

第六周作业

1. 证明一维谐振子的升降算符满足

- (1) $\hat{a}_-^\dagger = \hat{a}_+$,
- (2) $[\hat{a}_-, \hat{a}_+] = 1$,
- (3) $[\hat{a}_+, \hat{a}_+ \hat{a}_-] = -\hat{a}_+$,
 $[\hat{a}_-, \hat{a}_+ \hat{a}_-] = \hat{a}_-$.

证(1)

$$\begin{aligned}\hat{a}_+ &= \frac{1}{\sqrt{2\hbar m\omega}}(-i\hat{p} + m\omega\hat{x}), \\ \hat{a}_- &= \frac{1}{\sqrt{2\hbar m\omega}}(i\hat{p} + m\omega\hat{x}), \\ \hat{a}_-^\dagger &= \frac{1}{\sqrt{2\hbar m\omega}}(-i\hat{p}^\dagger + m\omega\hat{x}^\dagger) \\ &= \frac{1}{\sqrt{2\hbar m\omega}}(-i\hat{p} + m\omega\hat{x}) = \hat{a}_+.\end{aligned}$$

$$\begin{aligned}(2) \quad [\hat{a}_-, \hat{a}_+] &= \left[\frac{1}{\sqrt{2\hbar m\omega}}(i\hat{p} + m\omega\hat{x}), \frac{1}{\sqrt{2\hbar m\omega}}(-i\hat{p} + m\omega\hat{x}) \right] \\ &= \frac{1}{2\hbar m\omega} [i\hat{p} + m\omega\hat{x}, -i\hat{p} + m\omega\hat{x}] \\ &= \frac{1}{2\hbar m\omega} (im\omega[\hat{p}, \hat{x}] - im\omega[\hat{x}, \hat{p}]) \\ &= \frac{i}{2\hbar} (-i\hbar - i\hbar) = 1.\end{aligned}$$

$$\begin{aligned}(3) \quad [\hat{a}_+, \hat{a}_+ \hat{a}_-] &= [\hat{a}_+, \hat{a}_+] \hat{a}_- + \hat{a}_+ [\hat{a}_+, \hat{a}_-] \\ &= \hat{a}_+ [\hat{a}_+, \hat{a}_-], \\ [\hat{a}_+, \hat{a}_-] &= -1, \\ [\hat{a}_+, \hat{a}_+ \hat{a}_-] &= -\hat{a}_+, \\ [\hat{a}_-, \hat{a}_+ \hat{a}_-] &= [\hat{a}_-, \hat{a}_+] \hat{a}_- + \hat{a}_+ [\hat{a}_-, \hat{a}_-] \\ &= [\hat{a}_-, \hat{a}_+] \hat{a}_-, \\ [\hat{a}_-, \hat{a}_+] &= 1, \\ [\hat{a}_-, \hat{a}_+ \hat{a}_-] &= \hat{a}_-.\end{aligned}$$

2. 对于谐振子的能量本征态 ψ_n ,

- (1) 计算 \hat{x} 、 \hat{p} 的平均值,
- (2) 计算 \hat{x}^2 、 \hat{p}^2 的平均值,
- (3) 计算 $\Delta x = (\langle \hat{x}^2 \rangle - \langle \hat{x} \rangle^2)^{\frac{1}{2}}$ 、 $\Delta p = (\langle \hat{p}^2 \rangle - \langle \hat{p} \rangle^2)^{\frac{1}{2}}$.

解(1)

$$\hat{x} = \sqrt{\frac{\hbar}{2m\omega}}(\hat{a}_+ + \hat{a}_-),$$

$$\begin{aligned}
\hat{p} &= i \sqrt{\frac{m\omega\hbar}{2}} (\hat{a}_+ - \hat{a}_-), \\
\hat{a}_- \psi_n &= \sqrt{n} \psi_{n-1}, \\
\hat{a}_+ \psi_n &= \sqrt{n+1} \psi_{n+1}, \\
\hat{x} \psi_n &= \sqrt{\frac{\hbar}{2m\omega}} (\hat{a}_+ \psi_n + \hat{a}_- \psi_n) \\
&= \sqrt{\frac{\hbar}{2m\omega}} (\sqrt{n+1} \psi_{n+1} + \sqrt{n} \psi_{n-1}), \\
\hat{p} \psi_n &= i \sqrt{\frac{m\omega\hbar}{2}} (\hat{a}_+ \psi_n - \hat{a}_- \psi_n) \\
&= i \sqrt{\frac{m\omega\hbar}{2}} (\sqrt{n+1} \psi_{n+1} - \sqrt{n} \psi_{n-1}),
\end{aligned}$$

由本征态的正交归一化条件 $\langle \psi_n, \psi_{n'} \rangle = \delta_{nn'}$,

$$\langle \hat{x} \rangle = \langle \psi_n, \hat{x} \psi_n \rangle = 0,$$

$$\langle \hat{p} \rangle = \langle \psi_n, \hat{p} \psi_n \rangle = 0,$$

这个结论也可以利用波函数 $\psi_n(x)$ 的宇称性而得出。

$$\begin{aligned}
(2) \quad \hat{x}^2 &= \frac{\hbar}{2m\omega} (\hat{a}_+ + \hat{a}_-)^2 \\
&= \frac{\hbar}{2m\omega} (\hat{a}_+^2 + \hat{a}_-^2 + \hat{a}_+ \hat{a}_- + \hat{a}_- \hat{a}_+) \\
&= \frac{\hbar}{2m\omega} (\hat{a}_+^2 + \hat{a}_-^2 + 2\hat{N} + 1), \\
\hat{p}^2 &= -\frac{m\omega\hbar}{2} (\hat{a}_+ - \hat{a}_-)^2 \\
&= -\frac{m\omega\hbar}{2} (\hat{a}_+^2 + \hat{a}_-^2 - \hat{a}_+ \hat{a}_- - \hat{a}_- \hat{a}_+) \\
&= \frac{m\omega\hbar}{2} (2\hat{N} + 1 - \hat{a}_+^2 - \hat{a}_-^2),
\end{aligned}$$

由于

$$\hat{a}_+^2 \psi_n = \sqrt{(n+1)(n+2)} \psi_{n+2},$$

$$\hat{a}_-^2 \psi_n = \sqrt{n(n-1)} \psi_{n-2},$$

根据正交条件

$$\langle \psi_n, \hat{a}_+^2 \psi_n \rangle = 0, \quad \langle \psi_n, \hat{a}_-^2 \psi_n \rangle = 0,$$

因此

$$\langle \hat{x}^2 \rangle = \langle \psi_n, \hat{x}^2 \psi_n \rangle = \frac{\hbar}{2m\omega} \langle \psi_n, (2\hat{N} + 1) \psi_n \rangle = \frac{\hbar}{m\omega} \left(n + \frac{1}{2} \right),$$

$$\langle \hat{p}^2 \rangle = \langle \psi_n, \hat{p}^2 \psi_n \rangle = \frac{m\omega\hbar}{2} \langle \psi_n, (2\hat{N} + 1) \psi_n \rangle = m\omega\hbar \left(n + \frac{1}{2} \right).$$

$$(3) \quad \Delta x = (\langle \hat{x}^2 \rangle - \langle \hat{x} \rangle^2)^{\frac{1}{2}} = \sqrt{\frac{\hbar}{m\omega} \left(n + \frac{1}{2} \right)},$$

$$\Delta p = (\langle \hat{p}^2 \rangle - \langle \hat{p} \rangle^2)^{\frac{1}{2}} = \sqrt{m\omega\hbar \left(n + \frac{1}{2} \right)},$$

$$\Delta x \cdot \Delta p = \hbar \left(n + \frac{1}{2} \right).$$

对于基态, $n = 0, \Delta x \cdot \Delta p = \hbar/2$, 刚好是测不准关系所规定的下限。