

# An Introduction to Discontinuous Galerkin Methods

## Module 3A: To Higher-Orders - Solution Approximation

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# Weak Form- Revisited: Approximation Space

$$\int_{\Omega} \frac{\partial q}{\partial t} \phi_j dx + \int_{\Omega} \frac{\partial f(q)}{\partial x} \phi_j dx = 0 \quad (1)$$

- ▶ Let's revisit the weak form in greater detail, we glossed over some important details.
- ▶ To get a flexible method, we'd like to be able to generate an arbitrary order solution approximation. How do we do this in code?
- ▶ We glossed over the construction of the bases, we can think in terms of nodal or modal

$$\tilde{q}(x) = \sum_{i=0}^M a_i \psi_i(x) \quad (2)$$

## $L^2$ Projection: Minimize norm

- ▶ Where does the weak form come from? Approximate sol'n doesn't satisfy the PDE exactly, left with some residual

$$\frac{\partial \tilde{q}}{\partial t} dx + \frac{\partial f(\tilde{q})}{\partial x} dx = R(x) \quad (3)$$

- ▶ A projection finds the "closest" fit on a vector space compared to the original. An orthogonal projection does just this.
- ▶ Minimizes the  $L^2$  norm wrt the space projected upon

# The Galerkin Condition

- ▶ An orthogonal projection means that the difference between the original and the projection should be orthogonal to the space projected upon, so for each basis  $\phi \in V_h$

$$\int (f - f_h) \phi \, dx = 0 = \int R(x) \phi \, dx \quad (4)$$

- ▶ The projected space doesn't have to be the same space as the solution approximation, but if we do choose them to be the same we have the Bubnov-Galerkin approach
- ▶ We could choose the two spaces to be different, this is a Petrov-Galerkin approach. As an example, a weight space of delta functions yields a collocation method

# Arbitrary Order Basis: Monomials?

- ▶ Let's continue with the choice of polynomials for our basis functions.
- ▶ The easy and obvious choice is the monomials  $\psi_i(x) = x^i$  so the arbitrary order solution approximation is

$$\tilde{q}(x) = \sum_{i=0}^M a_i \psi_i(x) = \sum_{i=0}^M a_i x^i \quad (5)$$

- ▶ Can you think of why this a poor choice for a vector basis? Remember a basis needs to be complete and linearly independent. Let's see two separate ways why this is a poor choice.

# Linear Dependence

- ▶ Normalize the basis wrt to the  $L^2$  norm  $|g| = \sqrt{\int g^2}$  so  $\bar{\psi}_i = \sqrt{2i+1}x^i$
- ▶ Consider the inner product of two adjacent basis functions

$$\int \bar{\psi}_{i-1} \bar{\psi}_i dx = \sqrt{1 - \frac{1}{4i^2}} \quad (6)$$

- ▶ Linearly independent vectors are orthogonal, that is their inner product is zero. But for our monomials, even moderate values of  $i$  give an inner product of close to 1
- ▶ Very nearly dependent basis vectors mean small perturbations can lead to trouble finding the correct solution

# The ill-conditioned Hilbert matrix

- ▶ The alternate way to consider this is to look at the mass matrix for a monomial basis
- ▶ 1
- ▶  $x$
- ▶  $x^2$
- ▶  $x^3$
- ▶ When we integrate from 0 to 1 it is easy to see the pattern of values, this is a Hilbert matrix
- ▶ It is ill-conditioned, let's see why this is a problem in Matlab



## [Recall] Lagrange Interpolation

- ▶ Instead of a monomial basis, what if we were to use a Lagrange basis (commonly used for interpolation)
- ▶ Recall, the Lagrange basis is zero at all interp nodes except the basis node, which it is 1 at
- ▶ It has the form

$$L_q(x) = \prod_{p=1, p \neq q}^{N_p} \frac{x - x_p}{x_q - x_p} \quad (7)$$

- ▶ We can compactly represent this as  

```
nn=elim(nd(1:N)',nd(1:N)',[1 3 2]);  
Lag= @(x,nv)  
prod(bsxfun(@divide,bsxfun(@minus,x,nn(nv,:,:)),...  
bsxfun(@minus,nd(nv),nn(nv,:,:))),3);
```

# Lagrange Spatial Approximation

- ▶ We can now represent our arbitrary order sol'n approximation

$$\tilde{q}(x) = \sum_{i=0}^M a_i \psi_i(x) = \sum_{i=0}^M q(x_i) L_i(x) \quad (8)$$

- ▶ Where we must choose a collection of points to use as our interpolation nodes

# Interpolation Points: Equispaced?

- ▶ It would seem logical that equispaced points are a good choice. They evenly cover the domain and are simple to calculate positions for.
- ▶ Indeed these points work well for low  $Np$ , so one would expect the accuracy to increase as  $Np$  increases
- ▶ Let's see how it behaves at large  $Np$

# Runge's Phenomenon

- ▶ The divergence of the interpolation is a result of Runge's phenomenon
- ▶ Why this fails is explainable in terms of two relations: the Bernstein and Markov Inequalities. For a polynomial  $P^N \in [-1, 1]$

$$|P'(x)| \leq \frac{N}{\sqrt{1-x^2}} \|P\|_{\infty} \quad (9)$$

- ▶ The singularity in the Bernstein inequality can be bounded by the Markov inequality

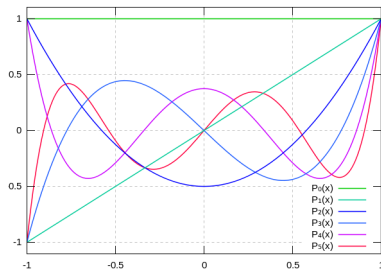
$$\|P'\|_{\infty} \leq N^2 \|P\|_{\infty} \quad (10)$$

- ▶ With the overall being

$$|P'(x)| \leq \min \left( \frac{N}{\sqrt{1-x^2}}, N^2 \right) \|P\|_{\infty} \quad (11)$$

# Interpolation Points: Legendre Roots

- ▶ The overall inequality implies we'd like quadratic spacing near the element boundaries.
- ▶ Though it may seem arbitrary, the roots of the Legendre polynomials have a favorable spacing that matches our needs
- ▶ They avoid Runge's phenomenon, and have several other very nice features we shall see in the next module



<http://en.wikipedia.org/wiki/File:Legendrepolynomials6.svg>