

# CS601: Software Development for Scientific Computing

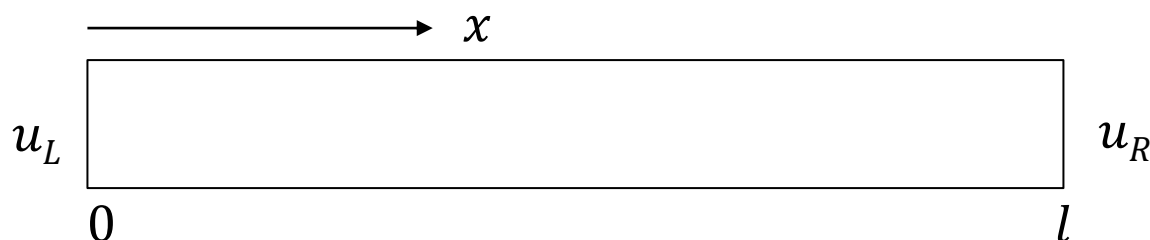
Autumn 2024

Week12: Structured Grids

# Recap

# Application: Heat Equation

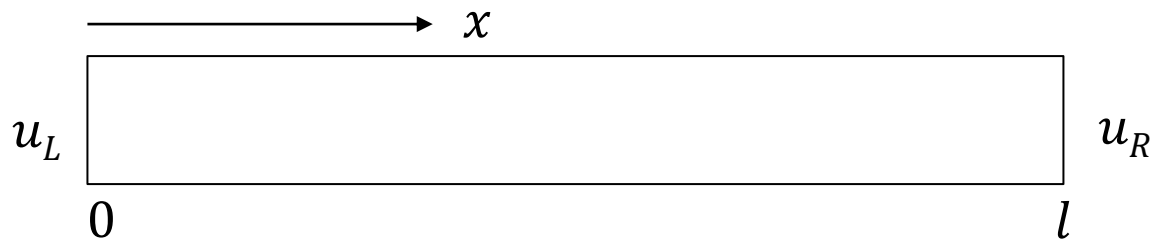
- Example: heat conduction through a rod



- $u = u(x, t)$  is the temperature of the metal bar at distance  $x$  from one end and at time  $t$
- Goal: find  $u$

# Initial and Boundary Conditions

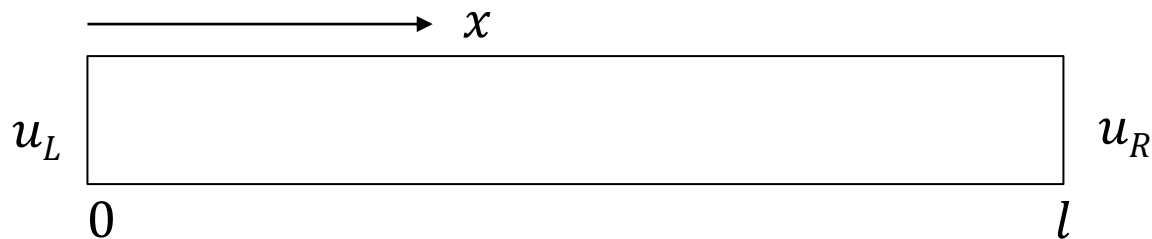
- Example: heat conduction through a rod



- Metal bar has length  $l$  and the ends are held at **constant temperatures**  $u_L$  at the left and  $u_R$  at the right
- Temperature distribution at the **initial time** is known  $f(x)$ , with  $f(0) = u_L$  and  $f(l) = u_R$

# Equations

- Example: heat conduction through a rod



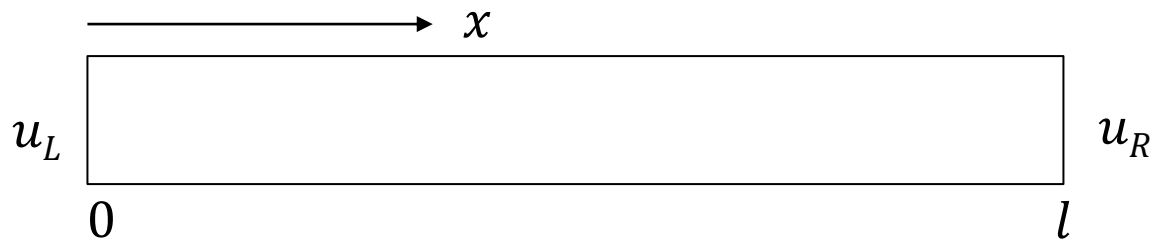
$$\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2} \quad (0 < x < l, t > 0)$$

$\alpha$  is thermal diffusivity

(a constant if the material is homogeneous and isotropic.  
copper =  $1.14 \text{ cm}^2 \text{ s}^{-1}$ , aluminium =  $0.86 \text{ cm}^2 \text{ s}^{-1}$ )

# Equations

- Example: heat conduction through a rod



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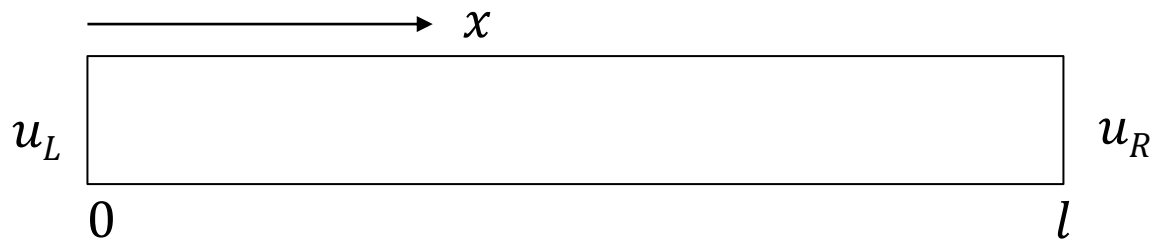
(a constant if the material is homogeneous and isotropic.

copper =  $1.14 \text{ cm}^2 \text{ s}^{-1}$ , aluminium =  $0.86 \text{ cm}^2 \text{ s}^{-1}$ )

- *Exercise: what kind of a PDE is this? (Poisson/Heat/Wave?)*

# Equations

- Example: heat conduction through a rod

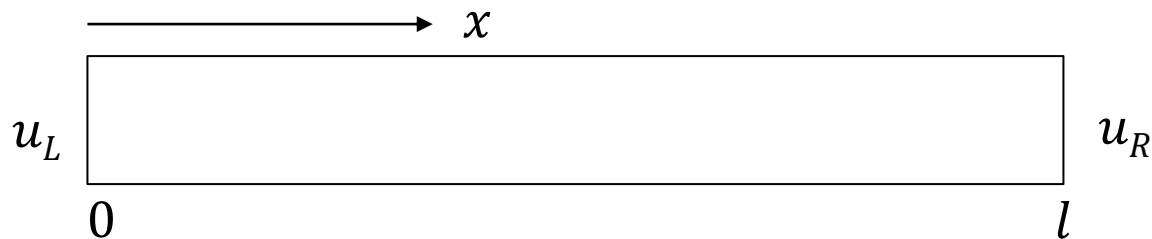


$$\partial_t u = \alpha \Delta u$$

as per the notation mentioned earlier

# Equations

- Example: heat conduction through a rod



$$\partial_t u = \alpha \Delta u$$

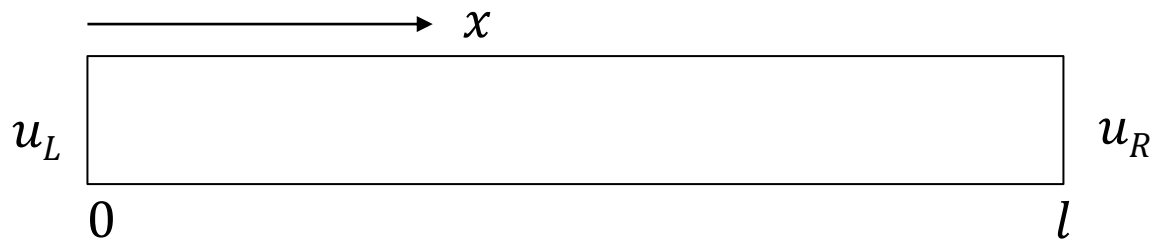
Can also be written as:

$$\partial_t u - \alpha \Delta u = 0$$



# Equations

- Example: heat conduction through a rod



$$\partial_t u - \alpha \Delta u = 0 ,$$

Based on initial and boundary conditions:

$$u(0, t) = u_L ,$$

$$u(l, t) = u_R ,$$

$$u(x, 0) = f(x)$$

# Equations

- Summarizing:

- $\partial_t u - \alpha \Delta u = 0, 0 < x < l, t > 0$

- $u(0, t) = u_L, t > 0$

- $u(l, t) = u_R, t > 0$

- $u(x, 0) = f(x), 0 < x < l$

- Solution:

$$u(x, t) = \sum_{m=1}^{\infty} B_m e^{-m^2 \alpha \pi^2 t / l^2} \sin\left(\frac{m\pi x}{l}\right),$$

$$\text{where, } B_m = 2/l \int_0^l f(s) \sin\left(\frac{m\pi s}{l}\right) ds$$

# Equations

- Summarizing:

1.  $\partial_t u - \alpha \Delta u = 0, 0 < x < l, t > 0$

2.  $u(0, t) = u_L, t > 0$

3.  $u(l, t) = u_R, t > 0$

4. *But we are interested in a numerical solution*

- Solution:

$$u(x, t) = \sum_{m=1}^{\infty} B_m e^{-m^2 \alpha \pi^2 t / l^2} \sin\left(\frac{m \pi x}{l}\right),$$

$$\text{where, } B_m = 2/l \int_0^l f(s) \sin\left(\frac{m \pi s}{l}\right) ds$$

# Approximating Partial Derivatives

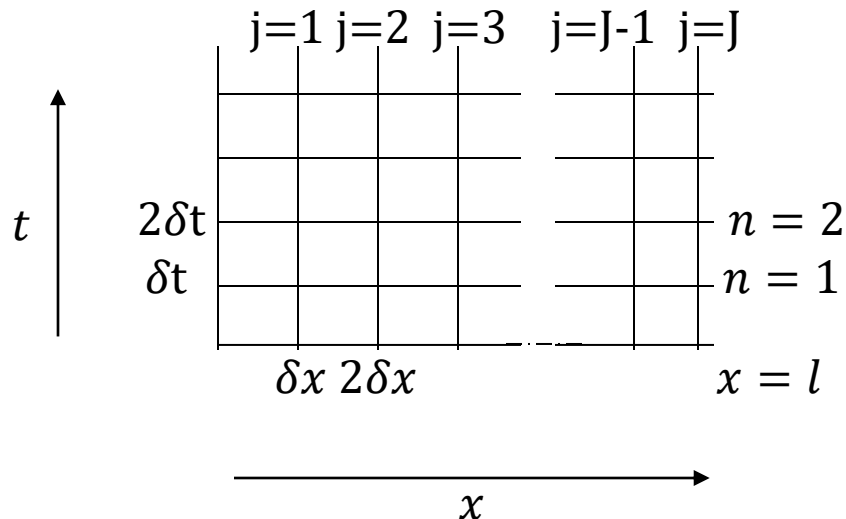
- Suppose  $y = f(x)$ 
  - Forward difference approximation to the first-order derivative of  $f$  w.r.t.  $x$  is:
$$\frac{df}{dx} \approx \frac{(f(x+\delta x) - f(x))}{\delta x}$$
  - Central difference approximation to the first-order derivative of  $f$  w.r.t.  $x$  is:
$$\frac{df}{dx} \approx \frac{(f(x+\delta x) - f(x-\delta x))}{2\delta x}$$
  - Central difference approximation to the second-order derivative of  $f$  w.r.t.  $x$  is:
$$\frac{d^2f}{dx^2} \approx \frac{(f(x+\delta x) - 2f(x) + f(x-\delta x))}{(\delta x)^2}$$

# Approximating Partial Derivatives

- In example heat application  $f = u = u(x, t)$  and
$$\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2}$$
  - First, approximating  $\frac{\partial u}{\partial t}$ :
$$\frac{\partial u}{\partial t} \approx \frac{(u(x, t + \delta t) - u(x, t))}{\delta t},$$
 where  $\delta t$  is a small increment in time
  - Next, approximating  $\frac{\partial^2 u}{\partial x^2}$ :
$$\frac{\partial^2 u}{\partial x^2} \approx \frac{(u(x + \delta x, t) - 2u(x, t) + u(x - \delta x, t))}{(\delta x)^2},$$
 where  $\delta x$  is a small increment in space (along the length of the rod)

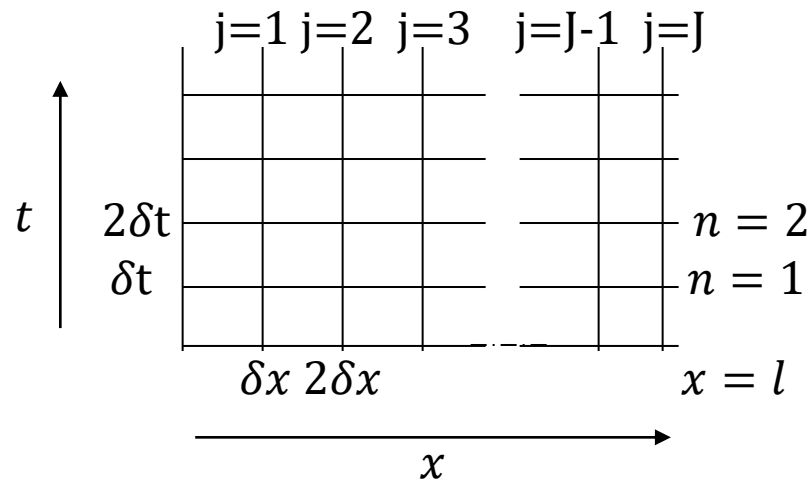
# Approximating Partial Derivatives

- Divide length  $l$  into  $J$  equal divisions:  $\delta x = l/J$  (space step)
- Choose an appropriate  $\delta t$  (time step)



# Approximating Partial Derivatives

- Find sequence of numbers which approximate  $u$  at a sequence of  $(x, t)$  points (i.e. at the intersection of horizontal and vertical lines below)



- Approximate the exact solution  $u(j \times \delta x, n \times \delta t)$  using the approximation for partial derivatives mentioned earlier

# Approximating Partial Derivatives

$$\begin{aligned}\frac{\partial u}{\partial t} &\approx \frac{(u(x, t + \delta t) - u(x, t))}{\delta t} \\ &= \frac{(u_j^{n+1} - u_j^n)}{\delta t}\end{aligned}$$

where  $u_j^{n+1}$  denotes taking  $j$  steps along  $x$  direction and  $n + 1$  steps along  $t$  direction

$$\begin{aligned}\text{Similarly, } \frac{\partial^2 u}{\partial x^2} &\approx \frac{(u(x + \delta x, t) - 2u(x, t) + u(x - \delta x, t))}{(\delta x)^2} \\ &= \frac{(u_{j+1}^n - 2u_j^n + u_{j-1}^n)}{(\delta x)^2}\end{aligned}$$



# Approximating Partial Derivatives

Plugging into  $\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2}$  :

$$\frac{(u_j^{n+1} - u_j^n)}{\delta t} = \alpha \frac{(u_{j+1}^n - 2 u_j^n + u_{j-1}^n)}{(\delta x)^2}$$

This is also called as difference equation because you are computing difference between successive values of a function involving discrete variables.

# Approximating Partial Derivatives

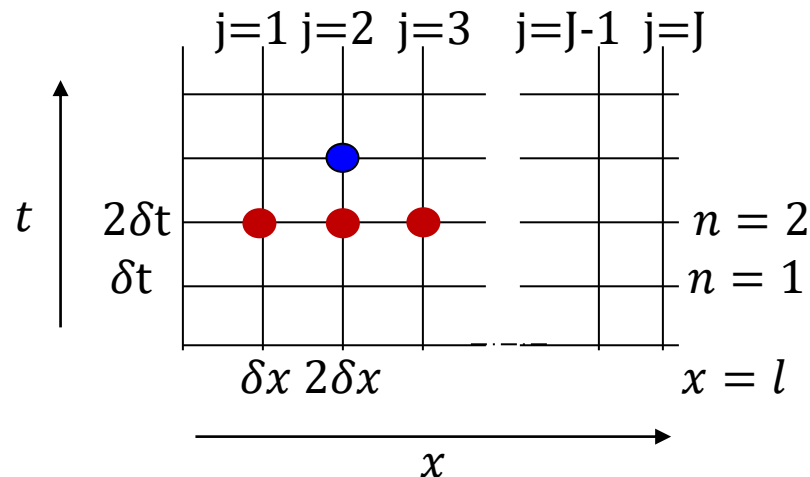
Simplifying:

$$\begin{aligned}u_j^{n+1} &= u_j^n + r(u_{j+1}^n - 2u_j^n + u_{j-1}^n) \\&= ru_{j-1}^n + (1 - 2r)u_j^n + ru_{j+1}^n, \\&\text{where } r = \alpha \frac{\delta t}{(\delta x)^2}\end{aligned}$$

# Approximating Partial Derivatives

visualizing,

$$u_j^{n+1} = ru_{j-1}^n + (1 - 2r)u_j^n + ru_{j+1}^n$$



*To compute the value of function at blue dot, you need 3 values indicated by the red dots – 3-point stencil*

# Approximating Partial Derivatives

- Initial and boundary conditions tell us that:

$$u(0, t) = u_L ,$$

$$u(l, t) = u_R ,$$

$$u(x, 0) = f(x)$$

- $u_0^0, u_1^0, u_2^0, \dots, u_j^0$  are known (at time  $t=0$ , the temperature at all points along the distance is known as indicated by  $f(x) = f_j$ ).
- $u_0^1$  is  $u_L$ ,  $u_j^1$  is  $u_R$
- Now compute points on the grid from left-to-right:

# Approximating Partial Derivatives

- Now compute points on the grid from left-to-right:

$$u_1^1 = u_1^0 + r(u_0^0 - 2u_1^0 + u_2^0)$$

$$u_2^1 = u_2^0 + r(u_1^0 - 2u_2^0 + u_3^0)$$

.

.

$$u_{j-1}^1 = u_{j-1}^0 + r(u_{j-2}^0 - 2u_{j-1}^0 + u_j^0)$$

- This constitutes the computation done in the first time step.
- Now do the second time step computation...and so on..

# Explicit Difference Method: Stability

- Given:  $l = 1$ ,  
 $u(0, t) = u_L = 0$ ,  
 $u(l, t) = u_R = 0$ ,  
 $u(x, 0) = f(x) = x(l - x)$   
 $\alpha = 1$ ,
- Choose:  $\delta x = 0.25, \delta t = 0.075$
- Solve.

# Explicit Difference Method: Stability

- Initialize  $u_j^0$  values from initial and boundary conditions i.e. *get time-step 0 values*

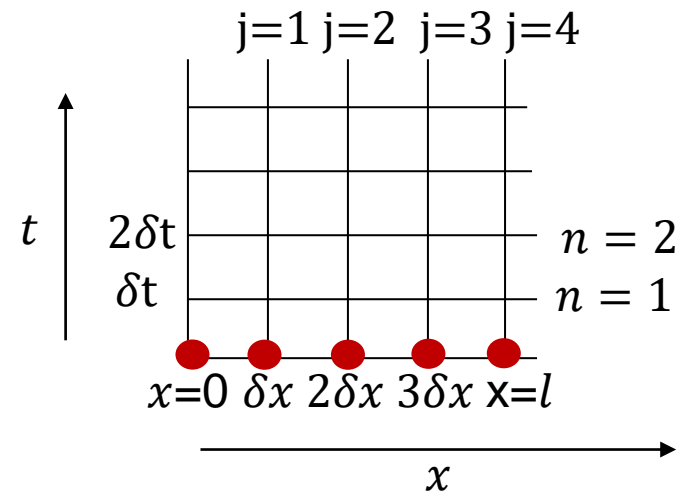
$$u_0^0 = 0$$

$$u_1^0 = f(\delta x) = \delta x(l - \delta x) = .1875$$

$$u_2^0 = f(2\delta x) = 2\delta x(l - 2\delta x) = .25$$

$$u_3^0 = f(3\delta x) = 3\delta x(l - 3\delta x) = .1875$$

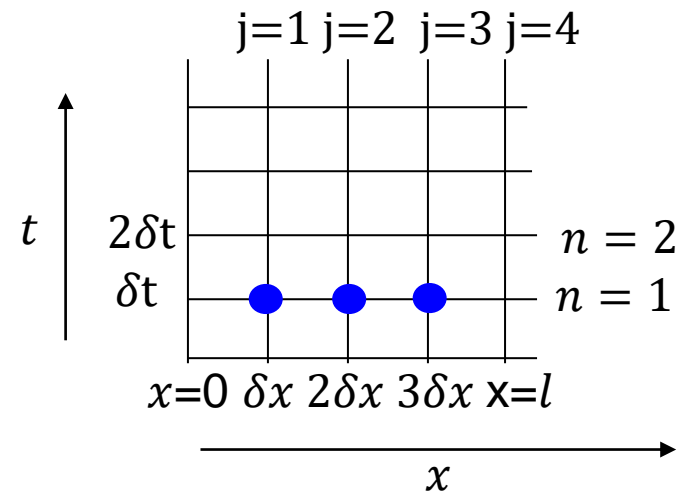
$$u_4^0 = 0$$



# Explicit Difference Method: Stability

- Compute time-step 1 values

$$u_j^{n+1} = ru_{j-1}^n + (1 - 2r)u_j^n + ru_{j+1}^n$$



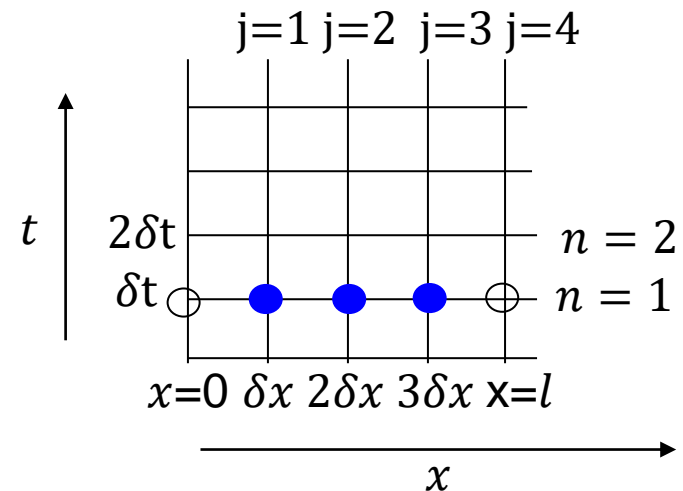


# Explicit Difference Method: Stability

- Compute time-step 1 values

$$u_j^{n+1} = ru_{j-1}^n + (1 - 2r)u_j^n + ru_{j+1}^n$$

What about values of  $u(x, t)$  at  $\bigcirc$  ?



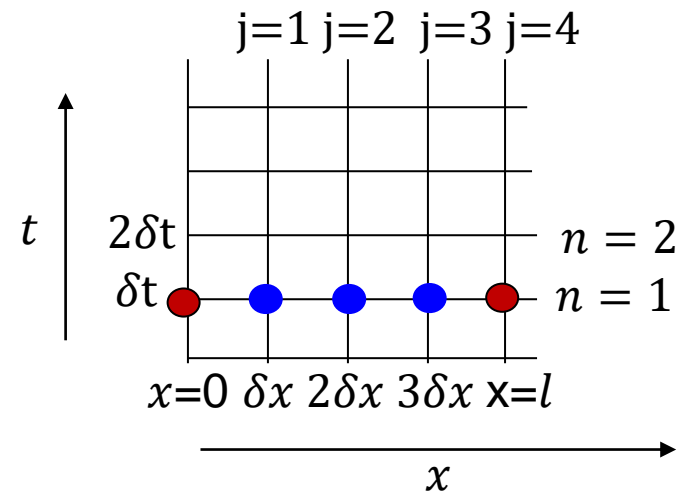
# Explicit Difference Method: Stability

- Compute time-step 1 values

$$u_j^{n+1} = ru_{j-1}^n + (1 - 2r)u_j^n + ru_{j+1}^n$$

What about values of  $u(x, t)$  at  $\circ$  ?

*Get it from boundary conditions*



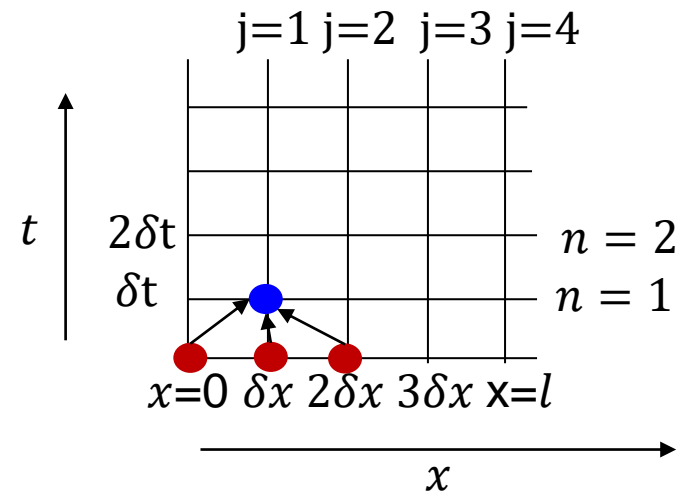
# Explicit Difference Method: Stability

- Compute time-step 1 values

$$u_j^{n+1} = ru_{j-1}^n + (1 - 2r)u_j^n + ru_{j+1}^n$$

$$r = \alpha \delta t / (\delta x)^2 = 1.2$$

$$u_1^1 = u_1^0 + r(u_0^0 - 2u_1^0 + u_2^0) = 0.03678$$



# Explicit Difference Method: Stability

- Compute time-step 1 values

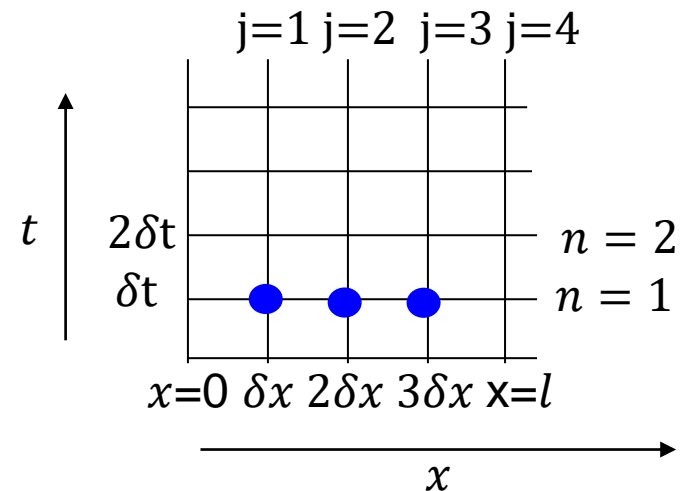
$$u_j^{n+1} = ru_{j-1}^n + (1 - 2r)u_j^n + ru_{j+1}^n$$

$$r = \alpha \delta t / (\delta x)^2 = 1.2$$

$$u_1^1 = u_1^0 + r(u_0^0 - 2u_1^0 + u_2^0) = 0.03678$$

$$u_2^1 = u_2^0 + r(u_1^0 - 2u_2^0 + u_3^0) = 0.1$$

$$u_3^1 = u_3^0 + r(u_2^0 - 2u_3^0 + u_4^0) = 0.03678$$



# Explicit Difference Method: Stability

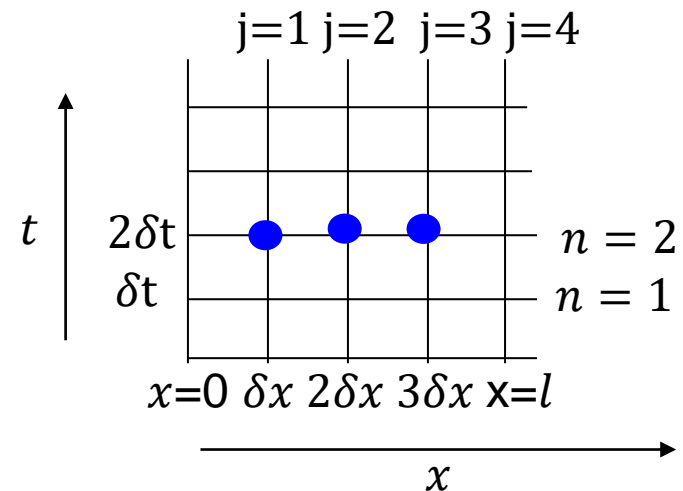
- Compute time-step 2 values

$$u_j^{n+1} = ru_{j-1}^n + (1 - 2r)u_j^n + ru_{j+1}^n$$

$$u_1^2 = u_1^1 + r(u_0^1 - 2u_1^1 + u_2^1) = 0.06851$$

$$u_2^2 = u_2^1 + r(u_1^1 - 2u_2^1 + u_3^1) = \mathbf{-0.05173}$$

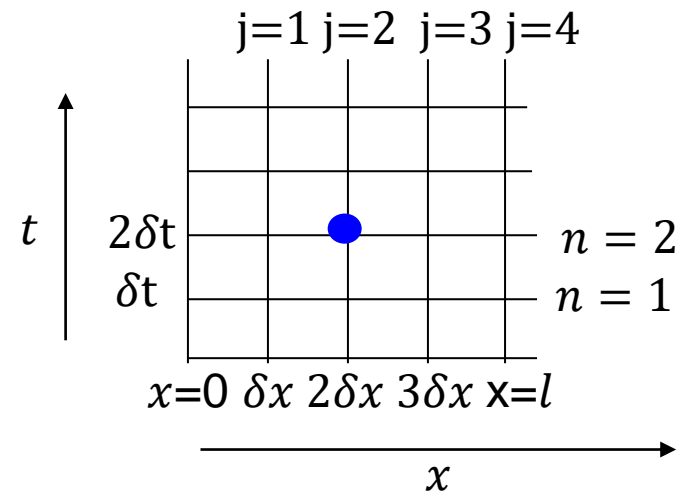
$$u_3^2 = u_3^1 + r(u_2^1 - 2u_3^1 + u_4^1) = 0.06851$$



# Explicit Difference Method: Stability

- Temperature at  $2\delta x$  after  $2\delta t$  time units went into negative! (when the boundaries were held constant at 0)
  - Example of *instability*

$$u_2^2 = u_2^1 + r(u_1^1 - 2u_2^1 + u_3^1) = \mathbf{-0.05173}$$



*The solution is stable (for heat diffusion problem) only if the approximations for  $u(x, t)$  do not get bigger in magnitude with time*

# Explicit Difference Method: Stability

- The solution for heat diffusion problem is stable only if:

$$r \leq \frac{1}{2}$$

Therefore, choose your time step in such a way that:

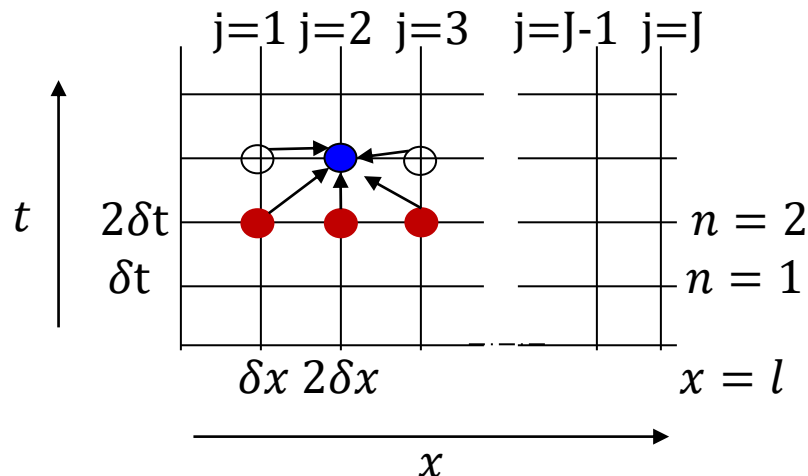
$$\delta t \leq \frac{\delta x^2}{2\alpha}$$

*But this is a severe limitation!*

# Implicit Method: Stability

- Overcoming instability:

$$u_j^{n+1} = u_j^n + \frac{1}{2} r (u_{j-1}^n - 2u_j^n + u_{j+1}^n + u_{j-1}^{n+1} - 2u_j^{n+1} + u_{j+1}^{n+1})$$



*To compute the value of function at blue dot, you need 6 values indicated by the red dots (known) and 3 additional ones (unknown) above*



# Implicit Method: Stability

- Overcoming instability:

$$u_j^{n+1} = u_j^n + 1/2 r( u_{j-1}^n - 2u_j^n + u_{j+1}^n + u_{j-1}^{n+1} - 2u_j^{n+1} + u_{j+1}^{n+1} )$$

- Extra work involved to determine the values of unknowns in a time step
  - Solve a system of simultaneous equations. Is it worth it?

# Suggested Reading

- *J.W. Thomas. Numerical Partial Differential Equations: Finite Difference Methods*
- *Parabolic PDEs:*  
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