) \\
&= \frac{-3}{\tau_1}\left(\frac{\rho}{\rho_0}\right)^{\frac{5}{3}} \left(\left(\frac{\rho}{\rho_0}\right)^{\frac{2}{3}} \delta_{in}\delta_{mj} + \left(\frac{\rho}{\rho_0}\right)^{\frac{2}{3}} \delta_{jn}\delta_{im} + \left(\frac{\rho}{\rho_0}\right)^{\frac{2}{3}} \delta_{im}\delta_{jn} - \frac{1}{3}\left(\frac{\rho}{\rho_0}\right)^{\frac{2}{3}} \delta_{im}\delta_{jn}\delta_{kl}\delta_{kl} - \frac{2}{3}\left(\frac{\rho}{\rho_0}\right)^{\frac{2}{3}} \delta_{ij}\delta_{mn}\right) \\
&\quad \times \left(A_{mn} - \left(\frac{\rho}{\rho_0}\right)^{\frac{1}{3}} \delta_{mn}\right) \\
&= \frac{-3}{\tau_1}\left(\frac{\rho}{\rho_0}\right)^{\frac{7}{3}} \left(\delta_{in}\delta_{mj} + \delta_{im}\delta_{jn} - \frac{2}{3}\delta_{ij}\delta_{mn}\right)\left(A_{mn} - \left(\frac{\rho}{\rho_0}\right)^{\frac{1}{3}} \delta_{mn}\right) \\
&= \frac{-3}{\tau_1}\left(\frac{\rho}{\rho_0}\right)^{\frac{7}{3}} \left(A_{mn}\left(\delta_{in}\delta_{mj} + \delta_{im}\delta_{jn} - \frac{2}{3}\delta_{ij}\delta_{mn}\right) - \left(\frac{\rho}{\rho_0}\right)^{\frac{1}{3}} \delta_{mn}\left(\delta_{in}\delta_{mj} + \delta_{im}\delta_{jn} - \frac{2}{3}\delta_{ij}\delta_{mn}\right)\right) \\
&= \frac{-3}{\tau_1}\left(\frac{\rho}{\rho_0}\right)^{\frac{7}{3}} \left(\left(A_{ji} + A_{ij} - \frac{2}{3}\operatorname{tr}\left(A\right)\delta_{ij}\right) - \left(\frac{\rho}{\rho_0}\right)^{\frac{1}{3}} \left(\delta_{ij} + \delta_{ij} - 2\delta_{ij}\right)\right) \\
&= \frac{-3}{\tau_1}\left(\frac{\rho}{\rho_0}\right)^{\frac{7}{3}} \left(A + A^T - \frac{2}{3}\operatorname{tr}\left(A\right) I\right)
\end{aligned} \tag{2.19}$$

The matrix for this system, in row-major form, is:

$$\frac{-3}{\tau_1}\left(\frac{\rho}{\rho_0}\right)^{\frac{7}{3}} \begin{pmatrix}
\frac{4}{3} & 0 & 0 & 0 & -\frac{2}{3} & 0 & 0 & 0 & -\frac{2}{3} \\
0 & 1 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 1 & 0 & 0 & 0 & 1 & 0 & 0 \\
0 & 1 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\
-\frac{2}{3} & 0 & 0 & 0 & \frac{4}{3} & 0 & 0 & 0 & -\frac{2}{3} \\
0 & 0 & 0 & 0 & 0 & 1 & 0 & 1 & 0 \\
0 & 0 & 1 & 0 & 0 & 0 & 1 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 1 & 0 & 1 & 0 \\
-\frac{2}{3} & 0 & 0 & 0 & -\frac{2}{3} & 0 & 0 & 0 & \frac{4}{3}
\end{pmatrix} \tag{2.20}$$

The eigenvalues and eigenvectors are:

$$\{0, 0, 0, 0, -2k, -2k, -2k, -2k, -2k\} \tag{2.21}$$

$$\begin{pmatrix}0\\1\\0\\-1\\0\\0\\0\\0\\0\end{pmatrix}\begin{pmatrix}0\\0\\1\\0\\0\\0\\-1\\0\\0\end{pmatrix}\begin{pmatrix}0\\0\\0\\0\\0\\1\\0\\-1\\0\end{pmatrix}\begin{pmatrix}1\\0\\0\\0\\1\\0\\0\\0\\1\end{pmatrix}\begin{pmatrix}-1\\0\\0\\0\\1\\0\\0\\0\\0\end{pmatrix}\begin{pmatrix}-1\\0\\0\\0\\0\\0\\0\\0\\1\end{pmatrix}\begin{pmatrix}0\\1\\0\\1\\0\\0\\0\\0\\0\end{pmatrix}\begin{pmatrix}0\\0\\1\\0\\0\\0\\1\\0\\0\end{pmatrix}\begin{pmatrix}0\\0\\0\\0\\0\\1\\0\\1\\0\end{pmatrix} \tag{2.22}$$

where $k = \frac{3}{\tau_1}\left(\frac{\rho}{\rho_0}\right)^{\frac{7}{3}}$. Thus, the solution is:

$$\begin{aligned}
&\frac{A_{12}-A_{21}}{2}\begin{pmatrix}0&1&0\\-1&0&0\\0&0&0\end{pmatrix}+\frac{A_{13}-A_{31}}{2}\begin{pmatrix}0&0&1\\0&0&0\\-1&0&0\end{pmatrix}+\frac{A_{23}-A_{32}}{2}\begin{pmatrix}0&0&0\\0&0&1\\0&-1&0\end{pmatrix} \\
&+\frac{A_{11}+A_{22}+A_{33}}{3}\begin{pmatrix}1&0&0\\0&1&0\\0&0&1\end{pmatrix} \\
&+\frac{2A_{22}-A_{11}-A_{33}}{3}e^{-2kt}\begin{pmatrix}-1&0&0\\0&1&0\\0&0&0\end{pmatrix}+\frac{2A_{33}-A_{11}-A_{22}}{3}e^{-2kt}\begin{pmatrix}-1&0&0\\0&0&0\\0&0&1\end{pmatrix} \\
&+\frac{A_{12}+A_{21}}{2}e^{-2kt}\begin{pmatrix}0&1&0\\1&0&0\\0&0&0\end{pmatrix}+\frac{A_{13}+A_{31}}{2}e^{-2kt}\begin{pmatrix}0&0&1\\0&0&0\\1&0&0\end{pmatrix}+\frac{A_{23}+A_{32}}{2}e^{-2kt}\begin{pmatrix}0&0&0\\0&0&1\\0&1&0\end{pmatrix}
\end{aligned} \tag{2.23}$$

This is equal to:

$$\frac{1}{2}\left(A-A^T\right)+\frac{\operatorname{tr}(A)}{3}I+e^{-2kt}\left(\frac{1}{2}\left(A+A^T\right)-\frac{\operatorname{tr}(A)}{3}I\right) \tag{2.24}$$

Results with Stokes' First Problem look good with this linearisation. The ODE step takes a negligible amount of time, meaning that if accuracy is maintained to second order, the solver is now fast enough.

2.5.2 Linearized Reduced Distortion ODE Solver

Taking system (**??**), note that the Jacobian of the system is given by:

$$J = -k\begin{pmatrix}4x_1-x_2-x_3 & -x_1 & -x_1\\ -x_2 & 4x_2-x_3-x_1 & -x_3\\ -x_3 & -x_3 & 4x_3-x_1-x_2\end{pmatrix} \tag{2.25}$$

Evaluated at stationary point $x_i = \sqrt[3]{c}$ we have:

$$J(\boldsymbol{x_0}) = -k\sqrt[3]{c}\begin{pmatrix}2&-1&-1\\-1&2&-1\\-1&-1&2\end{pmatrix} \tag{2.26}$$

Thus, the system is linearized to:

$$\begin{aligned}\frac{d\boldsymbol{x}}{dt} &\approx -k\sqrt[3]{c}\begin{pmatrix}2 & -1 & -1\\ -1 & 2 & -1\\ -1 & -1 & 2\end{pmatrix}\left(\begin{pmatrix}x_1\\ x_2\\ x_3\end{pmatrix}-\sqrt[3]{c}\begin{pmatrix}1\\ 1\\ 1\end{pmatrix}\right)\\ &= k\sqrt[3]{c}\begin{pmatrix}-2 & 1 & 1\\ 1 & -2 & 1\\ 1 & 1 & -2\end{pmatrix}\begin{pmatrix}x_1\\ x_2\\ x_3\end{pmatrix}\end{aligned} \tag{2.27}$$

The eigenvalues of this system matrix are $\{-3k\sqrt[3]{c}, -3k\sqrt[3]{c}, 0\}$ and the eigenvectors are:

$$\begin{pmatrix}-1\\ 0\\ 1\end{pmatrix},\begin{pmatrix}0\\ -1\\ 1\end{pmatrix},\begin{pmatrix}1\\ 1\\ 1\end{pmatrix} \tag{2.28}$$

Thus, the linearized solution is:

$$\boldsymbol{x}(t)=\frac{-2x_1+x_2+x_3}{3}e^{-3k\sqrt[3]{c}t}\begin{pmatrix}-1\\ 0\\ 1\end{pmatrix}+\frac{x_1-2x_2+x_3}{3}e^{-3k\sqrt[3]{c}t}\begin{pmatrix}0\\ -1\\ 1\end{pmatrix}+\frac{x_1+x_2+x_3}{3}\begin{pmatrix}1\\ 1\\ 1\end{pmatrix} \tag{2.29}$$

This may represent a faster way to calculate the evolution of the stretch terms of A. Note that some kind of normalization will probably be necessary, as:

$$\frac{x_1+x_2+x_3}{3}\geq(x_1x_2x_3)^{\frac{1}{3}} \tag{2.30}$$

with equality if and only if $x_1 = x_2 = x_3$.

2.6 Primitive WENO and DG Reconstruction

As suggested in [2], the WENO and DG can be performed in primitive variables, which is less computationally expensive than evaluating fluxes using conserved variables. Achieves around 20% speedup in DG step, at double cost in WENO step. Minimal speedup in FV step, as both primitive and conserved variables must be calculated for the flux updates. Not enough.

2.7 Change to Row-Major Ordering

The original GPR papers state the equations for A in column-major order, probably because the authors use Fortran. For C++ and Python implementations it is faster to work in row-major order. ~10% speedup was achieved by implementing this.

The GPR equations are:

$$\frac{\partial\rho}{\partial t}+\frac{\partial\left(\rho v_k\right)}{\partial x_k}=0 \tag{2.31a}$$

$$\frac{\partial\left(\rho E\right)}{\partial t}+\frac{\partial\left(\rho E v_k+\left(p\delta_{ik}-\sigma_{ik}\right)v_i+q_k\right)}{\partial x_k}=0 \tag{2.31b}$$

$$\frac{\partial\left(\rho v_i\right)}{\partial t}+\frac{\partial(\rho v_i v_k+p\delta_{ik}-\sigma_{ik})}{\partial x_k}=0 \tag{2.31c}$$

$$\frac{\partial A_{ij}}{\partial t}+\frac{\partial\left(A_{ik}v_k\right)}{\partial x_j}+v_k\left(\frac{\partial A_{ij}}{\partial x_k}-\frac{\partial A_{ik}}{\partial x_j}\right)=-\frac{\psi_{ij}}{\theta_1(\tau_1)} \tag{2.31d}$$

$$\frac{\partial\left(\rho J_i\right)}{\partial t}+\frac{\partial\left(\rho J_i v_k+T\delta_{ik}\right)}{\partial x_k}=-\frac{\rho H_i}{\theta_2\left(\tau_2\right)} \tag{2.31e}$$

Under row-major ordering, we have:

$$\boldsymbol{Q}=\begin{pmatrix}\rho & \rho E & \rho v_1 & \rho v_2 & \rho v_3 & A_{11} & A_{12} & A_{13} & A_{21} & A_{22} & A_{23} & A_{31} & A_{32} & A_{33} & \rho J_1 & \rho J_2 & \rho J_3\end{pmatrix}^T \tag{2.32a}$$

$$\boldsymbol{P}=\begin{pmatrix}\rho & p & v_1 & v_2 & v_3 & A_{11} & A_{21} & A_{31} & A_{12} & A_{22} & A_{32} & A_{13} & A_{23} & A_{33} & J_1 & J_2 & J_3\end{pmatrix}^T \tag{2.32b}$$

$$\boldsymbol{F_1}=\begin{pmatrix}\rho v_1\\ \rho v_1 E+v_1 p-\sigma_{1m}v_m+q_1\\ \rho v_1^2+p-\sigma_{11}\\ \rho v_1 v_2-\sigma_{12}\\ \rho v_1 v_3-\sigma_{13}\\ A_{1m}v_m\\ 0\\ 0\\ A_{2m}v_m\\ 0\\ 0\\ A_{3m}v_m\\ 0\\ 0\\ \rho J_1 v_1+T\\ \rho J_2 v_1\\ \rho J_3 v_1\end{pmatrix}\quad \boldsymbol{F_2}=\begin{pmatrix}\rho v_2\\ \rho v_2 E+v_2 p-\sigma_{2m}v_m+q_2\\ \rho v_1 v_2-\sigma_{21}\\ \rho v_2^2+p-\sigma_{22}\\ \rho v_2 v_3-\sigma_{23}\\ 0\\ A_{1m}v_m\\ 0\\ 0\\ A_{2m}v_m\\ 0\\ 0\\ A_{3m}v_m\\ 0\\ \rho J_1 v_2\\ \rho J_2 v_2+T\\ \rho J_3 v_2\end{pmatrix}\quad \boldsymbol{F_3}=\begin{pmatrix}\rho v_3\\ \rho v_3 E+v_3 p-\sigma_{3m}v_m+q_3\\ \rho v_1 v_3-\sigma_{31}\\ \rho v_2 v_3-\sigma_{32}\\ \rho v_3^2+p-\sigma_{33}\\ 0\\ 0\\ A_{1m}v_m\\ 0\\ 0\\ A_{2m}v_m\\ 0\\ 0\\ A_{3m}v_m\\ \rho J_1 v_3\\ \rho J_2 v_3\\ \rho J_3 v_3+T\end{pmatrix} \tag{2.33}$$

$$B_1 = \begin{pmatrix}
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&-v_2&-v_3&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&v_1&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&v_1&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&-v_2&-v_3&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&v_1&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&v_1&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&-v_2&-v_3&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&v_1&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&v_1&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0
\end{pmatrix} \tag{2.34}$$

$$B_2 = \begin{pmatrix}
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&v_2&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&-v_1&0&-v_3&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&v_2&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&v_2&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&-v_1&0&-v_3&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&v_2&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&v_2&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&-v_1&0&-v_3&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&v_2&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0
\end{pmatrix} \tag{2.35}$$

$$B_3 = \begin{pmatrix}
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&v_3&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&v_3&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&-v_1&-v_2&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&v_3&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&v_3&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&-v_1&-v_2&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&v_3&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&v_3&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&-v_1&-v_2&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0\\
0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0&0
\end{pmatrix} \tag{2.36}$$

$$\boldsymbol{S}=-\frac{1}{\theta_1\left(\tau_1\right)}\begin{pmatrix}0\\0\\0\\0\\0\\\psi_{11}\\\psi_{12}\\\psi_{13}\\\psi_{21}\\\psi_{22}\\\psi_{23}\\\psi_{31}\\\psi_{32}\\\psi_{33}\\0\\0\\0\end{pmatrix}-\frac{1}{\theta_2\left(\tau_2\right)}\begin{pmatrix}0\\0\\0\\0\\0\\0\\0\\0\\0\\0\\0\\0\\0\\0\\\rho H_1\\\rho H_2\\\rho H_3\end{pmatrix} \tag{2.37}$$

$$\Psi_{ij}=\rho v_i v_j-\sigma_{ij} \tag{2.38a}$$

$$\Phi_{ij}^k=\rho v_k \psi_{ij}-v_m\frac{\partial\sigma_{mk}}{\partial A_{ij}} \tag{2.38b}$$

$$\Omega_i=v_i\left(E+\rho E_\rho\right)-\frac{\sigma_{im}v_m}{\rho}+T_\rho H_i \tag{2.38c}$$

$$\Upsilon=\frac{\left\|\boldsymbol{v}\right\|^2+\boldsymbol{H}\cdot\boldsymbol{J}-E-\rho E_\rho}{\rho E_p} \tag{2.38d}$$

$$\tilde{\boldsymbol{H}}=E_{\boldsymbol{JJ}} \tag{2.38e}$$

$$\boldsymbol{S_p}=\frac{1}{\theta_1\left(\tau_1\right)}\begin{pmatrix}0\\\left(\gamma-1\right)\rho\left\|\psi\right\|_F^2\\0\\0\\0\\-\psi_{11}\\-\psi_{12}\\-\psi_{13}\\-\psi_{21}\\-\psi_{22}\\-\psi_{23}\\-\psi_{31}\\-\psi_{32}\\-\psi_{33}\\0\\0\\0\end{pmatrix}+\frac{1}{\theta_2\left(\tau_2\right)}\begin{pmatrix}0\\\left(\gamma-1\right)\rho\left\|\boldsymbol{H}\right\|^2\\0\\0\\0\\0\\0\\0\\0\\0\\0\\0\\0\\0\\-H_1\\-H_2\\-H_3\end{pmatrix} \tag{2.47}$$

$$
\frac{\partial \boldsymbol{Q}}{\partial \boldsymbol{P}} =
\begin{pmatrix}
1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
E+\rho E_\rho & \rho E_p & \rho v_1 & \rho v_2 & \rho v_3 & \rho\psi_{11} & \rho\psi_{12} & \rho\psi_{13} & \rho\psi_{21} & \rho\psi_{22} & \rho\psi_{23} & \rho\psi_{31} & \rho\psi_{32} & \rho\psi_{33} & \rho H_1 & \rho H_2 & \rho H_3 \\
v_1 & 0 & \rho & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
v_2 & 0 & 0 & \rho & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
v_3 & 0 & 0 & 0 & \rho & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\
J_1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \rho & 0 & 0 \\
J_2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \rho & 0 \\
J_3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \rho
\end{pmatrix}
\tag{2.39}
$$

$$
\frac{\partial \boldsymbol{P}}{\partial \boldsymbol{Q}} =
\begin{pmatrix}
1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
\Upsilon & \frac{1}{\rho E_p} & -\frac{v_1}{\rho E_p} & -\frac{v_2}{\rho E_p} & -\frac{v_3}{\rho E_p} & -\frac{\psi_{11}}{E_p} & -\frac{\psi_{12}}{E_p} & -\frac{\psi_{13}}{E_p} & -\frac{\psi_{21}}{E_p} & -\frac{\psi_{22}}{E_p} & -\frac{\psi_{23}}{E_p} & -\frac{\psi_{31}}{E_p} & -\frac{\psi_{32}}{E_p} & -\frac{\psi_{33}}{E_p} & -\frac{H_1}{\rho E_p} & -\frac{H_2}{\rho E_p} & -\frac{H_3}{\rho E_p} \\
-\frac{v_1}{\rho} & 0 & \frac{1}{\rho} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
-\frac{v_2}{\rho} & 0 & 0 & \frac{1}{\rho} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
-\frac{v_3}{\rho} & 0 & 0 & 0 & \frac{1}{\rho} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\
-\frac{J_1}{\rho} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{1}{\rho} & 0 & 0 \\
-\frac{J_2}{\rho} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{1}{\rho} & 0 \\
-\frac{J_3}{\rho} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{1}{\rho}
\end{pmatrix}
\tag{2.40}
$$

$$\frac{\partial \boldsymbol{F_1}}{\partial \boldsymbol{P}} = \begin{pmatrix}
v_1 & 0 & \rho & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
\Omega_1 & \left(v_1\left(\rho E_p+1\right)+T_p H_1\right) & \left(\Psi_{11}+\rho E+p\right) & \Psi_{12} & \Psi_{13} & \Phi^1_{11} & \Phi^1_{12} & \Phi^1_{13} & \Phi^1_{21} & \Phi^1_{22} & \Phi^1_{23} & \Phi^1_{31} & \Phi^1_{32} & \Phi^1_{33} & \left(\rho v_1 H_1 + T\tilde{H}_1\right) & \rho v_1 H_2 & \rho v_1 H_3 \\
\frac{\Psi_{11}}{\rho} & 1 & 2\rho v_1 & 0 & 0 & -\frac{\partial\sigma_{11}}{A_{11}} & -\frac{\partial\sigma_{11}}{A_{12}} & -\frac{\partial\sigma_{11}}{A_{13}} & -\frac{\partial\sigma_{11}}{A_{21}} & -\frac{\partial\sigma_{11}}{A_{22}} & -\frac{\partial\sigma_{11}}{A_{23}} & -\frac{\partial\sigma_{11}}{A_{31}} & -\frac{\partial\sigma_{11}}{A_{32}} & -\frac{\partial\sigma_{11}}{A_{33}} & 0 & 0 & 0 \\
\frac{\Psi_{12}}{\rho} & 0 & \rho v_2 & \rho v_1 & 0 & -\frac{\partial\sigma_{12}}{A_{11}} & -\frac{\partial\sigma_{12}}{A_{12}} & -\frac{\partial\sigma_{12}}{A_{13}} & -\frac{\partial\sigma_{12}}{A_{21}} & -\frac{\partial\sigma_{12}}{A_{22}} & -\frac{\partial\sigma_{12}}{A_{23}} & -\frac{\partial\sigma_{12}}{A_{31}} & -\frac{\partial\sigma_{12}}{A_{32}} & -\frac{\partial\sigma_{12}}{A_{33}} & 0 & 0 & 0 \\
\frac{\Psi_{13}}{\rho} & 0 & \rho v_3 & 0 & \rho v_1 & -\frac{\partial\sigma_{13}}{A_{11}} & -\frac{\partial\sigma_{13}}{A_{12}} & -\frac{\partial\sigma_{13}}{A_{13}} & -\frac{\partial\sigma_{13}}{A_{21}} & -\frac{\partial\sigma_{13}}{A_{22}} & -\frac{\partial\sigma_{13}}{A_{23}} & -\frac{\partial\sigma_{13}}{A_{31}} & -\frac{\partial\sigma_{13}}{A_{32}} & -\frac{\partial\sigma_{13}}{A_{33}} & 0 & 0 & 0 \\
0 & 0 & A_{11} & A_{12} & A_{13} & v_1 & v_2 & v_3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{21} & A_{22} & A_{23} & 0 & 0 & 0 & v_1 & v_2 & v_3 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{31} & A_{32} & A_{33} & 0 & 0 & 0 & 0 & 0 & 0 & v_1 & v_2 & v_3 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
v_1 J_1 + T_\rho & T_p & \rho J_1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \rho v_1 & 0 & 0 \\
v_1 J_2 & 0 & \rho J_2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \rho v_1 & 0 \\
v_1 J_3 & 0 & \rho J_3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \rho v_1
\end{pmatrix} \tag{2.41}$$

$$
\frac{\partial \boldsymbol{F_2}}{\partial \boldsymbol{P}} = \begin{pmatrix}
v_2 & 0 & 0 & \rho & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
\Omega_2 & \left(v_2\left(\rho E_p + 1\right) + T_p H_2\right) & \Psi_{21} & \left(\Psi_{22} + \rho E + p\right) & \Psi_{23} & \Phi^2_{11} & \Phi^2_{12} & \Phi^2_{13} & \Phi^2_{21} & \Phi^2_{22} & \Phi^2_{23} & \Phi^2_{31} & \Phi^2_{32} & \Phi^2_{33} & \rho v_2 H_1 & \left(\rho v_2 H_2 + T\tilde{H}_2\right) & \rho v_2 H_3 \\
\frac{\Psi_{21}}{\rho} & 0 & \rho v_2 & \rho v_1 & 0 & -\frac{\partial\sigma_{21}}{A_{11}} & -\frac{\partial\sigma_{21}}{A_{12}} & -\frac{\partial\sigma_{21}}{A_{13}} & -\frac{\partial\sigma_{21}}{A_{21}} & -\frac{\partial\sigma_{21}}{A_{22}} & -\frac{\partial\sigma_{21}}{A_{23}} & -\frac{\partial\sigma_{21}}{A_{31}} & -\frac{\partial\sigma_{21}}{A_{32}} & -\frac{\partial\sigma_{21}}{A_{33}} & 0 & 0 & 0 \\
\frac{\Psi_{22}}{\rho} & 1 & 0 & 2\rho v_2 & 0 & -\frac{\partial\sigma_{22}}{A_{11}} & -\frac{\partial\sigma_{22}}{A_{12}} & -\frac{\partial\sigma_{22}}{A_{13}} & -\frac{\partial\sigma_{22}}{A_{21}} & -\frac{\partial\sigma_{22}}{A_{22}} & -\frac{\partial\sigma_{22}}{A_{23}} & -\frac{\partial\sigma_{22}}{A_{31}} & -\frac{\partial\sigma_{22}}{A_{32}} & -\frac{\partial\sigma_{22}}{A_{33}} & 0 & 0 & 0 \\
\frac{\Psi_{23}}{\rho} & 0 & 0 & \rho v_3 & \rho v_2 & -\frac{\partial\sigma_{23}}{A_{11}} & -\frac{\partial\sigma_{23}}{A_{12}} & -\frac{\partial\sigma_{23}}{A_{13}} & -\frac{\partial\sigma_{23}}{A_{21}} & -\frac{\partial\sigma_{23}}{A_{22}} & -\frac{\partial\sigma_{23}}{A_{23}} & -\frac{\partial\sigma_{23}}{A_{31}} & -\frac{\partial\sigma_{23}}{A_{32}} & -\frac{\partial\sigma_{23}}{A_{33}} & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{11} & A_{12} & A_{13} & v_1 & v_2 & v_3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{21} & A_{22} & A_{23} & 0 & 0 & 0 & v_1 & v_2 & v_3 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{31} & A_{32} & A_{33} & 0 & 0 & 0 & 0 & 0 & 0 & v_1 & v_2 & v_3 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
v_2 J_1 & 0 & 0 & \rho J_1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \rho v_2 & 0 & 0 \\
v_2 J_2 + T_\rho & T_p & 0 & \rho J_2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \rho v_2 & 0 \\
v_2 J_3 & 0 & 0 & \rho J_3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \rho v_2
\end{pmatrix} \tag{2.42}
$$

$$
\frac{\partial \boldsymbol{F_3}}{\partial \boldsymbol{P}} = \begin{pmatrix}
v_3 & 0 & \rho & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
\Omega_3 & (v_3(\rho E_p + 1) + T_p H_3) & \Psi_{31} & \Psi_{32} & (\Psi_{33} + \rho E + p) & \Phi^3_{11} & \Phi^3_{12} & \Phi^3_{13} & \Phi^3_{21} & \Phi^3_{22} & \Phi^3_{23} & \Phi^3_{31} & \Phi^3_{32} & \Phi^3_{33} & \rho v_3 H_1 & \rho v_3 H_2 & \left(\rho v_3 H_3 + T\tilde{H}_3\right) \\
\frac{\Psi_{31}}{\rho} & 0 & \rho v_3 & 0 & \rho v_1 & -\frac{\partial\sigma_{31}}{A_{11}} & -\frac{\partial\sigma_{31}}{A_{12}} & -\frac{\partial\sigma_{31}}{A_{13}} & -\frac{\partial\sigma_{31}}{A_{21}} & -\frac{\partial\sigma_{31}}{A_{22}} & -\frac{\partial\sigma_{31}}{A_{23}} & -\frac{\partial\sigma_{31}}{A_{31}} & -\frac{\partial\sigma_{31}}{A_{32}} & -\frac{\partial\sigma_{31}}{A_{33}} & 0 & 0 & 0 \\
\frac{\Psi_{32}}{\rho} & 0 & 0 & \rho v_3 & \rho v_2 & -\frac{\partial\sigma_{32}}{A_{11}} & -\frac{\partial\sigma_{32}}{A_{12}} & -\frac{\partial\sigma_{32}}{A_{13}} & -\frac{\partial\sigma_{32}}{A_{21}} & -\frac{\partial\sigma_{32}}{A_{22}} & -\frac{\partial\sigma_{32}}{A_{23}} & -\frac{\partial\sigma_{32}}{A_{31}} & -\frac{\partial\sigma_{32}}{A_{32}} & -\frac{\partial\sigma_{32}}{A_{33}} & 0 & 0 & 0 \\
\frac{\Psi_{33}}{\rho} & 1 & 0 & 0 & 2\rho v_3 & -\frac{\partial\sigma_{33}}{A_{11}} & -\frac{\partial\sigma_{33}}{A_{12}} & -\frac{\partial\sigma_{33}}{A_{13}} & -\frac{\partial\sigma_{33}}{A_{21}} & -\frac{\partial\sigma_{33}}{A_{22}} & -\frac{\partial\sigma_{33}}{A_{23}} & -\frac{\partial\sigma_{33}}{A_{31}} & -\frac{\partial\sigma_{33}}{A_{32}} & -\frac{\partial\sigma_{33}}{A_{33}} & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{11} & A_{12} & A_{13} & v_1 & v_2 & v_3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{21} & A_{22} & A_{23} & 0 & 0 & 0 & v_1 & v_2 & v_3 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{31} & A_{32} & A_{33} & 0 & 0 & 0 & 0 & 0 & 0 & v_1 & v_2 & v_3 & 0 & 0 & 0 \\
v_3 J_1 & 0 & 0 & 0 & \rho J_1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \rho v_3 & 0 & 0 \\
v_3 J_2 & 0 & 0 & 0 & \rho J_2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \rho v_3 & 0 \\
v_3 J_3 + T_\rho & T_p & 0 & 0 & \rho J_3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \rho v_3
\end{pmatrix} \tag{2.43}
$$

$$M_1 = \begin{pmatrix}
v_1 & 0 & \rho & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & v_1 & \gamma p & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \Gamma\alpha^2 T & 0 & 0 \\
-\frac{\sigma_{11}}{\rho^2} & \frac{1}{\rho} & v_1 & 0 & 0 & -\frac{1}{\rho}\frac{\partial\sigma_{11}}{\partial A_{11}} & -\frac{1}{\rho}\frac{\partial\sigma_{11}}{\partial A_{12}} & -\frac{1}{\rho}\frac{\partial\sigma_{11}}{\partial A_{13}} & -\frac{1}{\rho}\frac{\partial\sigma_{11}}{\partial A_{21}} & -\frac{1}{\rho}\frac{\partial\sigma_{11}}{\partial A_{22}} & -\frac{1}{\rho}\frac{\partial\sigma_{11}}{\partial A_{23}} & -\frac{1}{\rho}\frac{\partial\sigma_{11}}{\partial A_{31}} & -\frac{1}{\rho}\frac{\partial\sigma_{11}}{\partial A_{32}} & -\frac{1}{\rho}\frac{\partial\sigma_{11}}{\partial A_{33}} & 0 & 0 & 0 \\
-\frac{\sigma_{12}}{\rho^2} & 0 & 0 & v_1 & 0 & -\frac{1}{\rho}\frac{\partial\sigma_{12}}{\partial A_{11}} & -\frac{1}{\rho}\frac{\partial\sigma_{12}}{\partial A_{12}} & -\frac{1}{\rho}\frac{\partial\sigma_{12}}{\partial A_{13}} & -\frac{1}{\rho}\frac{\partial\sigma_{12}}{\partial A_{21}} & -\frac{1}{\rho}\frac{\partial\sigma_{12}}{\partial A_{22}} & -\frac{1}{\rho}\frac{\partial\sigma_{12}}{\partial A_{23}} & -\frac{1}{\rho}\frac{\partial\sigma_{12}}{\partial A_{31}} & -\frac{1}{\rho}\frac{\partial\sigma_{12}}{\partial A_{32}} & -\frac{1}{\rho}\frac{\partial\sigma_{12}}{\partial A_{33}} & 0 & 0 & 0 \\
-\frac{\sigma_{13}}{\rho^2} & 0 & 0 & 0 & v_1 & -\frac{1}{\rho}\frac{\partial\sigma_{13}}{\partial A_{11}} & -\frac{1}{\rho}\frac{\partial\sigma_{13}}{\partial A_{12}} & -\frac{1}{\rho}\frac{\partial\sigma_{13}}{\partial A_{13}} & -\frac{1}{\rho}\frac{\partial\sigma_{13}}{\partial A_{21}} & -\frac{1}{\rho}\frac{\partial\sigma_{13}}{\partial A_{22}} & -\frac{1}{\rho}\frac{\partial\sigma_{13}}{\partial A_{23}} & -\frac{1}{\rho}\frac{\partial\sigma_{13}}{\partial A_{31}} & -\frac{1}{\rho}\frac{\partial\sigma_{13}}{\partial A_{32}} & -\frac{1}{\rho}\frac{\partial\sigma_{13}}{\partial A_{33}} & 0 & 0 & 0 \\
0 & 0 & A_{11} & A_{12} & A_{13} & v_1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & v_1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & v_1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{21} & A_{22} & A_{23} & 0 & 0 & 0 & v_1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_1 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{31} & A_{32} & A_{33} & 0 & 0 & 0 & 0 & 0 & 0 & v_1 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_1 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_1 & 0 & 0 & 0 \\
-\frac{T}{\rho^2} & \frac{T}{\rho(p+p_\infty)} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_1 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_1 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_1
\end{pmatrix} \tag{2.44}$$

$$M_2 = \begin{pmatrix}
v_2 & 0 & 0 & \rho & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & v_2 & 0 & \gamma p & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \Gamma\alpha^2 T & 0 \\
-\frac{\sigma_{21}}{\rho^2} & 0 & v_2 & 0 & 0 & -\frac{1}{\rho}\frac{\partial\sigma_{21}}{\partial A_{11}} & -\frac{1}{\rho}\frac{\partial\sigma_{21}}{\partial A_{12}} & -\frac{1}{\rho}\frac{\partial\sigma_{21}}{\partial A_{13}} & -\frac{1}{\rho}\frac{\partial\sigma_{21}}{\partial A_{21}} & -\frac{1}{\rho}\frac{\partial\sigma_{21}}{\partial A_{22}} & -\frac{1}{\rho}\frac{\partial\sigma_{21}}{\partial A_{23}} & -\frac{1}{\rho}\frac{\partial\sigma_{21}}{\partial A_{31}} & -\frac{1}{\rho}\frac{\partial\sigma_{21}}{\partial A_{32}} & -\frac{1}{\rho}\frac{\partial\sigma_{21}}{\partial A_{33}} & 0 & 0 & 0 \\
-\frac{\sigma_{22}}{\rho^2} & \frac{1}{\rho} & 0 & v_2 & 0 & -\frac{1}{\rho}\frac{\partial\sigma_{22}}{\partial A_{11}} & -\frac{1}{\rho}\frac{\partial\sigma_{22}}{\partial A_{12}} & -\frac{1}{\rho}\frac{\partial\sigma_{22}}{\partial A_{13}} & -\frac{1}{\rho}\frac{\partial\sigma_{22}}{\partial A_{21}} & -\frac{1}{\rho}\frac{\partial\sigma_{22}}{\partial A_{22}} & -\frac{1}{\rho}\frac{\partial\sigma_{22}}{\partial A_{23}} & -\frac{1}{\rho}\frac{\partial\sigma_{22}}{\partial A_{31}} & -\frac{1}{\rho}\frac{\partial\sigma_{22}}{\partial A_{32}} & -\frac{1}{\rho}\frac{\partial\sigma_{22}}{\partial A_{33}} & 0 & 0 & 0 \\
-\frac{\sigma_{23}}{\rho^2} & 0 & 0 & 0 & v_2 & -\frac{1}{\rho}\frac{\partial\sigma_{23}}{\partial A_{11}} & -\frac{1}{\rho}\frac{\partial\sigma_{23}}{\partial A_{12}} & -\frac{1}{\rho}\frac{\partial\sigma_{23}}{\partial A_{13}} & -\frac{1}{\rho}\frac{\partial\sigma_{23}}{\partial A_{21}} & -\frac{1}{\rho}\frac{\partial\sigma_{23}}{\partial A_{22}} & -\frac{1}{\rho}\frac{\partial\sigma_{23}}{\partial A_{23}} & -\frac{1}{\rho}\frac{\partial\sigma_{23}}{\partial A_{31}} & -\frac{1}{\rho}\frac{\partial\sigma_{23}}{\partial A_{32}} & -\frac{1}{\rho}\frac{\partial\sigma_{23}}{\partial A_{33}} & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & v_2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{11} & A_{12} & A_{13} & 0 & v_2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & v_2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{21} & A_{22} & A_{23} & 0 & 0 & 0 & 0 & v_2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_2 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_2 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{31} & A_{32} & A_{33} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_2 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_2 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_2 & 0 & 0 \\
-\frac{T}{\rho^2} & \frac{T}{\rho(p+p_\infty)} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_2 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_2
\end{pmatrix} \tag{2.45}$$

$$M_3 = \begin{pmatrix}
v_3 & 0 & 0 & 0 & \rho & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & v_3 & 0 & 0 & \gamma p & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \Gamma\alpha^2 T \\
-\frac{\sigma_{31}}{\rho^2} & 0 & v_3 & 0 & 0 & -\frac{1}{\rho}\frac{\partial\sigma_{31}}{\partial A_{11}} & -\frac{1}{\rho}\frac{\partial\sigma_{31}}{\partial A_{12}} & -\frac{1}{\rho}\frac{\partial\sigma_{31}}{\partial A_{13}} & -\frac{1}{\rho}\frac{\partial\sigma_{31}}{\partial A_{21}} & -\frac{1}{\rho}\frac{\partial\sigma_{31}}{\partial A_{22}} & -\frac{1}{\rho}\frac{\partial\sigma_{31}}{\partial A_{23}} & -\frac{1}{\rho}\frac{\partial\sigma_{31}}{\partial A_{31}} & -\frac{1}{\rho}\frac{\partial\sigma_{31}}{\partial A_{32}} & -\frac{1}{\rho}\frac{\partial\sigma_{31}}{\partial A_{33}} & 0 & 0 & 0 \\
-\frac{\sigma_{32}}{\rho^2} & 0 & 0 & v_3 & 0 & -\frac{1}{\rho}\frac{\partial\sigma_{32}}{\partial A_{11}} & -\frac{1}{\rho}\frac{\partial\sigma_{32}}{\partial A_{12}} & -\frac{1}{\rho}\frac{\partial\sigma_{32}}{\partial A_{13}} & -\frac{1}{\rho}\frac{\partial\sigma_{32}}{\partial A_{21}} & -\frac{1}{\rho}\frac{\partial\sigma_{32}}{\partial A_{22}} & -\frac{1}{\rho}\frac{\partial\sigma_{32}}{\partial A_{23}} & -\frac{1}{\rho}\frac{\partial\sigma_{32}}{\partial A_{31}} & -\frac{1}{\rho}\frac{\partial\sigma_{32}}{\partial A_{32}} & -\frac{1}{\rho}\frac{\partial\sigma_{32}}{\partial A_{33}} & 0 & 0 & 0 \\
-\frac{\sigma_{33}}{\rho^2} & \frac{1}{\rho} & 0 & 0 & v_3 & -\frac{1}{\rho}\frac{\partial\sigma_{33}}{\partial A_{11}} & -\frac{1}{\rho}\frac{\partial\sigma_{33}}{\partial A_{12}} & -\frac{1}{\rho}\frac{\partial\sigma_{33}}{\partial A_{13}} & -\frac{1}{\rho}\frac{\partial\sigma_{33}}{\partial A_{21}} & -\frac{1}{\rho}\frac{\partial\sigma_{33}}{\partial A_{22}} & -\frac{1}{\rho}\frac{\partial\sigma_{33}}{\partial A_{23}} & -\frac{1}{\rho}\frac{\partial\sigma_{33}}{\partial A_{31}} & -\frac{1}{\rho}\frac{\partial\sigma_{33}}{\partial A_{32}} & -\frac{1}{\rho}\frac{\partial\sigma_{33}}{\partial A_{33}} & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & v_3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & v_3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{11} & A_{12} & A_{13} & 0 & 0 & v_3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{21} & A_{22} & A_{23} & 0 & 0 & 0 & 0 & 0 & v_3 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_3 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_3 & 0 & 0 & 0 & 0 \\
0 & 0 & A_{31} & A_{32} & A_{33} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_3 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_3 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_3 & 0 \\
-\frac{T}{\rho^2} & \frac{T}{\rho(p+p_\infty)} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & v_3
\end{pmatrix} \tag{2.46}$$

Chapter 3

Slow Flow

3.1 Studying numerical smearing with slow flow past a barrier

A checkerboard pattern appears around the corner of the barrier, leading to a crash, using reflective boundary conditions (in velocity) for the barrier. Do we need a staggered grid?

Chapter 4

RGFM

The RGFM does nothing without a temperature fix when applied to the heat conduction test. The linearisation upon which it is based results in a stationary solution when $q_L = q_R$ and $\sigma_L = \sigma_R$ initially. Barton's RGFM is similar. Should q be fixed? Maybe use analytical solution to heat equation at $t = \Delta t$?

Bibliography

[1] Doron Levy. A Third Order Central WENO Scheme for 2D Conservation Laws 1 Introduction 2 Description of the scheme. 33:415–421, 1998.

[2] Olindo Zanotti and Michael Dumbser. Efficient conservative ADER schemes based on WENO reconstruction and space-time predictor in primitive variables. *Computational Astrophysics and Cosmology*, 3(1):1, 2016.